
HABILITATION à **D**IRIGER des **R**ECHERCHES

Mardi 08 Juillet 2003
Amphi Blandin , 15h

PHILIPPE BOURGES

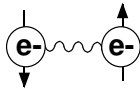
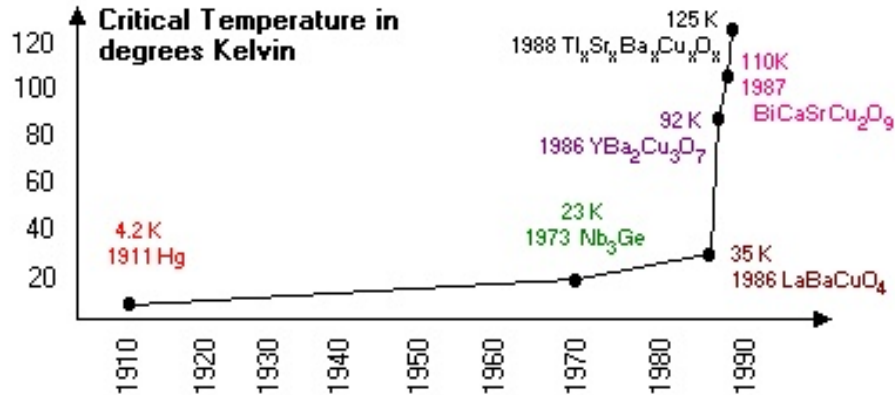
**Dynamique de spins dans les oxydes de cuivre
supraconducteurs à haute température critique**

Jury constitué de : **H. Alloul**
C. Berthier
B. Hennion
H. Raffy
L.P. Regnault
C. Varma
C. Vettier

Supraconductivité à haute température critique

1987 Découverte supraconductivité à haute température critique

$$T_C^{max} \sim 135 \text{ K}$$

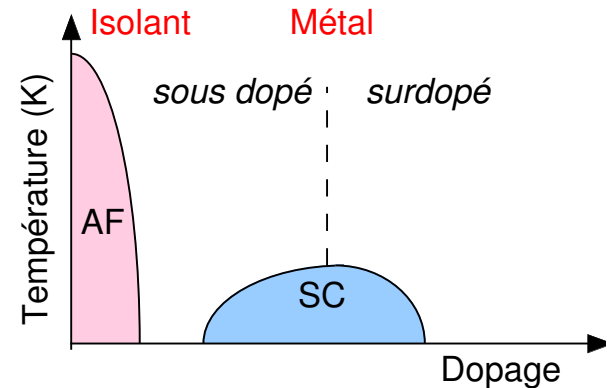
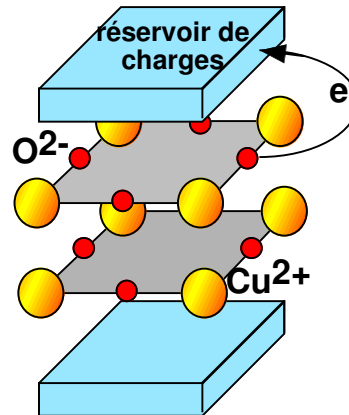


Mécanisme d'appariement ?

- Gap supraconducteur : $2\Delta_{max} \sim 9k_B T_C$
(Théorie BCS : $2\Delta = 3.5k_B T_C$)
- Supraconductivité non-conventionnelle : symétrie d
- Proximité d'une phase magnétique ordonné
(fermions lourds, Sr₂RuO₄, ferromagnétiques UGe₂,...)

⇒ Mécanisme d'appariement magnétique

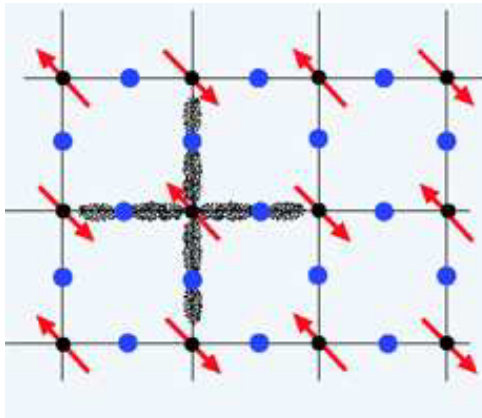
Oxydes de cuivre supraconducteurs : diagramme de phase



1. Dopage en trous des plans CuO_2 par transfert de charges
2. Rôle des corrélations antiferromagnétiques
3. Nature de l'état sous-dopé : Pseudo-gap et ordre caché ?

Cuprates : Système à électrons fortement corrélés

- cuprates : structure bidimensionnelle



$$H = U \sum_i n_{i\uparrow} n_{i\downarrow} - t \sum_{i,j,\sigma} c_{i\sigma}^\dagger c_{j\sigma}$$

- plan CuO_2 dopage nul : isolant Mott-Hubbard
Etat de Néel, $\text{Cu}^{2+} \text{ S} = \frac{1}{2}$,
couplage antiferromagnétique Cu-O-Cu : $J = 120 \text{ meV}$
- dopage en trous : singulet Zhang-Rice
- Autres interactions électroniques
- Couplage électron-phonon

Plan

1. Mode magnétique résonant (phase SC)
2. Corrélations magnétiques de l'état normal
3. Poids spectral \Rightarrow mécanisme de la supraconductivité

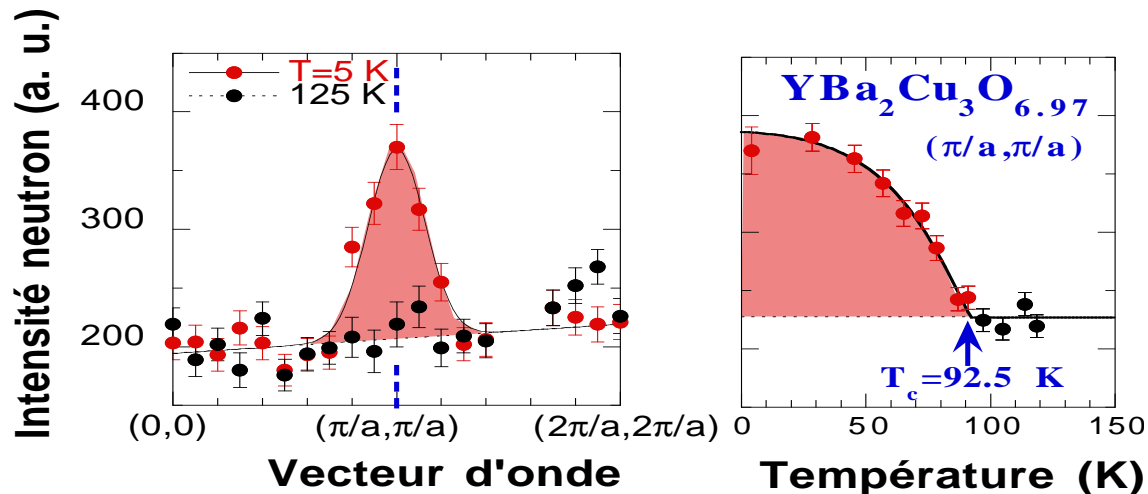
Diffusion inélastique de neutrons \Rightarrow Fonction de diffusion:

$$S^{\alpha\beta}(Q, \omega) = \frac{1}{2\pi\hbar} \int_{-\infty}^{+\infty} dt \exp(-i\omega t) \langle S_Q^\alpha S_{-Q}^\beta(t) \rangle$$

