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# Local approximation by linear systems and Almost-Riemannian Structures on Lie groups and Continuation method in rolling problem with obstacles 

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## Local approximation by linear systems

 and Almost-Riemannian Structures on Lie groups and Continuation method in rolling problem with obstacles Approximation locale par des systèmes linéaires et structures presque-Riemanniennes et Méthode de continuation dans un problème de roulement avec obstacles
## Thèse de doctorat de l'université Paris-Saclay

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## ÉCOLE DOCTORALE

## université PARIS-SACLAY

Titre: Approximation locale par des systèmes linéaires et structures presque-Riemanniennes et Méthode de continuation dans un problème de roulement avec obstacles.

Mots clés: Géométrie presque riemannienne, approximation nilpotente, approximation résoluble, problème des corps roulants, obstacles, planification nonholonomique du mouvement.

Résumé: L'objectif de cette thèse est d'étudier deux sujets en géométrie sub-riemannienne. D'une part, I'approximation locale d'une structure presque riemannienne aux points singuliers, et d'autre part, le système cinématique d'une variété à 2 dimensions roulant (sans tourner ni glissement) sur le plan euclidien avec des régions interdites.

Une structure presque Riemannienne de dimension $n$ peut être définie localement par $n$ champs vectoriels satisfaisant la condition de rang de l'algèbre de Lie, jouant le rôle d'un cadre orthonormé. L'ensemble des points où ces champs vectoriels sont colinéaires est appelé l'ensemble singulier $(\mathcal{Z})$. Aux points de tangence, c'est-àdire aux points où l'espace linéaire engendré par champs vectoriels est égale à l'espace tangent de $\mathcal{Z}$, l'approximation nilpotente peut être remplacée par l'approximation résoluble. Dans cette thèse, sous des conditions génériques, nous établissons l'ordre d'approximation de la distance originale par $\widetilde{d}$ (la distance induite par l'approximation solvable) et nous prouvons que $\widetilde{d}$ est plus proche que la distance induite par l'approximation nilpotente de la
distance originale. En ce qui concerne les structures des systèmes d'approximation, l'algèbre de Lie générée par cette nouvelle famille de champs vectoriels est de dimension finie et solvable (dans le cas générique). De plus, l'approximation solvable est équivalente à un ARS linéaire sur un espace homogène ou un groupe de Lie.

D'autre part, les systèmes nonholonomes ont attiré l'attention de nombreux auteurs de différentes disciplines pour leurs applications variées, principalement en robotique. Le problème du corps roulant (sans glissement ni rotation) d'une variété riemannien bidimensionnel sur une autre variété peut être écrit comme un système nonholonomique. De nombreuses méthodes, algorithmes et techniques ont été développés pour le résoudre. Une implémentation numérique de la méthode de continuation pour résoudre le problème dans lequel une surface convexe roule sur le plan euclidien avec des régions interdites (ou obstacles) sans glisser ou tourner est effectuée. Plusieurs exemples sont illustrés.

Sciences et technologies
de l'information et de
la communication (STIC)

Title: Local approximation by linear systems and Almost-Riemannian Structures on Lie groups and Continuation method in rolling problem with obstacles.
Keywords: Almost-Riemannian geometry, nilpotent approximation, solvable approximation, rolling body problem, obstacles, nonholonomic motion planning.


#### Abstract

The aim of this thesis is to study two topics in sub-Riemannian geometry. On the one hand, the local approximation of an almostRiemannian structure at singular points, and on the other hand, the kinematic system of a 2 dimensional manifold rolling (without twisting or slipping) on the Euclidean plane with forbidden regions.

A $n$-dimensional almost-Riemannian structure can be defined locally by $n$ vector fields satisfying the Lie algebra rank condition, playing the role of an orthonormal frame. The set of points where these vector fields are colinear is called the singular set $(\mathcal{Z})$. At tangency points, i.e., points where the linear span of the vector fields is equal to the tangent space of $\mathcal{Z}$, the nilpotent approximation can be replaced by the solvable one. In this thesis, under generic conditions, we state the order of approximation of the original distance by $\widetilde{d}$ (the distance induced by the solvable approximation), and we prove that $d$ is closer than the distance induced


by the nilpotent approximation to the original distance. Regarding the structure of the approximating system, the Lie algebra generated by this new family of vector fields is finite-dimensional and solvable (in the generic case). Moreover, the solvable approximation is equivalent to a linear ARS on a homogeneous space or a Lie group.

