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Angela Digulescu

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THÈSE

Pour obtenir le grade de

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Communauté Université Grenoble Alpes et
l'Académie Technique Militaire de Bucarest**

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Présentée par

Angela DIGULESCU

Thèse dirigée par **Cornel IOANA** et **Alexandru ȘERBĂNESCU**
préparée au sein des **Grenoble Images Parole Signal
Automatique (GIPSA-Lab)** et **Département des
Communications et Systèmes Electroniques Militaires**
dans les **Écoles Doctorales d'Electronique, Electrotechnique,
Automatique et Traitement du Signal**, respectivement
**Systèmes Electroniques, Informatiques et des
Communications pour la Défense et la Sécurité**

Caractérisation des phénomènes dynamiques à l'aide de l'analyse du signal dans les diagrammes des phases

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devant le jury composé de :

Charles L. WEBBER, Jr.

Professeur, Université Loyola Chicago, Etats Unis (Rapporteur)

Marie CHABERT

Professeur, INP-ENSEEIH Toulouse, France (Rapporteur)

Corneliu BURILEANU

Professeur, Université "Politehnica" de Bucarest, Roumanie (Président)

Gabriel CIOCAN

Professeur Associé, Université de Laval, Canada (Membre)

Jérôme MARS

Professeur, Université Grenoble Alpes, France (Membre)

Alexandru ȘERBĂNESCU

Professeur, Académie Technique Militaire de Bucarest, Roumanie (Directeur
de thèse)

Cornel IOANA

Maitre de conférences HDR, Université Grenoble Alpes, France (Directeur de
thèse)

Alexandre GIRARD

ING EDF R&D, France (Membre invité)



THESIS

CHARACTERIZATION OF DYNAMIC PHENOMENA BASED ON THE SIGNAL ANALYSIS IN PHASE DIAGRAM REPRESENTATION DOMAIN

Presented and defended by

Angela DIGULESCU

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*Faculty of Military Electronic and Informatics Systems
Specialization: Electronic Engineering and Telecommunications*

Thesis directed by Cornel IOANA and Alexandru ȘERBĂNESCU.

Prepared in the Grenoble Image Parole Signal Automatique laboratory (GIPSA-lab) and at the Faculty of Military Electronic and Informatics Systems.

Defended in Grenoble, on 17th January 2017, in front of the jury:

Reviewers: Charles L. WEBBER, Jr. - Professor, Loyola University Chicago, U.S.A.

Marie CHABERT - Professor, INP-ENSEEIH T Toulouse, France

Examiners: Corneliu BURILEANU - Professor, University "Politehnica" Bucharest, Romania

Gabriel CIOCAN - Associated Professeur, University Laval, Canada

Jérôme MARS - Professor, Université Grenoble Alpes, France

Cornel IOANA - Associate Professor, Université Grenoble Alpes

Alexandru ȘERBĂNESCU - Professor, Military Technical Academy, Romania

Invited: Alexandre GIRARD - ING EDF R&D, France

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Contents

- 1. Introduction: Physical parameters exploring in heterogeneous environments 1
 - 1.1 Sensing in heterogeneous environments – physical modeling and particularities 1
 - 1.1.1. Active configurations 2
 - 1.1.2. Passive configurations 6
 - 1.2. Thesis outline..... 9
- 2. Thesis work motivation and positioning..... 11
 - 2.1. Emission 12
 - 2.1.1. Narrowband pulses 12
 - 2.1.2. Wideband pulses..... 13
 - 2.2. Reception 14
 - 2.2.1. Energetic techniques..... 14
 - 2.2.2. Projective techniques..... 14
 - 2.2.3. Data-driven techniques 16
 - 2.3. Chapter summary 19
- 3. Active sensing: phase diagram-based adaptive waveform 21
 - 3.1. Introduction..... 21
 - 3.2. Phase diagram-based waveform design..... 23
 - 3.3. Emission 23
 - 3.4. Reception 27
 - 3.2.1. Phase space-based IFL estimation..... 34
 - 3.2.2. “Marked” lobes..... 38
 - 3.3. Chapter summary 41
- 4. Passive sensing based on the RQA concept..... 43
 - 4.1. Time-distributed recurrence..... 43
 - 4.2. Multi-lag phase-space analysis 46
 - 4.3. Diagonal lines quantification 50
 - 4.4. Chapter summary 54
- 5. Applications in active configurations 55
 - 5.1. Cavitating environments 55
 - 5.1.1. Experimental configuration and results 56
 - 5.2. Underwater object tracking..... 64
 - 5.3. High speed sensing using phase diagram-based adaptive waveform 69
 - 5.4. Chapter summary 72

6. Applications in passive configurations	75
6.1. Measuring pressure transient in water pipes.....	75
6.1.1. Water hammer parameters.....	75
6.1.2. Experimental setup and results	76
6.2. Electrical arcs.....	80
6.2.1. Detection performance	81
6.3. Chapter summary.....	85
7. Conclusions and perspectives	87
7.1. Work summary	87
7.2. Perspectives	88
Bibliography	97

List of Figures

Fig. 1-1: The experimental configuration for the cavitation generation	3
Fig. 1-2 : Signal's multi-path propagation for the cavitating phenomena.....	3
Fig. 1-3: Emitted vs. received signal in a multi-path propagation defined by the solid-monophasic water flow-solid.	4
Fig. 1-4: Illustration of the Doppler deformation due to the cavitation	4
Fig. 1-5 : The experimental configuration of the underwater object tracking	5
Fig. 1-6 : Signal's propagation in the heterogeneous environment defined by the presence of an object in the water flow, when sensing with an active acoustic paths.....	5
Fig. 1-7: Example of signals used for the object tracking application	6
Fig. 1-8 : Experimental setup for the water hammer (CV – closing valve)	6
Fig. 1-9: Signal's propagation in a heterogeneous pressure field	7
Fig. 1-10: The acoustic signals recorded by the ultrasonic transducers and the pressure recorded by the intrusive pressure transducer	7
Fig. 1-11 : Experimental configuration for the detection and the localization of the electrical arcs.....	8
Fig. 1-12: Signal's propagation for a discharge of air particle.....	8
Fig. 1-13: Examples of electrical arcs recorded with the acoustic sensors	9
Fig. 2-1: General non-intrusive acoustic sensing system	11
Fig. 2-2: Example of a narrowband pulse used for emission (left) and its received version (right) in a heterogeneous environment.....	12
Fig. 2-3: Example of wideband signal used for emission (left) and its received version (right)	13
Fig. 2-4: The matched filter applied for the signals presented in Fig. 2-2 (left), respectively in Fig. 2-3 (right)	14
Fig. 2-5: Transient signals considered for studying the wavelet transform	16
Fig. 2-6: The wavelet transform (using the Mexican Hat mother-function) applied on the three synthetic signals	16
Fig. 2-7: Example of phase space trajectory implementation (Ioana, et al., 2014).....	18
Fig. 2-8: DM and RM when applying the Euclidean distance (Ioana, et al., 2014).....	18
Fig. 3-1: General principle in radar and acoustic sensing	21
Fig. 3-2: Favorable scenario using the classical approach for the acoustic inspection (left) and de unfavorable scenario using the same approach (right)	22
Fig. 3-3 : Our proposed approach for the high-speed sensing.....	22
Fig. 3-4: The overall diagram of the proposed algorithm of waveform design in phase space	23

Fig. 3-5: Director vector example in 51 points: $a = 1, b = 1.5, c = 2$	24
Fig. 3-6: The auxiliary signal aux with the following parameters: $a_1 = 1.5, b_1 = 1.1, c_1 = 1.9,$ $a_2 = 1.5, b_2 = 1.2, c_2 = 1.8$, $k = 50, N = 100, f_s = 5 MHz$ and $t_0 = 1 / f_s = T_s$	25
Fig. 3-7: The phase space lobe obtained using the signal aux from Fig. 3-6.....	25
Fig. 3-8: The phase space lobe and its corresponding aux signal based on eq. (3.5) and (3.7) with the following parameters: $N = 100, a_1 = 1, b_1 = 1.2, c_1 = 1.3, n_1 = 50, n_2 = 50, n_3 = 60$	26
Fig. 3-9: The phase space lobe and its corresponding aux signal based on eq. (3.5) and (3.7) with the following parameters: $N = 100, a_1 = 1, b_1 = 1.2, c_1 = 1.3, n_1 = 40, n_2 = 50, n_3 = 60$	27
Fig. 3-10: The amplitude modulated signals (left) which have as envelope the auxiliary signal defined in eq. (3.8) shown in phase space representation (right)	27
Fig. 3-11: The recovered signal's envelopes (normalized values) for: $SNR \rightarrow \infty$ and $SNR = 20dB$ (left); $SNR \rightarrow \infty$ and $SNR = 15dB$ (right);	28
Fig. 3-12: The frequency modulated signals whose IFLs are auxiliary signals from eq. (3.9)	29
Fig. 3-13: The spectrogram of the modulated signals s_1 and s_2 from Fig. 3-12.....	29
Fig. 3-14: The extracted IFLs using the spectrogram for the modulated signals s_1 (left) and s_2 (right).....	30
Fig. 3-15: The phase space lobes provided by the IFLs from s_1 and s_2	30
Fig. 3-16: The correlation of the modulated signals s_1 and s_2 (left) and their amplitude spectrum (right)	31
Fig. 3-17: The correlation of the auxiliary signals aux_1 and aux_2 (left) and their amplitude spectrum (right)	31
Fig. 3-18: Phase space lobes in noise-free conditions: original IFL (triangles) and the spectrogram recovered IFL (stars)	32
Fig. 3-19: Phase space lobes for $SNR = 10$ dB: original IFL (triangles) and the spectrogram recovered IFL (stars)	32
Fig. 3-20: Phase space lobes generated from the set of 6 auxiliary signals used IFLs	33
Fig. 3-21: The original IFLs (continuous line) and the spectrogram-based recovered IFLs (stars) – noise -free case	33
Fig. 3-22: The original IFLs (continuous line) and the spectrogram-based recovered IFLs (stars) for $SNR = 10$ dB.....	34
Fig. 3-23: The original IFLs (continuous line) and the spectrogram-based recovered IFLs (stars) for $SNR = 5$ dB.....	34
Fig. 3-24: The time evolution of a sine wave during one period (left) and its trajectory in 2D phase space (right): $f_0 = 20; f_s = 400; D = 1 / f_0; d = 3$	35

Fig. 3-25: The time evolution of an interpolated sine wave during one period (left) and its trajectory in 2D phase space (right): $f_0 = 20$; $f_s = 400$; $D = 1 / f_0$; $d = 3$ and the interpolation factor $f = 5$	35
Fig. 3-26: The time evolution of a sine wave during two periods (left) and its trajectory in 2D phase space (right): $f_0 = 20$; $f_s = 400$; $D = 2 / f_0$; $d = 3$	36
Fig. 3-27: The original IFL (in blue) and the phase space-based IFL estimation (in red – $f = 5, d = 10$)	37
Fig. 3-28: The phase space-based IFL for $d = 5$ and $SNR \rightarrow \infty$ (left) and the MLPA phase space-based IFL $d = 15$ 15 13 14 15 15 (according to the colors: clockwise from blue to green) and $SNR \rightarrow \infty$ (right); continuous line – original IFL; starred line – estimated IFL.....	37
Fig. 3-29: The phase space-based IFL for $d = 5$ and $SNR = 5dB$ (left) and the MLPA phase space-based IFL $d = 21$ 27 19 30 21 25 (according to the colors: clockwise from blue to green) and $SNR = 5dB$ (right); continuous line – original IFL; starred line – estimated IFL..	38
Fig. 3-30: The ideal phase space representation of the “marked” lobe (auxiliary signal)	38
Fig. 3-31: The phase diagram space approach versus classical approach; in the phase diagram, the estimation of TOF is done using the distance between the marks of emitted and received signals, respectively.....	39
Fig. 3-32: The IFL (auxiliary signal) (left) that forms the phase space marked lobe (right): $x_1 = 0.8, x_2 = 0.6, x_3 = 1$ and $r = 1.2$	40
Fig. 3-33: The multiple lobes with signature obtained based on the lobe from Fig. 3-32 and rotated with an angle of $\pi / 6$	40
Fig. 3-34: Example of correlation between the signals that contain the IFLs blue and green from Fig. 3-33 (left) and their corresponding amplitude spectral representation (right)	40
Fig. 3-35: The original IFLs (continuous line) and the phase space-based estimated IFLs (starred lines) for $SNR = 20dB$ (left) and their corresponding phase space lobes (right)	41
Fig. 3-36: The original IFLs (continuous line) and the phase space-based estimated IFLs (starred lines) for $SNR = 7dB$ (left) and their corresponding phase space lobes (right)	41
Fig. 4-1: The transient detection using the TDR^* measure: (up) the analyzed transient signal; (left-down) the recurrence matrix of the analyzed signal; (right-down) the analyzed signal and its TDR^* based detection curve.....	44
Fig. 4-2: Transient signal detection using the TDR measure for a transition of only 11 samples where $m = 3$, $d = 2$ and $\alpha = 0.8$	45
Fig. 4-3: Transient signals considered for the multi-lag phase space analysis.....	47
Fig. 4-4: The phase space trajectory evolution using different delays	48
Fig. 4-5: The evolution of the area according to the lag (delay): for the signals presented in Fig. 2-5 (left) and for the normalized signals (right).....	48
Fig. 4-6: The polar coordinates representation for the transient signals considered in our study: $s_1 : d_1 = 5, s_2 : d_2 = 5$ and $s_3 : d_3 = 6$	49

Fig. 4-7: The schematic presentation of the multi-lag tools.....	50
Fig. 4-8: Example of DLQ: (left) test signal – a sine wave with the frequency $f = 10\text{Hz}$; (center) the DLQ measure; (right) zoom on the DLQ measure;	51
Fig. 4-9: The time evolution of the test signal modulated in amplitude	52
Fig. 4-10: The spectral representation of the diagonal lines quantification (up - left), the classical amplitude spectrum representation (up – right) and the envelope amplitude spectrum (down – center); for the computation of the distance matrix, it is used the squared Euclidean distance, $m = 3$ and $d = 2$	52
Fig. 4-11: The frequency modulated test signal	53
Fig. 4-12: The classical amplitude spectrum representation (left) and the spectral representation of the diagonal lines quantification (right); for the computation of the distance matrix, it is used the squared Euclidean distance, $m = 3$ and $d = 2$	54
Fig. 5-1: The installation for the first set of measurements.....	56
Fig. 5-2: The on-site setup for the first set of measurements	57
Fig. 5-3: The installation for the second set of measurements	57
Fig. 5-4: The on-site setup for the second set of measurements	57
Fig. 5-5: Cavitating vortex stages: A1 no vortex (left), A2 incipient vortex (center), A3 developed vortex (right)	57
Fig. 5-6: The emitted signal for each measurement, a received signal during a trial for the incipient cavitation and a received signal during a trial for the cavitation flow	58
Fig. 5-7: The amplitude spectrum of the envelope for a received signal when the cavitation vortex is fully developed	59
Fig. 5-8: The amplitude spectrum for a received signal when the cavitation vortex is fully developed.....	59
Fig. 5-9: The amplitude DLQ (eq. (4.13)) for an acquired signal when the cavitation vortex is fully developed	59
Fig. 5-10: The DLQ frequency position (left) and its mean (for 20 samples - right) using the diagonal lines quantification approach for the signals acquired at each operating point (the number in red box indicates the configuration); the RPA parameters are $m = 5$ and $d = 3$ using the Euclidean distance	60
Fig. 5-11: The cavitation evolution characterized in terms of frequency standard deviation (in red) and pressure standard deviation (in black).....	61
Fig. 5-12: DLQ absolute value for an acquisition of 2.6ms in A1 region.....	62
Fig. 5-13: DLQ absolute value for an acquisition of 2.6ms in A3 region.....	62
Fig. 5-14: Time variation of cavitation vortex component (in blue) and the flow rate evolution (in red); the RPA parameters are $m = 5$ and $d = 15$ using the Euclidean distance.....	63
Fig. 5-15: The experimental configuration: the dash-dot line shows the trajectory of the object	65

Fig. 5-16: The recorded signal arrived at Rx2 sensor	66
Fig. 5-17: The detection map obtained for the Rx2 signal for $m = 3$, $d = 13$ and $\alpha = 0.5$	66
Fig. 5-18: The TOAs (red) of each response of Rx2 signal and its signal envelope (black) ...	67
Fig. 5-19: The recorded signal arrived at Rx1 sensor	67
Fig. 5-20: The detection map obtained for the Rx1 signal for $m = 3$, $d = 7$ and $\alpha = 0.5$	68
Fig. 5-21: The TOAs (red) of each response of Rx1 signal and its signal envelope (black) ...	68
Fig. 5-22: The IFL used at emission (left) and its corresponding “marked” lobe (right)	69
Fig. 5-23: The positioning of the two transducers for the experimental test.....	70
Fig. 5-24: The signal used at emission and the received signal as well as their amplitude spectrum	70
Fig. 5-25: The IFLs of the emitted and received signals (left) and their representation in phase space (right).....	70
Fig. 5-26: The time shifted IFLs used for overlapping where $N = 1000$	71
Fig. 5-27: The overlapped signals which form the signal $s[n]$ (left) and its spectrogram (right)	71
Fig. 5-28: Phase space lobes of the IFLs used for overlapping.....	71
Fig. 5-29: The sum of the IFLs (left) and the phase space lobes of the IFL 2 (right)	72
Fig. 5-30: The sum of the IFLs (original – no noise, estimated – SNR=25 dB) (left) and the phase space lobes of the IFL 2 (original – no noise, estimated – SNR=25 dB) (right).....	72
Fig. 6-1: Experimental setup	76
Fig. 6-2: The on-site experimental setup of the ultrasonic transceivers performed at Power Engineering Faculty, ”Politehnica” University of Bucharest in December 2014	77
Fig. 6-3: The acoustic signals recorded by the two sensor and the simultaneous pressure time evolution: (left) acoustic signal arrived at sensor S2; (right) acoustic signal arrived at sensor S1	78
Fig. 6-4: The water hammer effect detection on the signals recorded by the two sensors.....	78
Fig. 6-5: The experimental configuration of the EA locator system.....	81
Fig. 6-6: The recorded EAs and the detection curve obtained with the TDR measure: $m = 3$, $d = 8$ and $\alpha = 0.7$	81
Fig. 6-7: The detection performances of each method; the ROCs where computed for the same white additive Gaussian noise (100 realizations) for each SNR level and method: (up-left) TDR^* measure for $m = 3$, $d = 8$ and $\alpha = 0.7$; (up-right) the spectrogram-based detection; (down) the wavelet-based detection using the Mexican Hat mother function	82
Fig. 6-8: The electrical arc recording s_1 and its reflections s_2, s_3	84
Fig. 6-9: The areas of the estimated ellipses that circumscribe the phase-space trajectory of the normalized multi-path signals	84

Fig. 6-10: The polar coordinates representation for the multi-path acoustic signals: $d_1 = 11$,
 $d_2 = 12$ and $d_3 = 8$ 85

List of Tables

Table 2-1: The requirements for a general sensing system	11
Table 5-1: System's parameters for each operating point.....	60
Table 5-2: Frequency content parameters	61
Table 5-3: Criteria limits for the cavitating vortex stages classification in increasing flow rate configuration	63
Table 5-4: Criteria validation for decreasing flow rate configuration.....	64
Table 6-1: Characteristic parameters of the pipe and fluid for theoretical computation of the pressure wave	79
Table 6-2: The pressure variations using different approaches.....	80
Table 6-3: The spatial localization accuracy for the electrical arcs	83

Table of acronyms and abbreviations

CV – Closing Valve,	6
DLQ - Diagonal Lines Quantification,	50
DM – Distance Matrix,	17
EA – Electrical Arc,	8
FFT - Fast Fourier Transform,	33
IFL - Instantaneous Frequency Law,	13
LFM - Linear Frequency Modulations,	28
MF – Matched Filter,	14
MLPA - Multi-Lag Phase Analysis,	37
PRF - Pulse Repetition Frequency,	11
RM – Recurrence Matrix,	17
ROC – Receiving Operating Characteristic,	80
RPA - Recurrence Plot Analysis,	16
RQA - Recurrence Quantification Analysis,	17
SNR – Signal-to-Noise Ratio,	28
TDR - Time-Distributed Recurrence,	43
TOA – Time Of Arrival,	12
TOF – Time Of Flight,	11

1. Introduction: Physical parameters exploring in heterogeneous environments

Signals' deformation along their propagation path is among the most important aspect which has to be taken into account at reception. These effects are caused by phenomena like attenuation, reflection, dispersion and noise. Whereas the first two are rather easy to monitor, because they affect the amplitude, respectively the delay, the latter two are more difficult to control, because they change signals' parameters (amplitude, frequency and phase) in an environment-dependent manner.

Generally, the consideration of the signal deformations at the reception is motivated by the requirement to compensate their effects in order to ensure satisfactory receiving performances. In this thesis, the main objective is to contribute to the deep analysis of the propagation effects in order to better understand them as well as to estimate their parameters that are interesting from application point of view. In this thesis, the studied phenomena take place in dynamic systems which, during their time evolution, rapidly (sometimes suddenly) change their parameters.

1.1 Sensing in heterogeneous environments – physical modeling and particularities

The purpose of this part is to present the problem of sensing in heterogeneous environments from signal analysis point of view, by proposing a general model, adaptable to different applications we have the occasion to treat. The place of signal analysis in these applications is central, as it concerns the analysis of signals received by a monitoring system devoted to the measurement of different parameters that will be discussed further.

Let $s(t)$ be the transmitted signal:

$$s(t) = A \cdot \exp(j \cdot \phi(t)) \quad (1.1)$$

where A is the amplitude and $\phi(t)$ is the instantaneous phase.

We define a general model of the received signal after the propagation in heterogeneous environment:

$$r(t) = \sum_k \mathcal{D}_k \{s(t)\} + n(t) \quad (1.2)$$

where $\mathcal{D}_k \{\cdot\}$ is a general deformation operator associated to the k^{th} component of the received signal and $n(t)$ is an additive white Gaussian noise.

In a static homogenous environment (e.g. unbounded free-space), the attenuation is time invariant, it is only range and frequency dependent and the delay is only given by the propagation path, if the propagation velocity is assumed to be constant.

Let us consider, as an example, two omnidirectional antennas which transmit in their line of sight (i.e. only their direct path is received). Then, the received signal is:

$$r(t) = \alpha \cdot s(t - \tau) + n(t) \quad (1.3)$$

where $\alpha = \frac{\lambda}{4\pi R_0}$ is the path loss, c is the electromagnetic wave speed ($\cong 3 \cdot 10^8 \text{ m/s}$), R_0 is the distance between the antennas, λ is the wavelength of the transmitted signal and $\tau = \frac{R_0}{c}$.

In this case, the deformation operator \mathcal{D}_k can be written as:

$$\mathcal{D}_k \{s(t)\} = \alpha \cdot s(t - \tau), k = 1 \quad (1.4)$$

Yet, in several environments (e.g. water, cables, etc.) the propagated waves are dispersive which means that the initial signal is forced to transform such that its energy is spread out in a larger bandwidth. In dynamic systems, things become even more complicated when the studied phenomena have complex models (e.g. cavitation, water hammer, etc.). The main environments of interest for our work are discussed in the next part of this chapter.

1.1.1. Active configurations

Cavitating environments

The cavitation phenomena, widely encountered in various underwater applications and hydro power engineering, are subject of high interest both for understanding complex phenomena, as well as for the monitoring the equipment in the presence of cavitation (Takahashi & Miyazaki, 2006). The effect of the cavitation consists in the variation of the pressure in the fluid and, then, the apparition of gas bubbles. That is, the fluid becomes di-phasic as long the both states – liquid and gas- exist. In terms of acoustic signals propagating in fluid under cavitation, the main effects are the amplitude attenuation, since the presence of gas is not favorable for acoustic propagation, and Doppler frequency-shifts, due to the motion of bubbles (Rubin, 2000).

As an example, let us consider the case of a reduced scale hydraulic facility, namely a hydraulic installation provided with a system where the flow rate can be manually changed (Candel, et al., 2014a). The installation has a stator inside its pipeline whose purpose is to produce cavitation vortex while the flow rate increases. Fig. 1-1 presents the experimental configuration.

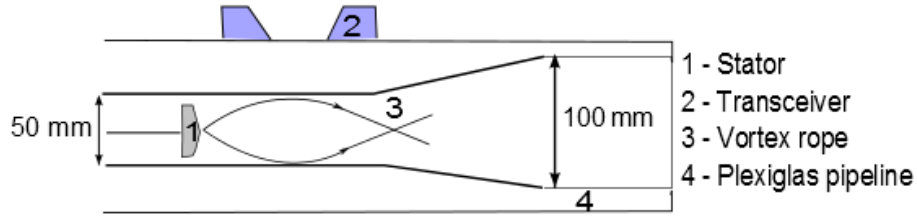


Fig. 1-1: The experimental configuration for the cavitation generation

In this facility, the cavitation parameters, in terms of rotational vortex speed and the density of bubbles, are controlled by the flow rate parameters. The purpose of our work, which will be detailed further in the thesis, is to define an active acoustic non-intrusive sensing system based on the measurement of signal's deformations when propagating in this dynamic diphasic environment. Actually, because of the presence of the two aggregation states, liquid and gas, we can consider that the propagation will be done in a heterogeneous environment, characterized by different zones with different propagation parameters. In addition, the conic shape of the acoustic sensors contributes to the multi-path behavior of the propagation in this environment. Another difficulty is the dynamic aspect of the configuration while the parameters of the flow and cavitation change over the time.

The next figure illustrates the specificity of this application described above.

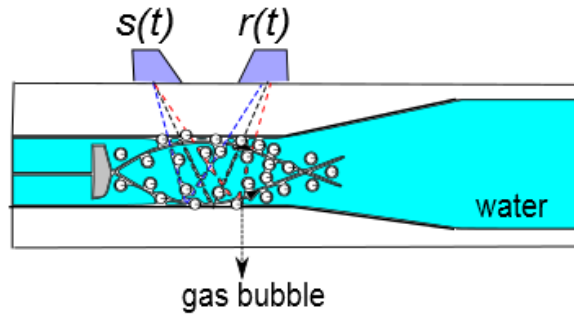


Fig. 1-2 : Signal's multi-path propagation for the cavitating phenomena

From signal's model point of view, we can derive two special cases:

- When the water is monophasic: the deformation operator \mathcal{D}_k is defined by the attenuation and delay (Fig. 1-2 - similar to eq. (1.4)):

$$\mathcal{D}_k \{s(t)\} = \alpha_k \cdot s(t - \tau_k) \quad (1.5)$$

where k is the index of the propagation path. Therefore, the received signal is:

$$r(t) = \sum_k \alpha_k \cdot s(t - \tau_k) + n(t) \quad (1.6)$$

Since the propagating paths are very close, the effect of the signal is the deformation due to the superposition of different arrivals so that the received signal has a different duration and a completely different amplitude shape than the emitted signal. This is illustrated in Fig. 1-3.

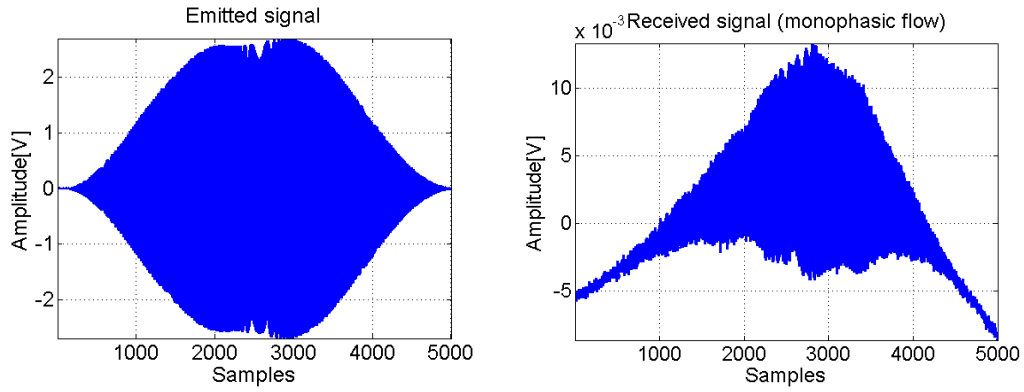


Fig. 1-3: Emitted vs. received signal in a multi-path propagation defined by the solid-monophasic water flow-solid.

- When the cavitation appears: then, the deformation operator \mathcal{D}_k is characterized by attenuation, delay as well as the Doppler effect that acts as a deformation of the instantaneous phase of the signal:

$$\mathcal{D}_k \{s(t)\} = \alpha_k \cdot s(t - \tau_k(t)) = \alpha_k \cdot A \cdot \exp(j \cdot \phi_k(t)) \quad (1.7)$$

where $\phi_k(t) = \phi \left(t - \frac{\bar{r}_k(t)}{c} \right)$ and $\bar{r}_k(t)$ is the time (Doppler) dependent position vector.

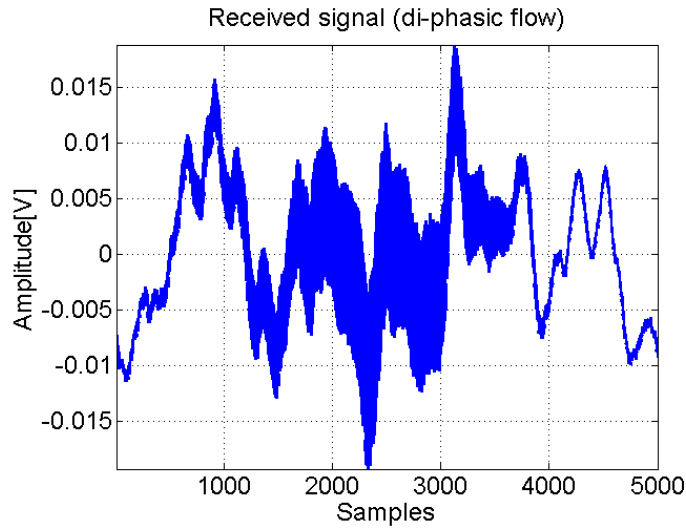


Fig. 1-4: Illustration of the Doppler deformation due to the cavitation

Fig. 1-4 illustrates the deformation of the transmitted signal in the presence of cavitation phenomena generated in our experimental setup. As it will be discussed in the devoted part of this thesis, this deformation is fast varying in time, while the cavitation structure changes in time.

Underwater object tracking

The underwater object tracking is an essential task in the marine community with uncountable number of applications – intrusion detection in protected area, underwater mammal surveillance, fish detection, pollution monitoring, to just mention few of possible applications (Yuan, et al., 2000).

The general context of this application field can be summarized by the presence of a moving object and the definition of two active acoustic paths, as indicated in Fig. 1-5.

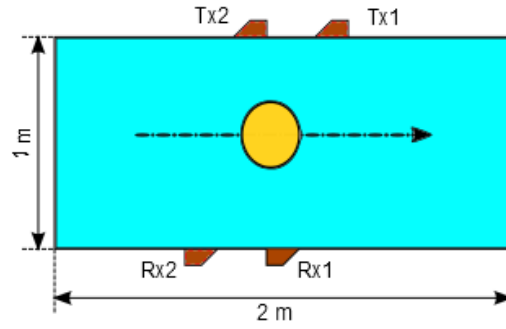


Fig. 1-5 : The experimental configuration of the underwater object tracking

In this case, the heterogeneity of the environment is given by the set water – object – water (Fig. 1-6).

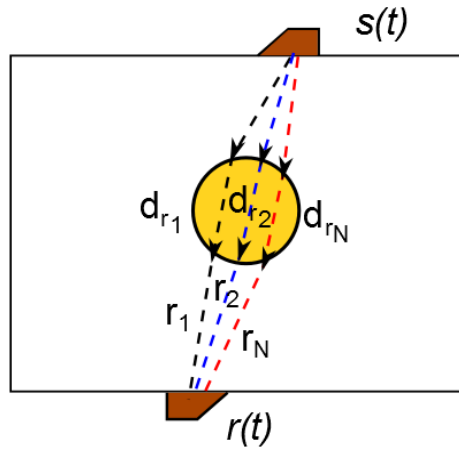


Fig. 1-6 : Signal's propagation in the heterogeneous environment defined by the presence of an object in the water flow, when sensing with an active acoustic paths

It goes that, the deformation operator is:

$$\mathcal{D}_k \{s(t)\} = \alpha_k \cdot A \cdot \exp \left(j \cdot \phi \left(t - \frac{r_k}{c_{water}} - \frac{d_{r_k}}{c_{oil}} - \tau_0 \right) \right), k = \overline{1, N} \quad (1.8)$$

where r_k is the k^{th} water propagation path, d_{r_k} is the k^{th} oil propagation path and $\tau_0 = \frac{2l_b}{c_{rubber}}$ where l_b is the thickness of the rubber wall of the balloon and c_{rubber} is the acoustic propagation velocity in rubber.

The next figure (Fig. 1-7) illustrates both transmitted and the received signal, corresponding to an impulse emission – propagation couple.

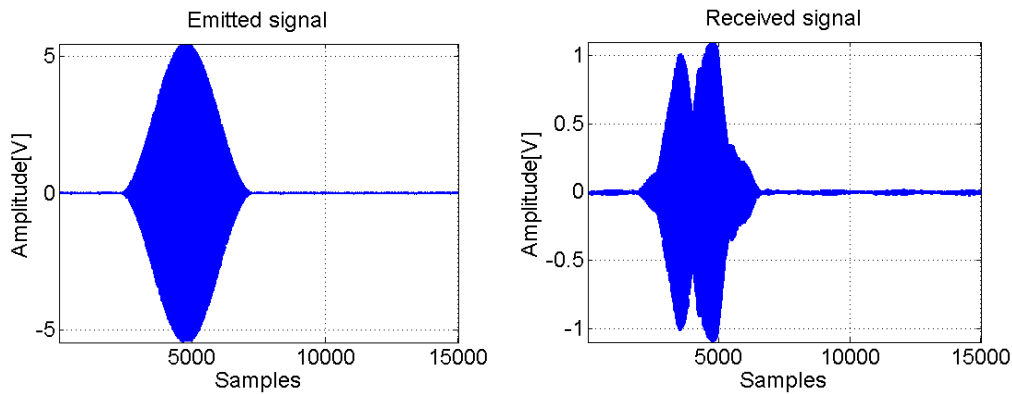


Fig. 1-7: Example of signals used for the object tracking application

The purpose of this application is to find a suitable method for the detection of an underwater object using bi-static acoustic paths and without synchronization between sensors. Actually, in some real application, when the transmitting and the receiving sensors need to be far, in order to cover a wider zone, the synchronization becomes an important issue, especially in the case of immersed sensors. As we will see further, the theoretical contributions of this thesis can be a potential solution in this context.

1.1.2. Passive configurations

Measuring pressure transient in water pipes

Pressurized pipes sometimes present unsteady flows which happen because of a sudden hydraulic change in the system (eg. closing of a valve). Among the effects of these changes, a commonly met phenomenon is the water hammer which usually is caused by a suddenly closed valve (Hachem & Schleiss, 2012). Thus, the fluid changes its direction or stops its motion. The major and dramatic effect of the water hammer phenomenon is that the system exhibits a pressure variation which may cause from vibrations to pipe collapse (Fig. 1-8).

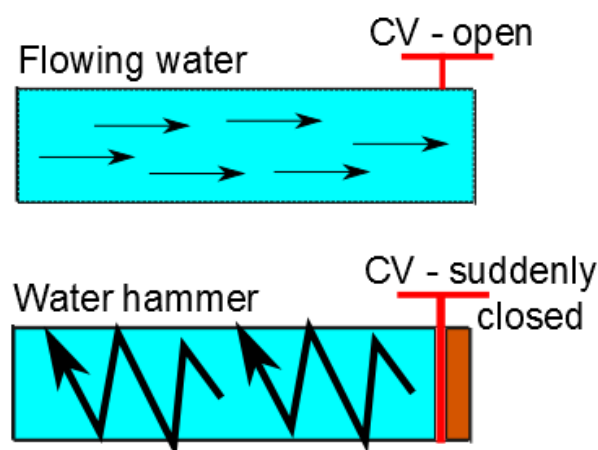


Fig. 1-8 : Experimental setup for the water hammer (CV – closing valve)

The heterogeneity of the environment, in this application, it is caused by the phenomenon itself: the environment is no longer homogenous, because the pressure field is no longer the same. Therefore, the environment is mainly characterized by the water and the

dynamic pressure field (Fig. 1-9). The main purpose in this application is to measure this dynamic pressure field from the interference between the acoustic propagating wave and the field, as illustrated in the next figure.

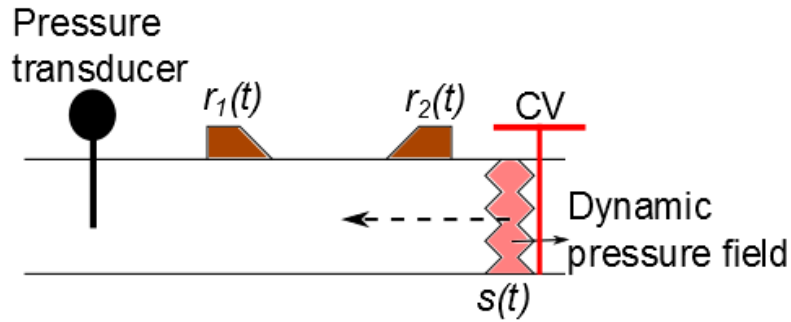


Fig. 1-9: Signal's propagation in a heterogeneous pressure field

The deformation operator \mathcal{D}_k becomes:

$$\mathcal{D}_k \{s(t)\} = \alpha_k \cdot A \cdot \exp \left(j \cdot \phi \left(t - \frac{L_k}{c_{pressure}} \right) \right) \quad (1.9)$$

where $c_{pressure}$ is the propagation speed of the pressure field and L_k is the distance between the CV and the ultrasonic transducers.

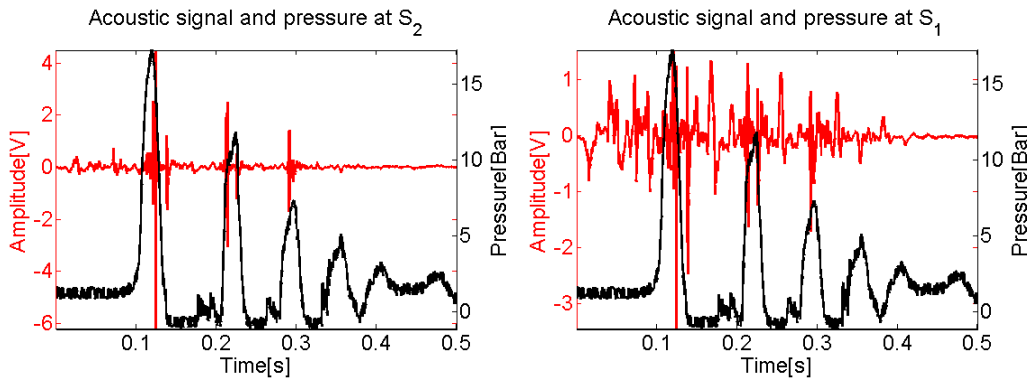


Fig. 1-10: The acoustic signals recorded by the ultrasonic transducers and the pressure recorded by the intrusive pressure transducer

This application has the purpose to evaluate, using a non-intrusive method, the water hammer effect, namely the pressure variation inside the pipeline. This can be achieved based on an accurate estimation of the pressure field propagation speed (Houghtalen, Akan, & Hwang, 2010) which is almost the same as the speed of the acoustic wave, Fig. 1-10.