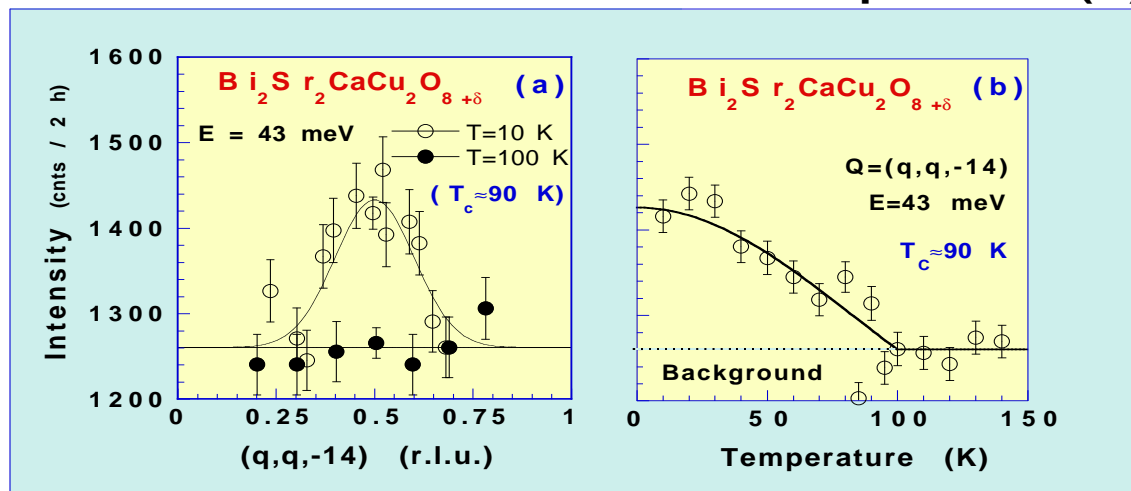
Susceptibilité magnétique:

$$S^{\alpha\beta}(Q, \omega) = \frac{1}{(g\mu_B)^2} \frac{1}{\pi} \{1 + n(\omega)\} \text{Im} \chi^{\alpha\beta}(Q, \omega)$$

Resonance peak in the high- T_C superconductors

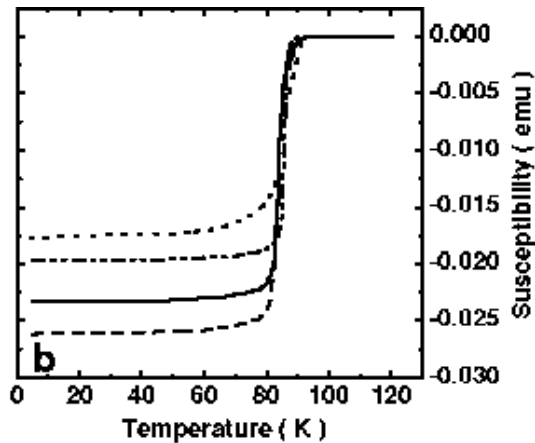
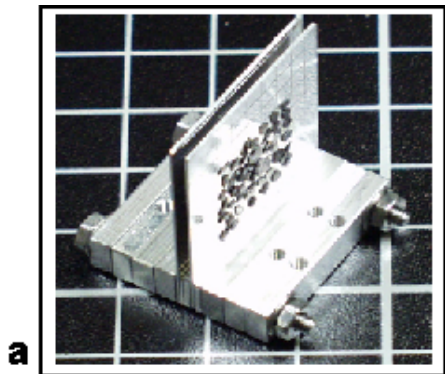


PRB, 53 876
(1996).

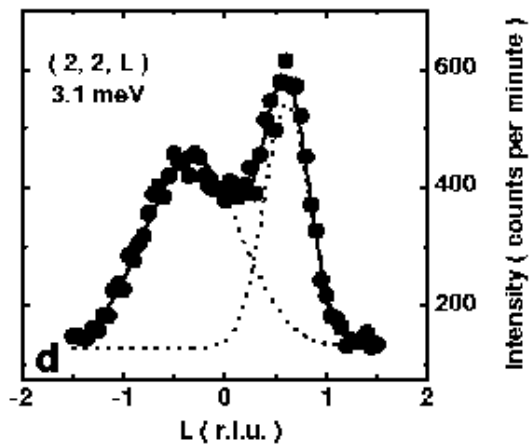
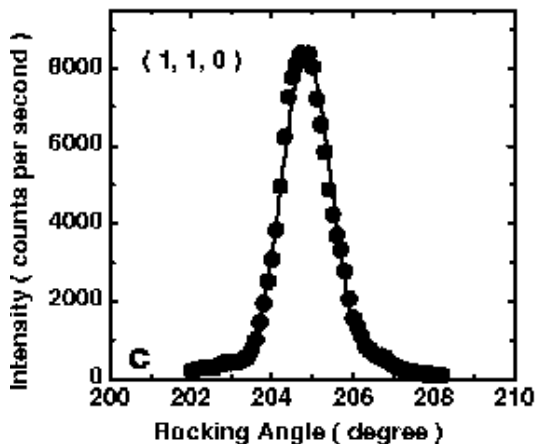


Nature, 398
588 (1999).

Resonance peak in $Tl_2Ba_2CuO_{6+\delta}$

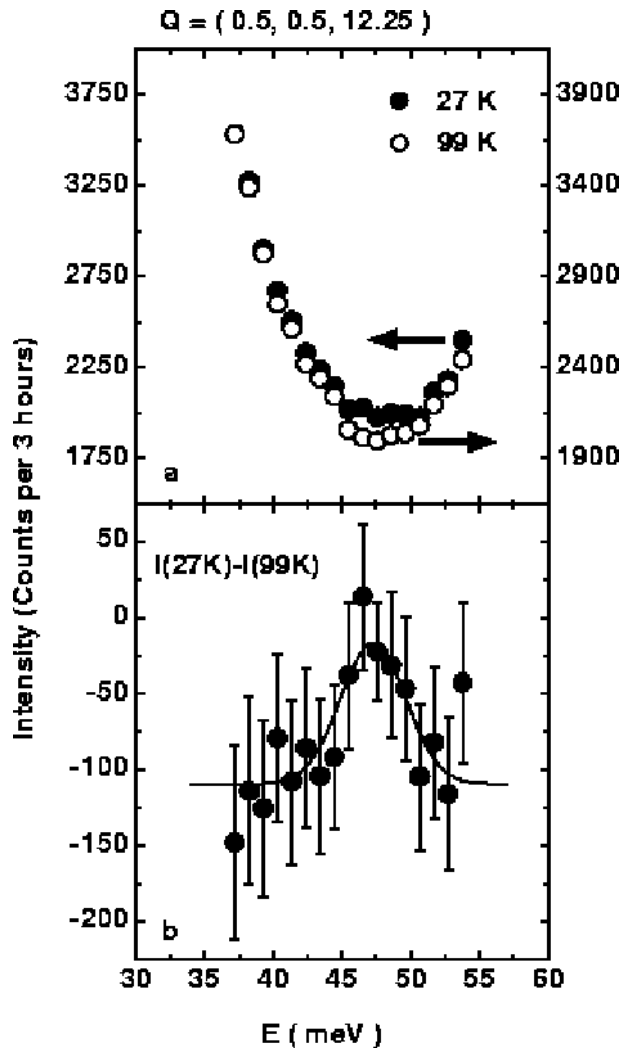


~ Samples 0.5-2 mg
 $T_C \simeq 90$ K
N.S. Berzigiarova,
N.N.Kolesnikov
(Chernogolovka,
Russia)