On the other hand, nonholonomic systems have attracted the attention of many authors from different disciplines for their varied applications, mainly in robotics. The rolling-body problem (without slipping or spinning) of a 2-dimensional Riemannian manifold on another one can be written as a nonholonomic system. Many methods, algorithms, and techniques have been developed to solve it. A numerical implementation of the Continuation Method to solve the problem in which a convex surface rolls on the Euclidean plane with forbidden regions (or obstacles) without slipping or spinning is performed. Several examples are illustrated.

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## General Introduction

The sub-Riemannian geometry is a generalization of the Riemannian geometry (geometry which was born with the work of Bernhard Riemann to generalize some of Gauss's results concerning the curvature). Its origin is in the paper [Strichartz, 1986]. However, some elements, examples, or applications of the sub-Riemannian geometry can be situated before, for instance, in works related to optimal control problems [Brockett, 1982], thermodynamic [Carathéodory, 1909], or even in papers related to the Riemannian geometry [Hermann, 1973] and the study of the geometry of the Heisenberg group, which is a famous example of the subRiemannian geometry [Gaveau, 1977]. This geometry exists thanks to the idea of non-integrable constraints; this is to say a constraint on the admissible direction of movements, and thus, it has received increasing attention in many disciplines such as control theory, robotics, classical mechanics, analysis of hypoelliptic operator, diffusion manifold, and even in other branches of geometry as Cauchy-Riemann geometry.

A particular sub-Riemannian structure is the so-called almost-Riemannian structure. Formally, a $n$-dimensional Almost-Riemannian Structure (ARS in short) is a rank-varying subRiemannian structure that can be locally defined by a set of $n$ smooth vector fields on a $n$-dimensional manifold, satisfying the Lie algebra rank condition. The set of points where the dimension of the linear span of the vector fields is not full is called the singular locus (or the singular set) and is denoted by $\mathcal{Z}$. Attractive models of ARS's can be described on Lie groups using invariant and linear vector fields. They are referred to as linear (or simple) ARSs on Lie groups (see [Ayala and Jouan, 2016]).

The aim of this thesis is to study two topics in sub-Riemannian geometry. On the one hand, the local approximation of an almost-Riemannian structure at singular points, where the nilpotent approximation lost the original structure, and on the other hand, the kinematic system of a manifold rolling on another manifold without twisting or slipping, particularly a numerical implementation of the Continuation Method when a 2-dimensional manifold rolling on the Euclidean plane with forbidden regions.

## Local approximation by linear systems and almost-Riemannian structures on Lie group

The purpose of this first part is to locally approximate almost-Riemannian structures at singular points by ARSs on Lie groups (or homogeneous spaces) and to show that this approximation is generally better than the nilpotent one.

Let $M$ be an $n$-dimensional differential manifold and consider $\left\{X_{1}, X_{2}, \ldots, X_{n}\right\}$ a set
of smooth vector fields on $M$. Locally, an ARS on $M$ can be defined by $\left\{X_{1}, X_{2}, \ldots, X_{n}\right\}$ satisfying the Lie algebra rank condition (Larc in short). This set of vector fields is considered as an orthonormal frame. We denote by $\Delta_{p}$ the linear span of the vector fields $\left\{X_{1}, X_{2}, \ldots, X_{n}\right\}$ at the point $p$. Recall that $\mathcal{Z}=\left\{p \in M: \operatorname{rank}\left(\Delta_{p}\right)<n\right\}$. If $\mathcal{Z}$ is empty, then the almostRiemannian structure is a Riemannian one (more details in [Agrachev et al., 2019]). In the 3-dimensional generic case, the singular set is formed by two types of points (see Figure 2.1): type-1 points where $\Delta_{p}$ has dimension 2 and is transversal to $\mathcal{Z}$, and type-2 points where $\Delta_{p}$ has dimension 2 and is tangent to $\mathcal{Z}$. Moreover, type-2 points are isolated (more details see [Boscain et al., 2015]). In the 2 -dimesional generic case also there are points where $\Delta_{p}$ is tangent to $\mathcal{Z}$ and isolated. Such points are called tangency points in [Agrachev et al., 2008].