The advantage of this approach is that the ultrasonic transceivers have just to be placed on the pipe without requiring any extra modification of the system and without direct contact with the fluid.

Electrical arcs

Another case of sensing in heterogeneous environment is the surveillance of electrical power systems - in the presence of transient phenomena. The surveillance of photovoltaic system in the presence of arc faults is the main application that is considered in our work.

One possible solution for the continuous surveillance of an equipment, in order to detect in real time the generation of faults such as electrical arcs (EA), is to listen, using a hybrid electromagnetic/acoustic array, the perimeter of interest (Fig. 1-11).

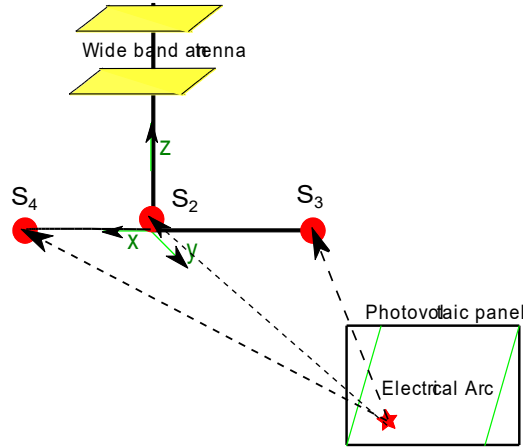


Fig. 1-11 : Experimental configuration for the detection and the localization of the electrical arcs

The particularity of this application is that, from acoustic point of view, the EA (discharges of air particles) creates different acoustic signatures, according to the direction of arrival with respect of each element of the acoustic array.

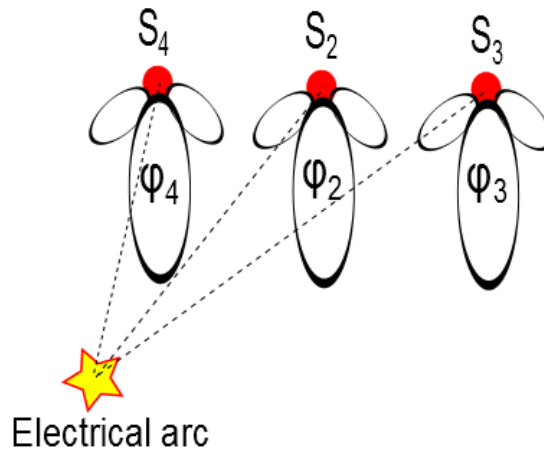


Fig. 1-12: Signal's propagation for a discharge of air particle

The heterogeneity is caused by the different angles of arrival, φ_k , at the sensing system, Fig. 1-12. The deformation operator can be written as:

$$\mathcal{D}_k \{s(t)\} = \alpha_k \cdot A \cdot \exp \left(j \cdot \phi \left(t - \frac{L_k}{c} \cos(\varphi_k) \right) \right) \quad (1.10)$$

Moreover, the use of the wide band antenna, namely the electromagnetic propagation mode accentuates the heterogeneity of the propagation of the transient electrical arc to the acoustic array (Fig. 1-13).

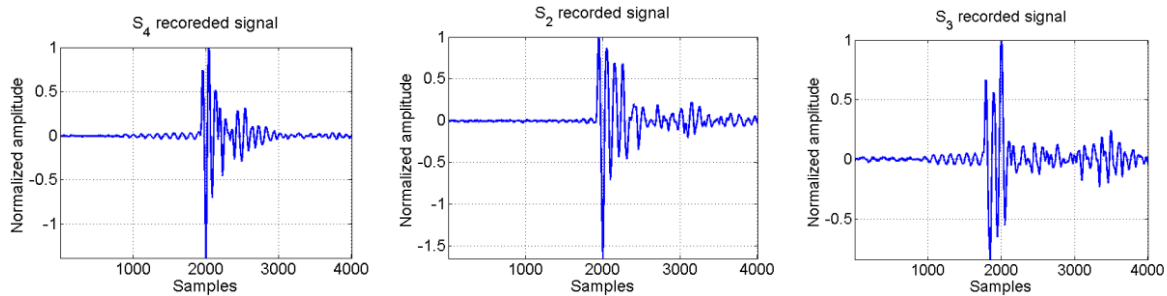


Fig. 1-13: Examples of electrical arcs recorded with the acoustic sensors

The applicative contexts briefly presented in this part are just few examples of problems consisting of the estimation of physical parameters in a heterogeneous propagation and dynamic environment. The main objective of the thesis is to propose new methods for signal's analysis that could provide, with respect of the state of the art, an interesting alternative for the sensing techniques required for these applications.

1.2. Thesis outline

Chapter 2 presents a state of the art of the techniques used in acoustic sensing field, from signal processing perspective, in a dynamic configuration. Then, we present a data-driven analysis technique based on the phase diagram concept that constitutes an interesting representation space for the problems we propose to solve.

Chapter 3 describes a new concept for the active sensing. The major part of this principle is the emission which is a phase-space-based waveform that presents a unique shape in the phase space and which can be easily translated into a signal. We start from the desired phase-space-based waveform and then we move backwards, to the time series concept in order to have a 1D signal which contains the desired information. This waveform design technique presents the interest to create separable waveforms, even if they share the same time-frequency plan. The analysis of such signals, using instantaneous frequency law tracking in phase diagram domain, is then proposed as the processing done at the receiving level.

Chapter 4 focuses on the analysis of the signals using new descriptors defined in the phase diagram/RPA domain. Firstly, the time-distributed recurrence measure is described and, then, used for the detection of transient signals. Next, the multi-lag phase analysis is introduced and applied on synthetic signals with close parameters in order to discriminate them. Finally, the diagonal lines quantification is presented as applied on modulated signals.

Chapter 5 presents the experimental results using proposed concepts previously presented in active configurations: the cavitating environment application is presented and the flow evolution is characterized using the diagonal lines quantification; the underwater object tracking is, then, presented and, with the use of the time-distributed measure, the object is tracked.

Chapter 6 describes the water hammer application and estimates the phenomenon parameters using the RQA descriptors. Then, the electrical arc application is presented for the

detection, localization and characterization of electrical faults using the phase diagram/ RPA concepts.

Finally, the conclusion lays out a synthesis of the thesis, our main contributions and proposes the perspectives for further research works.

2. Thesis work motivation and positioning

The general context of this thesis evolves around the sensing of transient and dynamic phenomena (Gao, et al., 2007), (Serbanescu, Cernaianu, & Ivan, 2009) such as cavitation, pressure transient, etc. as it was discussed in the Chapter 1.

In order to define the positioning of our work in this general context, we consider as an illustrative example, an active sensing system composed by two acoustic transducers: one for emission and the other for reception (Fig. 2-1).

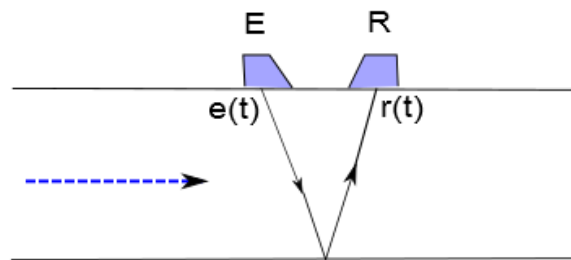


Fig. 2-1: General non-intrusive acoustic sensing system

The requirements for sensing (Emitter/Receiver) system are synthesized in Table 2-1.

	Requirements	Emitted waveform	Received signal processing
1.	Accurate <i>time</i> parameter estimation for flow velocity, object detection and localization	Good robustness face to the propagation phenomena	High resolution and robustness
2.	Accurate <i>spectral</i> parameters	Long duration signals to integrate Doppler effects	- High resolution n spectral domain - Sensing phenomena with weak amplitude
3.	High sped sensing for dynamic (<i>fast</i>) phenomena	Signals with high pulse repetition frequency (PRF)	Capacity of satisfying the first and second requirements

Table 2-1: The requirements for a general sensing system

The purpose of this system is to achieve a non-intrusive characterization of high dynamic phenomena such as the fluid parameters (mono/biphasic). In these applications, one of main parameters to estimate is the time of flight (TOF) which relates us to the flow velocity (Ioana, et al., 2014), as well as the spectral information related to the speed flow profil in bi-

phasic configuration. Such (general) applicative context requires new signal analysis and processing techniques. The contributions proposed by this thesis constitute a potential interesting sensing tool.

2.1. Emission

2.1.1. Narrowband pulses

The active context of the sensing of dynamic phenomena requires a special attention for the emission part. The received signal characterizes better the studied phenomenon if it is well adapted to the propagation environment.

In the domain of acoustic inspection or parameters estimation the main sensing techniques are built around the time of arrival measurements. The simplest way is to measure the travel time between an emitted waveform and the received one. The most used transmitted waveforms are the pulses, defined as:

$$s(t) = A \cdot \exp(j \cdot \omega_0 \cdot t), t \in [kT, kT + T_s], k = \overline{1, N} \quad (2.1)$$

where $\omega_0 = 2 \cdot \pi \cdot f_0$, f_0 is the central frequency of the transducers, T is the pulse period and T_s is the duration of the pulse.

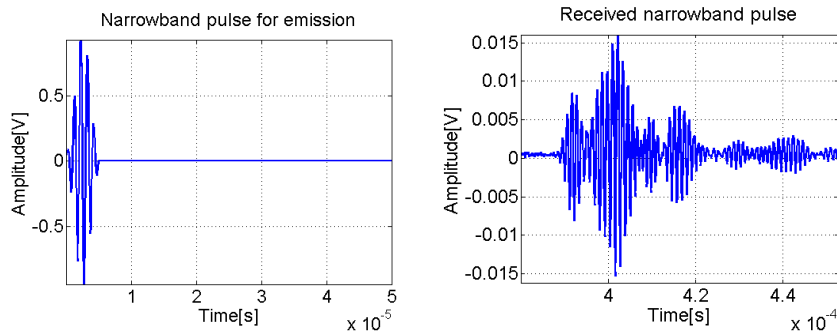


Fig. 2-2: Example of a narrowband pulse used for emission (left) and its received version (right) in a heterogeneous environment

The signals presented in Fig. 2-2 are recorded for an application related to the underwater object detection. The emitted signal is a windowed narrowband signal at 1MHz and a duration of $5 \mu\text{s}$. The configuration is static: both the water and the immersed object are not moving. As it is shown by Fig. 2-2, the received signal is strongly deformed, because of the heterogeneity of the environment - solid-fluid-solid.

This example illustrates the main limitation of narrow band pulses: the amplitude deformation as well as the complex behavior of the propagation environment makes difficult the problem of time of arrival (TOA) estimation.

As an alternative, the wideband signals can be used, starting from the fact that the frequency information is generally more robust than the amplitude's one.

2.1.2. Wideband pulses

The signals with high bandwidth-duration product are well known in radar and sonar applications (Daniels, 1999), thanks to their robustness in complex propagation contexts. The general definition of such signals is:

$$s(t) = A \cdot \exp(j \cdot \phi_{WB}(t)) \quad (2.2)$$

where $\phi_{WB}(t)$ is the instantaneous phase of the signal. From this variable, the Instantaneous Frequency Law (IFL) of the signal is defined as its first derivative:

$$\omega_{WB}(t) = \frac{d}{dt} \phi_{WB}(t) \quad (2.3)$$

The instantaneous phase can be characterized using an analytical polynomial form:

$$\phi_{WB}(t) = \sum_{k=0}^N a_k t^k \quad (2.4)$$

where a_k are the coefficients of the polynomial phase. Depending on these coefficients, the shape of the IFL can be flat ($k=1$), linear ($k=2$), cubic ($k=3$), etc.

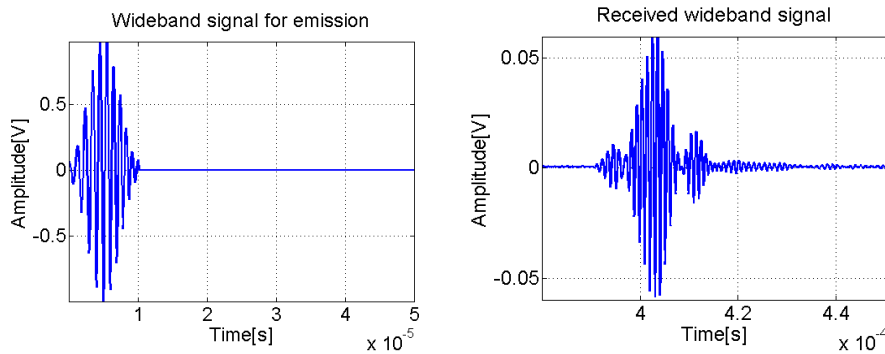


Fig. 2-3: Example of wideband signal used for emission (left) and its received version (right)

In Fig. 2-3, we present the signals emitted and received for the same application as the one conducting to the signal illustrated in Fig. 2-2. The emitted signal is a linear chirp from [800 kHz, 1200 kHz] of 10 μ s duration. Compared with the narrowband received signal, the wideband received one has higher amplitude and less interference, proving its better robustness to the propagation effects.

Still, a trade-off between the bandwidth of the signal and its duration must be considered: a higher frequency resolution implies a longer duration of the signal which usually leads to destructive interferences in multi-path configuration.

In order to achieve good parameters estimation capabilities in such experimental cases, the needing for new modulation concepts is obvious. The solution proposed in this thesis is to design different modulations, in the **same** time-frequency region, but with elements that will allow **their separation** at the reception level. Our contribution is to design these modulations in the phase diagram domain as we will present in the Chapter 3.

2.2. Reception

At the receiver level, either we are in active or passive configuration, the aim of our analysis is to estimate the physical parameters of the received signals. Independently of the waveform type, the received signals are strongly deformed. In order to place our contribution with respect of the state of the art, we will briefly present the major classes of methods traditionally used to estimate the signal's parameters.

2.2.1. Energetic techniques

In active configurations, one of most common analysis techniques is the matched filter (MF) method. Considering r the received signal, then the output of the matched signal is (Van Trees, 2001):

$$y(t) = \int_0^t r(\tau) \cdot h(t - \tau) d\tau \quad (2.5)$$

where t is the duration of the signal to be detected and h is the impulse response of the filter which is equal to the time reversed version of the transmitted signal, s , used as reference. Next, the output of the MF given in eq. (2.5) is compared with a threshold, and, if it is greater than the threshold, the signal s is considered to be present or not.

Multiple communication and physical parameters measurements equipments are based on this principle, especially for the detection part. The limitation of this method is that, in terms of characterization, it does not succeed to highlight similar (but not identical) deformations in the received signal since such deformations are not present in the reference signal, Fig. 2-4.

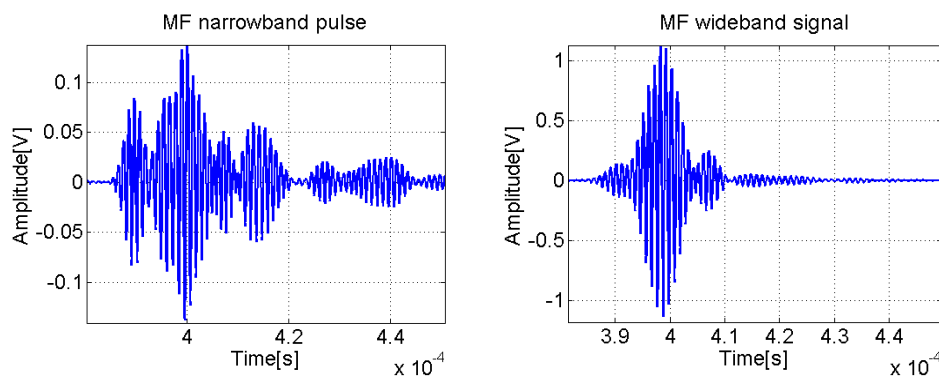


Fig. 2-4: The matched filter applied for the signals presented in Fig. 2-2 (left), respectively in Fig. 2-3 (right)

For this reason, a deeper analysis of the received signals is necessary, using the main techniques presented in the following part.

2.2.2. Projective techniques

The projective techniques are based on the principle of analyzing a signal using its projection on a set of elementary functions from a dictionary which are used as reference for the comparison with the studied signal.

There are various approaches for these techniques, we recall one among the most used technique which, for an adequate dictionary, highlights signal's time and scale evolution, being a powerful tool for the signals' characterization (Mallat, 1999). This is actually the representation using time-scale analysis and, particularly, the wavelet transform analysis.

A wavelet (Mallat, 1999) is a zero average function:

$$\int_{-\infty}^{\infty} \psi(t) dt = 0 \quad (2.6)$$

This function (often called mother wavelet) is dilated with the scale parameter s and translated with the parameter u . Hereby, the orthonormal basis can be built for $s = 2$.

$$\Psi_{s,u}(t) = \frac{1}{\sqrt{s}} \Psi\left(\frac{t-u}{s}\right) \quad (2.7)$$

As an example of well-known wavelets, we mention the families of functions presented in (Mallat, 1999), (Meyer, 1993), (Daubechies, 1992), (Daubechies, 1990), (Auger, et al.).

Therefore, the wavelet transform of a signal $x(t)$ at scale s and time position u is defined as:

$$Wx(u,s) = \frac{1}{\sqrt{s}} \int_{-\infty}^{\infty} x(t) \cdot \Psi^*\left(\frac{t-u}{s}\right) dt \quad (2.8)$$

It goes that a wavelet transform is able to quantify the time-frequency variations, but based on the scale provided by the wavelet basis. Hereby, in order to have a good wavelet representation, it is essential to determine an adequate mother function which resembles as much as possible the analyzed signal, $x(t)$ (Mallat & Zhong, 1992).

In order to illustrate the particularities of wavelet analysis in the case of transient signals having close features, we consider three transients with very close characteristics. These signals are defined by this generic formula:

$$s_{(a,f,b)}[n] = \begin{cases} a \cdot (\sin[2\pi \cdot f \cdot n] + b[n]), & n = \overline{1, N} \\ a \cdot b[n], & otherwise \end{cases} \quad (2.9)$$

where $N = \left\lceil \frac{f_s}{2f} \right\rceil$ and $f_s = 10 \text{ MHz}$ is the sampling frequency.

These signals, $s_{1(a_1, f_1, b_1)}$, $s_{2(a_2, f_2, b_2)}$, respectively $s_{3(a_3, f_3, b_3)}$ have the following relationships between their parameters: $a_1 / a_2 = 1 / 0.6 \cong 1.66$, $a_1 / a_3 = 1 / 0.3 \cong 3.33$, $f_1 / f_2 = (2 \cdot 10^5) / (1.9 \cdot 10^5) \cong 1.05$, $f_1 / f_3 = (2 \cdot 10^5) / (1.6 \cdot 10^5) = 1.25$, $b_1 = b_2 = b_3$ is an additive white Gaussian noise and $SNR_i = 20 \text{ dB}$, $i = \overline{1, 3}$. Fig. 2-5 presents the signals considered in eq. (2.9).

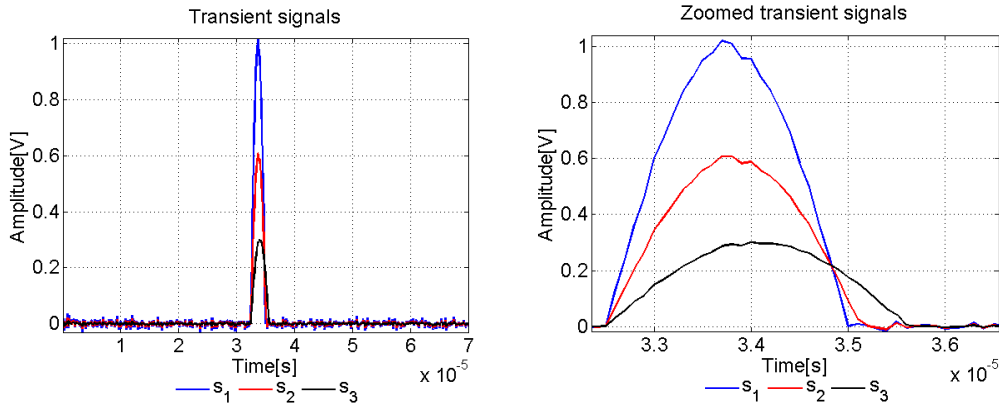


Fig. 2-5: Transient signals considered for studying the wavelet transform

These synthetic signals seem to be quite similar, but, at a closer look, slight differences appear. The results obtained with the wavelet analysis are shown in Fig. 2-6.

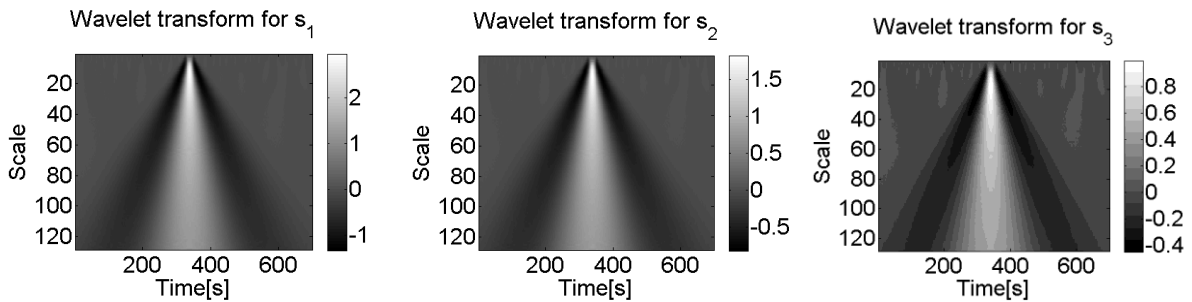


Fig. 2-6: The wavelet transform (using the Mexican Hat mother-function) applied on the three synthetic signals

It can be observed that the slight differences between the signals cannot be highlighted by the wavelet transform: their presence is detected and located in time, but their shape does not present any discriminating element (beside the amplitude differences).

Once again, even if the problem of detection and localization is solved (although, it frequently requires the managing of a large dictionary), the problem of signal characterization and discrimination between signals is present when it comes to analyze signals with close characteristics.

2.2.3. Data-driven techniques

While the projective techniques require the definition of the elementary functions, the interest of data-driven techniques is clear when this definition is not straightforward.

These techniques have the major advantage that they do not use any model to analyze the data. They analyze the signal based on the organization of samples in time. One of the current data-driven analysis techniques is based on the concept of Recurrence Plot Analysis (RPA) which has been firstly introduced by Eckmann et al. (Eckmann, Kamphorst, & Ruelle, 1987).

This concept is based on the phase space representation and, as its name suggests, it highlights the recurrences of trajectories from higher-dimension phase spaces. The term of “recurrence” - one main property of conservative dynamic systems (Zbilut & Webber Jr.,

1992), (Webber Jr. & Zbilut, 1994), (Webber Jr. & Zbilut, 2005), (Zbilut & Webber Jr., 2007), (Marwan, 2008), (Marwan, Schinkel, & Kurths, 2013) – means that the dynamical system under study returns in a state previously visited. The method is quantified using Recurrence Quantification Analysis (RQA).

The choice of the RQA based on RPA concept for the analysis of signals coming from heterogeneous environments is based on the fact that it is a data-driven method which does not require a priori information about the system, knowing that such information is not always available (Zbilut & Webber Jr., 2006), (Marwan, et al., 2007).

The recurrence information is very important, offering us new insights in the evolution of dynamic systems. In our work, we are interested of the system state changes that are not determined by random causes, but they are the results of a nonlinear input that causes them to change their state suddenly, exposing the system to major collapse.

The phase diagram concept starts by considering the following time series:

$$\mathbf{x} = \{x[1], x[2], \dots, x[N]\} \quad (2.10)$$

Then, this time series is represented in phase space. Its values become the coordinates of the m -dimensional space and, consequently, the vector sample is:

$$\bar{v}_i = \sum_{k=1}^m x[i + (k-1)d] \cdot \bar{e}_k, \quad i = 1, M \quad (2.11)$$

where m is the embedding dimension, d is the delay, \bar{e}_k is the unit vector of the axis that define the phase space, $M = N - (m-1)d$ and N is the length of the time series. Usually, the embedding dimension and the delay are chosen using the false nearest neighbor method and the mutual information method, respectively (Marwan, et al., 2007), (Kantz & Schreiber, 1997), (Thiel, et al., 2002), (Thiel, et al., 2004).

Fig. 2-7 illustrates the phase space construction algorithm. Each point of the phase space trajectory has as coordinates the time series values, namely signal's samples.

Next, the distances between the vectors in the phase space are represented on the distance matrix (DM), eq. (2.12). When compared with a threshold, the recurrence matrix (RM) is obtained, eq. (2.13).

$$D_{i,j} = \mathcal{D}(v_i, v_j) \quad (2.12)$$

$$R_{i,j} = \Theta(\varepsilon(i) - \mathcal{D}(v_i, v_j)) \quad (2.13)$$

where $\mathcal{D}(\cdot, \cdot)$ is a distance applied on the vectors from the phase space (Euclidean distance (Marwan, et al., 2007), (Marwan & Kurths, 2002), L1 norm (Kantz & Schreiber, 1997), angular distance (Birleanu, et al., 2012), (Ioana, et al., 2014), scalar dot product distance (Birleanu, et al., 2012), etc.), $\Theta(\cdot)$ is the Heaviside step function and $\varepsilon(i)$ is the threshold considered for recurrence. A graphical representation is presented in Fig. 2-8 when using the Euclidean distance.

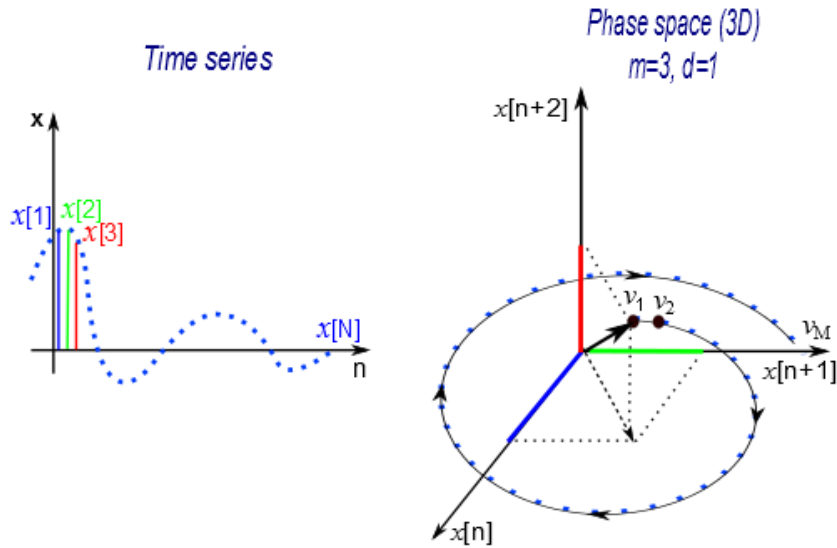


Fig. 2-7: Example of phase space trajectory implementation (Ioana, et al., 2014)

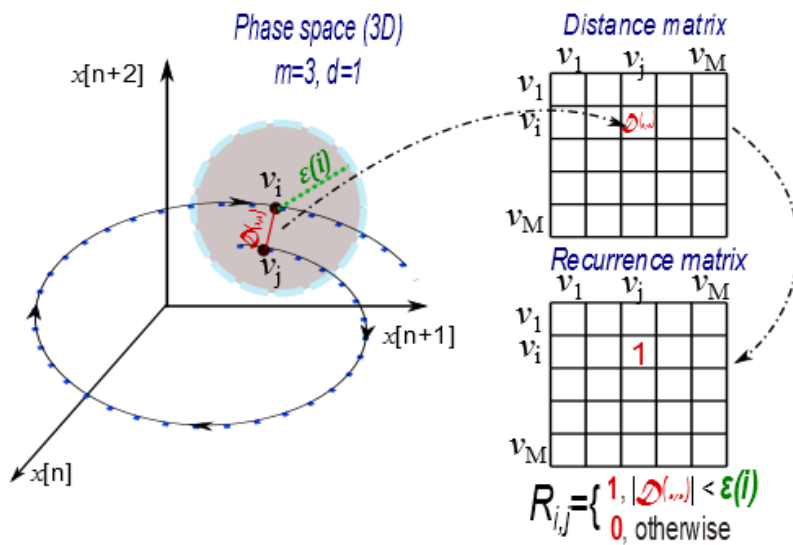


Fig. 2-8: DM and RM when applying the Euclidean distance (Ioana, et al., 2014)

Although the two methods, RPA and wavelet analysis, are based on totally different concepts (the recurrence of the system's trajectory in the phase space and the time-frequency representation, respectively), they are both able to highlight similar behaviors of a process like trends, sudden changes, discontinuities (Webber Jr. & Zbilut, 1994), (Popescu, et al., 2014), (Yang, 2011), (Chen & Yang, 2012), (Ramirez Avila, et al., 2013), (Webber Jr. & Zbilut, 2005), (Zbilut & Webber Jr., 2006), (Zbilut & Webber Jr., 2007), (Webber Jr., 2012), (Webber Jr., 2013), (Farge, 1992), (Torrence & Compo, 1998), (Digulescu, et al., 2013), (Larose, et al., 2015). In (Digulescu, et al., 2014c), several comparisons between the two methods along with other widely used signal processing methods (as matched filter, spectrogram, zero crossing) are shown.

The major advantage of the RPA method is that it does not need any orthogonal basis or any a priori information about the studied system (unlike the wavelet analysis: it requires the definition of the analyzing mother wavelet which must be close to parts of the signal to analyze). The method works in applications where a reference signal is not available (unlike

the matched filter method) (Digulescu, et al., 2014c), (Ioana, et al., 2014), (Digulescu, et al., 2016c), (Le Bot, Gervaise, & Mars, 2016).

In terms of RQA, we recall the most known measures as the recurrence rate, the determinism, the laminarity, etc. (Zbilut & Webber Jr., 1992), (Webber Jr. & Zbilut, 1994), (Webber Jr. & Zbilut, 2005), (Webber Jr., 2012), (Webber Jr., 2013), (Zbilut & Webber Jr., 2006), (Zbilut & Webber Jr., 2007) which measures the states of the system based on the RM.

2.3. Chapter summary

This chapter shortly presents the existing methods for the analysis of transient signals in terms of detection and characterization.

These methods present many advantages, but they also have limitations. In the next chapters, we present our approach based on the concept of phase diagram and the RQA based on the RPA method. For the active sensing part, we propose the use of phase space-based adaptive waveforms which are better suited for the propagation phenomena specific to the sensing of transient signals. The passive sensing part presents a new approach for the analysis and characterization of transient signals based on the concept of RQA.

3. Active sensing: phase diagram-based adaptive waveform

3.1. Introduction

The active sensing techniques are the heart of any system aimed to analyze physical parameters of an environment of interest, either by acoustic or electromagnetic waves. The sensing signal is always designed according to the application specificities: propagation phenomena, objects to detect, parameters to sense, etc. This is always the case in sonar/ radar domains, where the transmitted signal's parameters are chosen from physical considerations:

- the central frequency and the bandwidth are chosen to minimize the wave attenuation effect as well as to best fit to the sensors characteristics
- the signal's duration is chosen to meet the required resolution

Another parameter of great importance is the PRF (Pulse Repetition Frequency) which is related to the maximum range of sensed region. The classical approach considers the principle presented in Fig. 3-1.

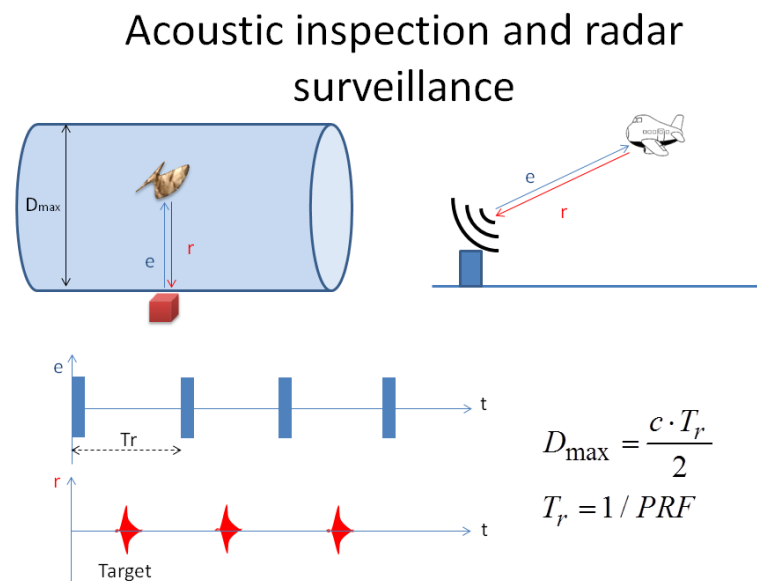


Fig. 3-1: General principle in radar and acoustic sensing

From this perspective, the environment inspection is done by respecting the relation between the maximum distance and the PRF.

Still, for sudden phenomena or fast changes, this approach could skip them, the PRF being too small to detect these transient phenomena. An example is presented in Fig. 3-2.

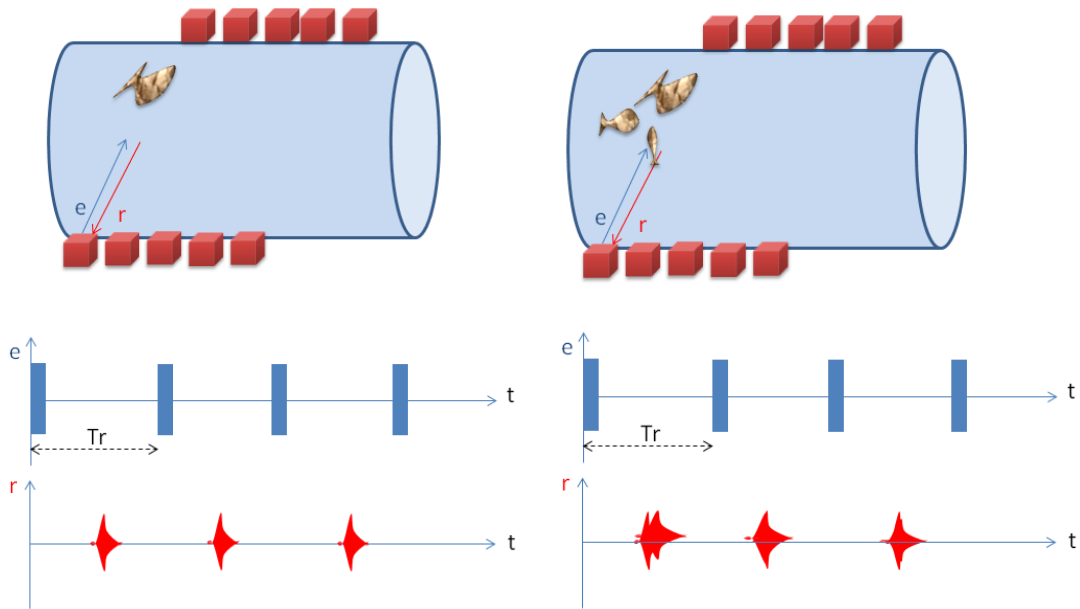


Fig. 3-2: Favorable scenario using the classical approach for the acoustic inspection (left) and de unfavorable scenario using the same approach (right)

In the first case (Fig. 3-2, left), the detection of the object passing in front of the sensors is done from the reflected wave, the range being calculated as a product between the wave velocity and the half of time of flight. With this technique, the characterization of the object is almost impossible as illustrated in the Fig. 3-2 (right). Two different objects are traduced, at the receiving sensor, by more or less similar reflected waves, making difficult their identification.

One solution is then to increase the PRF in order to allow a high speed sensing, as it is illustrated in Fig. 3-3.

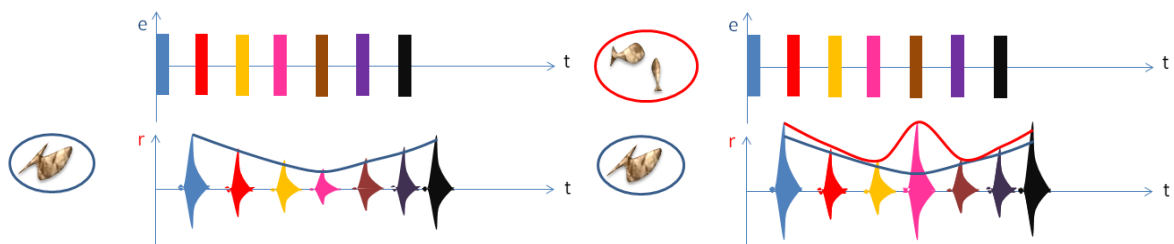


Fig. 3-3 : Our proposed approach for the high-speed sensing

These signals contain some unique information that allow us to increase the PRF disregarding the classical condition, in order to monitor the environment at higher sensing speed. The main difficulty is to add a unique identifier knowing that the amplitude or phase changes could interfere, from one emitted pulse to another, making difficult their separation at the receiving point. The appropriate design of the transmitted pulses is then the main objective of our contribution, presented in the next section.

3.2. Phase diagram-based waveform design

One way to generate separable pulses is to design them in diagram phase, in **different regions** which bring them “orthogonal” in this representation space.

Fig. 3-4 presents the overall diagram of the algorithm.

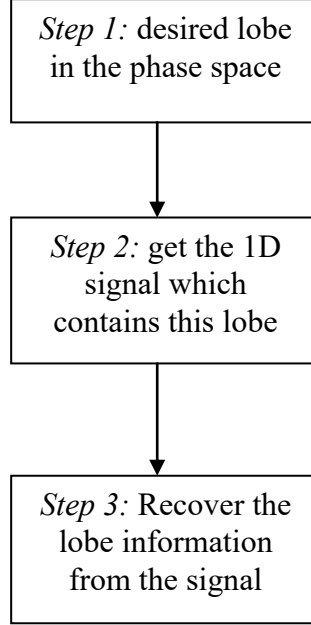


Fig. 3-4: The overall diagram of the proposed algorithm of waveform design in phase space

3.3. Emission

The starting point is the idea to design a pattern in the phase space which has a unique evolution and, afterwards, to translate it in a 1D signal into a one-to-one correspondence.

Step 1: we represent a lobe in a 3D phase diagram space. The core of this idea is suggested by the trajectory of a transient signal in phase space (Candel, et al., 2012a), (Candel, et al., 2012b).