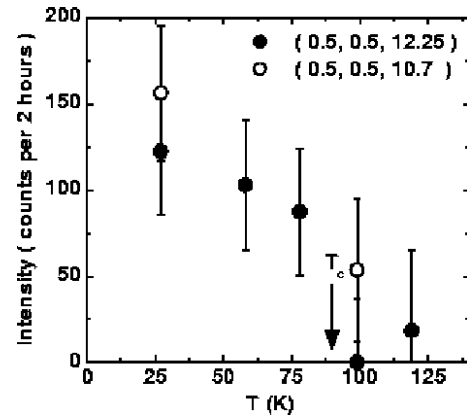
~ 300 samples
co-aligned (H.F. He)
 \Rightarrow
 $m \sim 0.7$ g ~ 0.1 cm³

Resonance peak in $Tl_2Ba_2CuO_{6+\delta}$



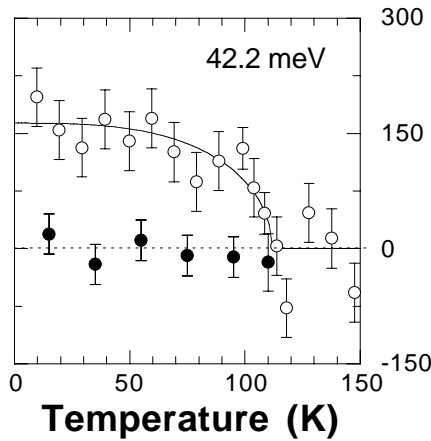
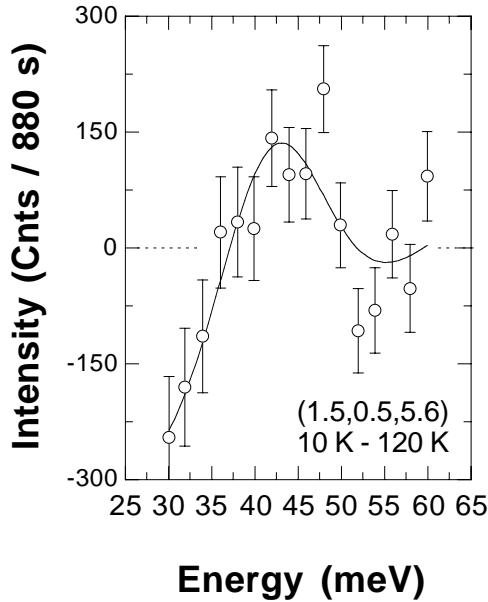
Energy scan
2T (LLB-Saclay)

Science [295](#) 1045 (2002)



□ same spectral weight as $YBCO_7$

resonance peak in Bi2223



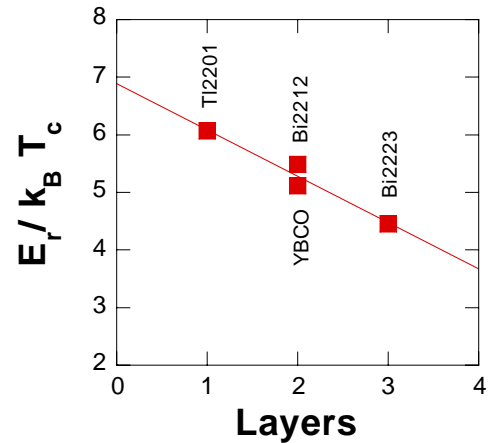
S. Bayrakci *et al.*, unpublished (2003).

3 layers

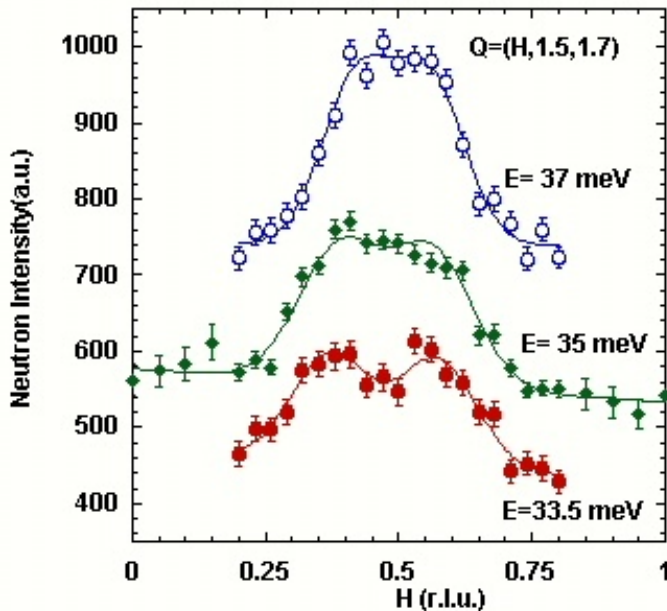
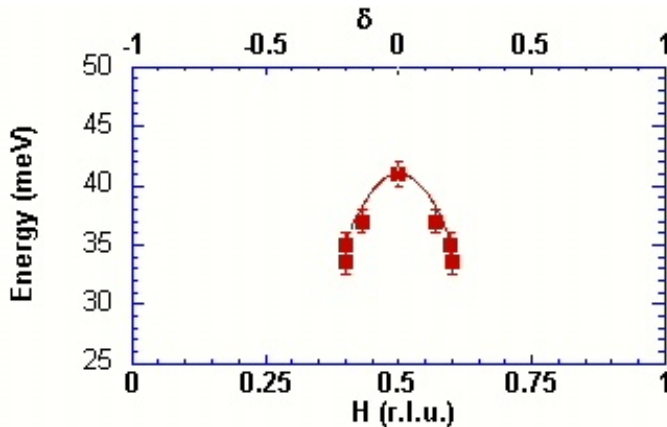
40 co-aligned single crystals

$T_C = 110$ K

$E_r = 42.2$ meV



Resonance peak dispersion



$\text{YBa}_2\text{Cu}_3\text{O}_{6.85}$ ($T_C = 89$ K)

For $E \leq E_r = 41$ meV

Energy dependent
Incommensurate peaks

at $Q_\delta = (\pi/a \pm \delta, \pi/a)$
 $\equiv (\pi/a, \pi/a \pm \delta)$

Only below T_C

\neq stripes

science, [288 1234](#) (2000).

Resonance peak is a generic feature

⇒ Magnetic collective mode in Optimally doped cuprates:

$$T_c^{\max} \geq 90 \text{ K}$$

	3-layers	2-layers		1-layer
Cuprates	Bi2223	YBCO	Bi2212	Tl-2201
Resonance energy (meV)	42	41	43	47
$E_r/k_B T_c$	4.4	5.1	5.4	6

□ $E_r < 2\Delta_{\max} \sim 70 \text{ meV} \sim 9k_B T_c$

1. Particle-hole scenario: "Spin-exciton" model

Onufrieva and Pfeuty PRB [65](#) 054515 (2002)

Abanov and Chubukov, Eschrig and Norman, PRL and PRB (2000-2002).