Figure 2.1: Two different points form $\mathcal{Z}$ in the 3-D generic case.
On the other hand, nilpotent approximations are used to locally study the behavior of almost-Riemannian structures due to their significant similarity to the original dynamics. However, there are cases where the nilpotent approximation of an ARS turns out to be a constant rank sub-Riemannian structure. That is, some vector fields may vanish. The above is precisely what happens in the generic 3 -dimensional case, dealt in [Boscain et al., 2015], where at type2 points (or tangency points), the nilpotent approximation is the Heisenberg sub-Riemannian structure and hence is not a 3 -ARS.

The solvable approximation was introduced in [Jouan and Manríquez, 2022] to recover the almost-Riemannian structure lost in the nilpotent approximation. In that paper was considered the case where only one of the vector fields vanishes in the nilpotent approximation and the other ones are independents. Then, the solvable approximation is a local approximation of an ARS at singular points where the nilpotent approximation is no longer an ARS but a constant rank sub-Riemannian structure. A generalization of this approximation is given in [?], where a complete description of the nilpotent and solvable approximations is addressed, including the generic ones.

We can identify, mainly, two axes of work in the research, related basically to the structures of the approximation systems and estimation of the different distances defined by the original structure, the nilpotent approximation and the solvable one. Concerning structures, the Lie algebra generated by this new family of vector fields is finite-dimensional and solvable (in the generic case). Thanks to the equivalence theorem of [Jouan, 2010] we know that the space $\mathbb{R}^{n}$ is diffeomorphic to some homogeneous space or Lie groups. Through this diffeomorphism, the
solvable approximation and the nilpotent one are equivalent to a linear ARS on a homogeneous space or a Lie group. Moreover, when we work with $n$-ARS, we can find different and complex structures. For this reason, we dealt with generic ARSs by determining the generic distributions on a $n$-dimensional connected manifold and used them to exhibit the generic nilpotent and solvable approximations. Regarding the axis of the distance estimation, the original system, the nilpotent and solvable approximations give rise to three different distances: $d, \widehat{d}$ and $\widetilde{d}$ respectively. We stated that generically, the distance $\widetilde{d}$ is closer than $\widehat{d}$ to $d$. To prove this result, it is important to determine two facts. First, to state the order of approximations of $d$ by $\widetilde{d}$, and second, to find translation directions such that the distance $\widetilde{d}$ of a pair of translated points is decreasing.

## Continuation Method in rolling problem with obstacles.

In the second part of this thesis, we address the motion planning problem of a strictly convex body $A_{2}$ rolling (without slipping nor spinning) on the Euclidean plane $\mathbb{R}^{2}$ with obstacles, with a numerical implementation of the Continuation Method.

A nonholonomic system on a $n$-dimensional manifold $M$ is a control system which is of the form

$$
\dot{x}=\sum_{i=1}^{m} u_{i} X_{i}(x), \quad x \in M,
$$

where $m>1$ is an integer and $X_{1}, X_{2}, \ldots, X_{m}$ are $C^{\infty}$ vector field on $M$ (cf. [Jean, 2014]). These systems have attracted the attention of many authors from different disciplines for their varied applications, mainly in robotics (see [Murray et al., 2017] and references therein). The rolling body problem (without slipping or spinning) of a 2-dimensional Riemannian manifold on another one (which is an excellent example of the fusion between the sub-Riemannian geometry and the (geometric) control theory), can be written as a nonholonomic system. More accurately, the rolling-body problem (without slipping nor spinning) is a control system $\Sigma$, which models the rolling of an embedded connected surface $A_{2}$ in $\mathbb{R}^{3}$ on another $A_{1}$ one. As a consequence of the rolling constraints, and given an absolutely continuous (a.c.) curve $\gamma_{1}$ on $A_{1}$, there exists a unique a.c. curve $\Gamma$ in the state space, which describes the rolling of the surface $A_{2}$ onto the surface $A_{1}$ along the curve $\gamma_{1}$. Thus, the admissible controls of $\Sigma$ correspond to the a.c. curves $\gamma_{1}$ of $A_{1}$ by their derivatives $\dot{\gamma}_{1}$. Then, the system of control can be written, in local coordinates, as a nonholonomic system