Hereby, the simplest way to represent the waveform’s trajectory in shape of a lobe is to compose it of two director vectors, as in Fig. 3-5:

$$\vec{v}_n = a_n \cdot \vec{i} + b_n \cdot \vec{j} + c_n \cdot \vec{k}, n \in \{1,2\}, \forall a_n, b_n, c_n > 0$$
$$\vec{v}_n(a_n, b_n, c_n) \tag{3.1}$$

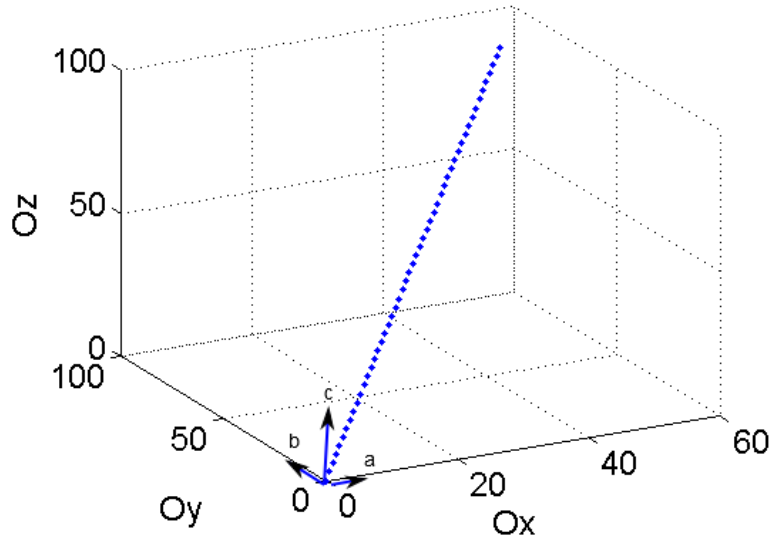


Fig. 3-5: Director vector example in 51 points: $a = 1, b = 1.5, c = 2$

For *step 2*, at a first look, the signal should be reconstructed based on the concept of phase space via Takens theorem (Takens, 1981). Using this approach, the coordinates of each point from phase space are the delayed samples of the signal.

This assumption leads us to the conclusion that $a_n \approx b_n \approx c_n$. For that matter, this contradicts our objective of creating a lobe anywhere in the first quadrant ($\forall a_n, b_n, c_n > 0$). Hereby, this approach is not feasible and the concept of phase space reconstruction must be reconsidered (without losing the meaning of system's evolution).

Let's consider that phase space lobe is characterized by the set of points $P_n(x_n, y_n, z_n)$, $x_n, y_n, z_n \in \mathbb{R}$ and $P_n \subset \mathbb{R}^3$. We define the 1D signal, called *auxiliary signal*, as follows:

$$\forall x_n, y_n, z_n \in \mathbb{R} \quad \mathbb{R} \quad \{x_1, \dots, x_N, y_1, y_2, \dots, y_N, z_1, z_2, \dots, z_N\} \quad (3.2)$$

Combining eq.(3.2) with eq.(3.1), it results in a phase space representation with a constant delay: $d = N$

The representation from eq.(3.2) is oscillations free for each part (x part, y part, z part). Using eq. (3.1) to identify the terms from eq.(3.2), a determined law appears to control signal's evolution:

$$\begin{cases} x_i = a_1 \cdot t_i, i = \overline{1, k} \text{ and } x_j = -a_2 \cdot t_j, j = \overline{k+1, N} \\ y_i = b_1 \cdot t_i, i = \overline{1, k} \text{ and } y_j = -b_2 \cdot t_j, j = \overline{k+1, N}, a_n, b_n, c_n > 0, n = \{1, 2\} \\ z_i = c_1 \cdot t_i, i = \overline{1, k} \text{ and } z_j = -c_2 \cdot t_j, j = \overline{k+1, N} \end{cases} \quad (3.3)$$

where k is the length of the first director vector and $N - k$ is the length of the second director vector. In other words, $P_i(x_i, y_i, z_i)$ are the points belonging to \vec{v}_1 , respectively $P_j(x_j, y_j, z_j)$ are the points belonging to \vec{v}_2 .

We recall that in *step 1*, it was imposed that $a_n, b_n, c_n > 0, n = \{1, 2\}$. Still, in eq. (3.3), the coordinates belonging to $\overline{v_2}$ appear to have negative coefficients. This is caused by two conditions that our lobe has to meet:

- The lobe must return to its original starting point
- The signal belongs to a continuous measurement. Therefore, we must accomplish the sampling condition: t_i and t_j are, in fact, the sampling moments of the signal with the sampling period T_s .

$$\begin{aligned} t_i &= t_0 + (i-1)T_s, \quad i = \overline{1, k} \\ t_j &= t_0 + (j-1)T_s, \quad j = \overline{k+1, N} \\ &= [t_0 + (k-1)T_s] + (j-k)T_s \end{aligned} \quad (3.4)$$

Based on the observations presented above, in Fig. 3-6, we illustrate such signal.

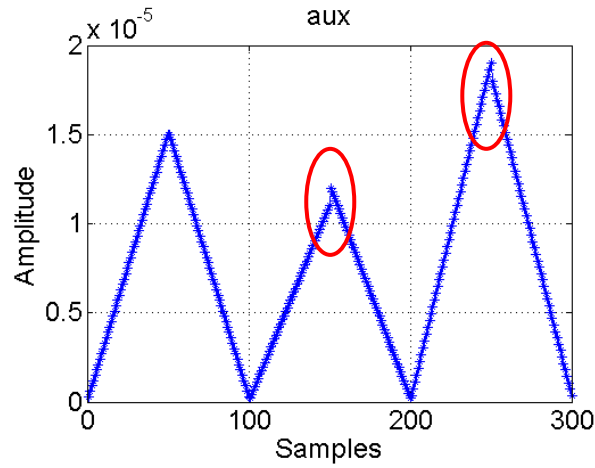


Fig. 3-6: The auxiliary signal **aux** with the following parameters: $a_1 = 1.5, b_1 = 1.1, c_1 = 1.9$,
 $a_2 = 1.5, b_2 = 1.2, c_2 = 1.8$, $k = 50, N = 100, f_s = 5 \text{ MHz}$ and $t_0 = 1/f_s = T_s$

Considering that $P_n(x_n, y_n, z_n), x_n, y_n, z_n \in \mathbb{R}$ is the evolution of the **aux** signal in the phase space, its trajectory is easily obtained (Fig. 3-7).

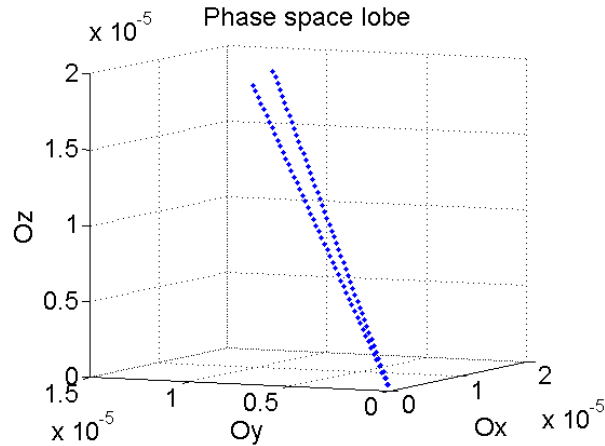


Fig. 3-7: The phase space lobe obtained using the signal **aux** from Fig. 3-6

The phase space lobe is obtained using the eq. (3.1), but it presents two limitations:

- the lobe is too tight, making this representation sensitive even to low levels of noise
- the corresponding signal of this phase space lobe presents discontinuities/ breaks in its evolution when the coefficient changes; in Fig. 3-6, $a_1 = a_2$, the signal is continuous, but $b_1 \neq b_2$ and $c_1 \neq c_2$, the signal is discontinuous when the coefficient changes (the effect is highlighted by the red circles)

In order to obtain a larger lobe, we impose different number of points for the two director vectors that create the phase space, so, the eq. (3.3) becomes:

$$\begin{cases} x_i = a_1 \cdot t_i, i = \overline{1, n_1} & \text{and} & x_j = -a_2 \cdot t_j, j = \overline{n_1 + 1, N} \\ y_i = b_1 \cdot t_i, i = \overline{1, n_2} & \text{and} & y_j = -b_2 \cdot t_j, j = \overline{n_2 + 1, N}, a_n, b_n, c_n > 0, n = \{1, 2\} \\ z_i = c_1 \cdot t_i, i = \overline{1, n_3} & \text{and} & z_j = -c_2 \cdot t_j, j = \overline{n_3 + 1, N} \end{cases} \quad (3.5)$$

The discontinuity issue of the **aux** signal corresponding to the phase space lobe is solved by imposing the continuity condition for the descending slopes from the **aux** signal.

$$\begin{cases} x_1[n_1] = x_2[n_1 + 1] \\ y_1[n_2] = y_2[n_2 + 1] \\ z_1[n_3] = z_2[n_3 + 1] \end{cases} \quad (3.6)$$

In fact, the slopes of the lines from the **aux** signal are the coefficients of \vec{v}_2 . Using the continuity condition, eq. (3.6) and considering that the lobe returns to its original point, meaning that $x_1[1] = x_2[N] = 0, y_1[1] = y_2[N] = 0, z_1[1] = z_2[N] = 0$, it results in a relationship between the two director vectors:

$$\begin{cases} a_2 = a_1 \cdot n_1 / (N - n_1) \\ b_2 = b_1 \cdot n_2 / (N - n_2) \\ c_2 = c_1 \cdot n_3 / (N - n_3) \end{cases} \quad (3.7)$$

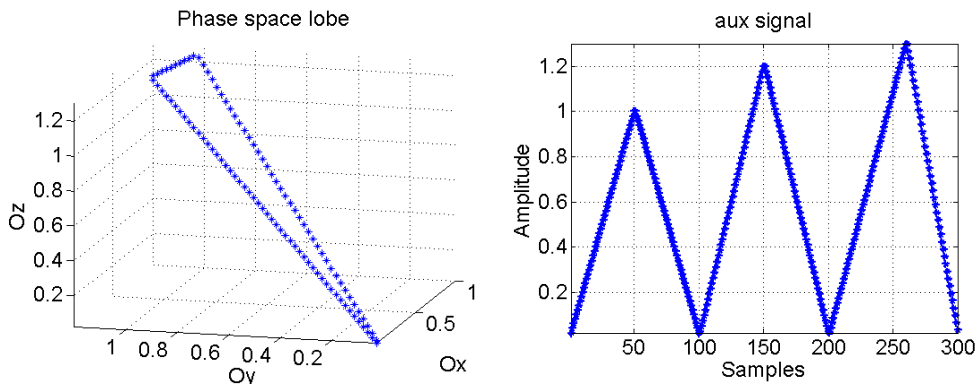


Fig. 3-8: The phase space lobe and its corresponding **aux** signal based on eq. (3.5) and (3.7) with the following parameters: $N = 100, a_1 = 1, b_1 = 1.2, c_1 = 1.3, n_1 = 50, n_2 = 50, n_3 = 60$

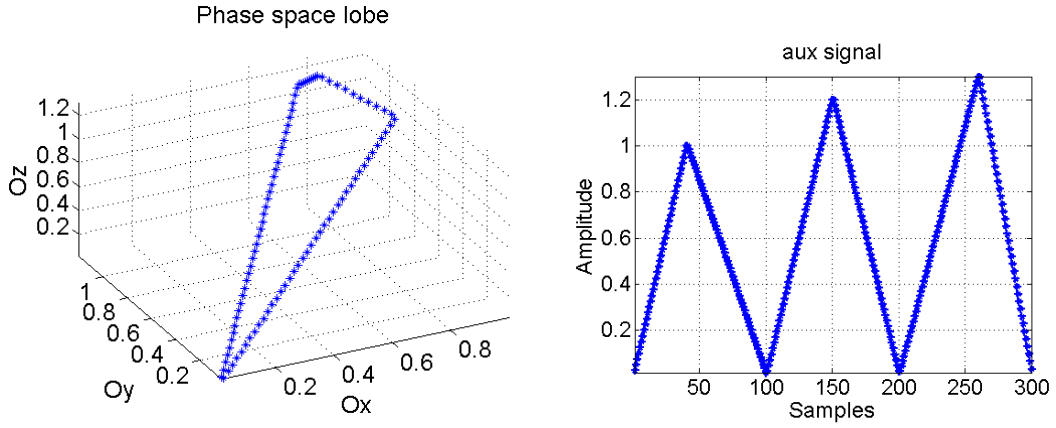


Fig. 3-9: The phase space lobe and its corresponding **aux** signal based on eq. (3.5) and (3.7) with the following parameters: $N = 100$, $a_1 = 1$, $b_1 = 1.2$, $c_1 = 1.3$, $n_1 = 40$, $n_2 = 50$, $n_3 = 60$

From Fig. 3-8 and Fig. 3-9, it can be observed that the parameters n_1, n_2 and n_3 reshape the phase space lobe such that it becomes larger and: it contains an extra vector (line) when two of them are equal (Fig. 3-8), or it contains two extra vectors (lines) when all three of them are different (Fig. 3-9).

3.4. Reception

Step 3: Recover the lobe information from the signal

Our first attempt was to transmit the phase space lobe into an amplitude modulation. Therefore, we define two auxiliary signals with the following parameters that are translated in amplitude modulations, Fig. 3-10:

$$\begin{aligned}
 aux_1 : & \begin{cases} a_1 = 1, b_1 = 1.5, c_1 = 2 \\ n_1 = 400, n_2 = 500, n_3 = 600, N = 1000 \end{cases} \\
 aux_2 : & \begin{cases} a_1 = 2, b_1 = 1, c_1 = 1.4 \\ n_1 = 400, n_2 = 500, n_3 = 600, N = 1000 \end{cases}
 \end{aligned} \tag{3.8}$$

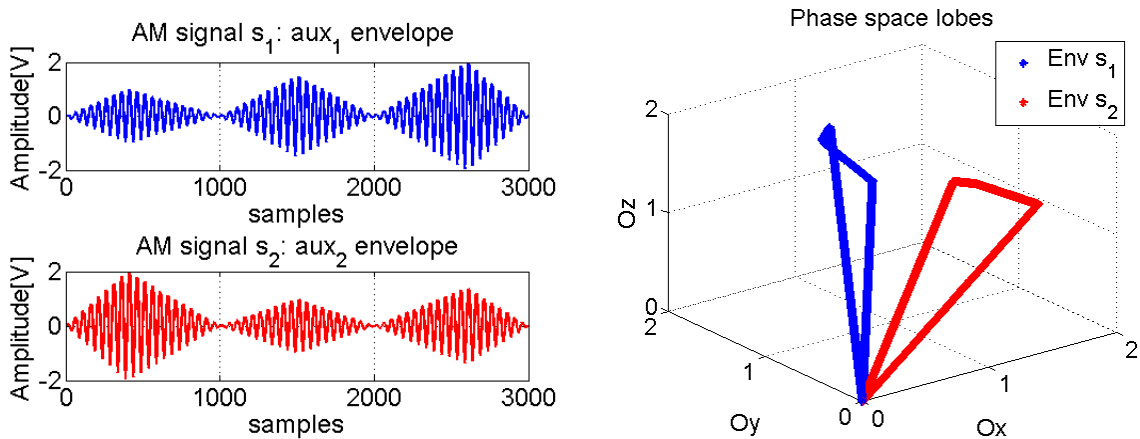


Fig. 3-10: The amplitude modulated signals (left) which have as envelope the auxiliary signal defined in eq. (3.8) shown in phase space representation (right)

At reception, the information held in signal's envelope is extracted in conditions free of noise and at different signal-to noise ratios (SNR), Fig. 3-11.

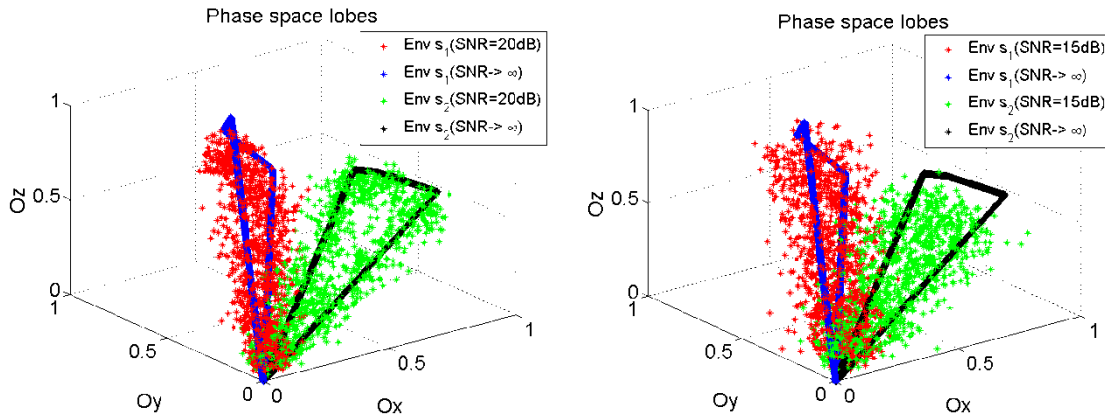


Fig. 3-11: The recovered signal's envelopes (normalized values) for: $SNR \rightarrow \infty$ and $SNR = 20dB$ (left); $SNR \rightarrow \infty$ and $SNR = 15dB$ (right);

Although the new representation space provides a new perspective in terms of adaptive waveform for active configurations, the modulation amplitude approach does not offer reliable results in real propagation conditions. It is obvious that even at 20 dB is hard to distinguish between the two lobes. Therefore, the alternative way is to introduce the lobe information in the frequency's evolution

Firstly, the frequency range must be established. In this new interpretation of phase diagram trajectories, the coefficients a_1, b_1 and c_1 become the maximum frequencies that the signals reach in time. In other words, as the auxiliary signals (that create the phase space lobes) contain only straight lines which vary in time, the modulated signal will contain just linear frequency modulations (LFM) that follow the trend given by the auxiliary signals. Hereby, these signals become from now the *IFLs* of the signals.

After defining the IFLs using the same algorithm as shown in eq. (3.5) and (3.7), this set is composed of LFM's which vary in the same time-frequency region in different manners, but their trajectory in phase space is UNIQUELY identifiable..

As an example, we consider a central frequency of 1 MHz and in the range of [800 kHz, 1200 kHz]. We define the auxiliary signals where the coefficients a_1, b_1 and c_1 correspond to the frequency parameters.

$$\begin{aligned} aux_1 : & \begin{cases} a_1 = 1 \cdot 10^6, b_1 = 1 \cdot 10^6, c_1 = 1.2 \cdot 10^6 \\ n_1 = 0.4N, n_2 = 0.5N, n_3 = 0.6N, N = 1000 \end{cases} \\ aux_2 : & \begin{cases} a_1 = 1.2 \cdot 10^6, b_1 = 1 \cdot 10^6, c_1 = 1 \cdot 10^6 \\ n_1 = 0.4N, n_2 = 0.5N, n_3 = 0.6N, N = 1000 \end{cases} \end{aligned} \quad (3.9)$$

The modulated signals with the IFLs defined in eq. (3.9) are presented in Fig. 3-12, as well as their spectrogram representation, Fig. 3-13.

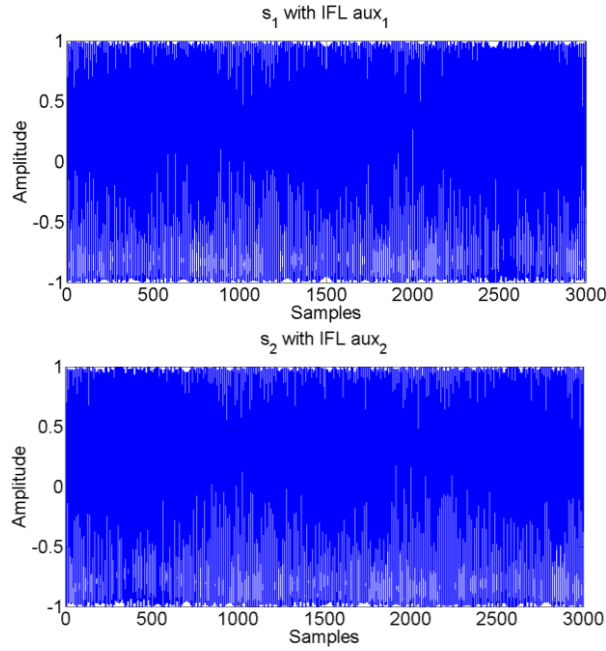


Fig. 3-12: The frequency modulated signals whose IFLs are auxiliary signals from eq. (3.9)

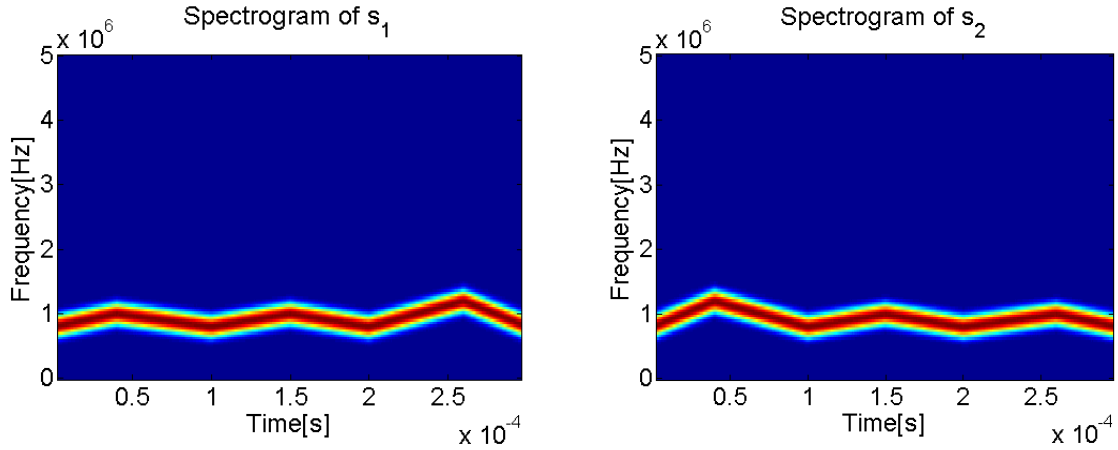


Fig. 3-13: The spectrogram of the modulated signals s_1 and s_2 from Fig. 3-12

The IFL spectrogram-based estimation (Fig. 3-14) determines the transmitted IFL in steps depending on the parameters of the spectrogram (given by the time-frequency resolution).

$$IFL_e = \max_f (|STFT(\tau, f)|) \quad (3.10)$$

where IFL_e is the estimated IFL and $STFT$ is the Short Time Fourier Transform.

Hereby, a linear fitting must be done by storing the moments (and their corresponding frequencies) when the spectrogram-based estimated frequency changes its value. Then, these stored values are linearly fitted. In order to have the same number of points as the original IFL, an interpolation is finally done on the fitted IFL.

$$IFL_f[i] = IFL_e[i_1] + \frac{IFL_e[i_2] - IFL_e[i_1]}{i_2 - i_1} i_1 \quad (3.11)$$

where IFL_f is the fitted IFL and IFL_e is the spectrogram estimated IFL (eq. (3.10)) where i_1 and i_2 are the time bins where the value of IFL_e changes.

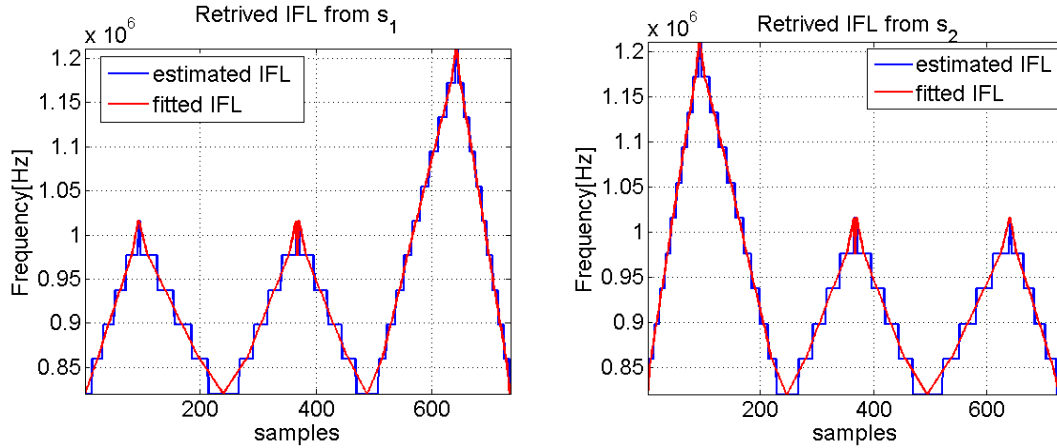


Fig. 3-14: The extracted IFLs using the spectrogram for the modulated signals s_1 (left) and s_2 (right)

In Fig. 3-15, the fitted IFL and the original IFL are plotted in the phase space.

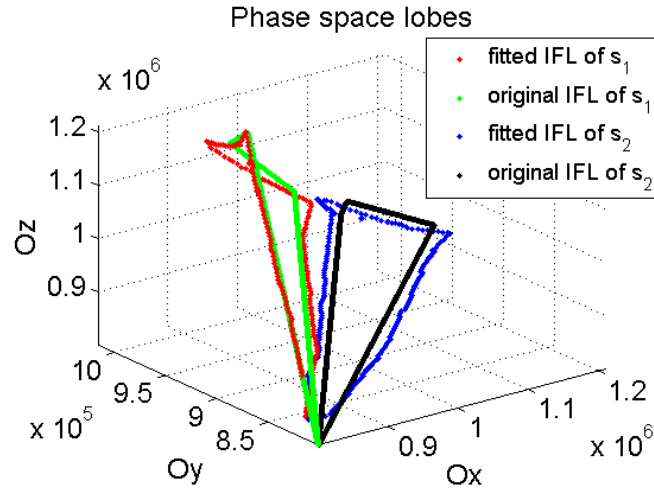


Fig. 3-15: The phase space lobes provided by the IFLs from s_1 and s_2

As shown in Fig. 3-15, the IFL shape is preserved in the phase space domain. Moreover, the phase space provides a new representation space for the discrimination of these signals.

This new waveform design concept has a great advantage in terms of separation of signals. In time and frequency domain, the new modulations are not separable (Fig. 3-16) as it is shown by their correlation and spectrum.

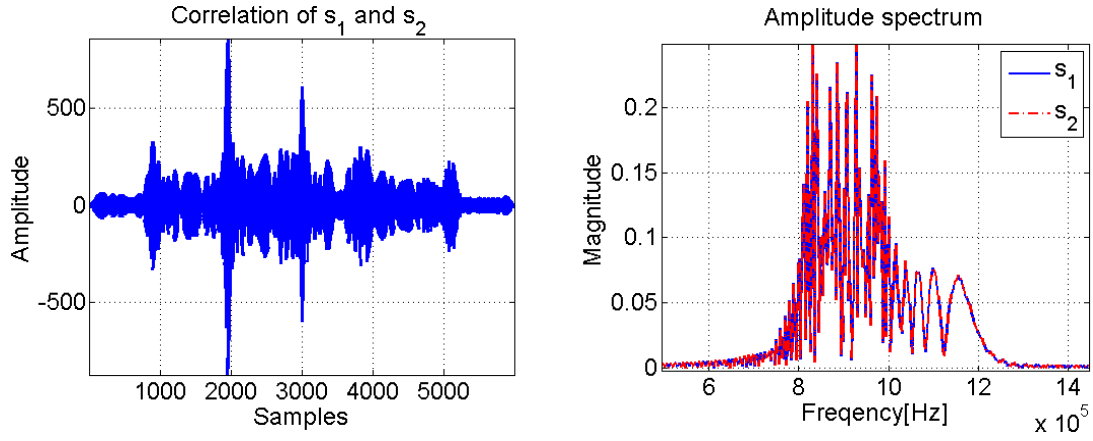


Fig. 3-16: The correlation of the modulated signals s_1 and s_2 (left) and their amplitude spectrum (right)

Obviously, the auxiliary signals also cannot be discriminated from time/frequency domain, Fig. 3-17.

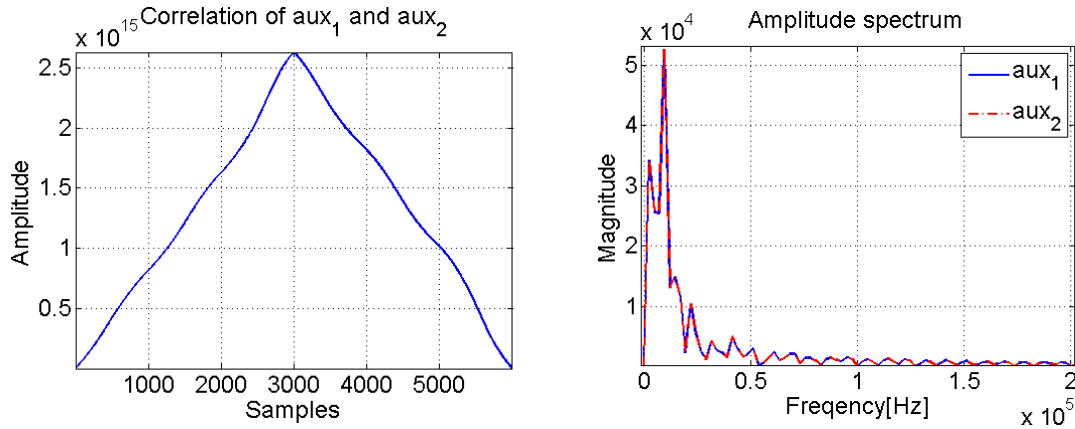


Fig. 3-17: The correlation of the auxiliary signals aux_1 and aux_2 (left) and their amplitude spectrum (right)

The capacity of phase diagram to encode separable signals is then used in order to generate multiple separable signals, that are overlap in time and in frequency but still separable in the phase diagram domain.

As an illustrative example, let us consider the following auxiliary signals:

$$\begin{aligned}
 aux_1 &: \begin{cases} a_1 = 1 \cdot 10^6, b_1 = 1.1 \cdot 10^6, c_1 = 1.2 \cdot 10^6 \\ n_1 = 0.5N, n_2 = 0.4N, n_3 = 0.6N, N = 500 \end{cases} \\
 aux_2 &: \begin{cases} a_1 = 1.1 \cdot 10^6, b_1 = 1.2 \cdot 10^6, c_1 = 1 \cdot 10^6 \\ n_1 = 0.5N, n_2 = 0.4N, n_3 = 0.6N, N = 500 \end{cases} \\
 aux_3 &: \begin{cases} a_1 = 1.2 \cdot 10^6, b_1 = 1 \cdot 10^6, c_1 = 1.1 \cdot 10^6 \\ n_1 = 0.5N, n_2 = 0.4N, n_3 = 0.6N, N = 500 \end{cases}
 \end{aligned} \tag{3.12}$$

Next, we study the frequency modulated signal's robustness to noise in terms of phase space lobe representation. From Fig. 3-18 and Fig. 3-19, it can be noticed that this approach is robust enough for a SNR of at least 10 dB.

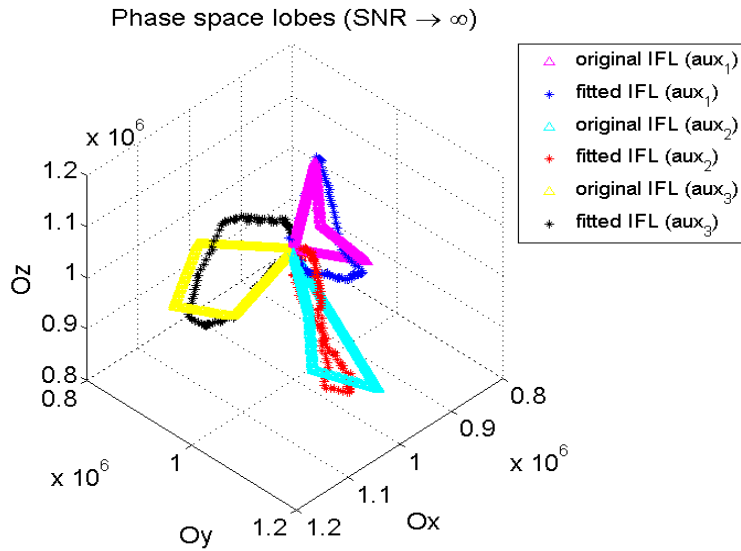


Fig. 3-18: Phase space lobes in noise-free conditions: original IFL (triangles) and the spectrogram recovered IFL (stars)

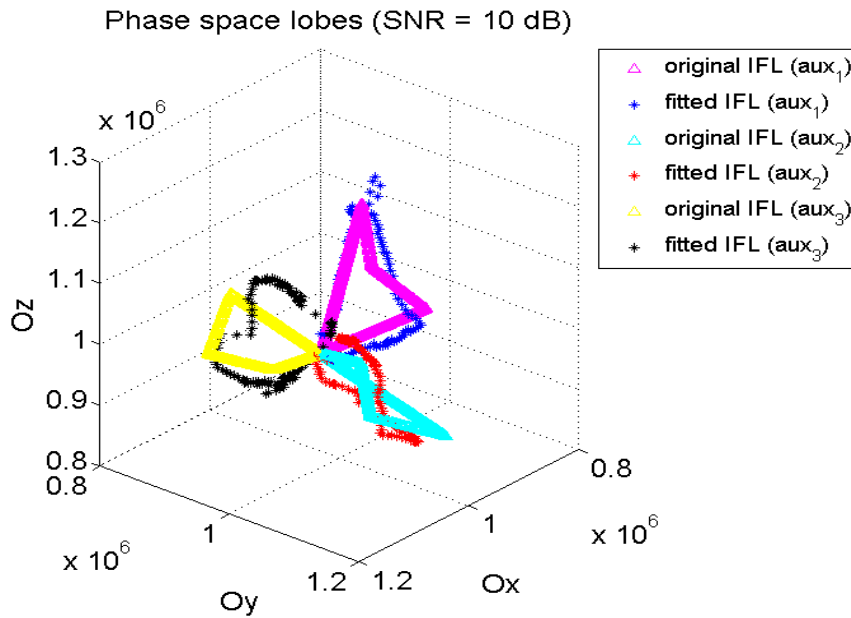


Fig. 3-19: Phase space lobes for SNR = 10 dB: original IFL (triangles) and the spectrogram recovered IFL (stars)

For SNRs below 10dB, there is a real problem with the choice of the parameters of the spectrogram that will conduct to satisfactory results. These parameters refer especially to the number of points for the Fourier transform computation and the overlapping of the windows. For example, for the noise-free case, the number of points is 512, whereas for the SNR = 10 dB, the number of points is 4096, these two setting conducting to similar quality of IFL estimate. Another limitation is the choice of these parameters which is done by several trials that are time consuming.

Considering more than 3 lobes in the phase space, we define them as all possible combinations of the coefficients $a_1 = 0.87 \cdot 10^6$, $b_1 = 1 \cdot 10^6$, $c_1 = 1.13 \cdot 10^6$. Hereby, we get a set of six signals, Fig. 3-20. The number of points are the same for the entire set: $n_1 = 0.5N$, $n_2 = 0.4N$, $n_3 = 0.6N$, $N = 500$.

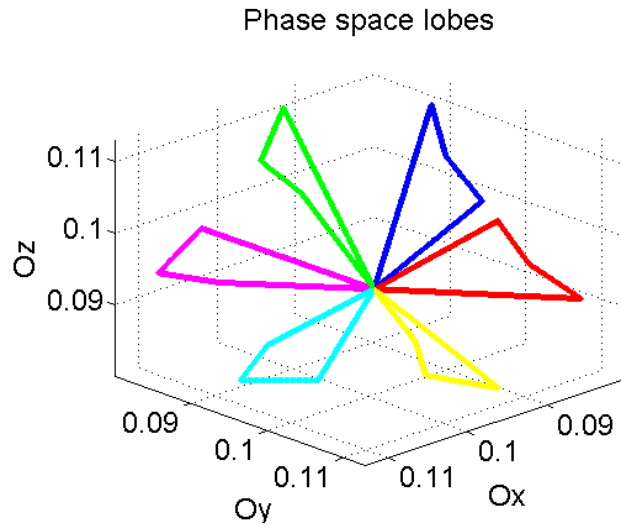


Fig. 3-20: Phase space lobes generated from the set of 6 auxiliary signals used IFLs

Next, these IFLs are used to modulate the transmitting signal. Fig. 3-21 presents the results of the phase space representation of the IFLs in the noise-free case.

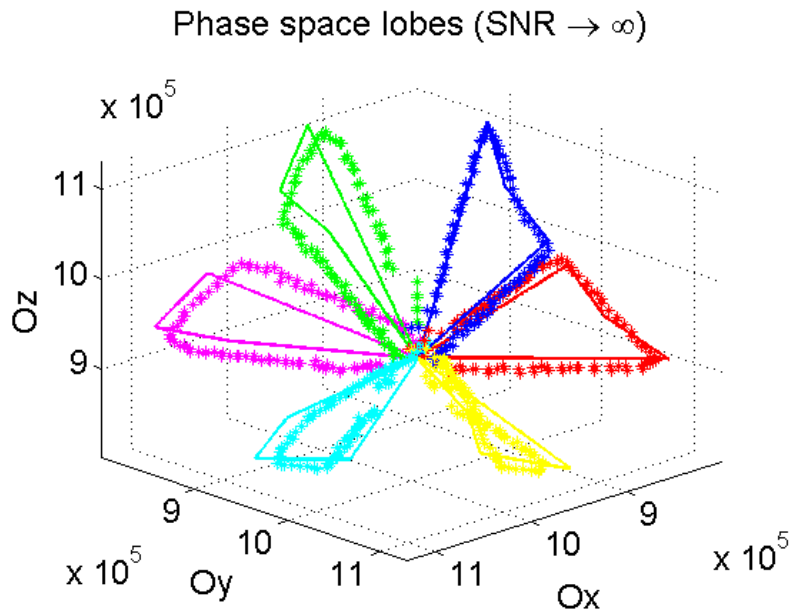


Fig. 3-21: The original IFLs (continuous line) and the spectrogram-based recovered IFLs (stars) – noise -free case

On the other hand, when white additive Gaussian noise corrupts the signal, the recovered signal is not very smooth. The results are presented in Fig. 3-22 where we reinforce that they are obtained for higher number of FFT (Fast Fourier Transform) points for the spectrogram (compared with the case of no noise (512 points), the number of points is eight times higher (4096 points)).