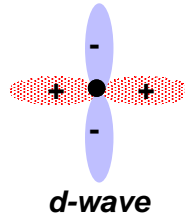
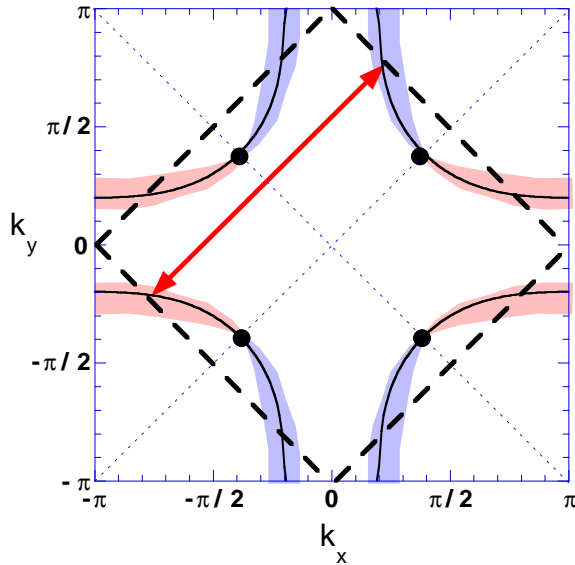
2. SO(5) (S.C. Zhang and Demler)

3. Nearly AF Liquid (D. Morr and D. Pines)

4. Stripes : Spin-waves emerging from incommensurate peaks

Batista, Ortiz, and Balatsky, PRB [64](#) 172508 (2001).

Particle-hole scenario



- Fermi surface
- BCS spin susceptibility (Superconducting state)
- SC order parameter: Δ_k
 \Rightarrow **d-wave** symmetry:
 $\Delta_k \Delta_{k+Q_{AF}} < 0$
- $E_k = \sqrt{\Delta_k^2 + \epsilon_k^2}$

$$\chi^0(q, \omega) \propto \lim_{\epsilon \rightarrow 0} \sum_k \left[1 - \frac{\Delta_k \Delta_{q+k} + \epsilon_{q+k} \epsilon_k}{E_{q+k} E_k} \right] \frac{1 - f_{q+k} - f_k}{E_{q+k} + E_k - \hbar\omega - i\epsilon}$$

↖

Coherence factor

↗

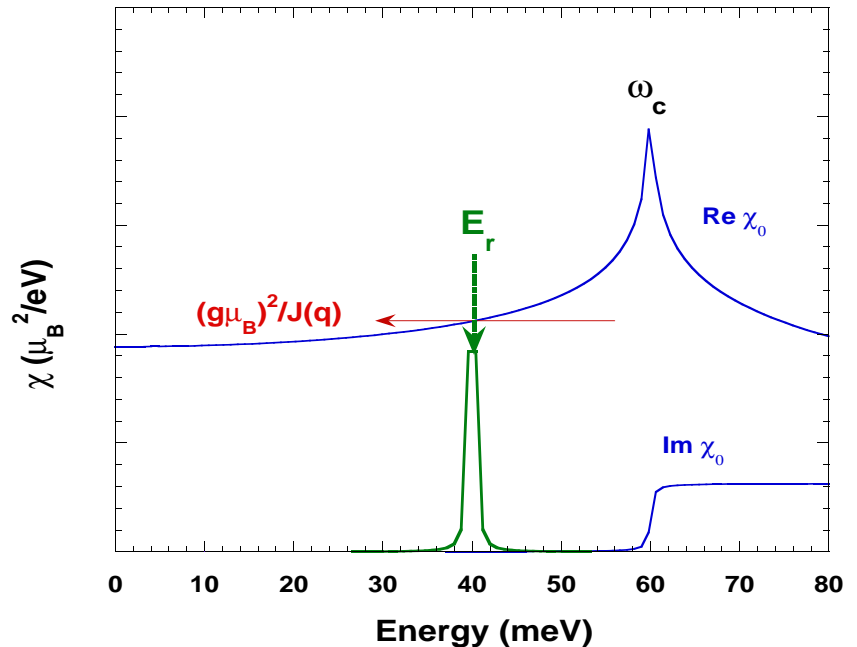
Threshold
Particle-hole
continuum

Collective mode: Spin exciton

Interactions $J(q)$ or U
(RPA)

$$\chi(q, \omega) = \frac{\chi^0(q, \omega)}{1 - J(q)\chi^0(q, \omega)}$$

□ Resonance peak
when
 $\text{Re}\chi_0(\mathbf{Q}, E_r) = 1/J(\mathbf{Q})$



In bilayer: $\mathbf{J}(\mathbf{Q}_{AF}) = \mathbf{J}_{//} + \mathbf{J}_{\perp} \Rightarrow$

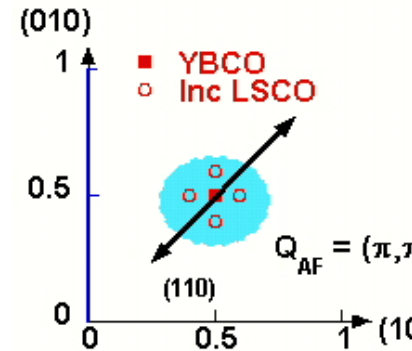
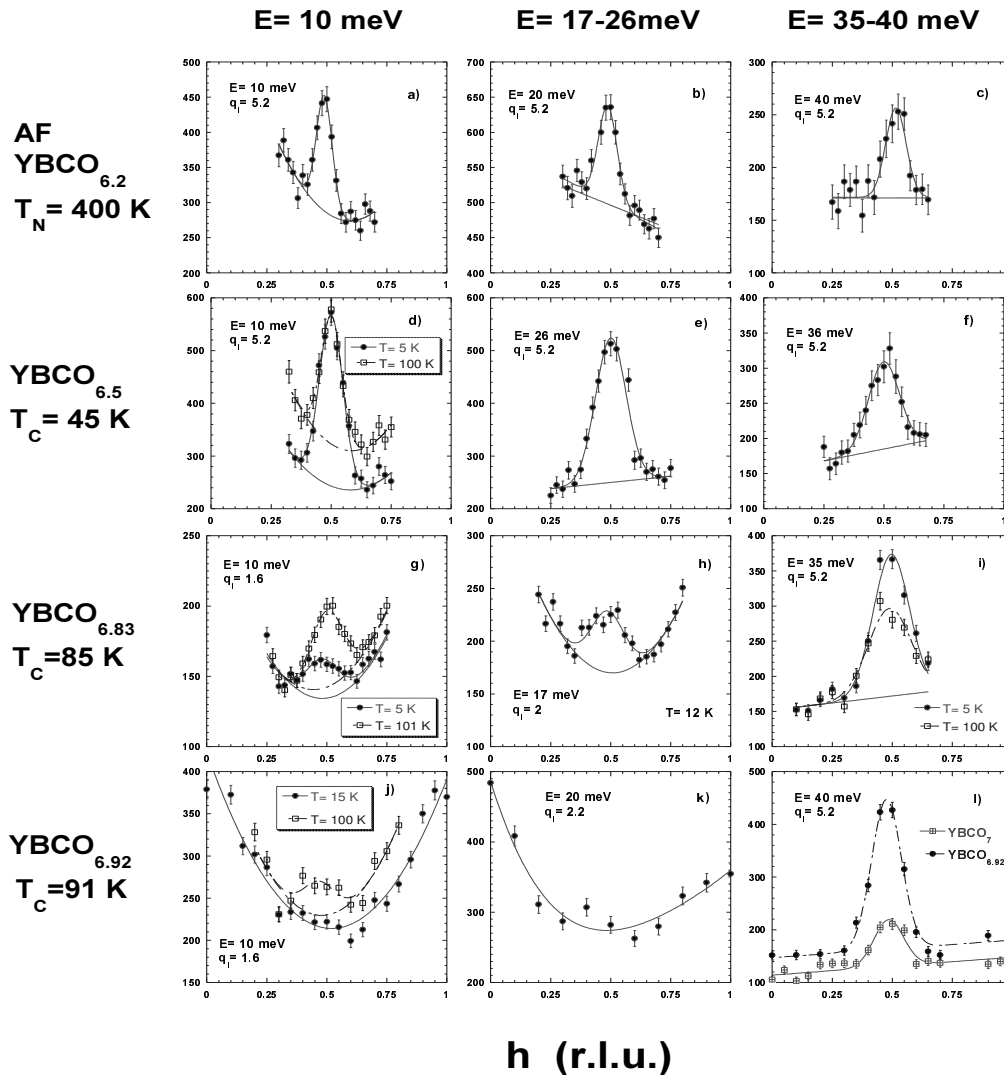
$$E_r(\text{1-layer}) > E_r(\text{2-layers})$$