$$
\Sigma: \quad \dot{x}=u_{1} X_{1}+u_{2} X_{2},
$$

where $\left(u_{1}, u_{2}\right) \in \mathbb{R}^{2}$ is the control, and $X_{1}$ and $X_{2}$ are vector fields.
The Continuation Method, was introduced in [Sussmann, 1992] and [Sussmann, 1993] and widely developed in [Chitour and Sussmann, 1998], [Chitour, 2002], and [Chitour, 2006]. It is used to solve nonlinear equations of the form $\mathcal{L}(x)=y$, where $x$ is the unknown, and $\mathcal{L}: X \rightarrow$ $Y$ is surjective. In the context of motion planning, $\mathcal{L}$ is the endpoint map (associated with some fixed point $p$ ) from the space of admissible inputs to the state space, that is $E_{p}: \mathcal{H} \rightarrow M$.

The convex body $A_{2}$, which can be embedded as a convex surface in $\mathbb{R}^{3}$, is assumed to have a stable periodic geodesic and it is defined by the function $a(x, y, z)$ such that $A_{2}=a^{-1}(0)$. We denote the state space simply by $M$. An obstacle $\mathcal{W}$ in $\mathbb{R}^{2}$ is a nonempty compact subset of $\mathbb{R}^{2}$. W maps in $M$ a region $C$. Thus, an obstacle in $M$ is a nonempty closed subset of $M$ such that $\widehat{M}=M \backslash C$ is also nonempty. Let us consider the control system on $M$ defined by

$$
\dot{y}=v_{1} \bar{X}_{1}(y)+v_{2} \bar{X}_{2}(y)
$$

where $\bar{X}_{i}=\zeta X_{i}, i=1,2$ and $\zeta: M \rightarrow \mathbb{R}$ such that $\zeta>0$ on $\widehat{M}, \zeta=0$ on $C$. Then $\widehat{M}$ is invariant under the above control system. Thus, the motion planning problem with obstacles is reduced to a motion planning problem for each connected component of $\widehat{M}$.

Therefore, we provide a complete numerical implementation of the Continuation Method presented above to solve the motion planning problem with forbidden regions on the plane. We give details about the fundamental points for the numerical implementation, which are the discretization of the control space $\mathcal{H}$, the computation of $D E_{p}(u)$, and the lift of the curve $\widetilde{\gamma}_{1}$ on the convex body $A_{2}$. This part provides three examples of the bodies rolling on the Euclidean plane (with obstacles): the sphere, the flattened ball, and an egg, this is to say, with a stable periodic geodesic (see Figure A.2). However, when we relax this geometric property, the method implemented still works.


Figure 2.2: Different examples of bodies rolling on the plane with obstacles.

## Part I

## Local approximation by linear systems and Almost-Riemannian Structures on Lie groups

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An Almost-Riemannian Structure (ARS in short) on an $n$-dimensional differential manifold is a rank-varying sub-Riemannian structure that can be defined, at least locally, by a set of $n$ vector fields satisfying the Lie algebra rank condition (Larc in short). We denote by $\Delta_{p}$ the linear span of the vector fields at the point $p$. The set of points where $\operatorname{dim}\left(\Delta_{p}\right)<n$ is called the singular locus or the singular set and denoted by $\mathcal{Z}$. Many papers dedicated to the study of ARSs can be found in the literature, for instance [Agrachev et al., 2010], [Bonnard et al., 2009], [Bonnard et al., 2011], [Boscain et al., 2013a] [Boscain et al., 2013b].

In the generic 3 -dimensional case, in which we are particularly interested, the singular set is a codimension one embedded submanifold and the points where $\Delta_{p}=T_{p} \mathcal{Z}$ are isolated. Such points are called tangency points in [Agrachev et al., 2008] and type-2 points in [Boscain et al., 2015].