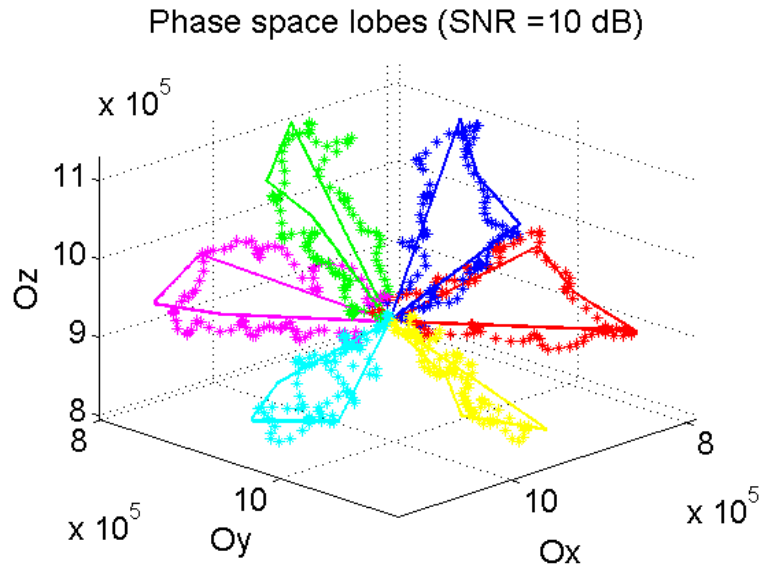


Fig. 3-22: The original IFLs (continuous line) and the spectrogram-based recovered IFLs (stars) for SNR = 10 dB

At lower SNRs (below 10 dB), the things become more complicated and it is difficult to separate the signals in the phase space, Fig. 3-23.

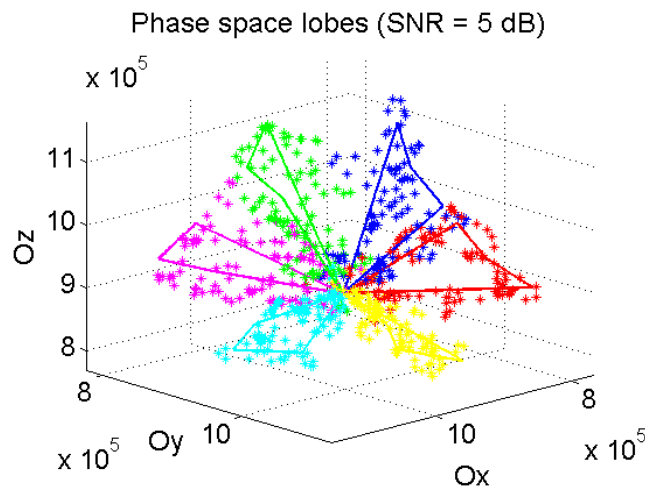


Fig. 3-23: The original IFLs (continuous line) and the spectrogram-based recovered IFLs (stars) for SNR = 5 dB

To improve the robustness to noise and to overcome the limitation of parameter dependence for the IFL's spectrogram-based estimation, we need to define a new approach for the IFLs estimation.

3.2.1. Phase space-based IFL estimation

In this section we introduce a new technique for IFL estimation using the phase space representation. This approach is, once again, suggested by the evolution of a signal – sine wave – in phase space.

Considering a sine wave of frequency f_0 and sampling frequency f_s , then, for a duration D , where $D = m / f_0$ and m is a positive integer, the sine wave contains $N_s = D \cdot f_s$ samples.

In phase space, the sine wave has a trajectory which forms an ellipse. It is known that a complete rotation of the ellipse is in fact the evolution of a sine wave during a complete period.

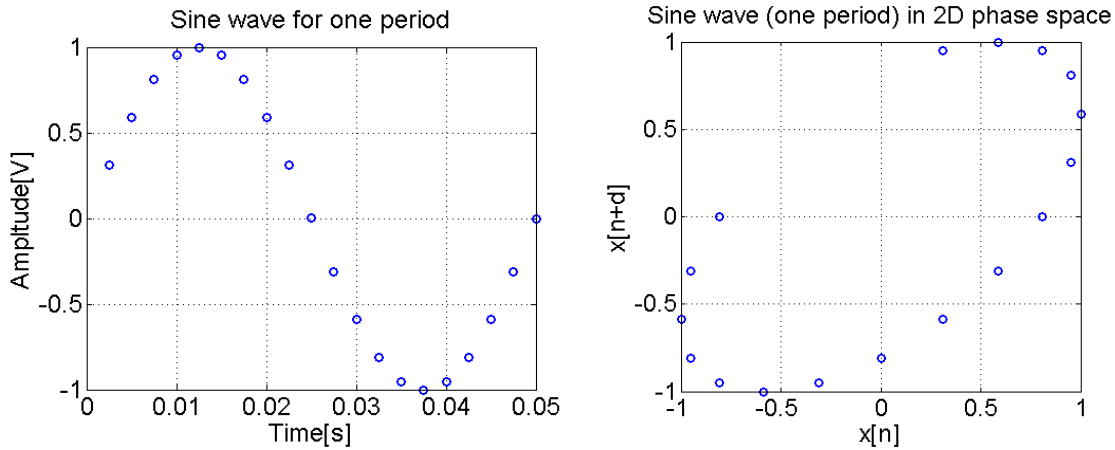


Fig. 3-24: The time evolution of a sine wave during one period (left) and its trajectory in 2D phase space (right):
 $f_0 = 20$; $f_s = 400$; $D = 1 / f_0$; $d = 3$

From Fig. 3-24, it can be noticed that the sine wave forms an ellipse in the phase space which is oriented along the first bisector, namely the ellipse is rotated with 45° . Still, for only one period, the ellipse is not complete. This is caused by the actual representation given by eq. (2.11). In other words, the number of samples from the phase space representation is $M = N_s - (m-1)d$. In our case, $M = 20 - (2-1)3 = 17$. In order to overcome this drawback, the considered signal can be interpolated so that the difference given by the factor $(m-1)d$ can be neglected (Fig. 3-25) and/ or the signal should contain two or more periods (Fig. 3-26).

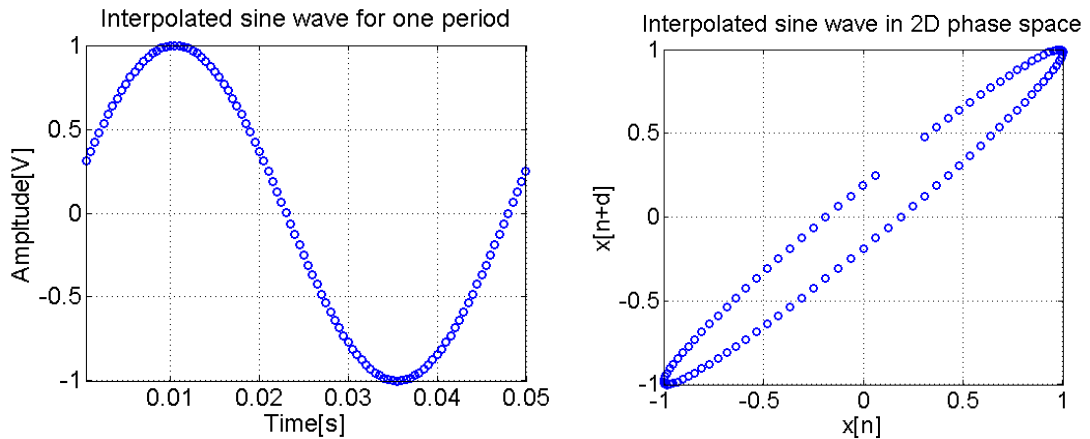


Fig. 3-25: The time evolution of an interpolated sine wave during one period (left) and its trajectory in 2D phase space (right): $f_0 = 20$; $f_s = 400$; $D = 1 / f_0$; $d = 3$ and the interpolation factor $f = 5$

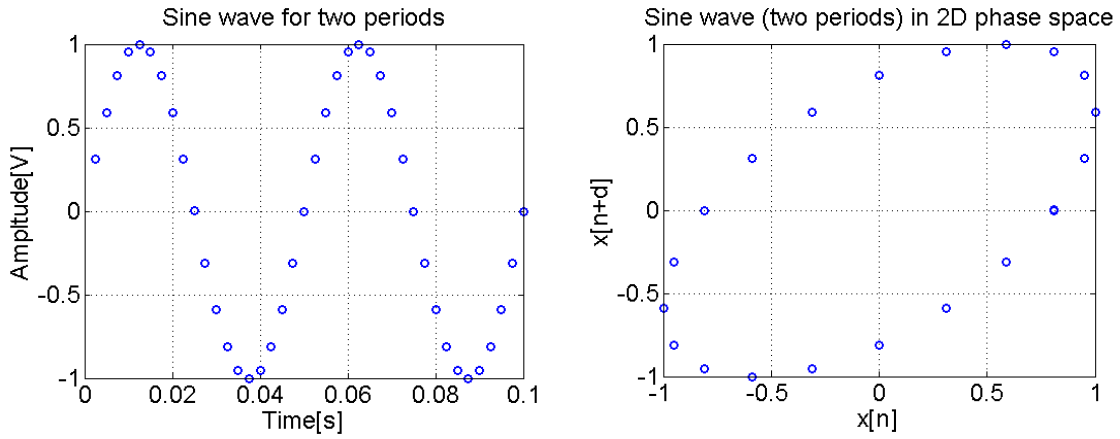


Fig. 3-26: The time evolution of a sine wave during two periods (left) and its trajectory in 2D phase space (right): $f_0 = 20$; $f_s = 400$; $D = 2 / f_0$; $d = 3$

Based on the observations presented above, for the phase space based IFL estimation, we have established the following algorithm:

- The trajectory of the analyzed signal is computed in 2D phase space (so that the factor $(m-1)d$ is minimized); previously, the signal can eventually be interpolated with an interpolation factor f
- The phase space is translated in the positive quadrant
- The polar angle of each point of phase space is computed as follows:

$$\theta[n] = \arctan\left(\frac{x[n+d]}{x[n]}\right) \quad (3.13)$$

- Knowing that the trajectory evolves along the first bisector, we eliminate this bias of 45° by subtracting this value from all polar angles
- Next, the ellipse's evolution is quantified: a counter is started when the first point with a positive polar angle is found and the counter is stopped when the trajectory returns in the same region (the angle passes from negative into positive). With the number of points counted in this interval, n_0 , we obtain the first f_0 estimation, where $f_0 = f_s / n_0$. The algorithm continues until the last point of the phase space is reached.

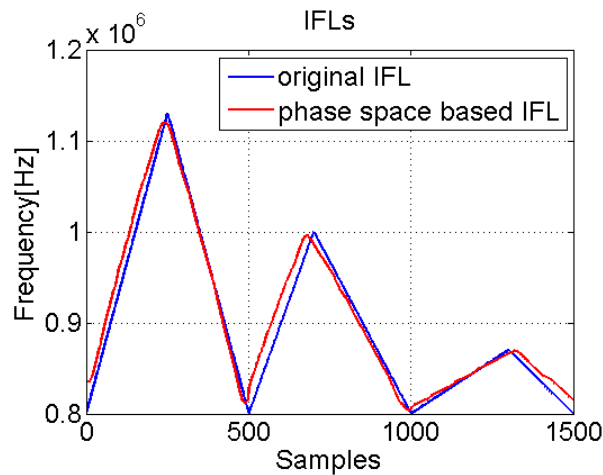


Fig. 3-27: The original IFL (in blue) and the phase space-based IFL estimation (in red – $f = 5, d = 10$)

For the signals presented in Fig. 3-27, we present the IFL estimation based on the phase space evolution.

Still, the choice of the delay d must be discussed. The optimal delay depends on the signal’s characteristics; therefore, we propose the use of the multi-lag phase analysis (MLPA) (Bernard, et al., 2016), (Digulescu, et al., 2016c), (Bernard, et al., 2014) for the IFL estimation. Besides the algorithm presented above, we also vary the delay and we choose the delay that gives the IFL estimate that is the closest to the transmitted signal. This decision is made by minimizing the cost function given by eq. (3.14).

$$C(d) = \sum |IFL_o - IFL_e(d)| \tag{3.14}$$

where IFL_o is the original IFL and IFL_e is the phase space-based IFL.

Next, a comparison between the use of a fixed delay and the MPLA is presented.

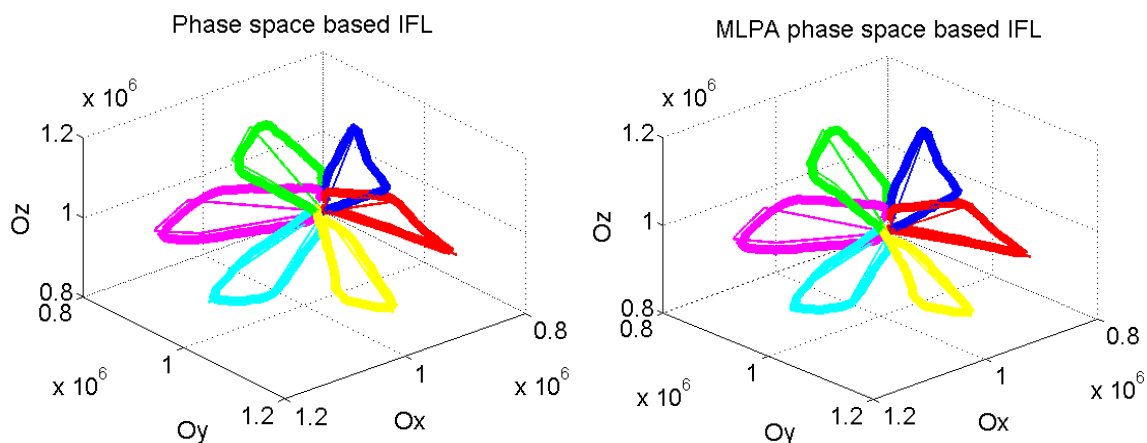


Fig. 3-28: The phase space-based IFL for $d = 5$ and $SNR \rightarrow \infty$ (left) and the MLPA phase space-based IFL $d = 15 \ 15 \ 13 \ 14 \ 15 \ 15$ (according to the colors: clockwise from blue to green) and $SNR \rightarrow \infty$ (right); continuous line – original IFL; starred line – estimated IFL

From Fig. 3-28, it goes that the IFL information do not depend on the delay when the signal does not contain any noise.

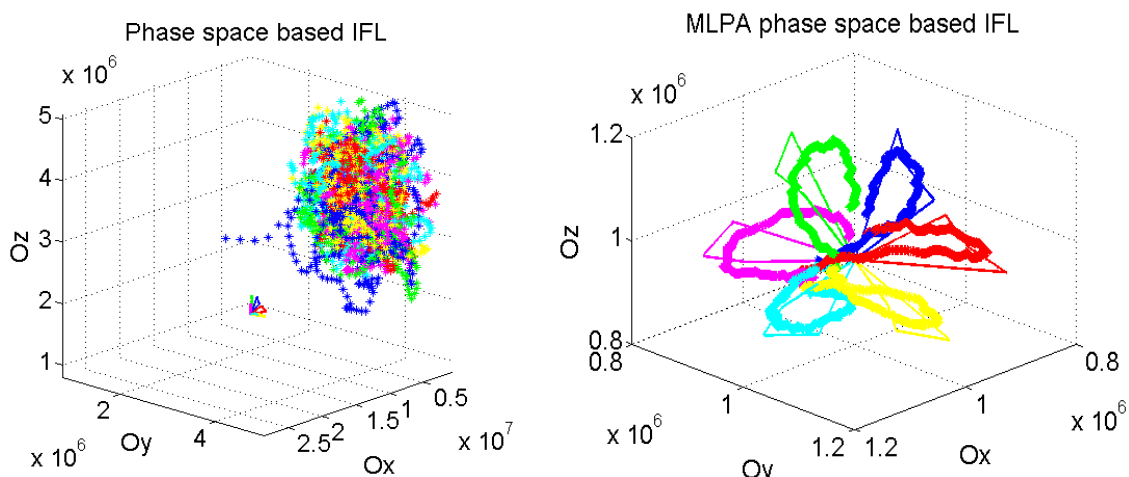


Fig. 3-29: The phase space-based IFL for $d = 5$ and $SNR = 5\text{ dB}$ (left) and the MLPA phase space-based IFL $d = 21\ 27\ 19\ 30\ 21\ 25$ (according to the colors: clockwise from blue to green) and $SNR = 5\text{ dB}$ (right); continuous line – original IFL; starred line – estimated IFL

When it comes to noisy signals, in Fig. 3-29, the $SNR = 5\text{ dB}$, a fixed delay does not succeed to track the IFL, (here $d = 5$ is just an example). Using the MLPA, it is obvious that the IFL is tracked, the optimal values of the delay varying from a signal to another.

Furthermore, although the estimated IFLs suffer a slight rotation from the original ones, the tracking of the IFL is successfully achieved, the signals being separated and distinguishable.

3.2.2. “Marked” lobes

This new waveform generation technique is based on the information coding via a frequency modulation defined by the lobes designed in the phase space. If the signals defined from these lobes are used in a time-of-flight measurement application, we propose to introduce a special mark that will help us measuring the time of flight as the distance, in time, between these marks are observed in the phase space.

Hereby, the two director vectors that define the lobe in the Cartesian system will present an extra point to mark the signature of the lobe, as presented in Fig. 3-30.

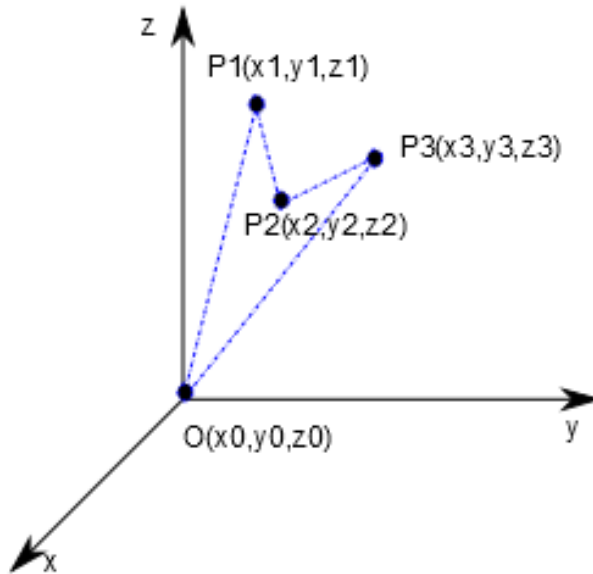


Fig. 3-30: The ideal phase space representation of the “marked” lobe (auxiliary signal)

From figure above, it can be observed that the “signature” of the lobe is given by the point $P_2(x_2, y_2, z_2)$ in the phase space. Based on this “signature”, we aim to determine the TOF (Fig. 3-31). The interest of this approach is that, even if the IFLs between transmitter and the receiver are not the same (due to propagation effect or event estimation error), if the marks are still visible, they allow a more accurate estimation of TOF than looking in time domain.

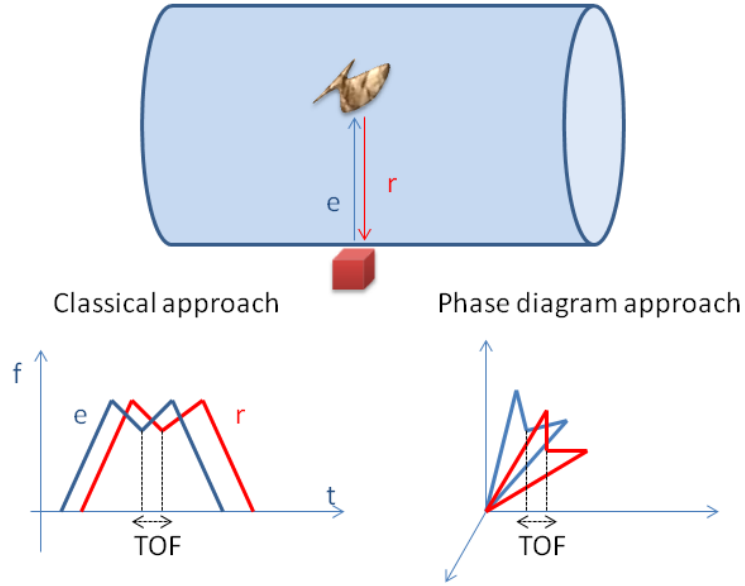


Fig. 3-31: The phase diagram space approach versus classical approach; in the phase diagram, the estimation of TOF is done using the distance between the marks of emitted and received signals, respectively

The *first condition* which we imposed is: the four points O, P_1, P_2 and P_3 should be coplanar. It goes that:

$$\begin{vmatrix} x_0 & y_0 & z_0 & 1 \\ x_1 & y_1 & z_1 & 1 \\ x_2 & y_2 & z_2 & 1 \\ x_3 & y_3 & z_3 & 1 \end{vmatrix} = 0 \quad (3.15)$$

Considering $x_0 = y_0 = z_0 = 0$, eq. (3.15) is equivalent to:

$$\begin{vmatrix} x_1 & y_1 & z_1 \\ x_2 & y_2 & z_2 \\ x_3 & y_3 & z_3 \end{vmatrix} = 0 \quad (3.16)$$

It is obvious that the points P_1, P_2 and P_3 should be different, therefore the lines of the determinant from eq. (3.16) must be different. In order to accomplish the condition of coplanarity for whatever points P_1, P_2 and P_3 , at least two columns of the determinant should be proportional. Therefore, the coplanarity condition (eq.(3.15)) restraints to the *second condition*:

$$\frac{x_i}{z_i} = r, \quad i \in \{1, 2, 3\}, \quad r \in (1, 2) \quad (3.17)$$

Then, the second column is defined as: $y_1 = x_3, y_2 = x_2$ and $y_3 = x_1$. With this condition, the phase space representation forms a lobe with a “signature” provided that $x_1 > x_2$ and $x_3 > x_2$.

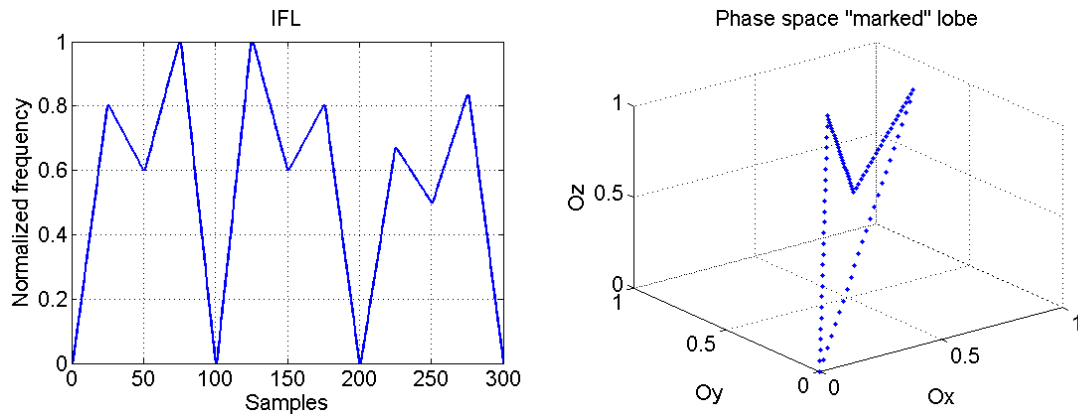


Fig. 3-32: The IFL (auxiliary signal) (left) that forms the phase space marked lobe (right):
 $x_1 = 0.8, x_2 = 0.6, x_3 = 1$ and $r = 1.2$

Next, the signals are generated in the same way as in the case of the simple lobe: the IFLs (auxiliary signals) are used to obtain the frequency modulated wide band signals

The multiple lobes approach results from the rotation in the phase space, with a given angle, of one “marked” lobe (for example, the one from Fig. 3-32). In this way, we get multiple “marked” lobes.

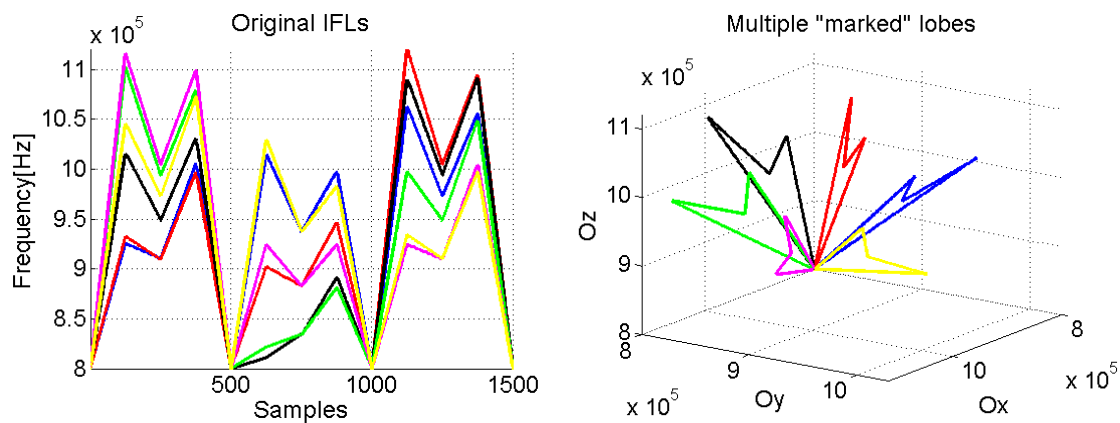


Fig. 3-33: The multiple lobes with signature obtained based on the lobe from Fig. 3-32 and rotated with an angle of $\pi / 6$

As expected, these signals are not separable in time and frequency domain, as it is shown in Fig. 3-34.

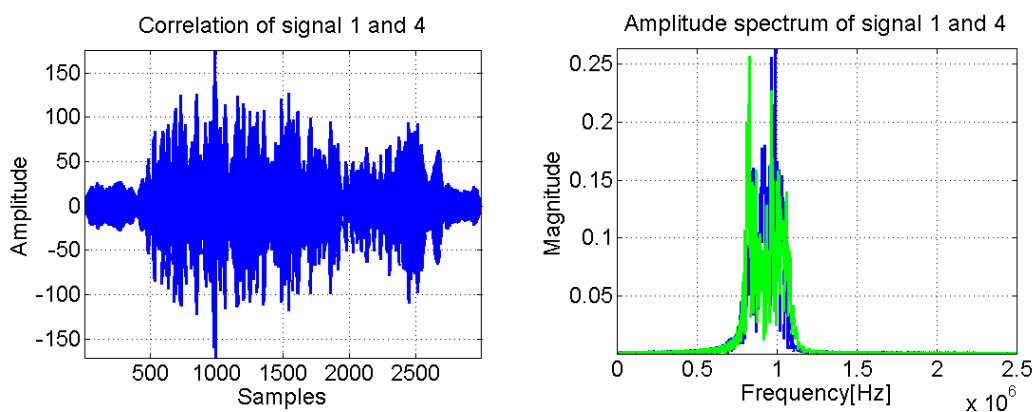


Fig. 3-34: Example of correlation between the signals that contain the IFLs blue and green from Fig. 3-33 (left) and their corresponding amplitude spectral representation (right)

Next, we study the robustness to noise of the multiple “marked” lobes.

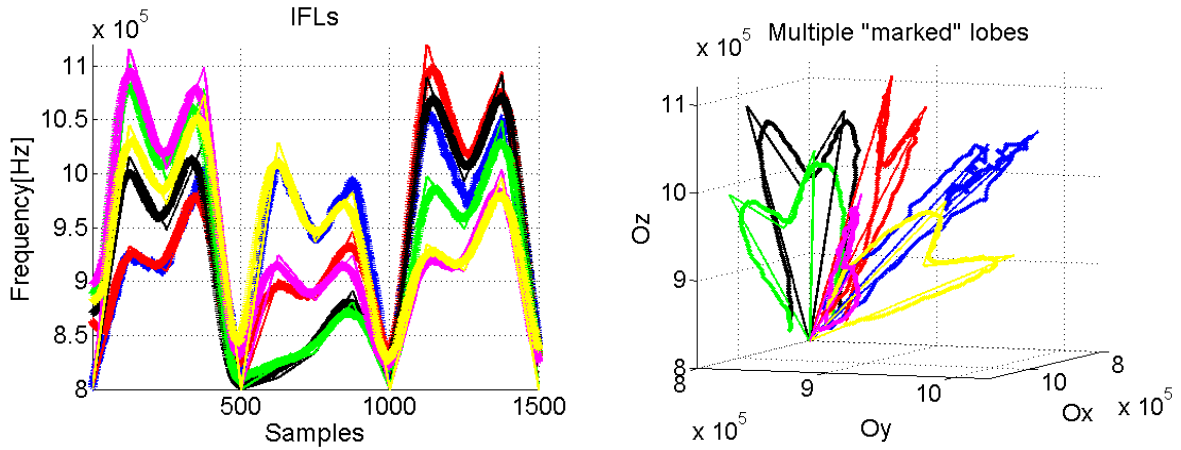


Fig. 3-35: The original IFLs (continuous line) and the phase space-based estimated IFLs (starred lines) for $SNR = 20dB$ (left) and their corresponding phase space lobes (right)

From Fig. 3-35, it can be noticed that the signals are discriminated in the phase space for a high SNR. Although, the recovered IFLs form a larger lobe in the phase space, than the original one, they do not overlap and present the approximately the same evolution.

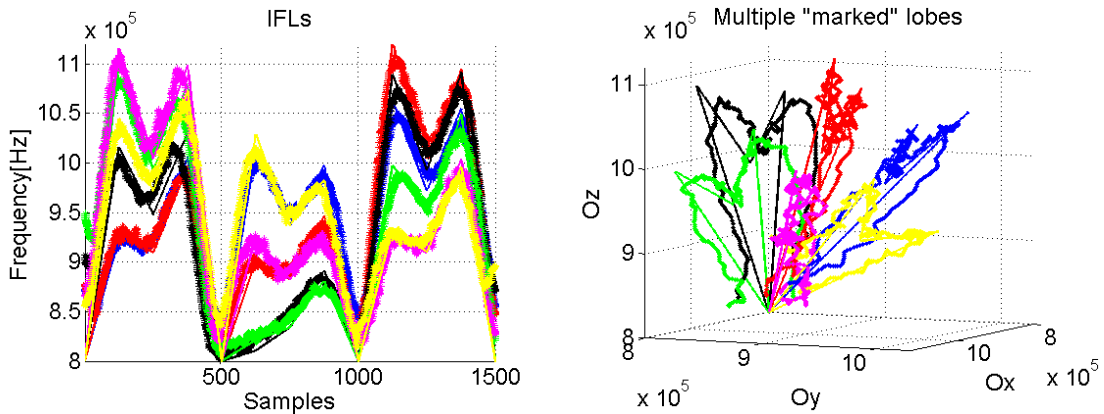


Fig. 3-36: The original IFLs (continuous line) and the phase space-based estimated IFLs (starred lines) for $SNR = 7dB$ (left) and their corresponding phase space lobes (right)

The limit of the SNR is around $7dB$, Fig. 3-36, in order that the lobes remain distinguishable in the phase space and without overlapping. We mention that the phase space-based IFL estimation is done using the MLPA and an interpolation factor of 10.

3.3. Chapter summary

This chapter illustrates a new concept for the waveform design, namely the phase space-based waveform generation, which can be used to create a set of orthogonal signals in phase space. This signal is created in the phase space and then its information is translated into 1D signal. Then, there is also presented a new IFL estimation method for the signal recovery, designed in the phase space.

These signals are tested and their feasibility is validated in our experimental setup presented in Chapter 5.

4. Passive sensing based on the RQA concept

Contrarily of the applications requiring an active sensing, there are a large number of applications where active sensing is useless. This is particularly the case of the application where the phenomena under study generate signals whose analysis allows us to understand (detection, classification, localization, etc.) them. Currently, a special attention is carried to the transient phenomena, characterized by very fast variation in time and frequency. The development of new analysis methods for such signals is a current work direction in the field of signal processing.

In this context, this chapter presents new methods for the detection, localization and characterization of transient signals based on the concept of RQA. As discussed in chapter 2, the RPA and RQA methods are the main approaches derived from the phase diagram and the information that could be retrieved from this analysis can bring a deeper understanding of the studied phenomena (Bernard, et al., 2016), (Marwan, Schinkel, & Kurts, 2013), (Webber Jr. & Zbilut, 2005), (Zbilut & Webber Jr., 2007).

Firstly, there is highlighted the measure used for the detection and localization of a transient signal, namely the time-distributed recurrence (TDR) measure. This measure exploits the fact that a sudden change in the signal under study is represented as a new state in the phase space which evolves in a totally different region than the rest of the signal's states. The initial idea has been presented in (Marwan, et al., 2007) and our contribution was to adapt this idea to the problem of detection of transient sources.

Then, the multi-lag phase-space analysis will be introduced. This concept is very useful for the characterization of transient signals (Bernard, et al., 2016), (Bernard, et al., 2014), (Digulescu, et al., 2016c). This approach allows us to discriminate between signals with close parameters based on their evolution in the phase diagram.

Finally, the diagonal lines quantification is presented. This approach is based on the fact that the diagonal lines from the distance/ recurrence matrix spectral content of any signal (Ioana, et al., 2014). With this property, we are able to highlight slight changes that appear in signal's spectral content, much weaker than the fundamental frequency, for instance.

4.1. *Time-distributed recurrence*

The core of the time-distributed recurrence (TDR) measure starts from the idea that a sudden change in a time series represents a new state of the dynamical system (Ioana, et al., 2014), (Birleanu, et al., 2012), (Marwan & Kurths, 2002), (Webber Jr. & Zbilut, 2005), (Webber Jr., 2012), (Zbilut & Webber Jr., 2007), (Ramirez Avila, et al., 2013), (Yang, 2011), (Digulescu, et al., 2016c), (Le Bot, Gervaise, & Mars, 2016). Namely, there is no recurrence

with the previous states. Therefore, when an appropriate distance is used, the distance/recurrence matrix presents a horizontal/ vertical band with much fewer recurrences. When the sum of the lines/ columns of the recurrence matrix is computed, we actually obtain the column average (Digulescu, et al., 2016c), (Ioana, et al., 2014), which, in the case of transient signals, significantly changes.

The TDR measure is defined as:

$$TDR[n] = \frac{1}{M} \sum_{i=1}^M R_{i,n} \quad (4.1)$$

where $R_{i,n}$ is the RM from eq. (2.13).

This measure can be interpreted as the column average recurrence of a given point i or the recurrence density heterogeneity in the point i . Hereby, a solitary position of a phase space vector changes significantly its average recurrence and it can be highlighted through the use of the complementary version of the TDR measure from eq. (4.1):

$$TDR^*[n] = 1 - \frac{TDR[n]}{\max\{TDR[n]\}} \quad (4.2)$$

Fig. 4-1 highlights the detection based on the TDR^* measure.

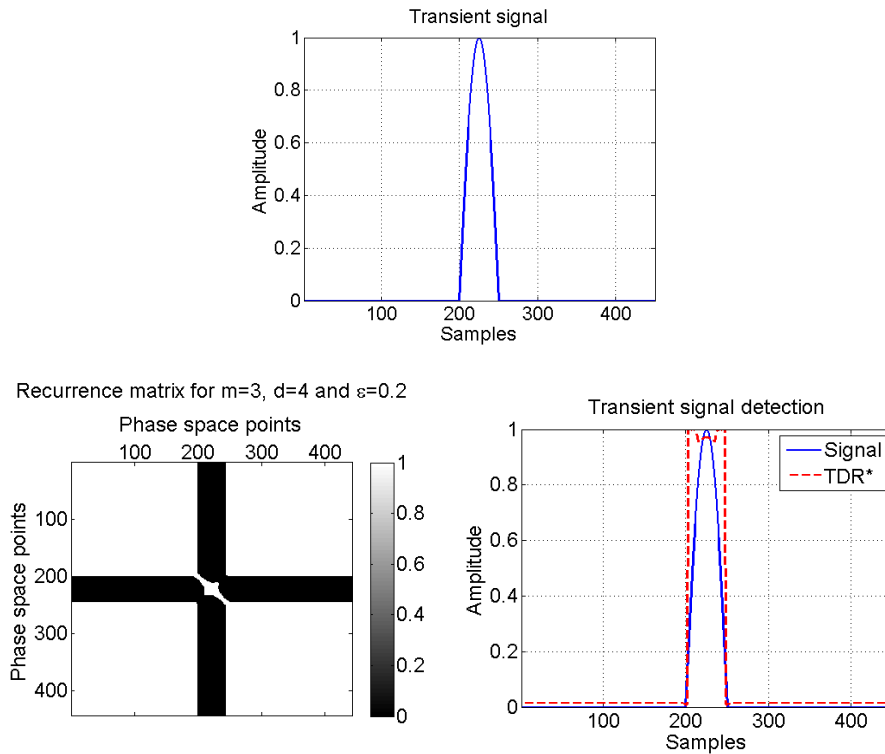


Fig. 4-1: The transient detection using the TDR^* measure: (up) the analyzed transient signal; (left-down) the recurrence matrix of the analyzed signal; (right-down) the analyzed signal and its TDR^* based detection curve

Moreover, in order to detect only the transient signal from the analyzed observation, the SNR must be computed. We computed it as follows: the last part of the acquired signal (when no transient occurs) is considered as noise, $z[i]$, (unwanted signal recording

environmental noise), whereas the part that has a different behavior is considered as the interest signal, $s[i]$. Both parts of the signal, $s[i]$ and $z[i]$ have the same length, N . Then, the SNR is computed as:

$$SNR = 10 \log_{10} \left(\frac{\left(\sum_{i=1}^N (s[i])^2 \right)}{\left(\sum_{i=1}^N (z[i])^2 \right)} \right) \quad (4.3)$$

Accordingly, the threshold ε of the recurrence matrix is chosen so that it includes a percentage, α , of the maximum value of the equivalent noise which has the same power as the interest signal:

$$\varepsilon = \alpha \cdot \max(|SNR \cdot z[i]|) \quad (4.4)$$

where α is a constant that is chosen to nonlinearly filter the noise. In our applications, α varies from 0.4 to 0.95.

Considering the eq.(4.4), the components of the noise (undesired parts of the signal) are considered as recurrences; therefore, the transient signal (useful part of the signal) is highlighted by the proposed measure.