BSCO-Optimal doping:

$$\Delta_{max} = 35 \text{ meV ARPES (J. Mesot et al PRL 1999)}$$

$$\omega_c = 1.8\Delta_{max} = 63 \text{ meV} \Leftrightarrow E_r = 43 \text{ meV} \simeq 1.2\Delta_{max}$$

Fluctuations magnétiques: vue d'ensemble

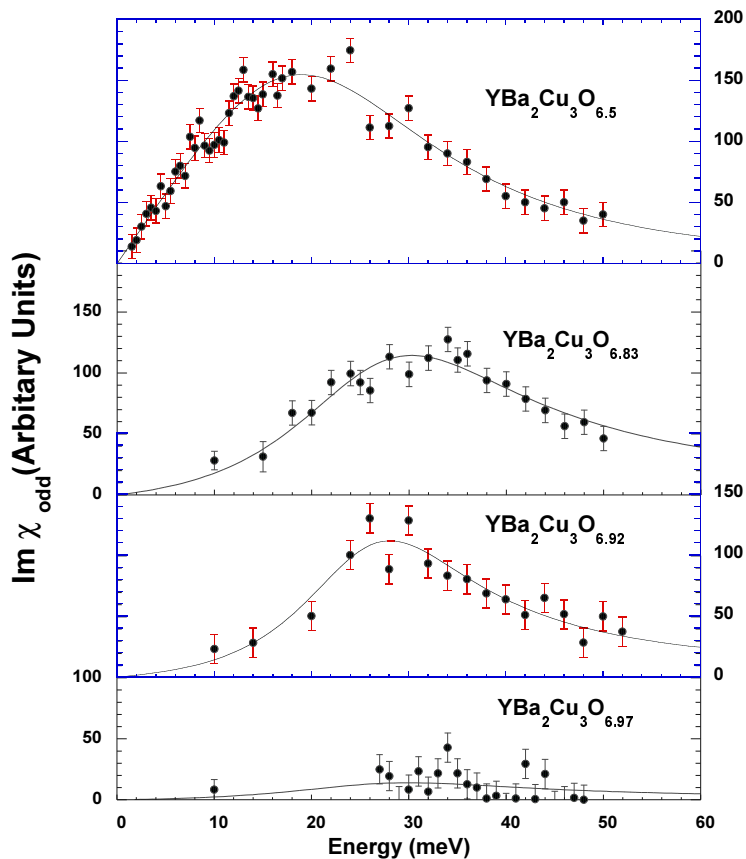


Même
échantillon
 $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$
(J.Y. Henry)

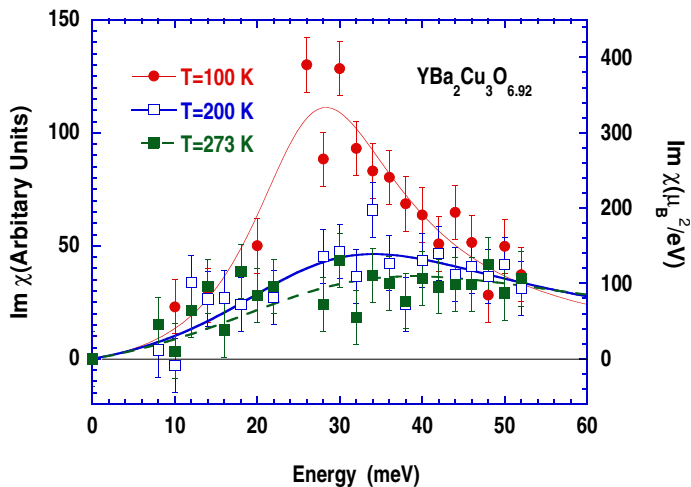
Même
spectromètre
2T (LLB)

Excitations magnétiques état normal à Q_{AF}

Normal State, 100 K, $Q=(\pi,\pi)$



Dépendance
en température $x=0.92$



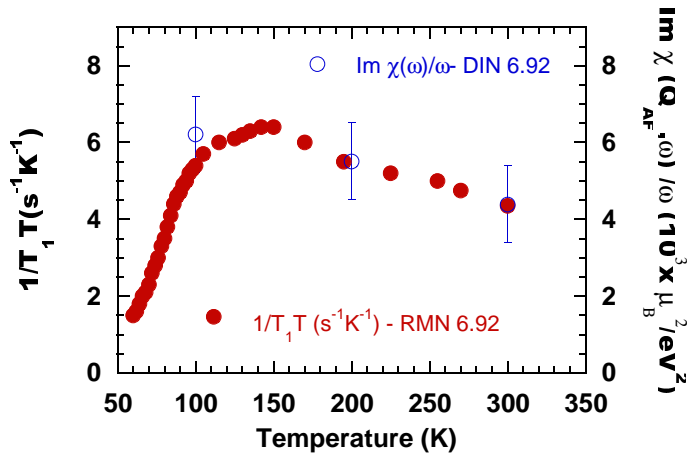
$$\text{Im } \chi^{AF} \sim 100\text{-}300 \mu_B^2/\text{eV}$$

\gg

$$\text{Im } \chi_0 \sim 1/t \sim 3\text{-}5 \mu_B^2/\text{eV}$$

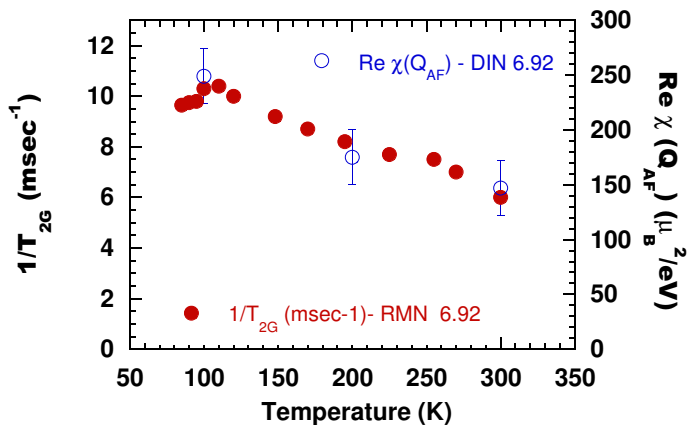
(susceptibilité uniforme)

Comparaison: Neutrons/RMN



J.A. Gillet, thèse Grenoble (1994).