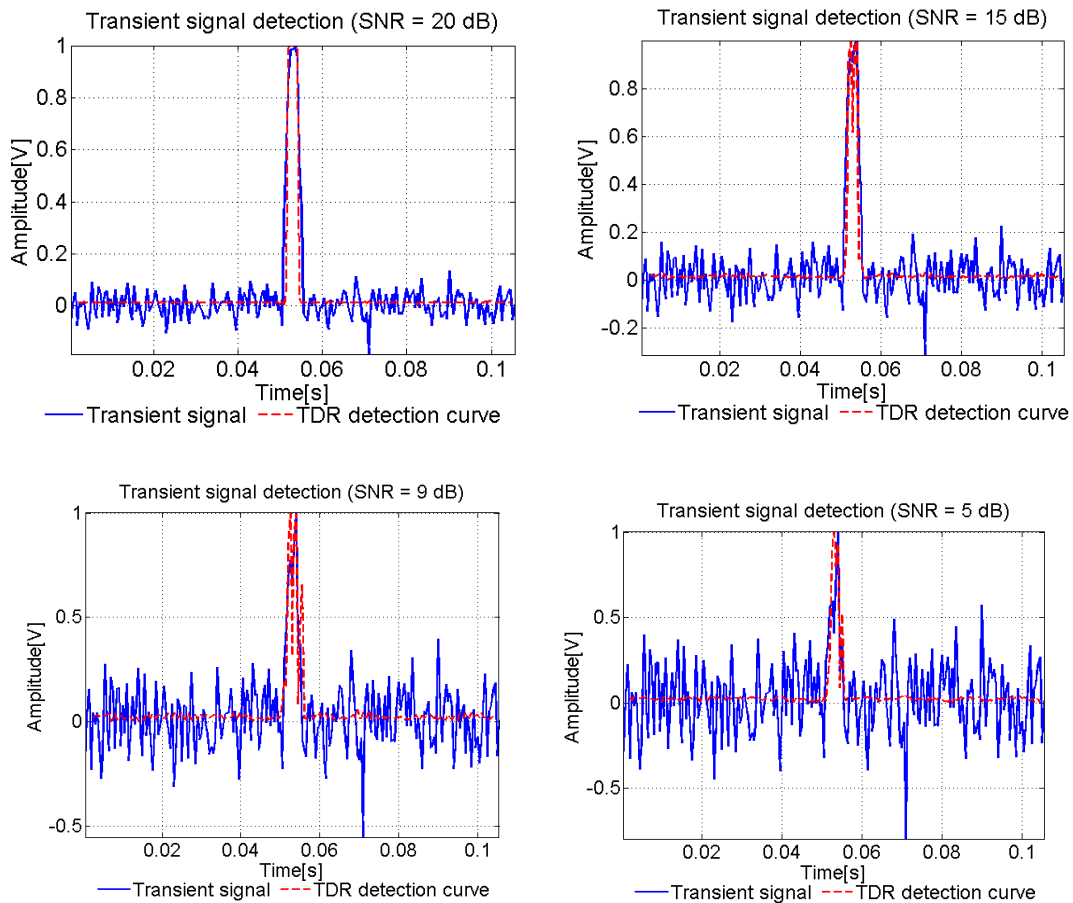


Fig. 4-2: Transient signal detection using the TDR measure for a transition of only 11 samples where $m = 3$, $d = 2$ and $\alpha = 0.8$

Fig. 4-2 emphasizes the advantage brought by the TDR measure, namely that the detection provided by the detection curve is more robust. In the examples above, the detection curve exhibits a SNR improvement of 30 dB .

The major advantage of this method is that, in applications where the SNR level of the acquired signal varies a lot, the TDR detection curve approach improves the SNR level after filtering the acquired signal with the detection curve. Through this approach, the signal's power after filtering is significantly increased (in our applications, at least 10 dB), which is very helpful for the classification and a further characterization.

4.2. Multi-lag phase-space analysis

The RPA method stands, as its name suggests it, on the concept of recurrences. Still, it is worth to pay more attention to the phase-space representation, in order to achieve richer characterization of signals with close parameters coming from the same source and having the same propagation and acquisition conditions.

The phase-space representation is very rich in information regarding the evolution of a dynamical system (Eckmann, Kamphorst, & Ruelle, 1987), (Kantz & Schreiber, 1997), therefore the analysis can be made in any dimension, but our approach restrains, for simplicity seeking, to a bi-dimensional representation.

Therefore, considering the Cartesian coordinate system, the phase-space representation is defined by plotting (x_n, y_n) defined as:

$$\begin{cases} x_n = s[n] \\ y_n = s[n + d] \end{cases} \quad (4.5)$$

where n is the n^{th} sample of the recorded time series. It goes that eq. (2.11) can be rewritten as: $\vec{v}_n = s[n] \cdot \vec{i} + s[n + d] \cdot \vec{j}$.

Firstly, the main attributes of this representation are recalled. Considering three signals s_1, s_2 and s_3 defined as:

$$\begin{aligned} s_1[n] &= s[n + \delta] \\ s_2[n] &= s[\alpha n] \\ s_3[n] &= \beta \cdot s[n] \end{aligned} \quad (4.6)$$

where α, β and δ are constant which modify the signal s through translation in time, scale and/or amplitude, the phase-space points present the following attributes (Marwan, et al., 2013), (Birleanu, et al., 2011):

$$\begin{aligned} \overline{v_{1[i]}} &= \overline{v_{i+\delta}} \\ \overline{v_{2[i]}} &= \overline{v_{\alpha i}} \\ \overline{v_{3[i]}} &= \beta \cdot \overline{v_{[i]}} \end{aligned} \quad (4.7)$$

The phase-space trajectory is invariant to translation and it points out the scale and amplitude change. In order to introduce the concept of multi-lag phase-space analysis, we recall the three transient signals with close characteristics given by eq. (2.9). These signals are given by the generic signal:

$$s_{(a,f,b)}[n] = \begin{cases} a \cdot (\sin[2\pi \cdot f \cdot n] + b[n]), n = \{1, 2, \dots, N\} \\ a \cdot b[n], \text{otherwise} \end{cases} \quad (4.8)$$

where $N = \left\lceil \frac{f_s}{2f} \right\rceil$ and $f_s = 10\text{MHz}$ is the sampling frequency.

The considered signals $s_{1(a_1, f_1, b_1)}$, $s_{2(a_2, f_2, b_2)}$ and $s_{3(a_3, f_3, b_3)}$, presented in Fig. 4-3, have the following relationships between their parameters: $a_1 / a_2 = 1 / 0.6 = 1.66$, $a_1 / a_3 = 1 / 0.3 = 3.33$, $f_1 / f_2 = (2 \cdot 10^5) / (1.9 \cdot 10^5) = 1.05$, $f_1 / f_3 = (2 \cdot 10^5) / (1.6 \cdot 10^5) = 1.25$, $b_1 = b_2 = b_3$ and $\text{SNR}_i = 20\text{dB}$, $i = \{1, 2, 3\}$.

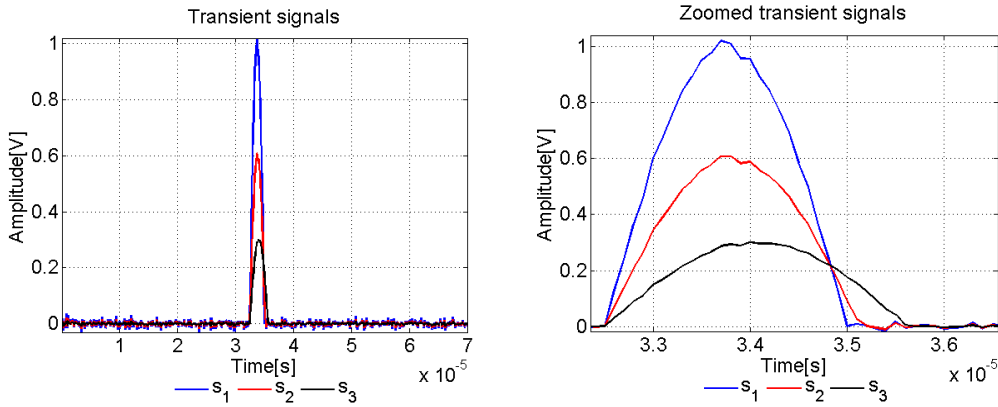


Fig. 4-3: Transient signals considered for the multi-lag phase space analysis

In order to distinguish between the synthetic transient signals considered in eq. (4.8), the multi-lag phase space analysis is considered (Digulescu, et al., 2016c), (Bernard, et al., 2016), (Bernard, et al., 2014). The study of the representation of the trajectory for multiple lags is done by two approaches: the elliptic modeling of the trajectory for the area estimation in the phase-space, respectively, the polar coordinate representation.

The elliptic modeling supposes to determine the ellipse that circumscribes the phase-space trajectory and, therefore, to estimate the area of the trajectory through the ellipse's area computation. The area of the modeling ellipse is a new descriptor of the transient signal.

The trajectory is modeled by considering the solution that minimizes the following system (Bernard, et al., 2016), (Bernard, et al., 2014), (Digulescu, et al., 2016c):

$$S = \sum_{i=1}^M [F(x_i, y_i)]^2 \quad (4.9)$$

where $F(x, y) = \Gamma x^2 + \Lambda y^2 + 1$. The least mean square estimation of eq. (4.9) determines the couple (Γ, Λ) . It goes that the major semi axes $a = 1 / \Gamma$ and the minor semi axes $b = 1 / \Lambda$.

The next step after the elliptic modeling is the area estimation and the determination of the optimal delay. The delay is considered to be optimal for the value that provides an average value of the area.

For example, in Fig. 4-4, we present this aspect for the synthetic transient signal shown in Fig. 4-1 (up).

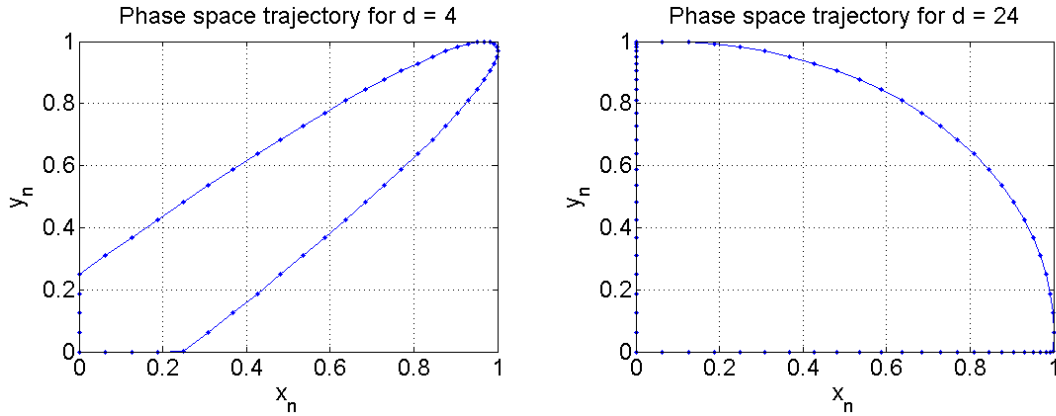


Fig. 4-4: The phase space trajectory evolution using different delays

This average value area ensures a suitable phase-space representation where the trajectory does not evolve too close to the main diagonal (case of redundancy) or its evolution is too complicated (case of irrelevance) (Casdagli, et al., 1991).

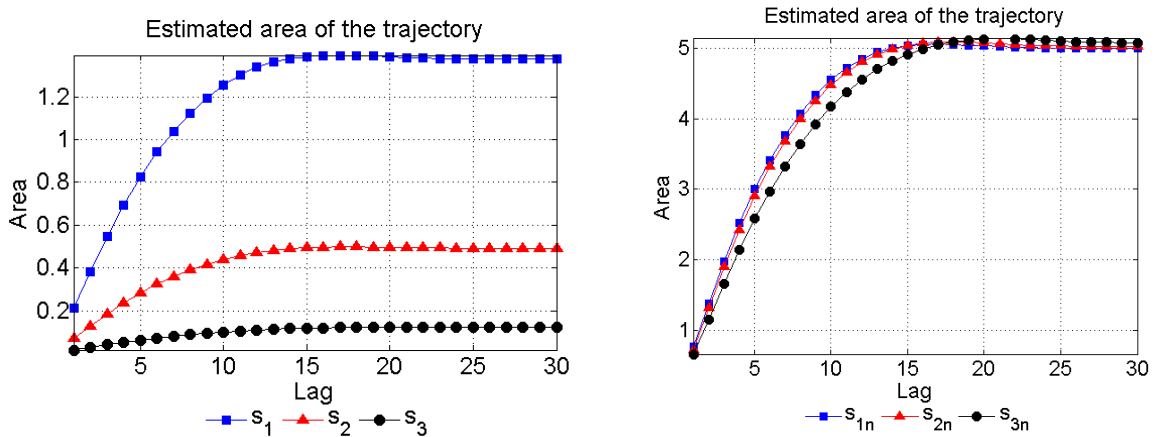


Fig. 4-5: The evolution of the area according to the lag (delay): for the signals presented in Fig. 2-5 (left) and for the normalized signals (right)

From Fig. 4-5, it can be noticed that the area of the signals has a similar trend with some differences for the normalized signals. Hereby, considering the average value of the area, $A=3$, the optimal delay for s_1 is $d_1 = 5$, for s_2 , $d_2 = 5$ and for s_3 , $d_3 = 6$.

With the chosen lags, the next step is to plot the phase-space trajectory into polar coordinates:

$$\begin{aligned}\varphi[n] &= \sqrt{x_n^2 + y_n^2} \\ \theta[n] &= \arctan \frac{y_n}{x_n}\end{aligned}\quad (4.10)$$

This representation is very useful for signals with different amplitudes, because it shows the evolution of the position vector regardless of its length (that is related to signal's amplitude).

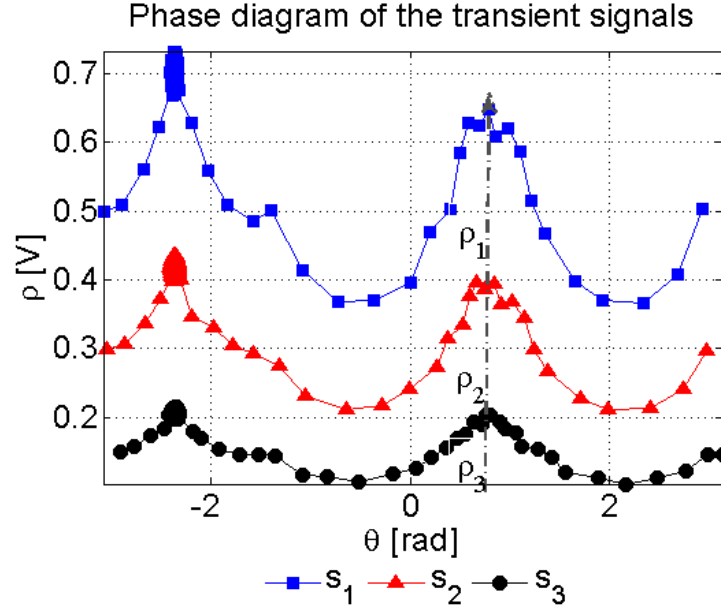


Fig. 4-6: The polar coordinates representation for the transient signals considered in our study:

$$s_1 : d_1 = 5, s_2 : d_2 = 5 \quad \text{and} \quad s_3 : d_3 = 6$$

Fig. 4-6 shows the evolution of the transient signals in the polar coordinates derived from the phase diagram. It can be observed that the noise is concentrated in a small region of the phase space which is translated in a small conglomeration of points around the peak corresponding to the angle $-3\pi/4$.

In this way, through the isolation of this region, the noise can be eliminated in the process of signal characterization. Moreover, the evolution of the trajectories is different for the peaks corresponding to the angle $\pi/4$. Measuring the values of the maximum length of the vector corresponding to the $\pi/4$ angle, it goes: $\rho_1/\rho_2 = 1.67$ and $\rho_1/\rho_3 = 3.43$. So, the ratio of the amplitude of the transient signals is then conserved in this new representation.

Furthermore, the number of points between the two angles $-3\pi/4$ and $\pi/4$ (after excluding the points corresponding to the noise) is directly related to the fundamental frequency of the signal. Let n_0 be this number of points from the phase diagram. The fundamental period of the signal, T_0 is:

$$T_0 = \frac{2n_0}{f_s} \quad (4.11)$$

From eq.(4.11), it goes that: $f_1 / f_2 = 1.04$ and $f_1 / f_3 = 1.24$ which are very close to the ratios proposed in eq. (2.9). These results allow us to discriminate between the proposed signals although their characteristics are very close.

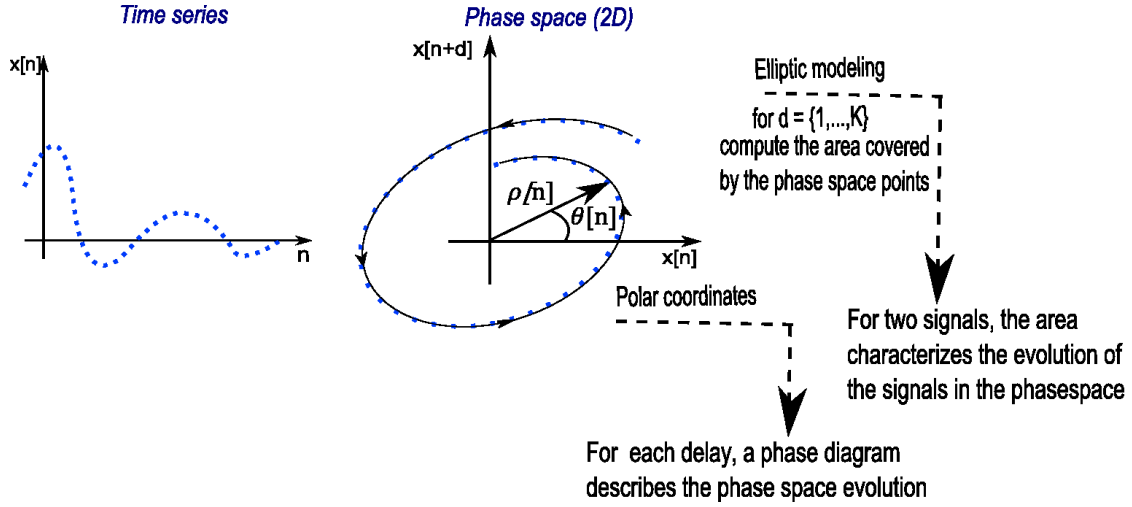


Fig. 4-7: The schematic presentation of the multi-lag tools

Fig. 4-7 presents the MPLA algorithm divided in two parts: the elliptic modeling and the polar coordinates representation, namely the phase diagram representation.

4.3. Diagonal lines quantification

The diagonal lines on the recurrence matrix indicate that the system returns in approximately the same state for a certain period of time. The choice of the diagonal lines quantification (DLQ - Diagonal Lines Quantification) is given by the fact that the distance between two successive diagonal lines provides the fundamental period of the signal (Digulescu, et al., 2016b), (Ioana, et al., 2014).

The quantification of these lines is done as follows:

$$dlq[n] = \frac{1}{M-n} \sum_{i=1}^{M-n} D^*_{i,i+n}, n = \overline{0, M-1} \quad (4.12)$$

where $D^*_{i,i+n} = 1 - D_{i,i+n} / \max\{D_{i,i+n}\}$ is the complement of the distance matrix.

Then, the Fourier transform is applied to this quantification:

$$DLQ[k] = \sum_{k=0}^{M-1} dlq[n] \cdot e^{-j2\pi kn/M} \quad (4.13)$$

where M is given in eq. (2.11). Fig. 4-8 presents the SD measure applied on a sine wave with the frequency $f = 10\text{Hz}$.

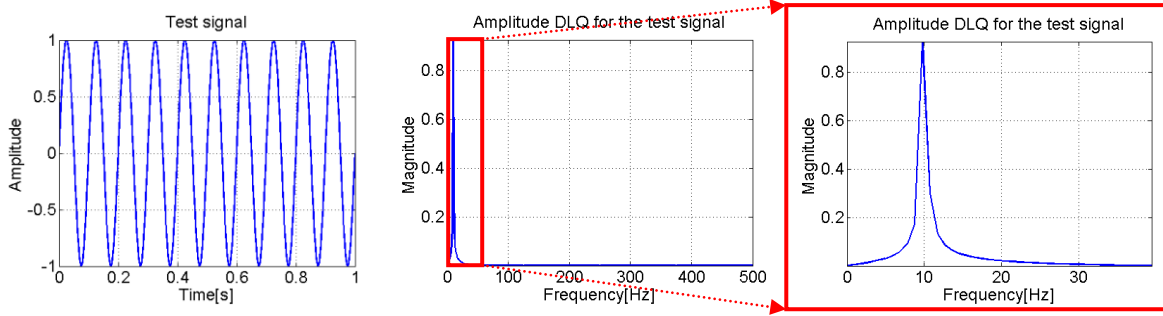


Fig. 4-8: Example of DLQ: (left) test signal – a sine wave with the frequency $f = 10\text{Hz}$; (center) the DLQ measure; (right) zoom on the DLQ measure;

This approach points out the frequency component in a different manner than the classical approach, where the Fourier transform is directly computed on the analyzed signal. In this case, the choice of the distance matrix is preferable to the recurrence matrix in order to follow the evolution of the signal entirely and not to separate it into a binary decision which usually brings in the spectral content various harmonics.

To illustrate this concept, let us consider a signal composed by the product of two cosines:

$$x[n] = \cos[\Omega n] \cdot \cos[\omega n], \Omega \gg \omega \quad (4.14)$$

To simplify the demonstration, we suppose that the embedding dimension $m = 2$ and the used distance is the squared Euclidean distance:

$$\vec{v}_i = x[i] \cdot \vec{i} + x[i+d] \cdot \vec{j} \quad (4.15)$$

Therefore the distance matrix depends on:

$$\begin{aligned} D_{i,i+n} &= (x[i] + x[i+n])^2 + (x[i+d] + x[(i+n)+d])^2 \\ &= E_i^2[n] + E_{i+d}^2[n] \end{aligned} \quad (4.16)$$

Replacing eq. (4.14) in eq.(4.16), it goes that:

$$E_i[n] = 2\alpha \cdot \sin \frac{\Omega[2i+n]}{2} \cos \frac{\omega[2i+n]}{2} \quad (4.17)$$

where $\alpha = \frac{\sin[(\Omega + \omega)n]}{2} \cong \frac{\sin[(\Omega - \omega)n]}{2}$ considering the relationship between the frequencies from eq. (4.14).

Using eq. (4.17) in eq. (4.16), leads to:

$$E_i^2[n] = \alpha^2 \left\{ 1 + \cos \left[2\omega \left(i + \frac{n}{2} \right) \right] - \cos \left[2\Omega \left(i + \frac{n}{2} \right) \right] - \cos[\Omega(2i+n)] \cos[\omega(2i+n)] \right\} \quad (4.18)$$

The eq. (4.18) develops the spectral representation of the diagonal lines quantification. Together with the $\Omega \pm \omega$ component that it is also found in the classical approach, in this representation there are present two other components: a component at 2ω and another at 2Ω .

To exemplify our approach, let $x[n]$ be a test signal:

$$x[n] = \cos[2\pi f_1 n] \cdot \cos[2\pi f_2 n] \quad (4.19)$$

where $f_1 = 500\text{Hz}$ and $f_2 = 39\text{Hz}$, conserving the condition $\Omega \gg \omega$. The test signal is shown in Fig. 4-9.

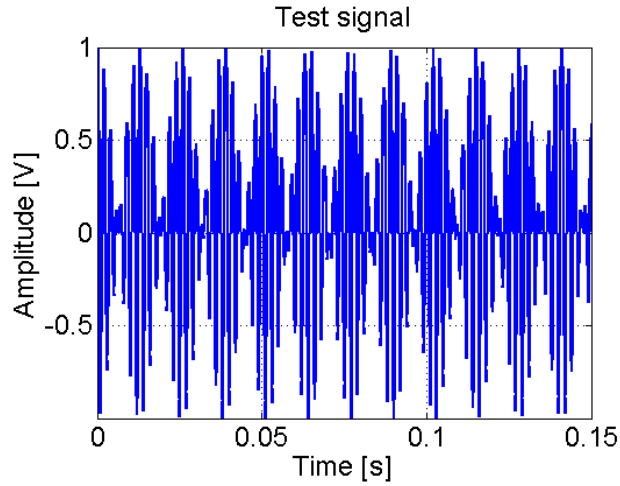


Fig. 4-9: The time evolution of the test signal modulated in amplitude

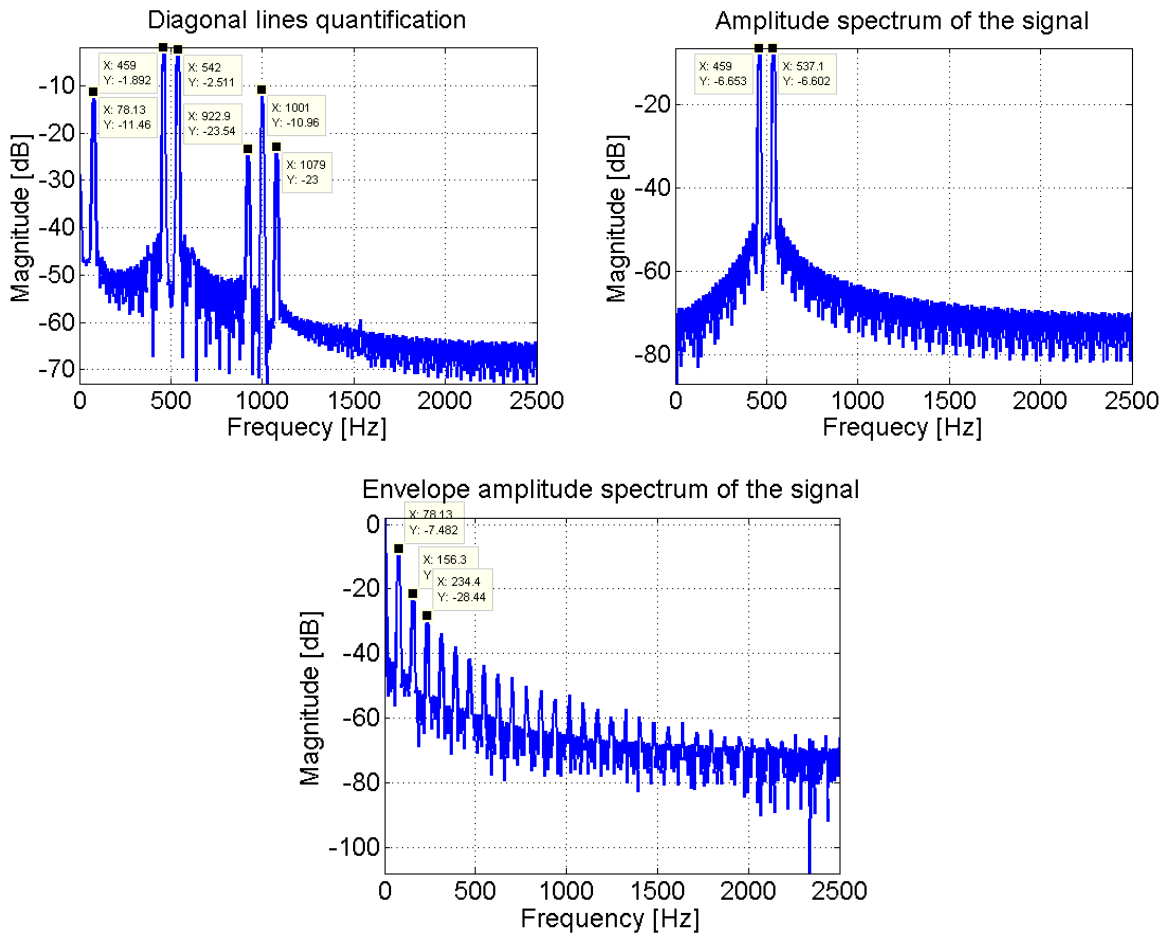


Fig. 4-10: The spectral representation of the diagonal lines quantification (up - left), the classical amplitude spectrum representation (up - right) and the envelope amplitude spectrum (down - center); for the computation of the distance matrix, it is used the squared Euclidean distance, $m = 3$ and $d = 2$

Fig. 4-10 presents the new information brought by this new approach: the diagonal lines representation contains three important components ($2\omega, \Omega \pm \omega, 2\Omega$), while the amplitude spectrum of the proposed signal presents only the $\Omega \pm \omega$ component and the envelope amplitude spectrum contains only the envelope frequency information, but also multiple harmonics of its envelope.

The displayed data tips show the values of the components: $f = 78\text{Hz}$, two components at $f = 500 \pm 39\text{Hz}$ and the third at $f = 2(500 \pm 39)\text{Hz}$ and $f = 2 \cdot 500\text{Hz}$. It is obvious that the classical amplitude spectrum contains only the frequency $f = 500 \pm 39\text{Hz}$ and the envelope amplitude spectrum the frequency $f = k \cdot 78\text{Hz}$, $k = \{1, 2, 3, \dots\}$.

This new quantification is very useful for applications where wide band signals interact with low frequencies which is the case of vortex detection and characterization: the wide band signals are suitable hydraulic system monitoring (Candel, et al., 2014a), (Digulescu, et al., 2014d), (Digulescu, et al., 2016a), (Digulescu, et al., 2016b) and the frequencies induced by the vortex are significantly smaller than the wide band signals (Digulescu, et al., 2016a).

It should be noticed that, when $\Omega \gg \omega$, the frequency modulation amplitude spectrum is quite similar with the amplitude spectrum of the amplitude modulation. Therefore, we apply the diagonal line quantification on the next test signal (Fig. 4-11):

$$x[n] = \cos[2\pi f_1 n + \beta \cdot \cos[2\pi f_2 n]] \quad (4.20)$$

where $f_1 = 500\text{Hz}$ and $f_2 = 39\text{Hz}$ and $\beta = 2$. The condition $\Omega \gg \omega$ stands valid.

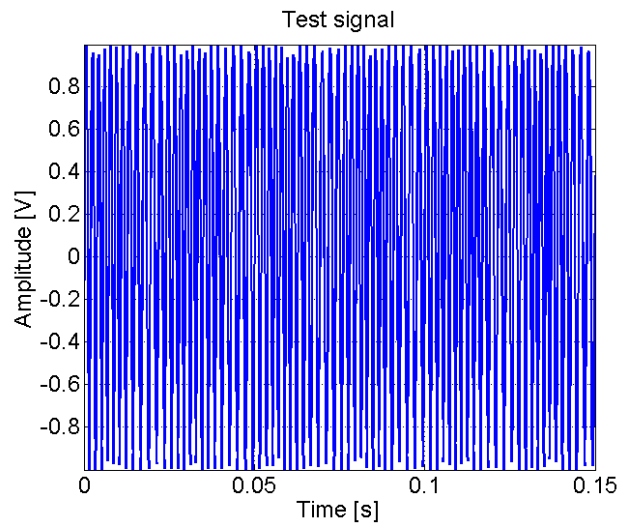


Fig. 4-11: The frequency modulated test signal

From Fig. 4-12, it can be noticed that in the case of a frequency modulated signal, the diagonal line quantification measure also emphasizes the classical components $\Omega \pm \omega$, but more important the extra component ω . Besides an extra component $2(\Omega \pm \omega)$.

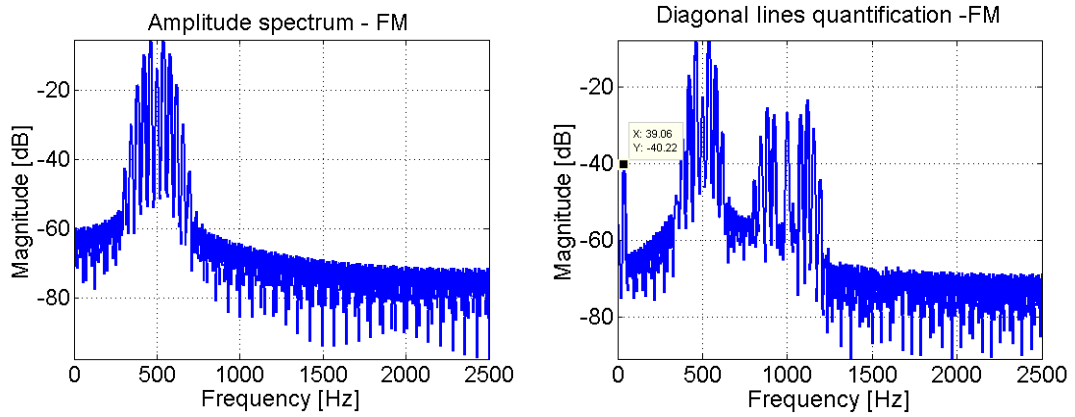


Fig. 4-12: The classical amplitude spectrum representation (left) and the spectral representation of the diagonal lines quantification (right); for the computation of the distance matrix, it is used the squared Euclidean distance, $m = 3$ and $d = 2$

This approach is very useful for applications where a phenomenon of interest is characterized by much lower spectral content than the investigating signal's one.

4.4. Chapter summary

This chapter presents new approaches for the detection, localization and characterization of transient signals using the RPA, RQA and the phase diagram analysis concept. The TDR measure is based in the evolution of the signal in phase space and it exploits the recurrence matrix properties in order to detect and localize transient signals.

Next, the MLPA is introduced based on the elliptic modeling and on the phase diagram representation. This type of analysis is very useful for the discrimination of signals with close parameters.

Finally, the diagonal lines quantification is introduced and, using synthetic signals, we emphasize the attributes that this type of quantification makes visible: the modulating component of an amplitude/ frequency modulated signal, even if they are much lower in frequency than the investigating signal's content.

5. Applications in active configurations

This chapter focuses on the experiments performed in active configuration. Firstly, we present the experimental configuration and, then, we show our results using the theoretical methods presented in the previous chapter. Moreover, these results are compared with classical methods.

5.1. Cavitating environments

The first applicative context is the detection of the cavitating vortex based on non-intrusive active sensing using an ultrasonic path defined by two wide band transducers. The objective of this application is to emphasize the effect of the cavitating vortex from the acoustic data in order to identify the different phases of the evolution of the phenomenon. Measurements were carried out in a laboratory facility which materializes the flow conditions downstream a hydraulic turbine.

The vortex flow and vortex spectrum

In literature the vortex vector is defined as $\nabla \times \mathbf{u} = 2 \cdot \boldsymbol{\omega}$, where $\boldsymbol{\omega}$ is the rotational speed of the fluid (namely the local spinning motion in a given point) and \mathbf{u} is the local flow velocity vector. The flow for an incompressible fluid of velocity is described by the Navier-Stokes (5.1) and continuity equations (5.2) (Emanuel, 1984), (Takahashi & Miyazaki, 2006):

$$\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} = -\frac{1}{\rho} \nabla p + \nu \nabla^2 \mathbf{u} \quad (5.1)$$

$$\nabla \cdot \mathbf{u} = 0 \quad (5.2)$$

where p is the fluid pressure, ρ is the density of the fluid, \mathbf{u} is the flow velocity vector, ν is the fluid viscosity and ∇ is del operator.

For the vortex analysis, the Rankine model can be used. Analytically, the Rankine vortex is defined by the components from eq. (5.3) (Takahashi & Miyazaki, 2006):

$$u_r = 0; u_z = 0; u_\theta = \begin{cases} \frac{r}{r_{\max}} v_{\max}, & r \leq r_{\max} \\ \frac{r_{\max}}{r} v_{\max}, & r > r_{\max} \end{cases} \quad (5.3)$$

where u_r , u_z and u_θ are the radial, axial and, respectively, tangential velocity components defined in a cylindrical coordinate system for the \mathbf{u} flow velocity. It should be mentioned

that, when $r \leq r_{\max}$, the tangential velocity u progresses in the forced vortex region, while, when $r > r_{\max}$, its progress happens in the free vortex region.

The frequency signature of the cavitating vortex is given by the solution (5.3) of equations (5.1) and (5.2) through the solution ω . Hereby, the relation between the vortex flow and its spectrum is:

$$u_r = 0; u_z = 0; u_\theta = \begin{cases} \omega \cdot r, & r \leq r_{\max} \\ \omega \cdot r_{\max}^2 / r, & r > r_{\max} \end{cases} \quad (5.4)$$

In (Rubin, 2000), there can be found more details about the spectral magnitude of Doppler velocity. For our experiments, this theoretical result shows that the vortex conduct to the generation of a low frequency component that our methodology aims to exploit in order to assess this phenomenon.

5.1.1. Experimental configuration and results

The experiments have been conducted in a reduced scale facility which consisted of a hydraulic installation where the flow rate can be manually changed (Candel, et al., 2014a). The acoustic transceivers operate around at 1 MHz and the emitted waveform is a wide band signal with the same central frequency. This signal, $v(\Omega)$ has a linear characteristic law.

$$v(\Omega, t) = \cos(\theta(t)), \text{ where } \theta(t) = 2\pi\alpha t^2 + 2\pi f_0 t + \phi(t) \text{ and } \Omega(t) = \frac{d\theta(t)}{dt} \quad (5.5)$$

where $f_0 = 800 \text{ kHz}$, $t \in [0, 100 \mu\text{s}]$, the coefficient α is chosen such that the final frequency of the sweep is $f = 1200 \text{ kHz}$.

Considering that the cavitating vortex is characterized by the unknown law $u(\omega)$ and the wide band signals (that are continuously emitted) are characterized by $v(\Omega)$, then, at the receiver, the signal carries the interaction between the emitted signal and the studied cavitating vortex in $w(\Omega, \omega)$. The purpose of the experiment is to highlight the signature of the cavitating vortex from the received signal.

The experiments are performed on two different pipelines; therefore, there are two configurations to consider. The pipeline of the installation has been placed inside a stator which produced the cavitation vortex as the flow rate increases, Fig. 5-1.

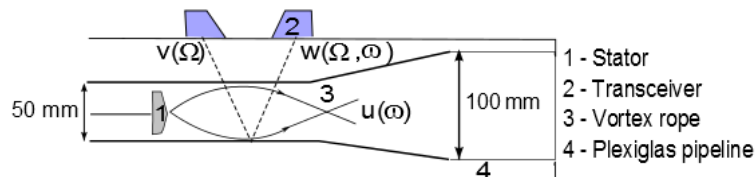


Fig. 5-1: The installation for the first set of measurements

There are two experimental setups that we performed: for the first one, the flow rate of the hydraulic system varies and, when the flow rate becomes stable, the acoustic active

system is started at this operating point. In the same time, the flow rate of the system is recorded. Moreover, on the tip of the stator a pressure sensor is placed so that we have a reference for the cavitation vortex (more the pressure is dropping, higher the intensity of the cavitation is). Fig. 5-2 presents the setup of the experiment.

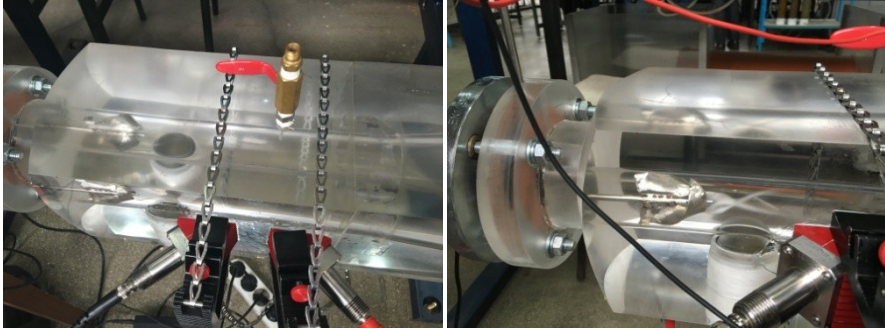


Fig. 5-2: The on-site setup for the first set of measurements

The second experimental setup is performed after the divergent region of the Plexiglas pipeline, as shown in Fig. 5-3 and Fig. 5-4.

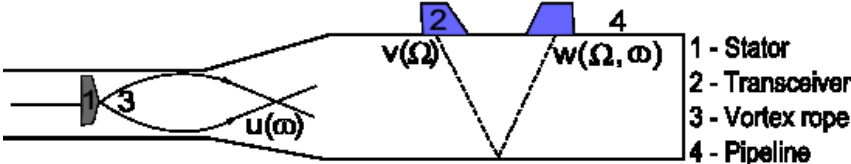


Fig. 5-3: The installation for the second set of measurements

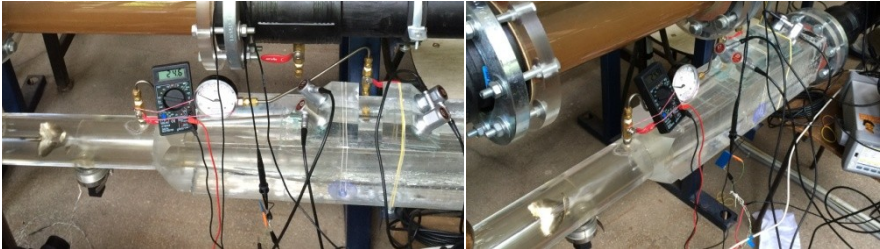


Fig. 5-4: The on-site setup for the second set of measurements

The experiment is done while slowly varying the flow rate of the water. In other words, during the experiment, the received signal is acquired and recorded in bursts of 2.6 ms simultaneously with the flow rate for approximately 90 s.



Fig. 5-5: Cavitating vortex stages: A1 no vortex (left), A2 incipient vortex (center), A3 developed vortex (right)

For both set of measurements, based on the visual observation (Fig. 5-5), three stages of the vortex evolution were identified: A1 region - no vortex presence (up to 6 l/s), A2 region incipient vortex (from 6 to 7.2 l/s) and A3 region developed vortex (more than 7.2 l/s).

Relating the cavitating vortex stages to the deformation operator, we recall that:

- when the water is monophasic: the deformation operator \mathcal{D}_k is defined by the attenuation and delay:

$$\mathcal{D}_k \{s(t)\} = \alpha_k \cdot s(t - \tau_k) \quad (5.6)$$

where k is the index of the propagation path. Therefore, the received signal is:

$$r(t) = \sum_k \alpha_k \cdot s(t - \tau_k) + n(t)$$

- when the cavitation appears: then, the deformation operator \mathcal{D}_k is characterized by attenuation, delay, as well as the Doppler effect that acts as a deformation of the instantaneous phase of the signal:

$$\mathcal{D}_k \{s(t)\} = \alpha_k \cdot s(t - \tau_k(t)) = \alpha_k \cdot A \cdot \exp(j \cdot \phi_k(t)) \quad (5.7)$$

where $\phi_k(t) = \phi \left(t - \frac{\bar{r}_k(t)}{c} \right)$ and $\bar{r}_k(t)$ is the time (Doppler) dependent position vector.

Based on this latter observation, we aim to detect and characterize this type of environment.

First experimental setup

This experiment was performed in March, 2016 at the Power Engineering Faculty, "Politehnica" University of Bucharest. The acoustic transducers are placed in "V" configuration for a longer interaction with the vortex. The emission transducer, the closest to the stator, generates 1 ms linear frequency modulations in a spectral bandwidth range of [0.8,1.2] MHz with a repetition rate of 2.5 ms. The sampling frequency is 5 MHz.

Fig. 5-6 presents the wide band acoustic signal used for the characterization of the cavitation vortex and two received signals in different states of the system: for the incipient cavitation, and for a fully developed cavitation, respectively.

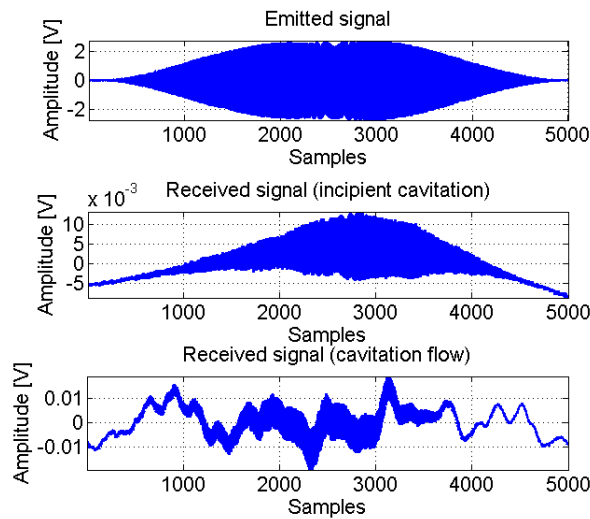


Fig. 5-6: The emitted signal for each measurement, a received signal during a trial for the incipient cavitation and a received signal during a trial for the cavitation flow

In Fig. 5-7 and Fig. 5-8, the limitations of the Fourier transform applied on the signal and its envelope are presented. The signal considered is acquired when the cavitation is clear eye visible (cavitation flow).

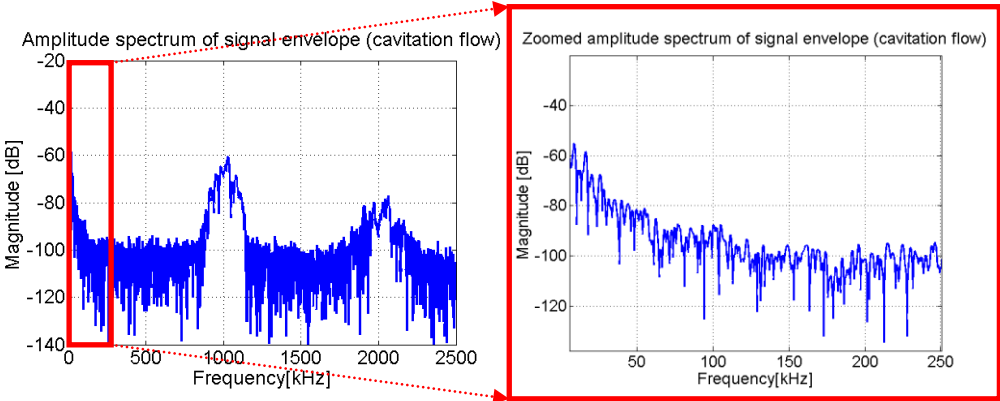


Fig. 5-7: The amplitude spectrum of the envelope for a received signal when the cavitation vortex is fully developed

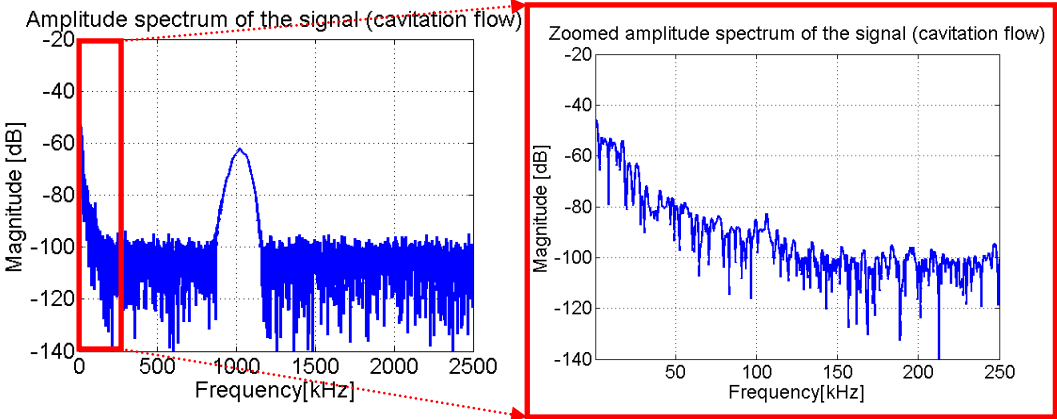


Fig. 5-8: The amplitude spectrum for a received signal when the cavitation vortex is fully developed

This approach is based on the assumption made in eq.(4.19), namely that the spectral content of the cavitation has a much lower frequency content than the wide band acoustic signal of the transceivers, $\Omega \gg \omega$. It can be seen that the associated phenomena (attenuation, dispersion, noise, etc.) hide the effect induced by the particle motion from the cavitation vortex, they only affect the low frequency content (around hundreds of Hz).

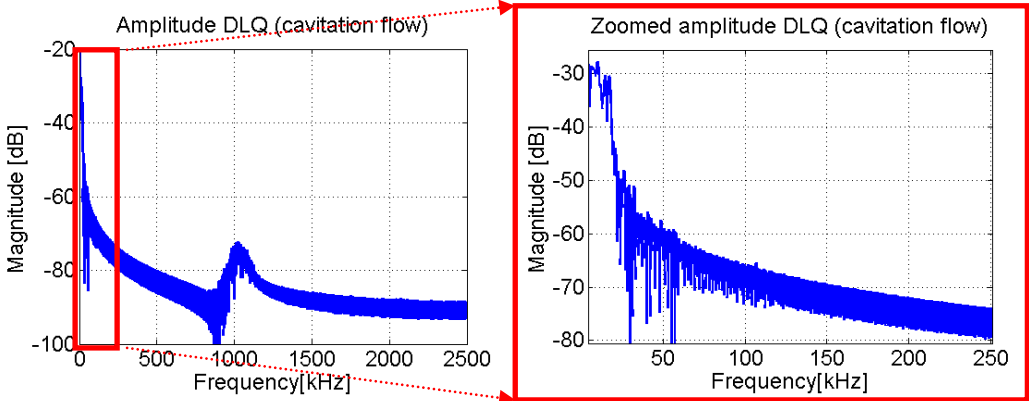


Fig. 5-9: The amplitude DLQ (eq. (4.13)) for an acquired signal when the cavitation vortex is fully developed

On the other hand, the use of the diagonal line quantification (eq.(4.13)) highlights, on the same signal, the cavitation vortex spectral content induced in the received signal when the cavitation is fully developed (Fig. 5-9).

Next, we applied our approach for all acquired signals in each configuration. Table 5-1 presents the system's characterizing parameters for each operating point of the system.

Operating point no.	Parameters		
	Flow rate (l/s)	Mean pressure (bar)	Pressure standard deviation (Pa)
3	1.14	1.23	1936
4	4.34	1.19	2437
5	9.2	1.05	3447
6	12.86	0.83	5784
7	14.77	0.64	6208
8	15.63	0.55	6985
9	16.01	0.51	7136
10	16.08	0.49	7039

Table 5-1: System's parameters for each operating point

It should be noticed that the greater the number of the configuration is, the greater the flow rate of the hydraulic system is. This has as a direct consequence the appearance of the cavitation which is emphasized by the decreasing of the pressure and the increasing of its standard deviation.

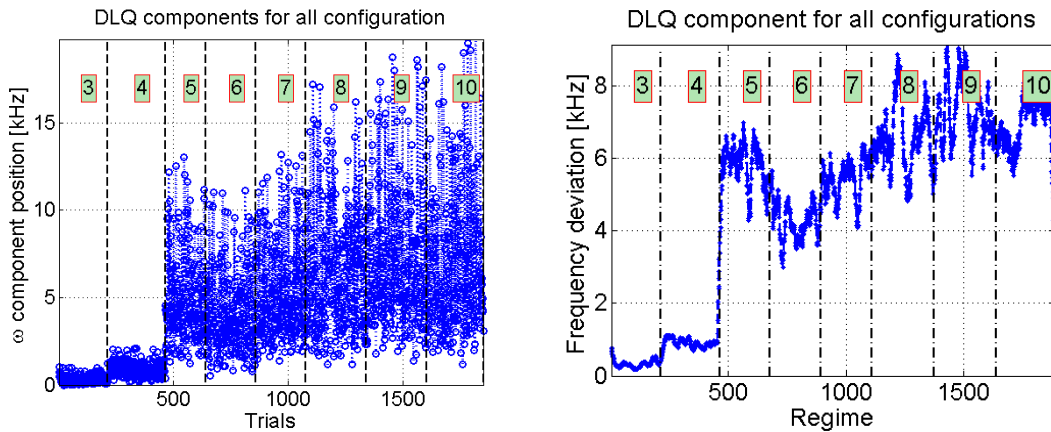


Fig. 5-10: The DLQ frequency position (left) and its mean (for 20 samples - right) using the diagonal lines quantification approach for the signals acquired at each operating point (the number in red box indicates the configuration); the RPA parameters are $m = 5$ and $d = 3$ using the Euclidean distance

As Fig. 5-10 presents, for the operating point no. 3, there is no cavitation. In operating point no. 4, the mean of the frequency deviation increases as well as its standard deviation, then, from operating point no. 5, the mean and the standard deviation of the frequency component induced by the particle motion increase gradually. Moreover, the evolution of the frequency component is strongly connected to the cavitating vortex stages shown in Fig. 5-5.

Therefore, for each operating point, we compute the mean of the frequency component and its standard deviation.

Operating point no.	Parameters	
	Mean of the frequency component (kHz)	Standard deviation of the frequency component (kHz)
3	0.84	0.23
4	1.42	0.39
5	5.65	2.6
6	4.7	2.3
7	5.9	2.7
8	7.07	3.42
9	7.75	3.74
10	7.48	3.85

Table 5-2: Frequency content parameters

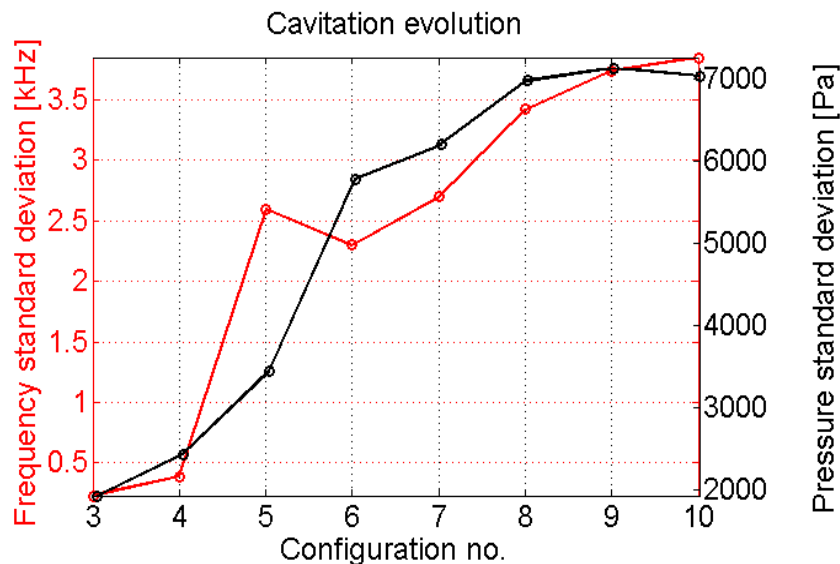


Fig. 5-11: The cavitation evolution characterized in terms of frequency standard deviation (in red) and pressure standard deviation (in black)

From Table 5-2 and Fig. 5-11, it can be seen that the mean of the frequency component induced by the particle motion has the same trend as the flow rate of the system, whereas the standard deviation of this component (non-intrusive technique) emphasized the cavitation phenomenon and has the same trend as the pressure standard deviation (intrusive technique).

We can then conclude that the cavitating vortex has a spectral component (deformation) which can be extracted using the diagonal lines quantification. This component describes the evolution of flow and of the cavitating vortex.

Second experimental setup

This experiment was performed in December, 2015 at the Power Engineering Faculty, "Politehnica" University of Bucharest. In this setup, the stator is placed 70 mm upstream the divergent section (previously, it was situated at 250 mm). For the acoustic analysis, two pairs of transceivers positioned in "V" configuration were used (Fig. 5-3 and Fig. 5-4) instead of only one. Moreover, the transducers were placed further downstream from the cavitating vortex area, on the divergent section (Fig. 5-3), while in the previous studies they were placed on top of the cavitating vortex area.

The first set of data is obtained for the appearance and developing of the cavitating vortex during discharge increasing. In order to cover the damping behavior of the cavitating vortex from a hydraulic turbine, the variation of the discharge is reversed in the setup, so the effect of the cavitating vortex disappearance was created.

We recall that the experiment is performed with a continuous slow variation of the water flow. During the experiment, the water flow is varied from $6\text{ to }19\text{ l/s}$ and backwards, in order to determine the critical flow rates for which the cavitating vortex appears/dampens.

Using the property of the diagonal line quantification, the position of the ω component (attributed to the presence of cavitating vortex) from the absolute value of SD measure is determined for each burst of 2.6 ms in each vortex region. Further, an example of ω component identification is presented in Fig. 5-12 and Fig. 5-13.

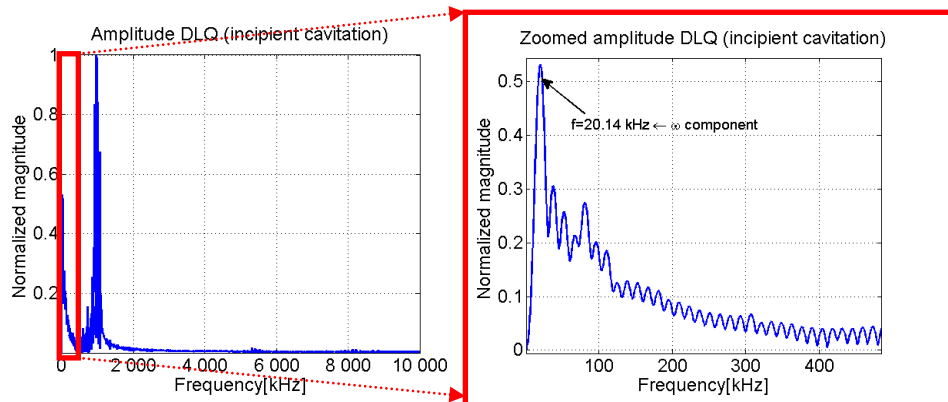


Fig. 5-12: DLQ absolute value for an acquisition of 2.6ms in A1 region

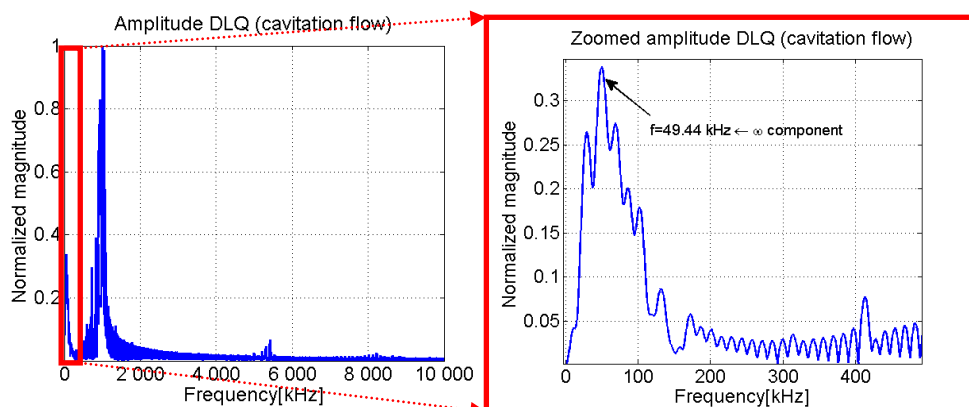


Fig. 5-13: DLQ absolute value for an acquisition of 2.6ms in A3 region

The main component at 1 MHz corresponds to the central frequency of the emitted signal at the acoustic sensors. Other significant components can be found in the range of 0–100 kHz and are attributed to ω component, as presented in the zoomed details.

The frequencies corresponding to the ω component are further presented as the time variation of the ω spectral component in Fig. 5-14. Also, the flow variation was represented simultaneously with time variation of ω spectral component in order to identify the three cavitating vortex regions.

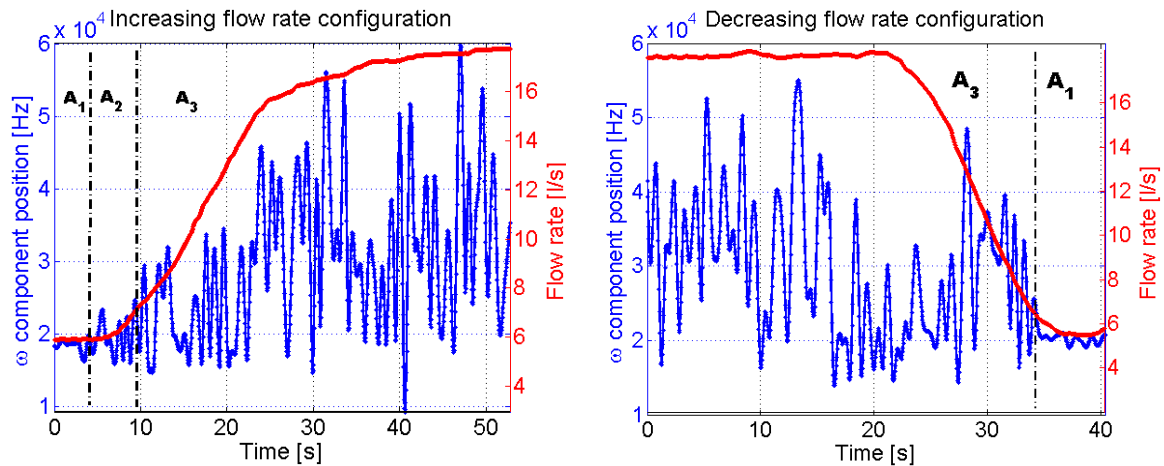


Fig. 5-14: Time variation of cavitation vortex component (in blue) and the flow rate evolution (in red); the RPA parameters are $m = 5$ and $d = 15$ using the Euclidean distance

For the increasing flow rate configuration, it can be seen that, while there is no vortex (A1 region), the cavitation vortex component has a reduced variation. When the vortex starts to appear in A2 region, the deviation of the frequency component is larger and is significantly increasing in A3 region for the developed vortex. The standard deviations of the cavitation vortex component, σ_ω , for each region are computed, and cavitating vortex characterization criteria are defined according to them (Table 5-3).

Region	σ_ω [kHz]	Limits[kHz]	Vortex stage
A1	0.76	$\sigma_\omega < 1$	No vortex
A2	2.26	$1 \leq \sigma_\omega \leq 3$	Incipient vortex
A3	9.71	$\sigma_\omega > 3$	Developed vortex

Table 5-3: Criteria limits for the cavitating vortex stages classification in increasing flow rate configuration

The validation of the criteria is done using recorded data during decreasing water flow configuration. The cavitation vortex component is determined and its time variation was plotted together with the water flow variation in Fig. 5-14 (left). The cavitating vortex regions are identified by visual observations and fitted with the defined criteria (Table 5-3).

Region	σ_{ω} [kHz]	Vortex stage
A1	0.74	No vortex
A2	-	-
A3	9.13	Developed vortex

Table 5-4: Criteria validation for decreasing flow rate configuration

In this case, the incipient vortex stage (A2 region) was not visible, which was also confirmed by the high value of standard deviation of the cavitation vortex frequency component computed. This may be caused by the inertia of the fluid while decreasing the flow, which was done in a shorter period than in the increasing case (about 15 s compared to 30 s).

Hereby, with this second setup, we observe that the cavitating vortex signature (deformation) is present in the spectral components of the diagonal lines quantification. This component describes the evolution of flow, as well as the cavitating vortex.

5.2. Underwater object tracking

This section presents a new method for the detection of an underwater object using bi-static acoustic paths. The method uses the dynamic evolution in the phase space of the received signals without any synchronization between the transmitted and received signal.

Generally, the object detection using acoustic barriers is based on the amplitude fading when the objects arrive between the emitting and receiving transducers (Digulescu, et al., 2014a), (Ioana, Candel, & Digulescu, 2015). Nevertheless, when the objects have close properties, in terms of density, to the water ones, the amplitude changes induced by the objects are very light and, for this reason, the objects cannot be always detected. This is the case of objects such as pockets of oil, fishes, etc.

In this chapter, an alternative detection method is proposed, based on the accurate estimation of the time of flight between the transmitter and the receiver. Actually, the aim is to take advantage of the fact that the waves traveling the objects have a different velocity than in water. At the receiving level, it comes that the wave will arrive with a different time of flight than in the case when the waves propagate only in water. This technique requires an accurate estimation of the time of flight and we will show that this information is more accurately estimated than from the classical envelope-based detection method.

The experimental setup

The experiment is performed in a reduced scale facility where the underwater object travels for 3 s on a distance of 1.5 m . The underwater object is an oil balloon entirely immersed, crossing the acoustic paths upon moving. The experimental configuration is presented in Fig. 5-15.

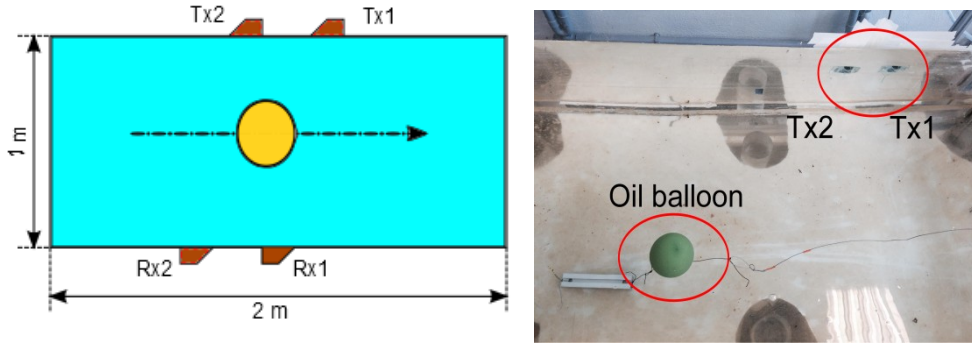


Fig. 5-15: The experimental configuration: the dash-dot line shows the trajectory of the object

All four acoustic transducers are identical and have a wide band reception centered on 1 MHz with a bandwidth that starts from 800 kHz and ends around 1200 kHz. The amplitude of the transmitted signal is 5 V. The emission transducers (Tx1 and Tx2) are placed on one side of the water tank and generate linear chirp burst in their bandwidth at an angle of 45°. On the opposite side of the water tank, the receiving sensors are placed at a minimum distance such that crosstalk between the acoustic paths is avoided in the clear water (when no object is in the water tank). The data acquisition is performed for duration of approximately 3.35 s in order to record the entire process.

In this case, the deformation operator that we look to estimate is:

$$\mathcal{D}_k \{s(t)\} = \alpha_k \cdot A \cdot \exp \left(j \cdot \phi \left(t - \frac{r_k}{c_{water}} - \frac{d_{r_k}}{c_{oil}} - \tau_0 \right) \right), k = \overline{1, N} \quad (5.8)$$

where r_k is the k^{th} water propagation path, d_{r_k} is the k^{th} oil propagation path and $\tau_0 = \frac{2l_b}{c_{rubber}}$ where l_b is the thickness of the rubber wall of the balloon and c_{rubber} is the acoustic propagation velocity in rubber.

Results

Taking into account that the passage of the object in front of the sensors, involves a change in the received time series, we aim to highlight this effect on the detection map based on the TDR^* measure (eq. (4.2)). Hereby, we define the detection map (Digulescu, et al., 2015a):

$$DM(l, \tau) = [TDR^*(l)]_{\tau=1, \frac{N}{s}} \quad (5.9)$$

where N is the length of the recorded signal and s is the size of the fixed time bin for which the TDR^* measure is computed.

The detection map contains both the amplitude (in the TDR^* measure) and time of arrival information (given by the time bin, τ). Therefore, this approach presents the time information, although the signals have no synchronization and their characterization is based only on the received signals by the two sensors.

For the detection part, we firstly present the signal recorded at Rx2 sensor (Fig. 5-15) which directly receives its bursts from Tx2. On this signal, Fig. 5-16, the detection map based on the TDR^* measure is applied, Fig. 5-17.

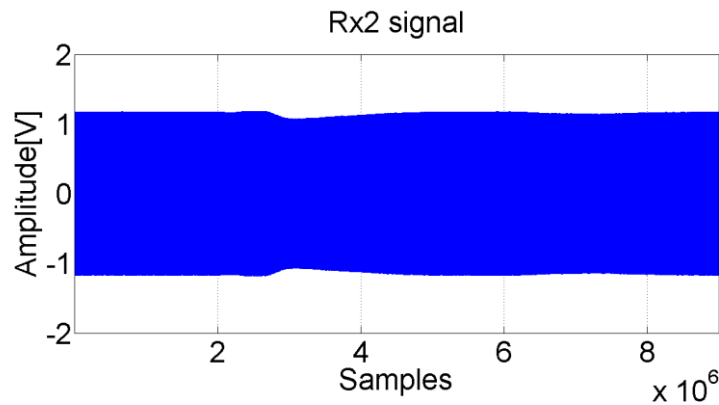


Fig. 5-16: The recorded signal arrived at Rx2 sensor

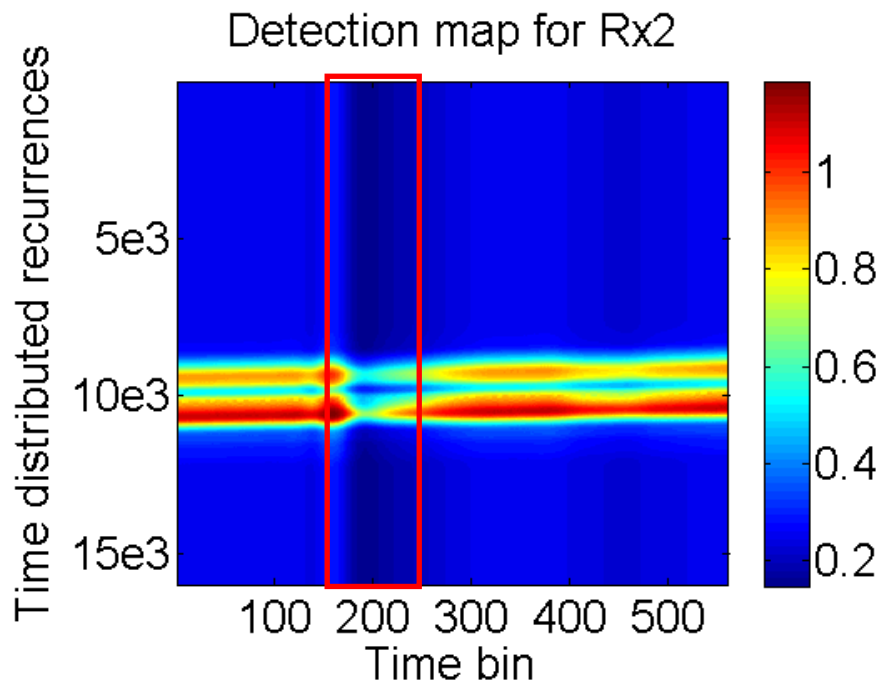


Fig. 5-17: The detection map obtained for the Rx2 signal for $m = 3$, $d = 13$ and $\alpha = 0.5$

It can be observed that the burst (the response of the object) is represented by the two brown horizontal lines which vary in amplitude (the average distance in the phase space changes), but also in time (the TOAs of the bursts in each time bin is not constant). Moreover, as the lines of the detection map represent the value of the TDR^* measure at each time bin, its columns contain the information regarding the TOAs of each burst at its corresponding time bin.

Therefore, in order to obtain the TOA information from the detection map (Fig. 5-17), we impose an appropriate threshold that discriminates between the part of the signal that has no burst and the one that has it. In this case, the threshold is given by the cyan region from the detection map. Hereby, the first point on each column of the detection map corresponding to the cyan region gives the TOA of the responses in each time bin. The results are presented in Fig. 5-18.

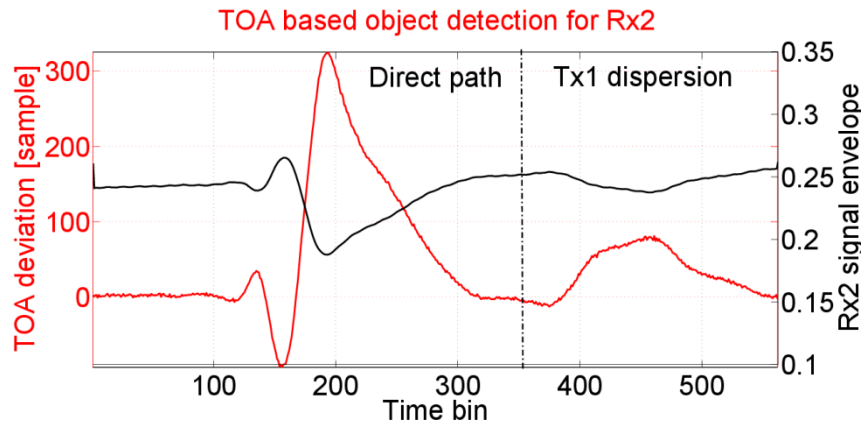


Fig. 5-18: The TOAs (red) of each response of Rx2 signal and its signal envelope (black)

Although, the signal's envelope does not offer reliable information regarding the object tracking (the amplitude variation can be caused by external vibrations, move of transducers, etc.), the TOA information is very robust, because it highlights the delay induced by the object. From Fig. 5-18, it can be seen that, even for a slight attenuation in signal's amplitude, the detection map approach is able to highlight a delay in the TOA of the burst for two time bin intervals.

This delay is explained by the fact that the speed of sound in the sunflower oil is smaller (around 1425 m/s) than the water speed of sound (around 1480 m/s) and its maximum corresponds to the delay induced by the oil ball diameter (15 cm) and the rubber wall (0.5 mm) of the balloon that folds the oil ball.

Moreover, in the TOA based object detection representation, the first interval corresponds to the object passing in front of the transducer Tx2 and the second corresponds to the crosstalk induced by the dispersion of the signals coming from Tx1. In this way, we suppose that the travel direction of the object is from Tx2 to Tx1. In order, to confirm our assumption, we also study the signals recorded by Rx1 (Fig. 5-19).

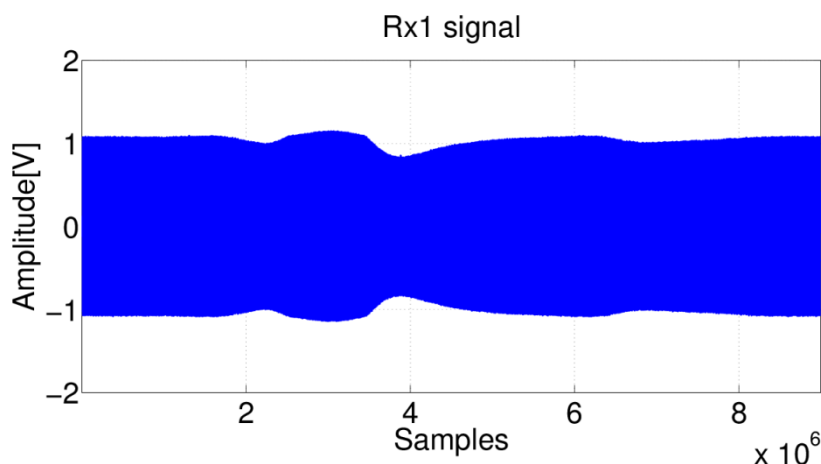


Fig. 5-19: The recorded signal arrived at Rx1 sensor

In this case, the signal is more disturbed in terms of amplitude. Fig. 5-20 and Fig. 5-21 present the detection map corresponding to the Rx1 signal and the TOAs of the responses

from each time bin. It seems that even if the amplitude of the signal is more disturbed, its TOAs present almost the same values for the positives peaks for both receivers.

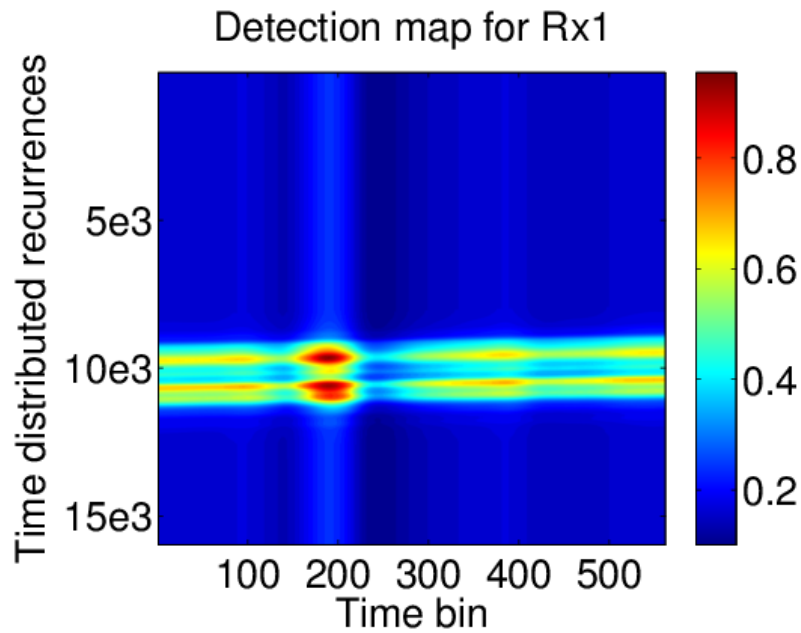


Fig. 5-20: The detection map obtained for the Rx1 signal for $m = 3$, $d = 7$ and $\alpha = 0.5$

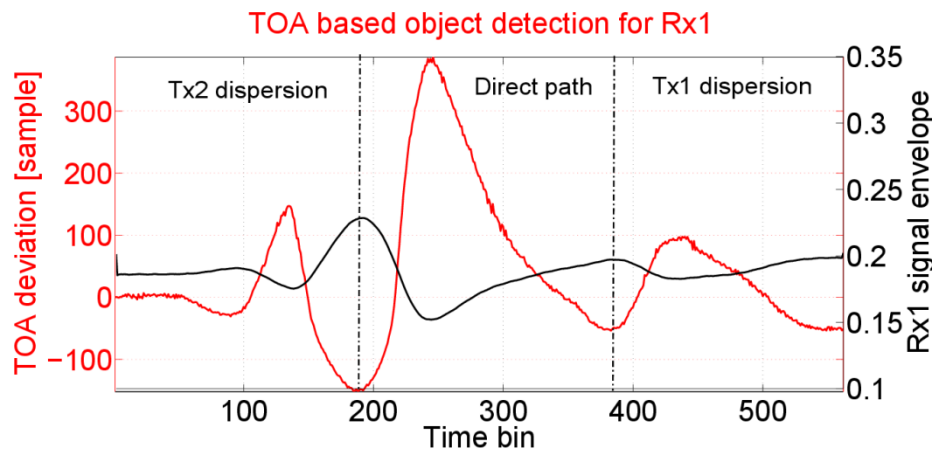


Fig. 5-21: The TOAs (red) of each response of Rx1 signal and its signal envelope (black)

The TOA representations from Fig. 5-18 and Fig. 5-21 point out approximately the same maximum delays induced by the object. This observation leads to the conclusion that one and the same object has passed in front of the sensors. Furthermore, the other variations of the TOAs are caused by the crosstalk that the signals' dispersion in the balloon induces to the sensors. We recall that, in the configuration with no object in the water, the sensors were placed at a minimum distance so that they record only the response coming from their corresponding transceiver (Tx2-Rx2 and Tx1-Rx1, respectively). But, in the presence of the object, its movement induces a cross-talk between the sensors.

In Fig. 5-18, two peaks can be observed: a maximum which automatically corresponds to the delay induced by the object that interferes with the direct path between the pair Tx2-Rx2, and the second peak corresponds to the cross-talk induced by Tx1 as the object travels towards its path. Also, in Fig. 5-21, three peaks are present: the maximum which corresponds

to the passing of the object through the direct path between Tx1- Rx1, a second peak present before the maximum which is caused by the reflections coming from Tx2 and a third peak caused by the reflections of Tx1. These observations lead to the conclusion that the object moves from left to right in the given configuration.