C. Berthier et al, J. Phys I, **6**, 2205 (1997).



RMN du noyau ^{63}Cu : sonde Q_{AF}

Temps de relaxation spin-réseau:

$$\frac{1}{^{63}T_1 T} \sim \lim_{\omega \rightarrow 0} \sum_Q |A(Q)|^2 \frac{\text{Im} \chi(Q, \omega)}{\omega}$$

Accord de l'évolution
en fonction de la température

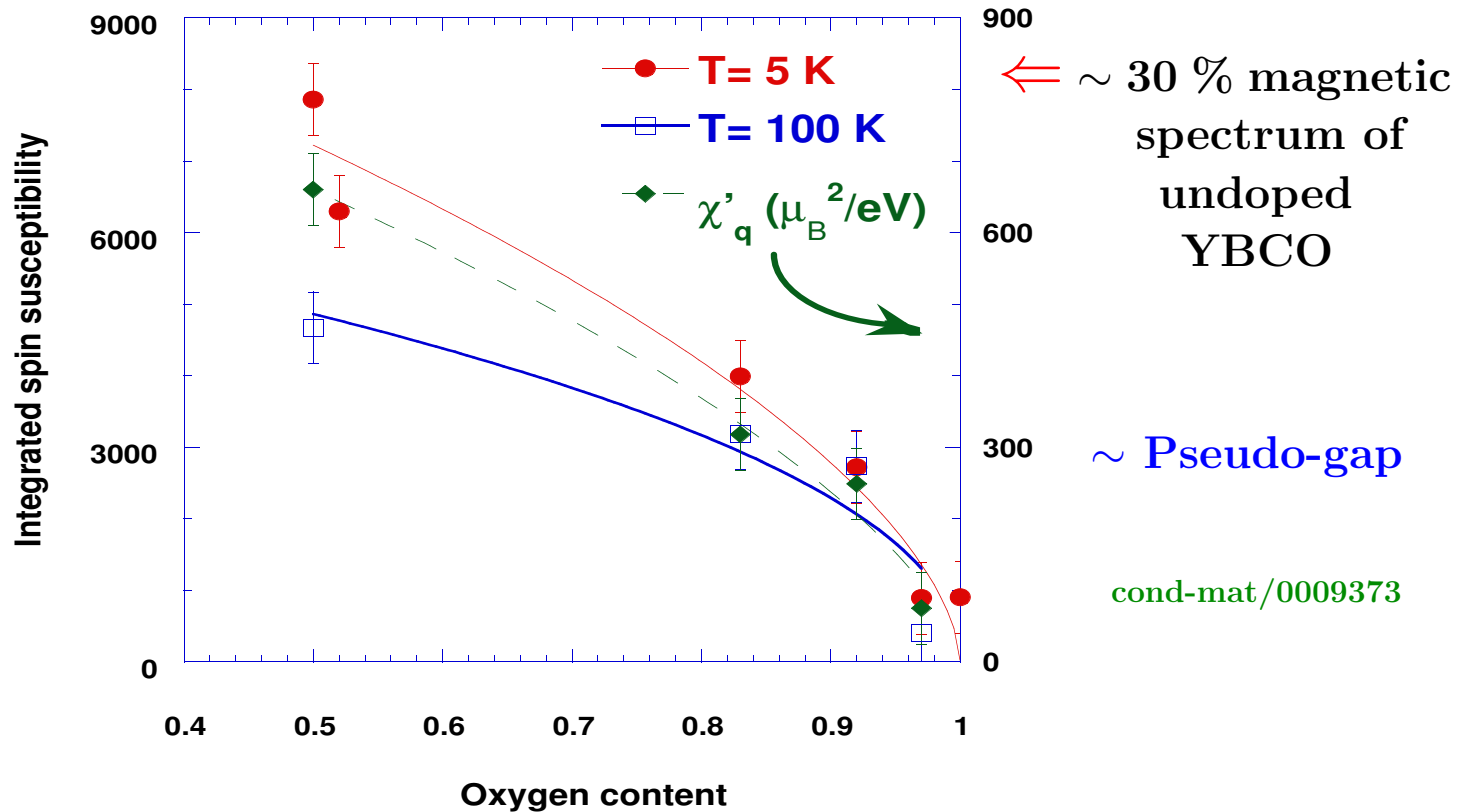
□ Facteur d'échelle 3-4 trop grand
⇒ problème dépendance en Q

Temps de relaxation spin-spin:

$$\frac{1}{^{63}T_{2G}^2} \sim \sum_Q \left[|A(Q)|^4 (\text{Re} \chi(Q))^2 - (|A(Q)|^2 \text{Re} \chi(Q))^2 \right]$$

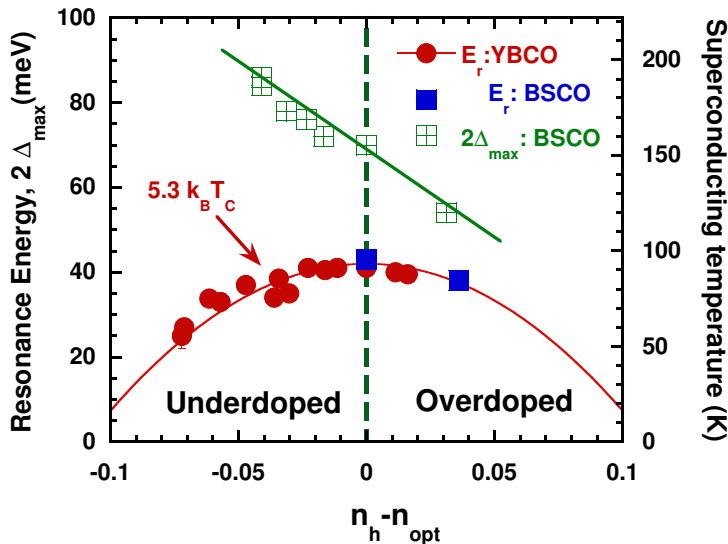
Total magnetic spectral weight

$$\int_0^{50\text{meV}} d\omega \text{Im}\chi(Q_{AF}, \omega)$$



Impact on electronic properties:

$$E_r \propto 5.3k_B T_C$$

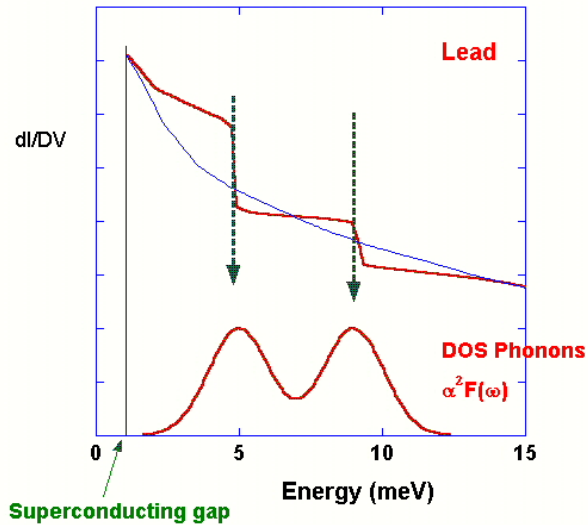


cond-mat/0211227.