5.3. High speed sensing using phase diagram-based adaptive waveform

The approaches proposed in the previous sections use the active sensing to detect the objects such as bubbles of the cavitation and pockets of oils. The detection and the measurement of object properties are based on the estimation of received signal parameters such as TOFs or Doppler deformation. Generally, the current sensing is able to provide these parameters just for few periods that could be enough for understating the phenomena. Furthermore, if we look to have a more detailed image of these phenomena, the estimation of the signal’s parameters for a large number of period is necessary. This is actually what we call high speed sensing and this section presents the feasibility of this concept using the phase diagram-based adaptive waveform discussed in the Chapter 3.

The experiment is the essential part which confirms if this concept can be applied in real – life. Fig. 5-22 presents the considered IFL and its corresponding phase space lobe.

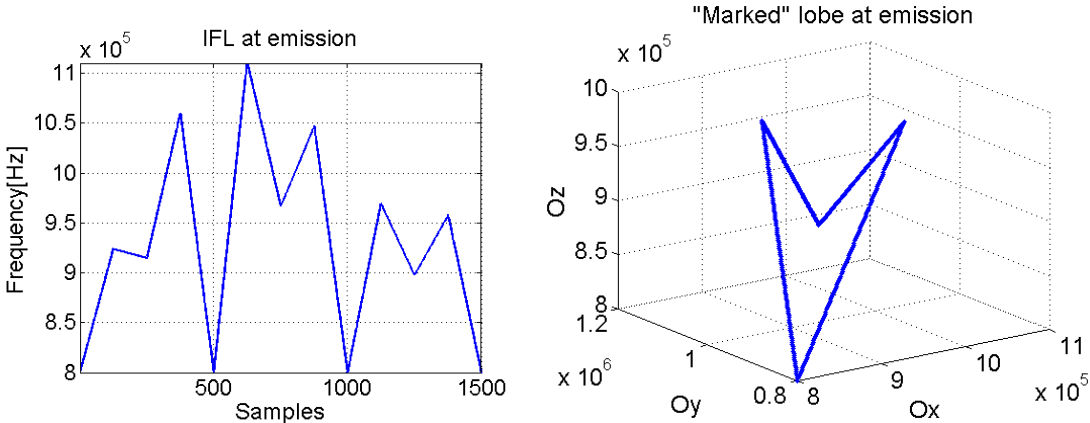


Fig. 5-22: The IFL used at emission (left) and its corresponding “marked” lobe (right)

The first test that we performed was to study if the characteristics of the signal containing this IFL remain unchanged when using two ultrasonic transducers for emission and, respectively, for reception. The two transducers are placed one in front of the other as shown in Fig. 5-23.

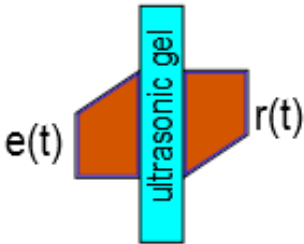


Fig. 5-23: The positioning of the two transducers for the experimental test

The signals used for emission and reception are shown in the next figure, Fig. 5-24.

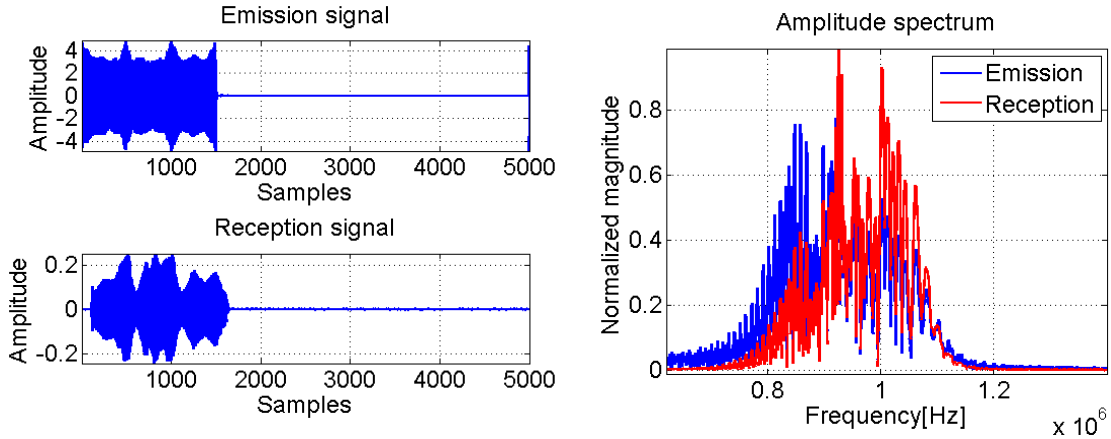


Fig. 5-24: The signal used at emission and the received signal as well as their amplitude spectrum

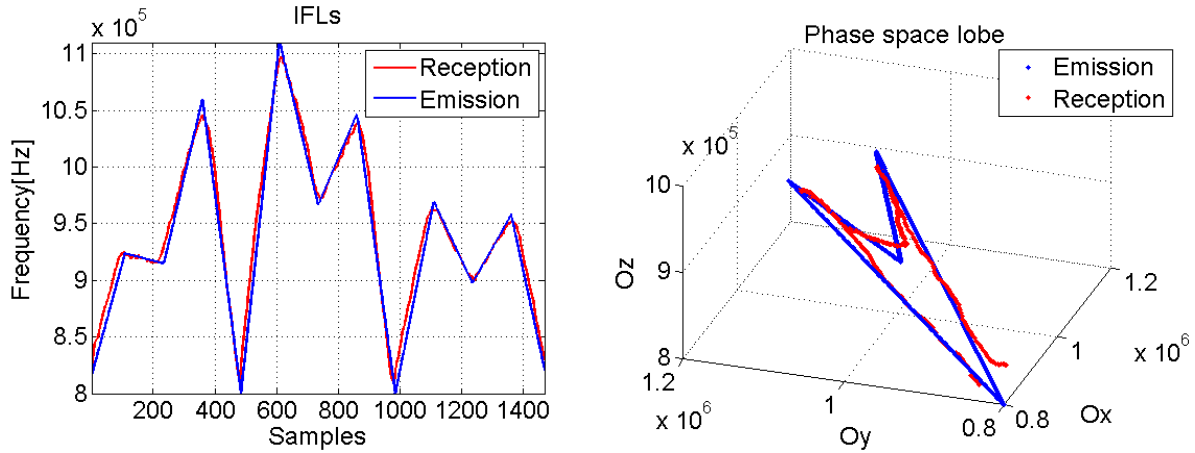


Fig. 5-25: The IFLs of the emitted and received signals (left) and their representation in phase space (right)

From Fig. 5-25, it seems that, although the transfer characteristic of the transducers eliminates part of the lower frequencies, the received signal maintains our proposed characteristics. Therefore, these new signals with “special signatures” in the phase space can be used in real-life applications.

The objective of the high-speed sensing using this waveform is to achieve a quasi-continuous emission which contains overlapped signals in time, but separable using their IFLs.

Hereby, considering the signal $s(t)$, where:

$$s[n] = s_1[n] + s_2[n - \beta] \quad (5.10)$$

where $s_1[n]$ and $s_2[n]$ are the two signals with a unique representation in phase space and β is a delay which provides an initial time separation between the signals. Considering the phase space representation of the phase space lobe, as shown in Fig. 5-26, β is equal with $2N$ where N is given in eq. (3.5).

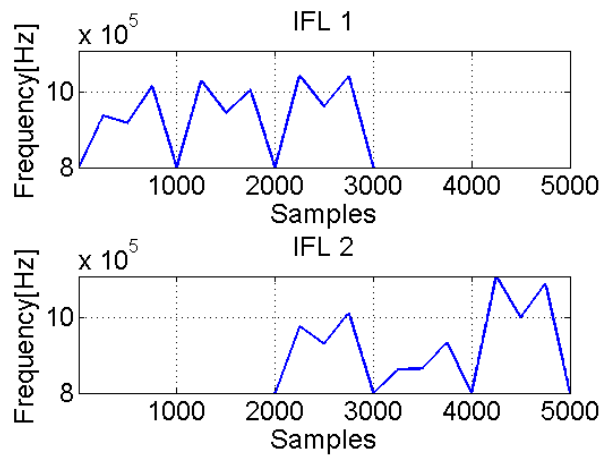


Fig. 5-26: The time shifted IFLs used for overlapping where $N = 1000$

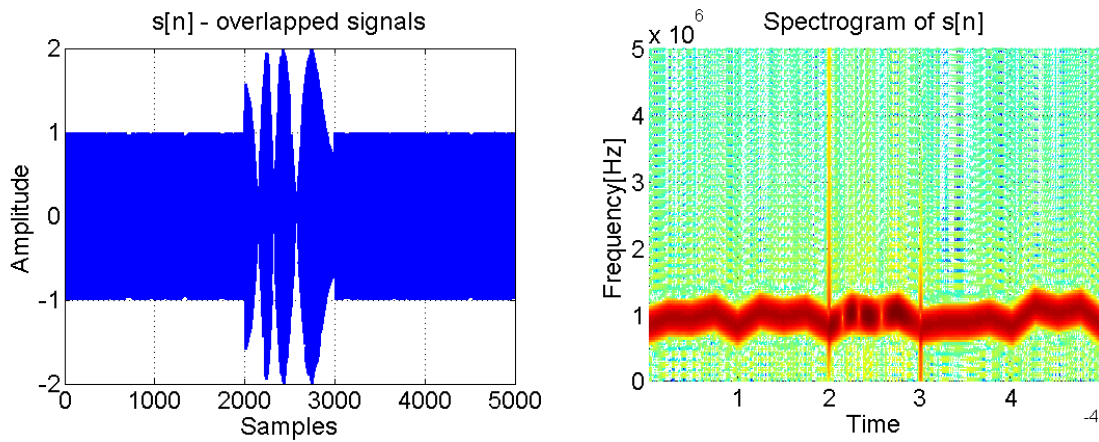


Fig. 5-27: The overlapped signals which form the signal $s[n]$ (left) and its spectrogram (right)

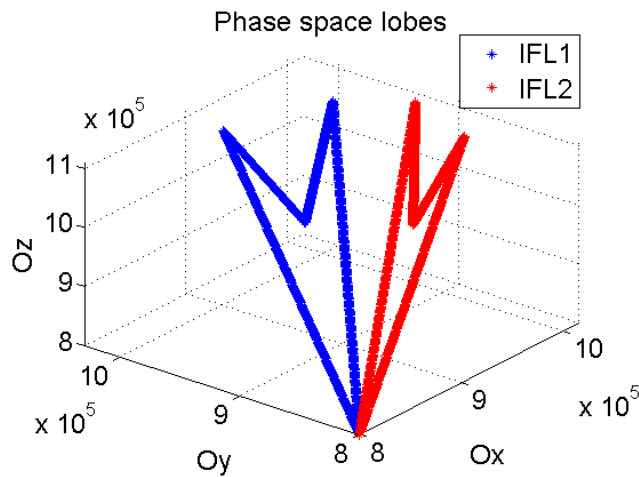


Fig. 5-28: Phase space lobes of the IFLs used for overlapping

Our objective is to separate the overlapped IFLs (Fig. 5-27) based on the algorithm presented in Chapter 3. Considering the first IFL known (IFL1), we want to separate the second IFL (IFL2) from the first, Fig. 5-28. Hereby, the algorithm is applied in the range $2N+1$ to $3N$.

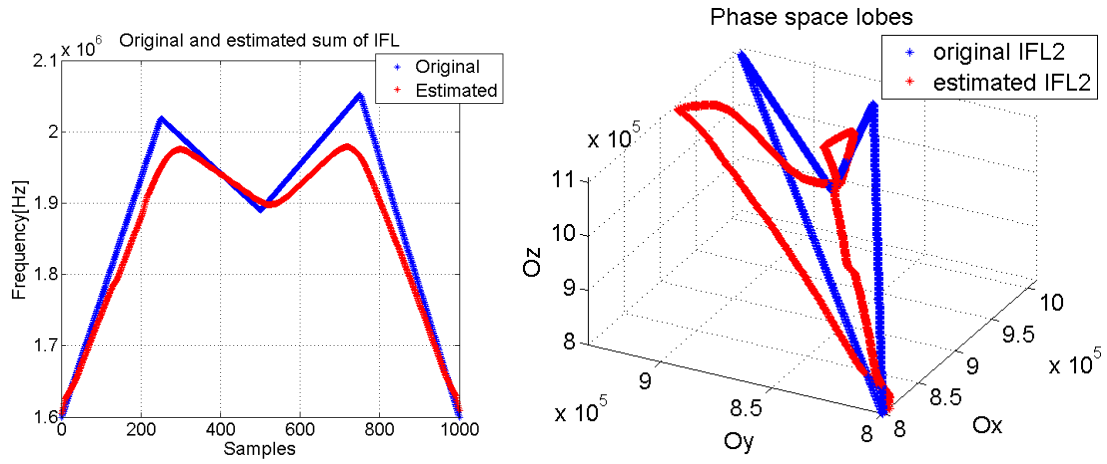


Fig. 5-29: The sum of the IFLs (left) and the phase space lobes of the IFL 2 (right)

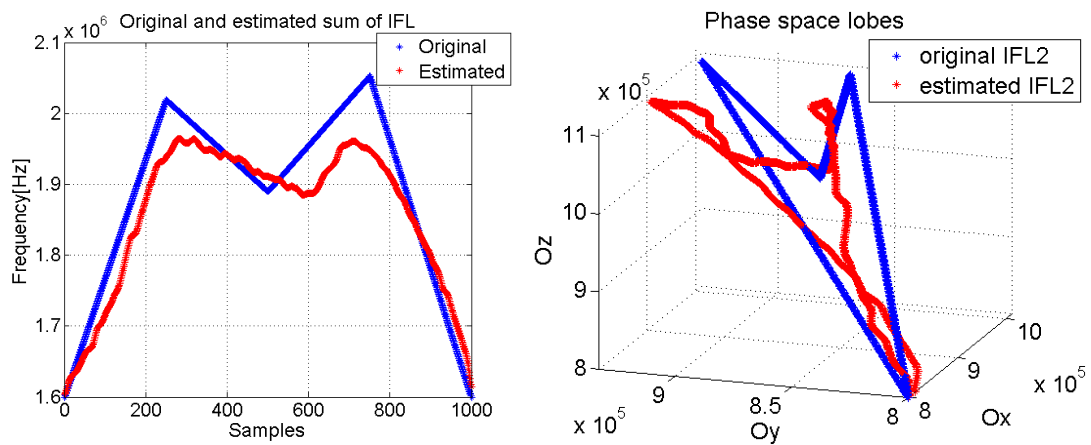


Fig. 5-30: The sum of the IFLs (original – no noise, estimated – SNR=25 dB) (left) and the phase space lobes of the IFL 2 (original – no noise, estimated – SNR=25 dB) (right)

This result, Fig. 5-29, shows that it is possible to separate two successive IFLs, which correspond to our high speed sensing sequence, by looking on their trajectory in phase space. Still, the method performs for quite high SNR (Fig. 5-30), the next research effort will attempt to improve its noise robustness.

5.4. Chapter summary

In this chapter, we described our experiments performed in active configurations using the theoretical methods introduced in this thesis.

For the cavitating environment application, we have shown that the generated vortex flow can be characterized in a non-intrusive manner using the diagonal lines quantification for of the received signal, when the phenomena are sensed with wide band signals.

The underwater object tracking is achieved using the TDR^* measure detection map. We point out the presence of the underwater object from the TOA parameter which is very robust because it highlights the delay induced by the object, although the signal's envelope is not that reliable.

These two applications can be extended to a high speed sensing concept presented after at the end of this Chapter, allowing to get the dynamic evolution of the cavitation process or the object's image, respectively. For this purpose, the phase diagram-based waveform design can be successfully used, as illustrated in the final part of this chapter.

6. Applications in passive configurations

This chapter presents the applications carried out in passive configuration. We present the experimental configurations and, then, we show the results obtained using methods proposed in the Chapter 4. Comparisons with classical method are performed.

6.1. *Measuring pressure transient in water pipes*

In this application, we present an acoustic signal processing approach for the surveillance of a water pipe and the characterization of the water hammer effect. The experiment is based on ultrasonic transceivers and the validation is made with intrusive pressure sensors. The major objective is to quantify the pressure variation into the pipe using a non-intrusive method that is usually used for the flow rate estimation (Wylie & Steerter, 1993). Therefore, the speed of the transient was determined using the time-of-flight (TOF) estimated by the RQA method. The estimation of the pressure using the same electronic device as for the flow rate estimation (Candel, et al., 2014), (Murgan, et al., 2016), (Petrut, et al., 2014) would be a great achievement since it offers to an existing device a second functionality.

In industrial applications, the water hammer must be supervised and controlled in order to minimize the destructive effects that it may cause to the hydraulic system. For this reason, the proposed approach that uses ultrasonic transceivers has the advantage to be suitable for most systems since it does not require any extra manipulations upon the system or pipeline.

6.1.1. Water hammer parameters

The water hammer is a periodic phenomenon caused by sudden maneuver of a closing vane (fully or partial). This forced stopping of the flowing fluid particles implies changing of their kinetic energy into pressure energy, so the pressure in the vane section will be increased with Δp .

The maximum overpressure in a pipe during water hammer due to sudden total closure of a vane can be determined using Joukowsky equation (Houghtalen, Akan, & Hwang, 2010), (Hachem & Schleiss, 2012), (Wylie & Steerter, 1993):

$$\Delta p = \rho \cdot c \cdot V_0 \quad (6.1)$$

where ρ is the fluid density, c is the pressure wave velocity and V_0 is the steady state flow velocity.

Some part of the energy is dissipated due to fluid viscosity (so to the hydraulic losses) which attenuates in time the water hammer phenomenon (Bucur, et al., 2014). The pressure waves travel inside the pipe along its length, L , between the closing vane and the reservoir, with high velocities, c , and have an oscillating period, T . Hereby, this period is given by:

$$c = \frac{4L}{T} \quad (6.2)$$

Also, the pressure wave speed, c , can be determined considering the characteristics of the fluid and of the pipe:

$$c = \frac{1}{\sqrt{\rho \left(\frac{1}{\varepsilon_w} + \frac{D}{e \cdot E} \right)}} \quad (6.3)$$

where ε_w is the bulk modulus of the fluid, D is the pipe diameter, e is the pipe wall thickness and E is the Young modulus of the pipe.

6.1.2. Experimental setup and results

The experimental setup (Fig. 6-1 and Fig. 6-2) consists in a horizontal pipe supplied by a tank (ST) with a total volume of 200 l . Together with an air compressor, it can create the constant supply pressures needed for the experiments.

The pipe has a length of $L=10.11 \text{ m}$, made entire of Plexiglas, with rectangular exterior section of $50 \times 50 \text{ mm}^2$ and circular interior section, $D=39 \text{ mm}$, used for high quality visualization of the process. The water is evacuated at the downstream end of the pipe through an 20 mm orifice diaphragm (d).

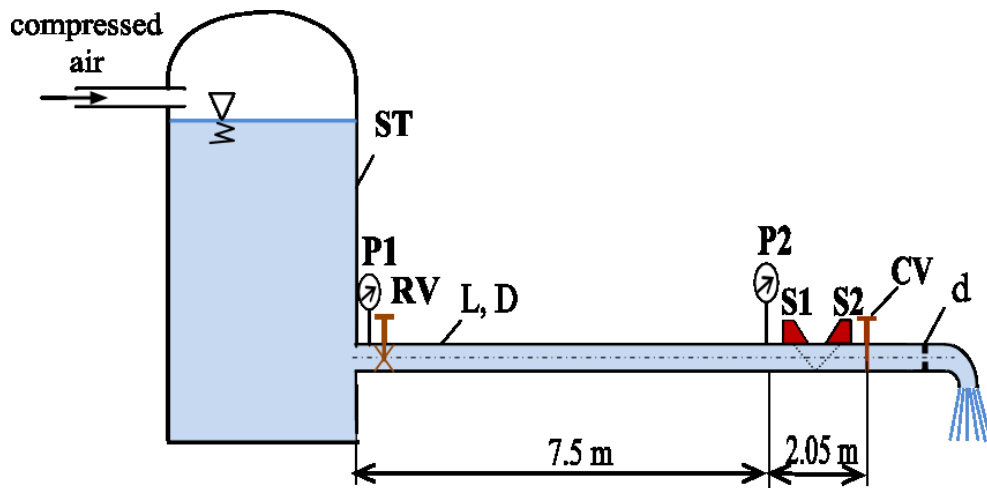


Fig. 6-1: Experimental setup

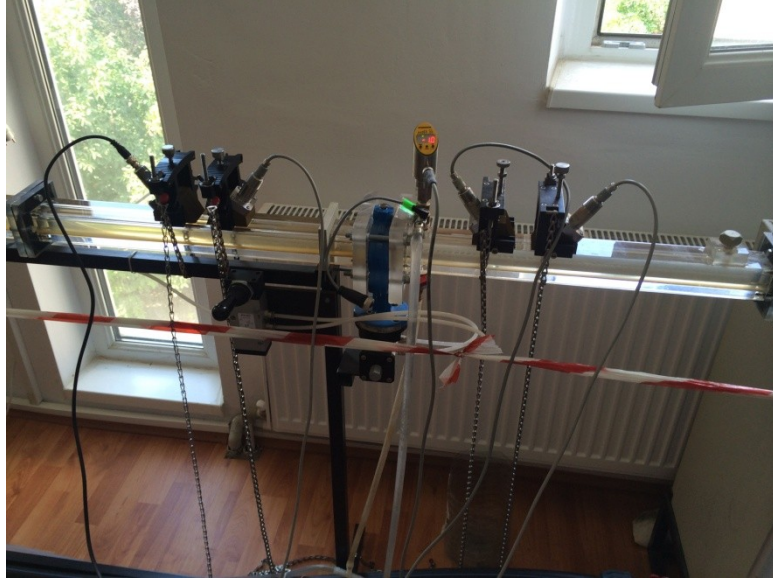


Fig. 6-2: The on-site experimental setup of the ultrasonic transceivers performed at Power Engineering Faculty, "Politehnica" University of Bucharest in December 2014

The discharge is regulated using a valve (RV) mounted at the entrance of the pipe, at 0.25 m downstream ST. The closing vane (CV) is a pneumatically driven butterfly with a pneumatic actuator mounted together with the valve. The valve closing time is approximately 20 ms obtained with 4 bar of the air pressure in the valves actuators pneumatic circuit.

The pressure transducer P2 is mounted upstream the fast closing vane, at 2.05 m . It has the measuring range between $0 - 40\text{ bar}$ relative pressure, with the output signal $0 - 10\text{ V}$, accuracy ± 0.25 and repeatability 0.1% of full scale (Fig. 6-1 and Fig. 6-2).

Also, two acoustic transceivers were placed upstream the closing vane at 8 cm and 16 cm respectively (Fig. 6-1 and Fig. 6-2). The acoustic transceivers have a wide band reception centered on 1 MHz , but being able to receive signals starting from 8 kHz . The output signal varies from -20 V to 20 V . Both transceivers were set in passive configuration and the acquisition was made at a sampling frequency of 50 MHz for a period of 0.67 s .

The transient regime (water hammer) during the tests was induced by the sudden closure of the vane CV. The initial conditions are 1.98 l/s flow rate and 1.4 bar supply pressure. The maximum pressure due to water hammer phenomenon recorded by P2 during tests was of 17.3 bar .

For this application, we recall that the deformation operator \mathcal{D}_k becomes:

$$\mathcal{D}_k \{s(t)\} = \alpha_k \cdot A \cdot \exp \left(j \cdot \phi \left(t - \frac{L_k}{c_{pressure}} \right) \right) \quad (6.4)$$

where $c_{pressure}$ is the propagation speed of the pressure field and L_k is the distance between the CV and the ultrasonic transducers.

This deformation operator translates, at the transducers, transient signals which offer the TOF information needed for the non-invasive pressure variation estimation.

Results using the RQA-based signal processing approach

The acquired acoustic signals at both sensors are presented in fig. 3. As the dashed blue ellipse highlights, it can be observed that the acoustic transient happen simultaneously with the pressure spikes. Moreover, as S_2 is closer to the CV (approx. 8 cm), the received acoustic transients are much more visible (fig. 3(a)) than those arrived at S_1 (fig. 3(b)). This is because the water hammer acoustic wave starts to diffuse. Still, the acoustic transient at S_1 has high energy (Digulescu, et al., 2015c).

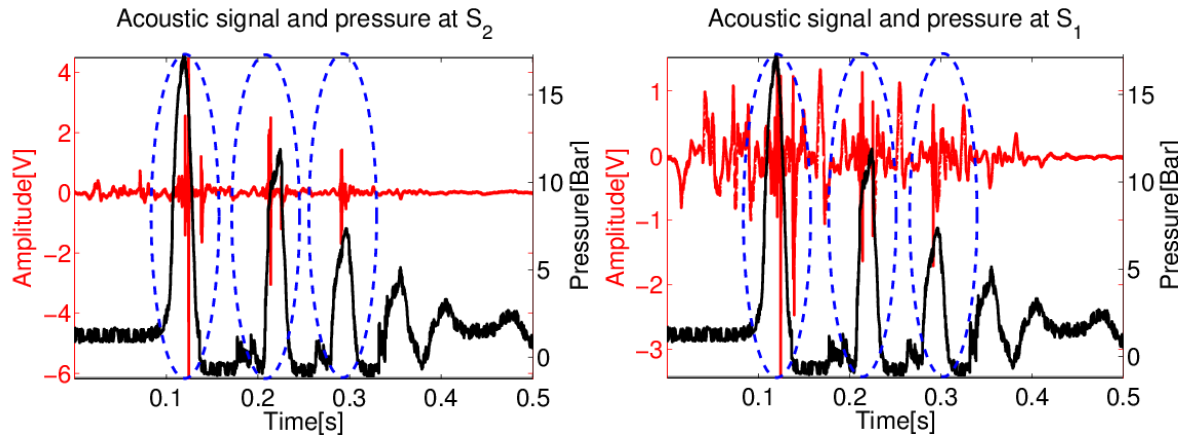


Fig. 6-3: The acoustic signals recorded by the two sensor and the simultaneous pressure time evolution: (left) acoustic signal arrived at sensor S_2 ; (right) acoustic signal arrived at sensor S_1

We applied the RPA method on these signals and computed the TDR^* measure for $m = 3, d = 10$ and $\varepsilon = 70$.

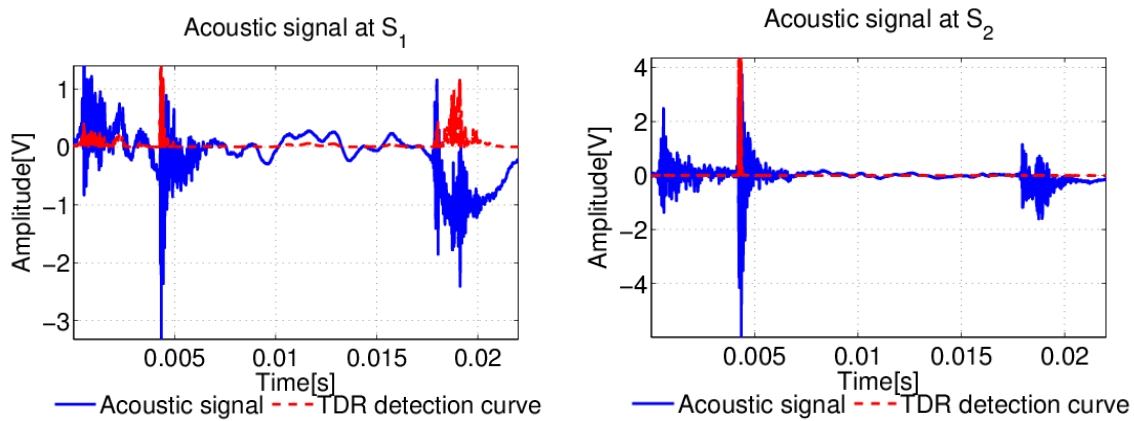


Fig. 6-4: The water hammer effect detection on the signals recorded by the two sensors

In Fig. 6-4, it can be seen that the TDR^* measure succeeds to detect the acoustic transients regardless of mechanical vibrations and environmental noise recorded by the sensors. In the computation of the TOF, we considered the first moments for which the TDR^* measure was 80% from the maxima, namely $TDR^* = 0.8$

Hereby, we determined that $TOF = 82\mu s$. Knowing that the distance between the two ultrasonic sensors was $d_s = 8cm$, we obtained the wave speed $c_{TDR} = 975 m/s$.

Alternatively, considering just the maxima given by the acoustic transients (result obtained typically from the envelope detection), we obtained $TOF = 99\mu s$, namely $c_{pk} = 805 m/s$. The difference between this approach and the RQA-based method is given by the diffusion effect of the acoustic wave generated by the water hammer.

Results using Joukowski equation

In the given configuration, the theoretical value of pressure wave speed in a Plexiglas pipe was computed using eq. (6.3) and considering the parameters of the water and the pipe as given by Table 6-1.

Parameter	Value
ρ	$1000 kg/m^3$
ϵ_w	$2.1 GPa$
e	$10 mm$
E	$5.66 GPa$

Table 6-1: Characteristic parameters of the pipe and fluid for theoretical computation of the pressure wave

Also, the presence of four butterfly valves with steel casing situated along the pipe was taken into account and an weighted elasticity modulus (E) of the setup was determined.

Hereby, the theoretical pressure wave speed value obtained is $c_{th} = 926 m/s$, value that is closest to the one determined by the RQA-based method. Using this value, the pressure variation during water hammer can be determined from eq. (6.1), and results a maximum increase of $\Delta p_{th} = 15.35 bar$.

Pressure variation

Using the maximum pressure difference measured during the tests after the sudden closure of the vane, $\Delta p = 15.9 bar$, the experimental value of pressure wave speed can be calculated with Joukowsky equation (6.1). In order to obtain this, the water velocity in the steady state flow is determined from the continuity equation, $V_0 = 4Q / (\pi D^2)$, as the flow rate before the closure is $Q = 1.98 l/s$. So, for the obtained value of $V_0 = 1.66 m/s$ the resulted experimental pressure wave speed is $c_e = 959 m/s$.

The results from Table 6-2 show that RQA approach gives very good results with an estimation error below 2%. This may be caused by the choice of the threshold of the TDR^* measure. Further investigations will be made for improvement. The next suitable approach is the one given by Joukowsky equation. Still, in industrial hydraulic systems, it may be difficult to determine a correct elasticity modulus of the materials that may be degraded and changed since their design, which will directly influence a correct evaluation of the pressure wave

speed and so, of the maximum pressure. The signal processing approach (pk) does not involve complex computational procedures using only the maxima of the acoustic transients but has a lower performance. The error may be due to the fact that the signal received by S1 is diffused, so the TOF may be altered. Also, it must be emphasized that the correlation, a state of the art method used in the computation of the TOF for acoustic signals, did not give any realistic result in this case.

Method	$c[m/s]$	$\Delta p[bar]$	Error[%]
c_e	959	15.91	-
c_{TDR}	975	16.18	1.69
c_{pk}	805	13.36	16.02
c_{th}	926	15.35	3.43

Table 6-2: The pressure variations using different approaches

The TDR^* measure has proved to be the most successful approach having an error of only 1.69%. The pressure variation computation using Joukowski equation for theoretical evaluation of pressure wave speed had an estimation error of 3.43%, so it can be useful with the condition that all the characteristics of the hydraulic system are known.

In conclusion, the RQA-based estimation technique, applied in acoustic non-intrusive passive sensing system, has a great potential for the monitoring of hydraulic transient phenomena. Further works will conduct to an operational tool that could be integrated in industrial equipments.

6.2. *Electrical arcs*

The permanent surveillance of electrical power systems must be accomplished in order to be aware and limit the effects of electrical faults. One of the most encountered electrical faults are materialized by the electrical arcs (EAs), characterized by wide band transient signals. Acoustic, electromagnetic and optics signatures are typical of such fault that is generally induced by physical defaults of electrical materials.

Since the presence of such phenomena is a preliminary stage in the continuous monitoring of distributed power systems (solar panels, transformers, etc.), we first studied the detection efficiency at different signal to noise ratios (SNRs) using the spectrogram and the wavelet transform than traditional-robust energy detectors (Mallat & Zhong, 1992). The performance of these detectors is presented in their corresponding receiver-operating-characteristic (ROC).

Afterward, these methods are applied to determine the time of arrival (TOA) of the electrical arcs in a multi-sensor configuration (Digulescu, et al., 2014c), (Digulescu, et al., 2013), (Digulescu, et al., 2016c), (Ioana, et al., 2014), (Digulescu, et al., 2013), (Digulescu, et al., 2014b), (Paun, et al., 2015). Then, the localization of the sources of the electrical arc is

obtained by solving the geometrical system which describes the experimental configuration. Finally, we presented the precision accuracy of each method.

Finally, we characterize these signals using the MLPA method in order to monitor the signal's changes, aiming to classify the electrical arc's signature (Digulescu, et al., 2015b), (Serbanescu, et al., 2015).

6.2.1. Detection performance

The experiment is performed with an electrical arc locator system composed of three acoustical sensors placed in a 3D configuration and a wide band antenna placed in the center of the system. The experimental configuration is presented in Fig. 6-5.

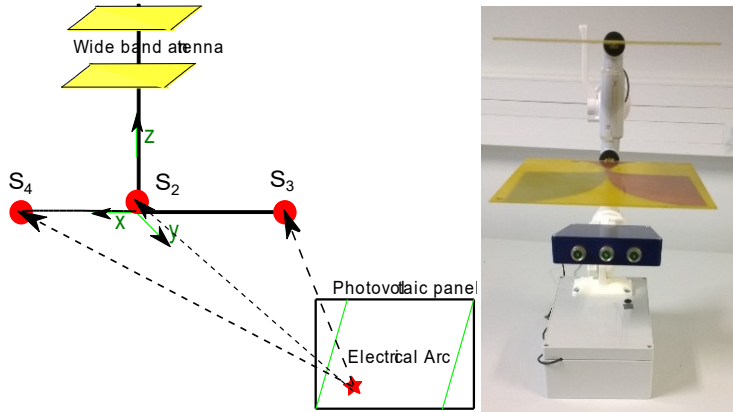


Fig. 6-5: The experimental configuration of the EA locator system

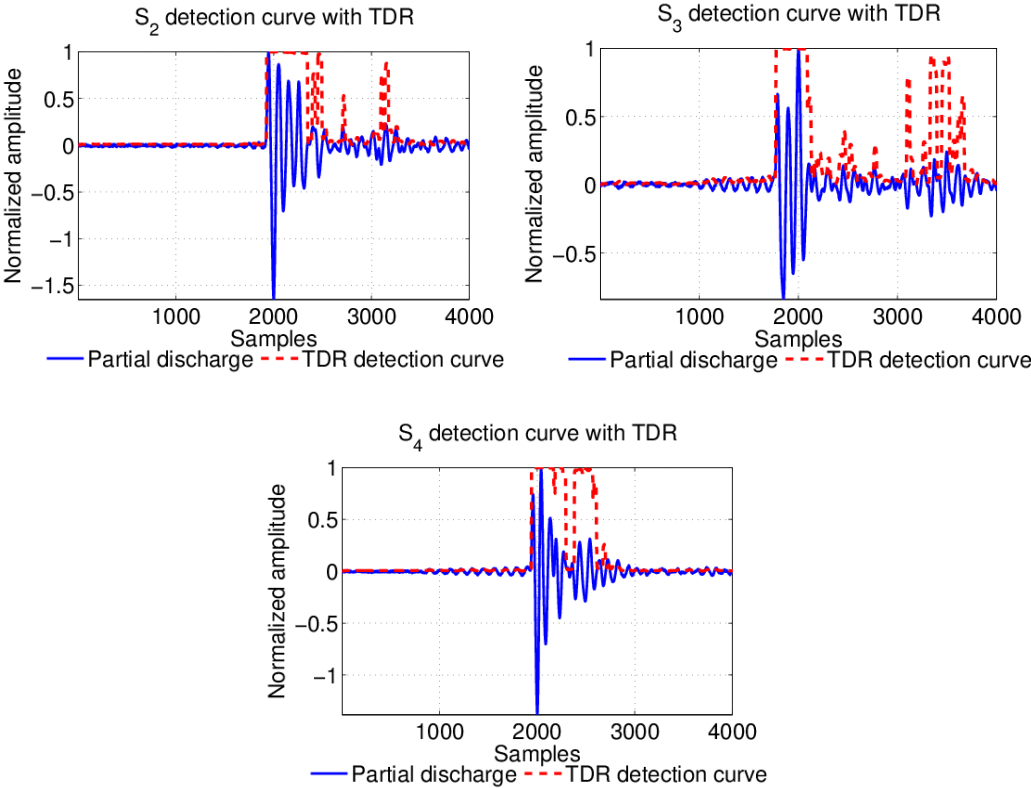


Fig. 6-6: The recorded EAs and the detection curve obtained with the TDR measure: $m = 3, d = 8$ and $\alpha = 0.7$

The detection and estimation of times of arrival of the transient signals generated by electrical arc at each sensor has been done using the TDR^* measure, the spectrogram and the wavelet transform. Fig. 6-6 presents an example of received signals and their detection curves obtained with the TDR^* measure.

For the ROCs computation, we considered the following algorithm for the SNR value:

- Let us consider that s is the test signal of length L
- For a given SNR , we have: $SNR = 10 \lg \frac{P_s + P_n}{P_n}$, where $P_s = \sum_{i=1}^L \frac{|s[i]|^2}{L}$
- Hereby, we obtain $P_n = \frac{P_s}{10^{10 \lg SNR - 1}}$
- Then, the imposed noise is: $n = \sqrt{P_n} \cdot n_0$, where n_0 is a random signal with normal distribution of mean zero and variance equal to 1
- The studied signal becomes: $s_n = s + n$

We compare the TDR^* measure with the spectrogram and wavelet methods by tracking the ROCs. The results are presented in Fig. 6-7.