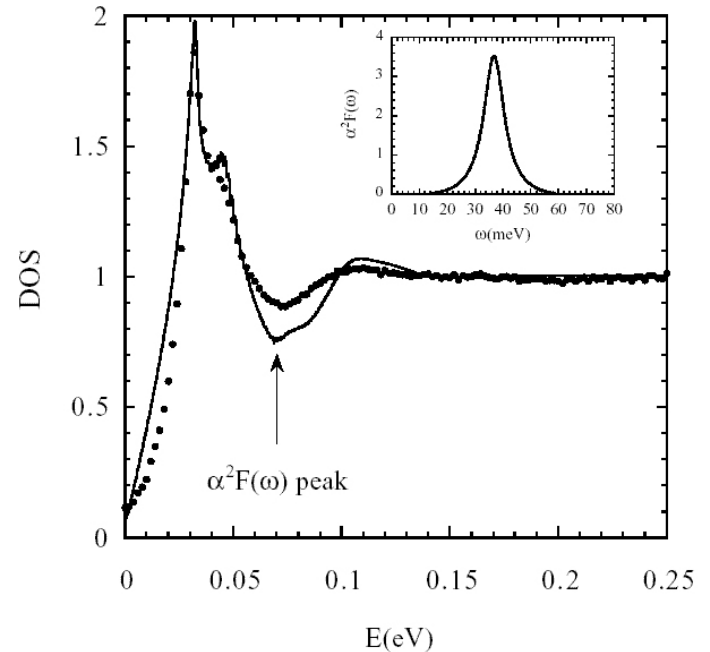
- Single particle electronic spectrum (ARPES)
J.C. Campuzano et al,
P.D. Johnson et al, J. Fink et al
- Tunneling
Zasadzinski et al,
PRL 87 067005(2001).
- Optical data
Carbote et al, Munzar et al
- Raman scattering

Tunneling Junction

□ Strong coupling analysis

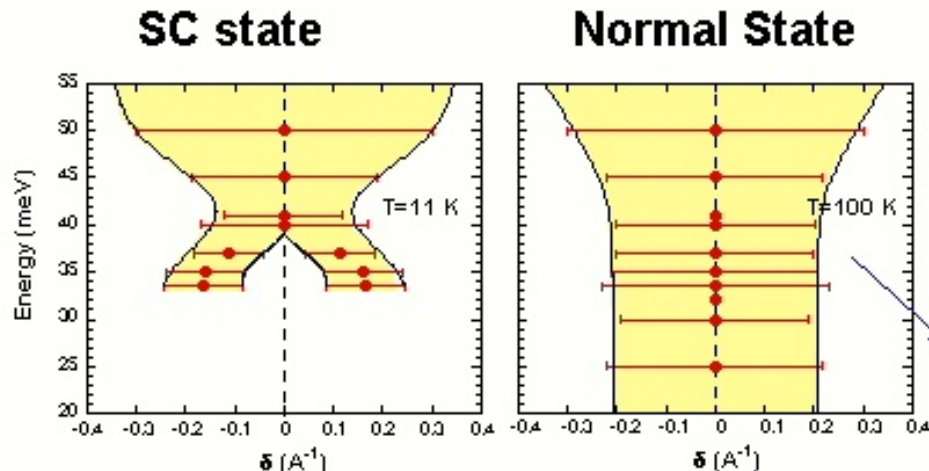


J.R. Schrieffer, *Theory of Superconductivity* (1964)



J.F. Zasadzinski et al, preprint.

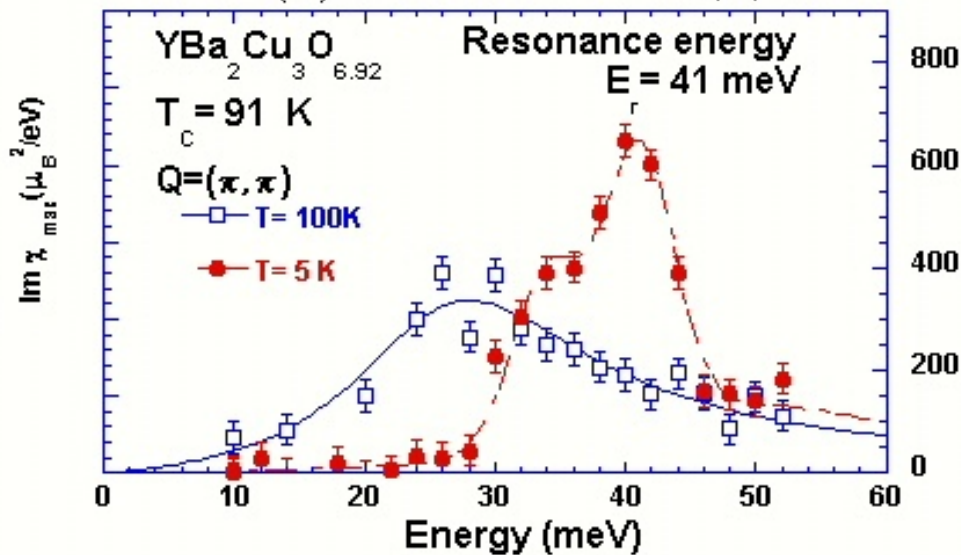
SC state *vs* Normal state



YBCO_{6.85}

$T_C = 89 \text{ K}$

$$\Delta_q = \alpha T_C$$



Spectral weight
Redistribution
in Q and ω
across T_C

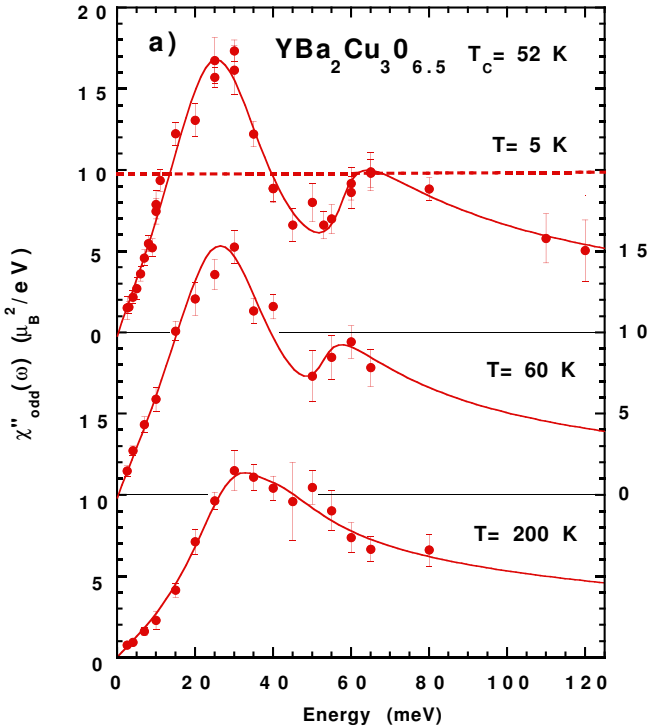
[cond-mat/9902067](https://arxiv.org/abs/cond-mat/9902067)

Conclusion/Perspectives

1. **Mode magnétique collectif (phase SC): $T_c^{\max} \geq 90$ K**
Pas présent dans $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$ ($T_c^{\max} \simeq 37$ K),
ni $\text{Bi}_2\text{Sr}_2\text{CuO}_{6+\delta}$ ($T_c^{\max} \simeq 27$ K)
2. **Fortes corrélations magnétiques de l'état normal**
 - Forte dépendance en dopage: $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$
quasiment pas détectable au dopage optimal:
 $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$, $\text{Tl}_2\text{Ba}_2\text{CuO}_{6+\delta}$
 - Accord $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$
(à part position en vecteur d'onde incommensurable)
 - Pas de signal détecté dans $\text{Bi}_2\text{Sr}_2\text{CuO}_{6+\delta}$
3. **Etat surdopé: YBCO-Ca**
4. **Role du magnétisme en comparaison avec Sr_2RuO_4 et $\text{Na}_{0.35}\text{CoO}_2, \text{H}_2\text{O}$**

Local Susceptibility: $Im\chi_{loc}(\omega) = \int_{ZB} d^2\mathbf{q} Im\chi(\mathbf{q}, \omega) / \int_{ZB} d^2\mathbf{q}$

YBa₂Cu₃O_{6.5}



□ Similar to La_{1.86}Sr_{0.14}CuO₄

S.M. Hayden et al, 76, 1344, PRL (1996)

Nearly AF Liquid

(Moriya et al, Pines et al)

$$\chi(\mathbf{q}, \omega) = \frac{\chi_{AF}}{1 + (q\xi)^2 - (\omega/\Delta)^2 - i(\omega/\omega_{sf})}$$



$$\chi_{NAFL}^{2D}(\omega) \simeq \frac{\chi_{AF}}{4(\xi/a)^2} \left[\frac{1}{2} + \frac{1}{\pi} \arctan \left[\omega_{sf} \frac{\omega^2 - \Delta^2}{\omega \Delta^2} \right] \right]$$

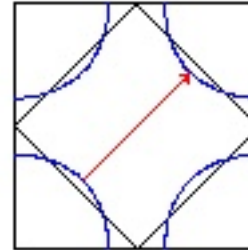
PRB, 56 11439 (1997).

Dépendance en vecteur d'onde

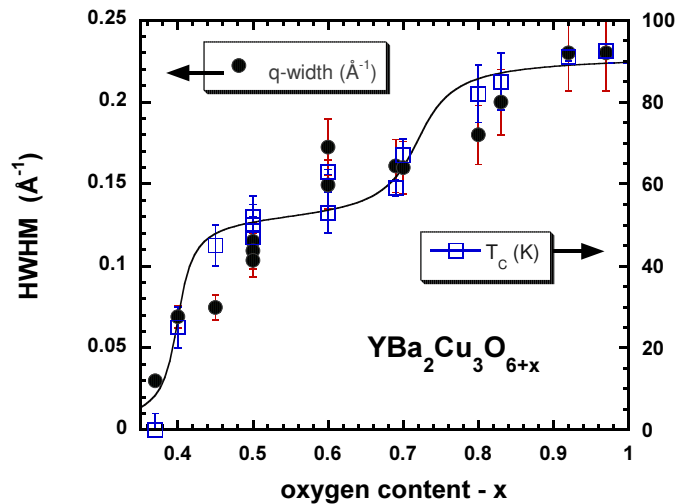
□ $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$:

Pics incommensurables

□ $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$: Largeur en Q



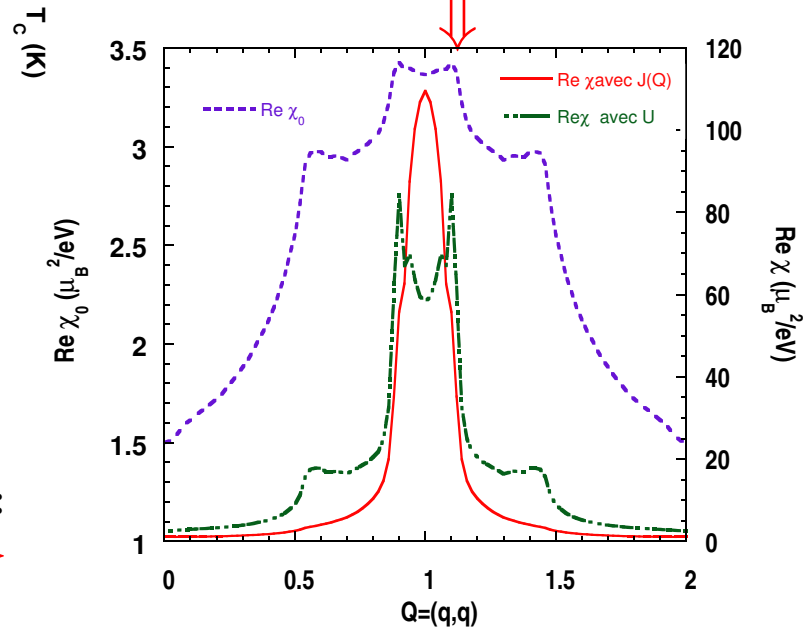
vecteur d'onde "node-node"



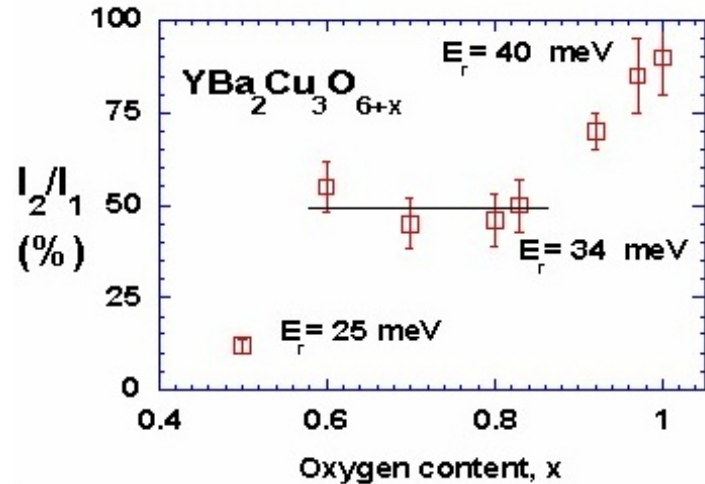
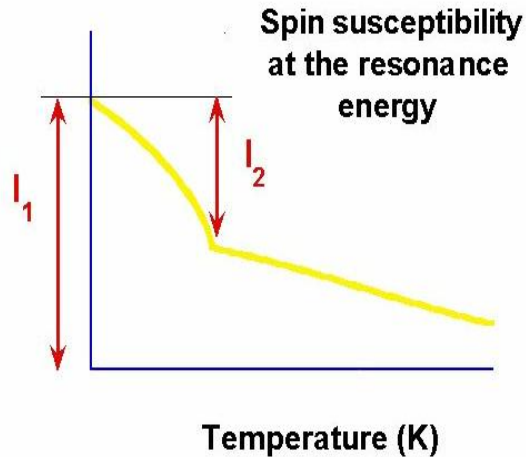
PRL, 82 5337 (1999).

susceptibilité RPA:

$$\chi(q) = \frac{\chi^{\circ}(q)}{1 - J(q)\chi^{\circ}(q)} \Rightarrow$$



Resonance peak spectral weight



□ For all doping: $\int d^3Q d\omega \text{Im}\chi(Q, \omega) \simeq 0.05 \mu_B^2$

⇒ about 2 % magnetic spectrum of undoped YBCO

□ Is it enough to explain electronic anomalies ?

Kee, Kivelson and Aeppli, PRL, 88, 257002 (2002) VS

Abanov, Chubukov, Eschrig, Norman and Schmalian, Phys. Rev. Lett., 89, 177002 (2002)