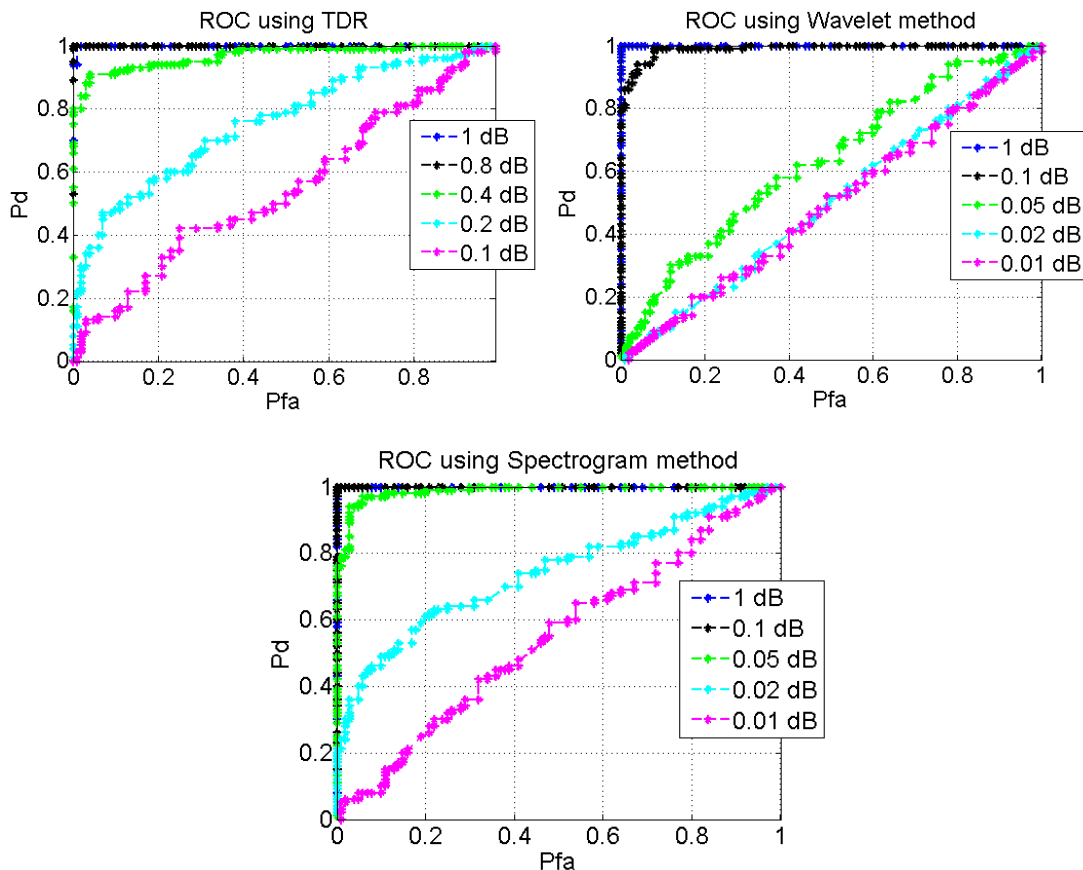


Fig. 6-7: The detection performances of each method; the ROCs where computed for the same white additive Gaussian noise (100 realizations) for each SNR level and method: (up-left) TDR^* measure for $m = 3, d = 8$ and $\alpha = 0.7$; (up-right) the spectrogram-based detection; (down) the wavelet-based detection using the Mexican Hat mother function

As expected, both dedicated detection methods have the best performances which are one size higher. Nevertheless, the RQA-based method, namely the TDR^* measure offers reliable performances, therefore, in the case of localization application, this approach must be considered in terms of localization accuracy.

The signals (Fig. 6-6) are recorded at a sampling frequency of $f_s = 5 \text{ MHz}$ for a period of 10 ms (50000 samples). In terms of location accuracy, the TDR^* measure is compared with the classical time-scale approaches.

The spatial localization is achieved by solving the geometrical system (eq.(6.5)) based on the time-of-arrival (TOA) of the EAs at each microphone. The TOA is obtained by imposing the same threshold (0.5) at the normalized detection curves based on each method.

$$\begin{cases} d_{PS_2} - d_{PS_1} = v \cdot t_{21} \\ d_{PS_3} - d_{PS_1} = v \cdot t_{31} \\ d_{PS_4} - d_{PS_1} = v \cdot t_{41} \end{cases} \quad (6.5)$$

where $d_{PS_i} = \sqrt{(x_{S_i} - x_P)^2 + (y_{S_i} - y_P)^2 + (z_{S_i} - z_P)^2}$, $i = \overline{1,4}$, each position of the acoustic sensor and antenna (considered S_1) is known $S_i(x_{S_i}, y_{S_i}, z_{S_i})$, $t_{(i+1)i}$ is the time of arrival of the EA at each microphone triggered by the S_1 wide band antenna. The localization accuracy using each method is presented in Table 6-3.

Method	Relative error [%]
TDR^* measure	6.2
Wavelet	11.2
Spectrogram	9.4

Table 6-3: The spatial localization accuracy for the electrical arcs

In this application, the signal is propagated on multiple paths, because the experiment is performed in a closed facility: an indoor facility, mainly characterized by three propagation paths.

Although, the classical techniques detect and localize the electrical arc source in an effective manner, when it comes to multi-path signal discrimination, these methods are limited.

But, the MLPA method provides better results. Using one of the signals arrived at S_2 microphone, the reflections s_2, s_3 of the electrical arc are compared with the direct path signal s_1 . These signals are presented in Fig. 6-8.

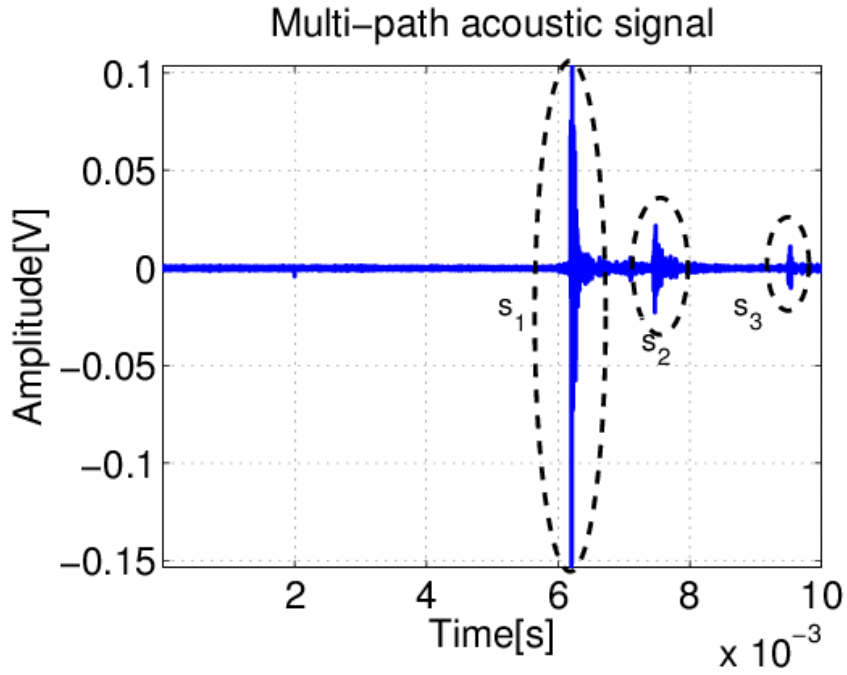


Fig. 6-8: The electrical arc recording s_1 and its reflections s_2, s_3

For the multi-path acoustic signals, the elliptic modeling (eq.(4.9)) is applied and the area of the estimated ellipse is determined (Fig. 6-9).

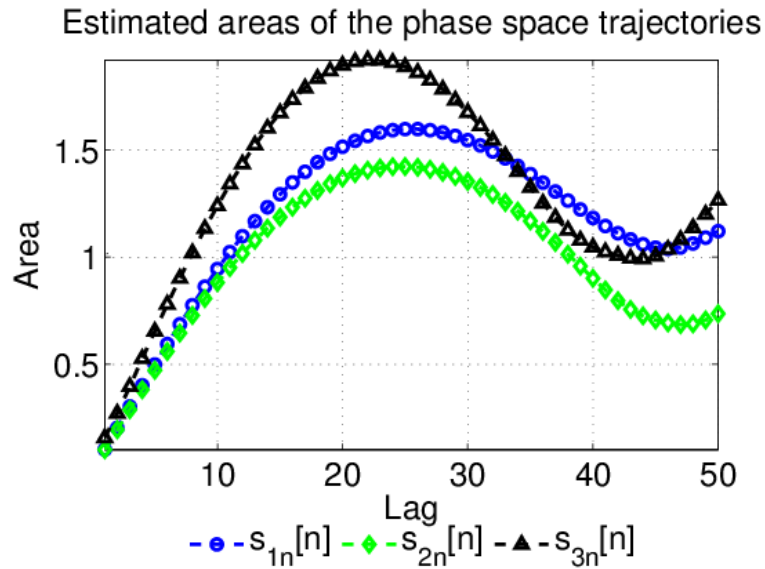


Fig. 6-9: The areas of the estimated ellipses that circumscribe the phase-space trajectory of the normalized multi-path signals

The estimated ellipses evolve in a different manner, but they have the same trend. After an average value of the area is chosen, $A=1$, an optimal lag is determined: $d_1=11$, $d_2=12$ and $d_3=8$. The results of the phase diagram representation are shown in Fig. 6-10.

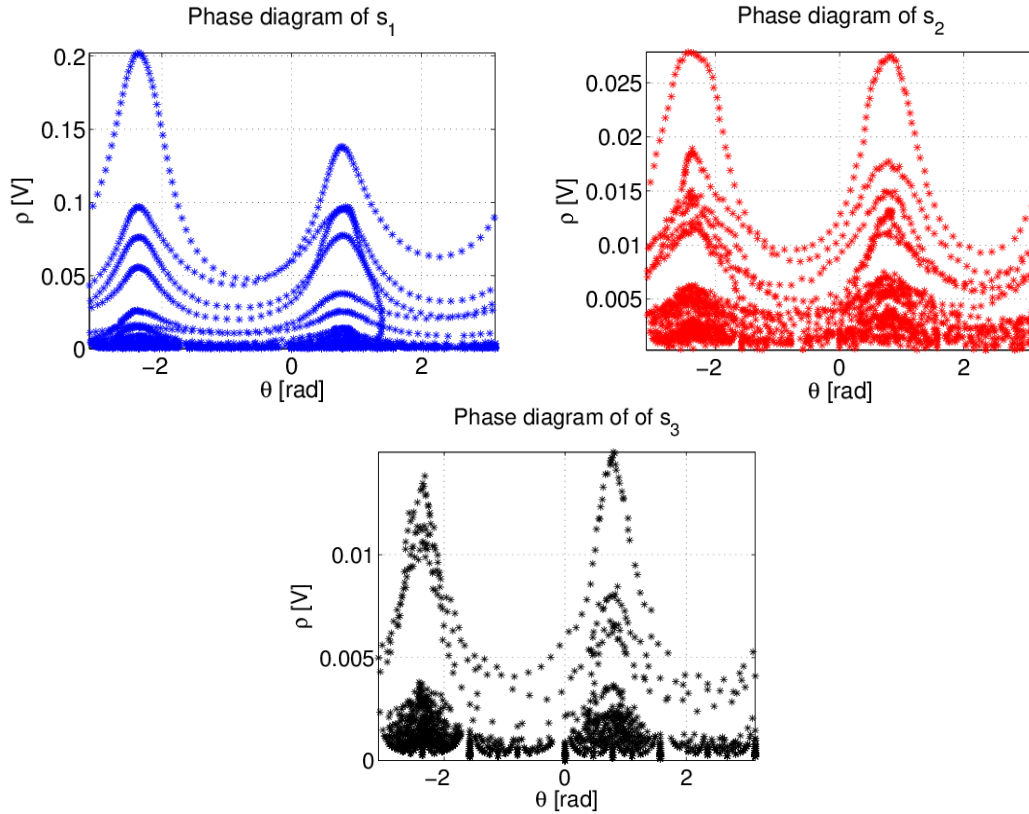


Fig. 6-10: The polar coordinates representation for the multi-path acoustic signals: $d_1 = 11$, $d_2 = 12$ and $d_3 = 8$

In terms of attenuation, the phase diagram points out the same information as the time evolution of the signals. Moreover, the reflections present fewer curves with peaks than the direct path of the signal. This means that the reflections contain fewer oscillations than the direct signal, hereby, the reflections are dispersed.

Using eq.(4.11), it goes that: $f_{s_1} / f_{s_2} = 0.94$ and $f_{s_1} / f_{s_3} = 0.92$, meaning that the reflections suffer a time dilatation. Because the frequency ratios are close to 1, it leads to the idea that the signals have the same source, but on their propagation path, the reflections are affected by multiple phenomena: diffusion, dispersion, attenuation, etc.

6.3. Chapter summary

This chapter presents two applications of passive sensing. For the pressure transient measurement, we have shown that the RQA-based method for transient detection and time localization is a method worth being considered for the phenomenon characterization.

Regarding the EA application, the RQA approach provides a better localization accuracy than the time-scale methods (but similar detection performances) and with the use of multi-lag phase-space analysis, it highlights new information regarding the characterization of transient signals and the changes that they suffer in a realistic multi-path propagation environment.

7. Conclusions and perspectives

7.1. Work summary

The main objective of this work was to contribute to the analysis of signals related to different physical phenomena, aiming to better understand them as well as to estimate their parameters that are interesting from application point of view. Different applicative contexts have been investigated in active and passive sensing configurations. For the active part, we mention the monitoring of cavitation phenomena and its characterization for hydraulic system surveillance. The second application of the active sensing is the underwater object detection and tracking without synchronization between sensors. For the passive configuration, we focus on the pressure transient analysis in water pipes investigation with a non-intrusive method and on the surveillance of electrical power systems in the presence of transient phenomena such as electrical arcs.

Despite the differences between the physical considerations, we propose a unique mathematical model of the signals issued from the active/passive sensing system used to analyze the considered phenomena.

Chapter 2 presents a state of the art of the techniques used for acoustic sensing methods, from signal processing perspective, in dynamic configurations. Starting with the well known analysis techniques based on energy methods and projective approaches (such as wavelet transform), we focus our research to the data-driven analysis technique based on the phase diagram concept.

Chapter 3 describes a new concept of active sensing system for analyzing dynamic transient phenomena. For the emission part we propose the phase-space-based waveform design that presents a unique shape in the phase space and which can be easily translated into a signal. Starting from this unique shape, called phase space lobe, we translate the representation into a time series (auxiliary signal). This time series carries the IFL (Instantaneous Frequency Law) of the further emission signal. The main advantage of this new concept is the capability to generate several transmitting signals, non-orthogonal in time/frequency domain but orthogonal in the phase diagram representation space. For the reception part, the IFL information is estimated in a phase space-based representation space. This waveform design technique presents the interest to conduct to a high speed sensing methods, very useful to monitor dynamic phenomena. This work will conduct to a patent software that will be pended before the end of this year.

Chapter 4 focuses on the signal analysis based on the major information that can be extracted from phase diagrams: the recurrence quantification analysis (RQA) based on the recurrence plot analysis (RPA). We propose new tools for the investigation of the signals characteristics. The first one is the TDR^* measure (Time Distributed Recurrences) that

quantifies the recurrence/ distance matrix and it is used for the detection of transient signals. The second one is the multi-lag phase analysis using multiple lags and it is successfully used to discriminate between signals with close parameters. Finally, the diagonal lines quantification of RPA matrix is proposed as an alternative for the analysis of modulated signals.

Chapter 5 and 6 present the experimental results using the proposed theoretical methods introduced by this thesis. The results are compared with classical techniques. The *cavitating environment* (a major problem in hydraulic turbines) evolution on a small scale facility is characterized using the *diagonal lines quantification* of the distance matrix (Digulescu, et al., 2016) (Digulescu, et al., 2016) (Digulescu, et al., 2014).

A new method for the detection of an *underwater object* using bi-static acoustic paths is presented. For this application of object tracking, the main advantage of the RQA and RPA methods is that it uses only the signals arrived at reception (so there is synchronization-free technique). Using the *TDR** measure, we define the detection map based on this measure which is more robust to noise than the signal's envelope-based detection. This map contains both amplitude and time information (Digulescu, et al., 2015a).

The *water hammer* phenomenon is pointed out through our non-intrusive approach that proves to be the closest to the reference method. The effect of the phenomenon is quantified using the *TDR* measure which is more robust to the effect of pressure wave diffusion that takes place inside the pipe during the experiment (Digulescu, et al., 2015c).

The *electrical arc* application points out that our approach presents another detection method based only on the RQA and RPA methods using the *TDR* measure (Bernard, et al., 2016), (Candel, et al., 2012a), (Digulescu, et al., 2013), (Digulescu, et al., 2016c), (Digulescu, et al., 2014b), (Digulescu, et al., 2014c). The localization accuracy outperforms the classical non-stationary signals processing methods. Moreover, the signals are characterized and discriminated using the concept of *multi-lag phase space analysis* (Bernard, et al., 2016) (Digulescu, et al., 2015b) (Digulescu, et al., 2016c).

These applications can be extended to a high speed sensing concept, allowing to get the dynamic evolution of the cavitation process or the object's image, respectively. For this purpose, the phase diagram-based waveform design can be successfully used, as illustrated in this work.

7.2. Perspectives

This thesis is structured in two main parts, one theoretical and another one applicative, where the theoretical contributions have been investigated in some applications of current interest. Consequently, the perspectives of my work will be concentrated on two main directions – theoretical and applicative. Since the both aspects are intimately related, I prefer to present them in a merged manner.

The **first theoretical** works will concern the phase diagram-based adaptive waveform, by investigating different partition of information in the phase diagram domain. The

organization of the information in lobes could be generalized in n -dimensional phase diagram domain, creating also the possibility to design better separable waveform carrying, in the same time, more information, bringing the modulations more robust and with better performances in terms of parameters of estimation.

At the receiver level, we will investigate the capability of instantaneous frequency law tracking in the phase diagram, trying to take advantage of the phase continuity of the phase trajectory. Namely, the IFL tracking can be successfully based on the association of the time-frequency points according to the continuity criteria that could be naturally defined in the phase diagram domain. In the case of the multi-lag phase diagram analysis, this continuity will be defined in more robust manner since the random noise effect will be minimized when looking to different lags.

In the applicative context, these new theoretical directions will conduct to the definition of a new imaging acoustic methodology where the high speed sensing will be allowed thanks to the construction of agile wide band waveforms, emitted as much lower PRF than the classical ones. In addition, the accurate estimation of the received signals characteristics will allow to form the acoustic images not only from the times of arrival, but also from the *relative deformation between the received signals*.

Another applicative extension of this new theoretical work will concern the definition of new communication modulations and multiple access technique, not only in acoustic but also electromagnetic waves.

The **second theoretical** direction is the development of new signal descriptors in the phase diagram plan. As a potential idea we mention here the dynamic of position vector in polar coordinates, but merged from *coherent* variations of this vector at different lags.

Another work idea is to generalize the definition of these vectors in n -D. The definition of a two-dimensional representation of the position vector parameters (calculated in n -dimensional phase diagram coordinates) – representation similar to RPA, would be very helpful for the intuitive data interpretation.

The relation between phase diagram analysis and the compressive sensing domain could be also another work direction worth to be followed.

In the applicative context, the aforementioned directions will conduct to new methods for signal analysis and characterization. Their using in the application we have investigated in this thesis is natural. Namely, these representation methods will conduct to a finer analysis of the interaction between acoustic waves and the hydrodynamic phenomena. The further contribution will be to help in monitoring hydraulic machines in non-stationary conditions that are nowadays under intensive multidisciplinary research works.

Another applicative work context of the further representation tools is the analysis of the electromagnetic transients, aiming to characterize and localize their sources. Applications in ultra wideband tracking, identification and communications as well as in monitoring of power equipments are just two domains where the further theoretical researches will serve.

Résumé étendu

Caractérisation des phénomènes dynamiques à l'aide de l'analyse du signal sur le diagramme des phases

La déformation des signaux au long de leur trajet de propagation est un des plus importants facteurs qui doivent être considérés à la réception. Ces effets sont dus à des phénomènes comme l'atténuation, la réflexion, la dispersion ou, encore, le bruit. Alors que les premiers deux phénomènes sont assez facile à surveiller, parce qu'elles affectent l'amplitude, respectivement le retard des signaux, les deux derniers phénomènes sont plus difficiles à contrôler, parce qu'elles changent les paramètres du signal (amplitude, fréquence et phase) de manière totalement dépendante de l'environnement.

Nous définissons un modèle générale pour le signal reçu après la propagation dans un environnement hétérogène :

$$r(t) = \sum_k \mathcal{D}_k \{s(t)\} + n(t) \quad (6)$$

où $\mathcal{D}_k \{\cdot\}$ est l'opérateur général de déformation associé à la composante k du signal reçu et $n(t)$ est le bruit additive blanc gaussien.

Dans cette thèse, l'objectif principal est de contribuer à l'analyse des signaux liés aux différents phénomènes physiques, en visant une meilleure compréhension de ces phénomènes, ainsi que l'estimation de leurs paramètres qui sont intéressants de point de vue applicatif. Plusieurs contextes applicatifs ont été investigués dans deux configurations de mesure: active et passive.

Pour la configuration active, le premier contexte applicatif consiste en l'étude du phénomène de cavitation dans le domaine de la surveillance de systèmes hydrauliques. La deuxième application de la configuration active est la détection et le suivi des objets immergés sans synchronisation entre les capteurs d'émission et de réception.

Pour la configuration passive, nous nous concentrons sur l'analyse des transitoires de pression dans les conduites d'eau en utilisant une méthode non-intrusive ainsi que sur la surveillance des réseaux d'énergie électrique en présence des phénomènes transitoires comme les arcs électriques.

Malgré les différences entre les considérations physiques spécifiques à ces applications, nous proposons un modèle mathématique unique pour les signaux issus des deux types de configuration.

Le Chapitre 2 présente les techniques les plus utilisés pour les méthodes de mesure acoustique en configurations dynamiques, dans la perspective du traitement du signal. En commençant avec les techniques d'analyse très connues basées sur l'approche énergétique ainsi que sur les approches projectives (comme la transformée en ondelettes), nous concentrons notre recherche sur les techniques d'analyse des signaux aiguillées par les données.

La méthode des récurrences commence en considérant la série de temps suivante :

$$\mathbf{x} = \{x[1], x[2], \dots, x[N]\} \quad (7)$$

Cette série est ensuite représentée dans l'espace des phases. Ses valeurs deviennent les coordonnées de l'espace m - dimensionnel et, par conséquence, les vecteurs de l'espace est :

$$\vec{v}_i = \sum_{k=1} x[i + (k-1)d] \cdot \vec{e}_k, \quad i = 1, M \quad (8)$$

où m est la dimension de l'espace, d est le retard entre les échantillons, \vec{e}_k est le vecteur unitaire des axes de l'espace, $M = N - (m-1)d$ et N est la longueur de la série de temps. Usuellement, la dimension de l'espace et le retard sont choisis en utilisant la méthode de faux voisins et, respectivement, la méthode de l'information mutuelle.

Par la suite, les distances entre les vecteurs de l'espace des phases sont représentées par la matrice des distances (DM), eq. (4). Comparées avec un seuil, on obtient la matrice des récurrences (RM), eq. (2.13).

$$D_{i,j} = \mathcal{D}(v_i, v_j) \quad (9)$$

$$R_{i,j} = \Theta(\varepsilon(i) - \mathcal{D}(v_i, v_j)) \quad (10)$$

où $\mathcal{D}(\cdot, \cdot)$ est la distance appliquée aux vecteurs de l'espace des phases (la distance euclidienne, la norme L1, la distance angulaire), $\Theta(\cdot)$ est la fonction Heaviside et $\varepsilon(i)$ est le seuil considéré pour la récurrence.

L'avantage majeur de la méthode des récurrences est qu'elle ne requiert aucune base orthogonale ou aucune information a priori du système étudié (contrairement à la méthode des ondelettes qui nécessite la définition d'un dictionnaire qui doit ressembler aux composantes du signal analysé). La méthode fonctionne pour des applications où la référence n'est pas disponible (contrairement à la méthode du filtrage adapté) et est plus précise pour la localisation.

Le Chapitre 3 décrit un nouveau concept d'un système actif pour l'analyse des phénomènes transitoires dynamiques. Pour la partie d'émission, nous proposons la construction des ondes d'émissions basées sur l'espace des phases. Les différentes ondes construites présentent une forme unique dans l'espace des phases.

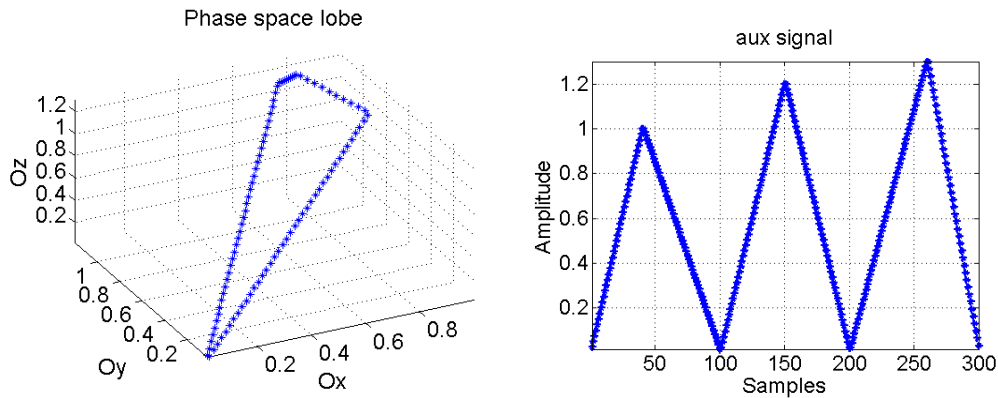


Fig. 1: Le lobe dans l'espace des phases et son signal auxiliaire correspondant

En commençant avec cette forme unique, appelée lobe dans l'espace des phases, nous translatons cette représentation dans une série temporelle (signal auxiliaire) (Fig. 1). Cette série temporelle porte la LFI (Loi de Fréquence Instantanée) du signal d'émission. L'avantage principal de ce concept est la capacité de générer plusieurs signaux à l'émission, non-orthogonaux dans le domaine de temps et de fréquence, mais qui est orthogonal dans l'espace de représentation constitué par le diagramme de phases.

Pour la partie de réception, l'information portée par la LFI est estimée dans une analyse des trajectoires dans l'espace des phases (Fig. 2).

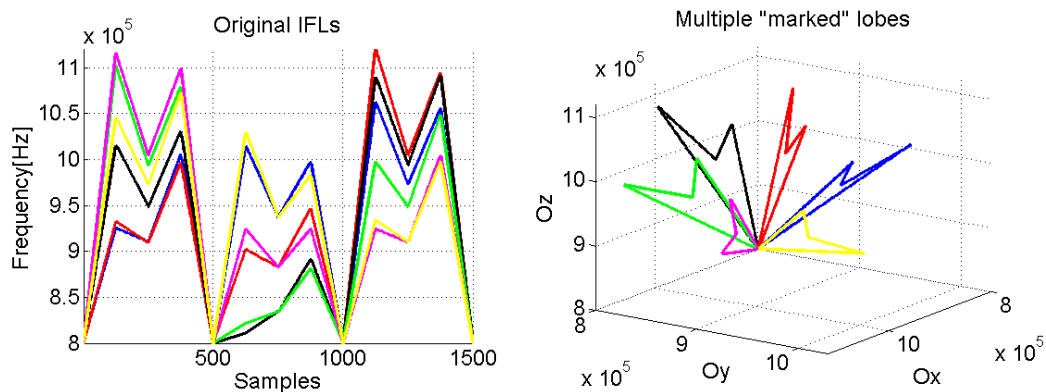


Fig. 2: Les LFI et leurs lobes multiples obtenus dans l'espace des phases

Cette technique de construction des formes d'ondes présente l'intérêt de conduire à des méthodes d'investigation actives à haute cadence, très utiles pour la surveillance des phénomènes dynamiques.

Le Chapitre 4 concerne l'analyse des signaux basée sur l'information majeure qui peut être extraite du diagramme des phases : l'analyse des récurrences. Nous proposons deux approches nouvelles pour l'investigation des caractéristiques des signaux. La première est la mesure TDR* (Time Distributed Recurrences) qui quantifie la matrice des récurrences/distances et qui est utilisée pour la détection des signaux transitoires.

La deuxième approche est l'analyse des phases à plusieurs retards et elle est utilisée pour la discrimination entre des signaux avec des paramètres très proches. Finalement, la quantification des lignes diagonales de la matrice des récurrences est proposée comme alternative pour l'analyse des signaux modulés en fréquence.

Les Chapitres 5 et 6 présentent les résultats expérimentaux en utilisant les méthodes théorétiques proposées dans cette thèse. Les résultats sont comparés avec des techniques classiques. L'évolution de l'environnement avec de la cavitation (un problème majeure dans la surveillance des installations hydrauliques) est caractérisée avec la quantification des lignes diagonales de la matrices des distances (Fig. 3).

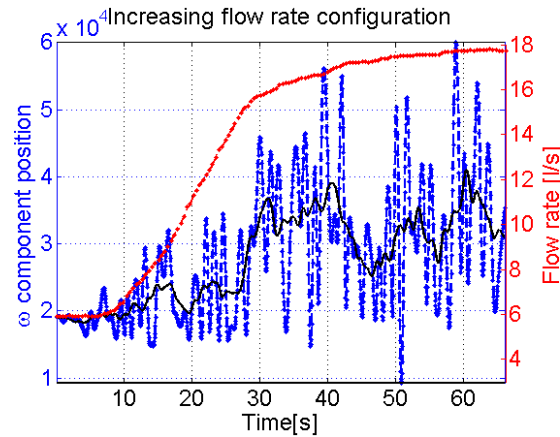


Fig. 3: La caractérisation du vortex avec des cavitations (générées par l'écoulement à l'interface avec un stator) en utilisant la quantification des lignes diagonales de la matrice des distances

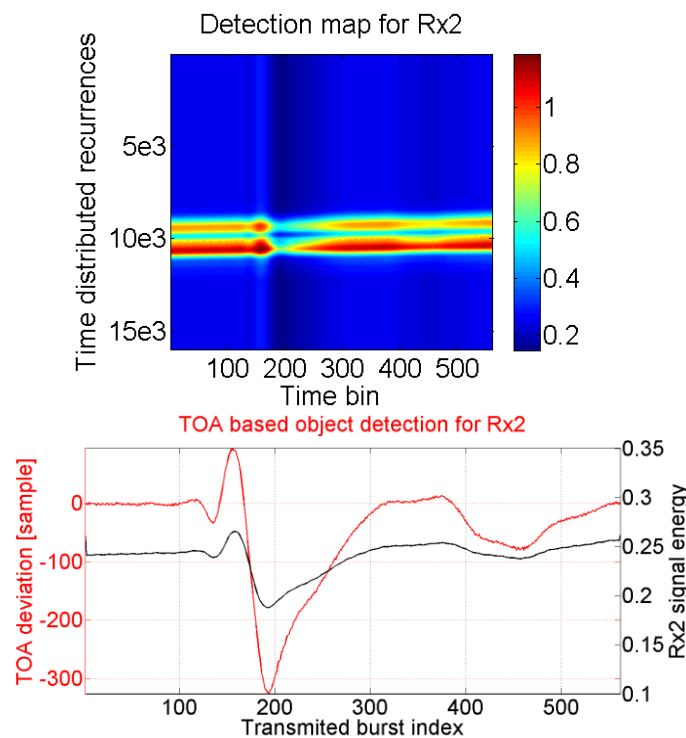


Fig. 4: La carte de détection basée sur la mesure TDR* et la détection de l'objet immergé en utilisant le temps d'arrivé de chaque pulse

Une méthode nouvelle pour la détection d'un objet immergé en utilisant deux trajets acoustiques statiques est également présentée. Pour l'application de suivi d'objet, l'avantage principal de la méthode des récurrences est qu'elle utilise seulement les signaux arrivés à la réception (par ailleurs, elle est une technique qui ne nécessite pas la synchronisation). En utilisant la mesure TDR*, nous définissons une carte de détection basée sur cette mesure qui

est plus robuste au bruit que la détection basée sur l'enveloppe du signal (Fig. 4). Cette carte contient les deux informations de temps et d'amplitude.

Le phénomène de coup de bélier est également un autre contexte applicatif étudié par notre approche non-intrusive qui montre qu'elle est la plus proche de la référence. L'effet du phénomène est quantifié en utilisant la mesure TDR* (Fig. 5) qui est plus robuste à l'effet de diffusion de l'onde de pression qui se passe à l'intérieur de la conduite pendant l'expérience (erreur d'estimation sous 2%).

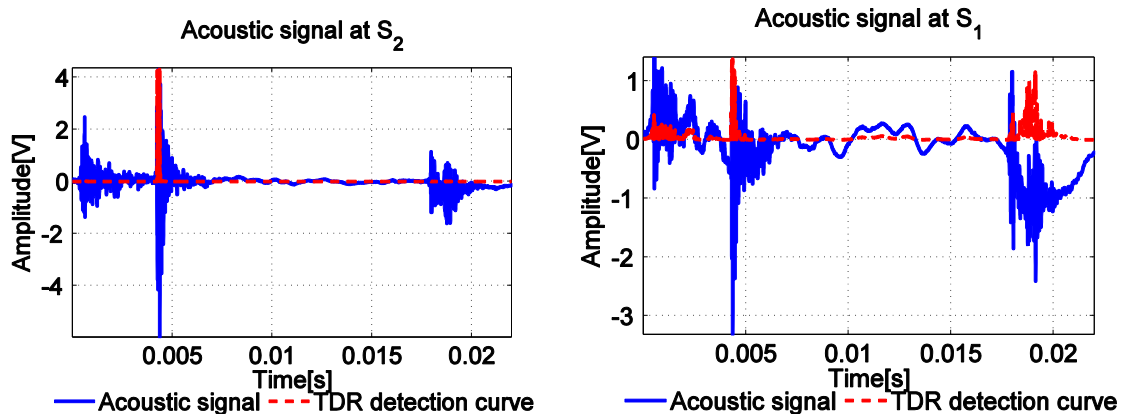


Fig. 5: La caractérisation du coup de bélier sur les signaux reçus par les capteurs

Méthodes basés sur l'estimation du temps d'arrivé	Erreur relative [%]
Mesure TDR*	6.2
Ondelettes	11.2
Spectrogramme	9.4

Table 1: Les résultats concernant la précision en termes de localisation

L'application concernant l'étude des arcs électriques met en évidence l'intérêt pour l'approche basée sur la nouvelle méthode de détection constituée par l'analyse des récurrences en utilisant la mesure TDR*. La précision en termes de localisation est démontrée en comparaison avec des techniques standard (Table 1). De plus, les signaux sont caractérisés et discriminés en utilisant le concept d'analyse dans l'espace des phases à plusieurs retards.

La méthode des récurrences offre des informations nouvelles concernant les systèmes réels industrielles. Des perspectives de ces travaux, tant dans les domaines théorique et qu'applicatif, sont discutés à la fin du mémoire.

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Résumé - La déformation des signaux au long de leur trajet de propagation est un des plus importants facteurs qui doivent être considérés à la réception. Ces effets sont dus à des phénomènes comme l'atténuation, la réflexion, la dispersion et le bruit. Alors que les premiers deux phénomènes sont assez facile à surveiller, parce qu'elles affectent l'amplitude, respectivement le retard des signaux, les deux derniers phénomènes sont plus difficiles à contrôler, parce qu'elles changent les paramètres du signal (amplitude, fréquence et phase) de manière totalement dépendante de l'environnement. Dans cette thèse, l'objectif principal est de contribuer à l'analyse des signaux liés aux différents phénomènes physiques, en visant une meilleure compréhension de ces phénomènes, ainsi que l'estimation de leurs paramètres qui sont intéressants de point de vue applicatif. Plusieurs contextes applicatifs ont été investigués dans deux configurations de : active et passive. Pour la configuration active, le premier contexte applicatif consiste en l'étude du phénomène de cavitation dans le domaine de la surveillance de systèmes hydrauliques. La deuxième application de la configuration active est la détection et le suivi des objets immergés sans synchronisation entre les capteurs d'émission et de réception. Pour la configuration passive, nous nous concentrons sur l'analyse des transitoires de pression dans les conduites d'eau en utilisant une méthode non-intrusive ainsi que sur la surveillance des réseaux d'énergie électrique en présence des phénomènes transitoires comme les arcs électriques. Malgré les différences entre les considérations physiques spécifiques à ces applications, nous proposons un modèle mathématique unique pour les signaux issus des deux types de configurations. Le modèle est basé sur l'analyse des récurrences. Avec ce concept, nous proposons une nouvelle approche pour les ondes basées sur l'espace des phases. Cette technique de construction des formes d'ondes présente l'intérêt de conduire à des méthodes de d'investigation active à haut cadence, très utiles pour la surveillance des phénomènes dynamiques. En plus, nous proposons des approches nouvelles pour l'investigation des caractéristiques des signaux. La première est la mesure *TDR** (Time Distributed Recurrences) qui quantifie la matrice des récurrences/ distances et qui est utilisée pour la détection des signaux transitoires. La deuxième approche est l'analyse des phases à plusieurs retards et elle est utilisée pour la discrimination entre des signaux avec des paramètres très proches. Finalement, la quantification des lignes diagonales de la matrice des récurrences est proposée comme alternative pour l'analyse des signaux modulés en fréquence. Les travaux présentent les résultats expérimentaux en utilisant les méthodes théorétiques proposées dans cette thèse. Les résultats sont comparés avec des techniques classiques. Des perspectives de ces travaux, tant dans les domaines théorique et qu'applicatif, sont discutés à la fin du mémoire.

Mots clés : Déformation du signal, Diagramme des phases, Analyse des récurrences, Ondes basées sur l'espace des phases, Détection et localisation, Caractérisation du signal

Abstract - Signals' deformation along their propagation path is among the most important aspect which has to be taken into account at reception. These effects are caused by phenomena like attenuation, reflection, dispersion and noise. Whereas the first two are rather easy to monitor, because they affect the amplitude, respectively the delay, the latter two are more difficult to control, because they change signals' parameters (amplitude, frequency and phase) in an environment-dependent manner. In this thesis, the main objective is to contribute to the analysis of signals related to different physical phenomena, aiming to better understand them as well as to estimate their parameters that are interesting from application point of view. Different applicative contexts have been investigated in active and passive sensing configurations. For the active part, we mention the monitoring of cavitation phenomena and its characterization for hydraulic system surveillance. The second application of the active sensing is the underwater object detection and tracking without synchronization between sensors. For the passive configuration, we focus on the pressure transient analysis in water pipes investigation with a non-intrusive method and on the surveillance of electrical power systems in the presence of transient phenomena such as electrical arcs. Despite the differences between the physical considerations, we propose a unique mathematical model of the signals issued from the active/passive sensing system used to analyze the considered phenomena. This model is based on the Recurrence Plot Analysis (RPA) and Recurrence Quantification Analysis (RQA) methods. With this concept, we propose the phase-space based waveform design. This waveform design technique presents the interest to conduct to a high speed sensing methods, very useful to monitor dynamic phenomena. Moreover, we propose new tools for the investigation of the signals characteristics. The first one is the *TDR** measure (Time Distributed Recurrences) that quantifies the recurrence/ distance matrix and it is used for the detection of transient signals. The second one is the multi-lag phase analysis using multiple lags and it is successfully used to discriminate between signals with close parameters. Finally, the diagonal lines quantification of RPA matrix is proposed as an alternative for the analysis of modulated signals. Our work presents the experimental results using the proposed theoretical methods introduced by this thesis. The results are compared with classical techniques. The perspectives of this thesis are presented at the end of this paper.

Keywords: Signal deformation, Phase diagram, Recurrence Plot Analysis, Phase-space based adaptive waveform, Detection and localization, Signal characterization