We are likewise interested in the so-called linear ARSs on Lie groups (or homogeneous space)(see [Ayala and Jouan, 2016]) because they will be used as approximating structures for general ARSs. Under some conditions, the singular set of such structures is a subgroup or an analytic, embedded, codimension one submanifold (for more details about these structures, see [Ayala and Jouan, 2016], and [Jouan et al., 2018]).

On the other hand, nilpotent approximations are used to locally study the behavior of almost-Riemannian structures due to their significant similarity to the original dynamics. However, in some cases the nilpotent approximation of an ARS degenerates, because it is no longer an ARS but a constant rank sub-Riemannian structure. In other words, it may happen that some of the vector fields of the nilpotent approximation vanish, changing the almost-Riemannian structure into a constant rank sub-Riemannian one. For instance, if

$$
X_{1}=\left(\begin{array}{l}
1 \\
0 \\
0
\end{array}\right), \quad X_{2}=\left(\begin{array}{l}
0 \\
1 \\
x
\end{array}\right) \quad \text { and } \quad X_{3}=\left(\begin{array}{c}
0 \\
0 \\
x^{2}
\end{array}\right)
$$

then its nilpotent approximation is

$$
\widehat{X}_{1}=\left(\begin{array}{l}
1 \\
0 \\
0
\end{array}\right), \quad \widehat{X}_{2}=\left(\begin{array}{l}
0 \\
1 \\
x
\end{array}\right) \quad \text { and } \quad \widehat{X}_{3}=0 .
$$

that is, the Heisenberg sub-Riemannian structure. The above is exactly what happens in the generic 3-dimensional case, dealt in [Boscain et al., 2015].

We assume that only $m$ elements of the $\widehat{X}_{1}, \widehat{X}_{2}, \ldots, \widehat{X}_{n}$ are linearly independent. Let us assume without loss of generality that these vector fields are $\widehat{X}_{1}, \widehat{X}_{2}, \ldots, \widehat{X}_{m}$ and that $\widehat{X}_{m+1}, \widehat{X}_{m+2}, \ldots, \widehat{X}_{n}$ vanish. Our aim consists in recovering the almost-Riemannian structure lost in the nilpotent approximation thanks to the vector fields, denoted by $\widetilde{X}_{i}$ for $i=$ $m+1, m+2, \ldots, n$, which are the homogeneous component of degree 0 of the Taylor expansion in privileged coordinates of the vector fields that vanish. The new family of vector fields composed by $\widehat{X}_{1}, \widehat{X}_{2}, \ldots, \widehat{X}_{m}$ and $\widetilde{X}_{m+1}, \widetilde{X}_{m+2}, \ldots, \widetilde{X}_{n}$ is called the solvable approximation. The Lie algebra generated by this new family of vector fields is finite-dimensional. Moreover, when only one of the vector fields vanishes and the other ones are independent, this Lie algebra is solvable (hence the name of this approximation). We are also interested in some nilpotent Lie group on which $\widetilde{X}_{m+1}, \widetilde{X}_{m+2}, \ldots, \widetilde{X}_{n}$ act as linear vector fields (see Definition 2.2.1).

## This first part is organized as follows.

Chapter 2 contains generalities about ARSs, nonholonomic order, privileged coordinates, the nilpotent approximation, linear vector fields, and linear ARS on Lie groups or homogeneous spaces.

In the first section of Chapter 3, we introduce the definition of a solvable approximation considering $s=1$; this is to say only one of the vector fields vanishes, and the other ones are independent in the nilpotent approximation; we analyze its algebraic structures and the distance defined by the solvable approximation ( $\widetilde{d}$ ), concluding that the Lie algebra generated by this new family of vector fields is finite-dimensional and solvable (Proposition 1.3.1), the solvable approximation is equivalent to a simple ARS on a homogeneous space or a Lie group (Theorem 1.3.1), and the distance $\widetilde{d}$ always satisfies $\widetilde{d} \leq \widehat{d}$ (Proposition 1.3.3), where $\widehat{d}$ is the distance induced by the nilpotent approximation.

In the second section, we addressed the 3 -dimensional generic case. Here, we show that the solvable approximation is a simple ARS on $\mathbb{R}^{3}$ diffeomorphic to a quotient of the 5 -dimensional group Heisenberg $\mathbb{H}^{2}$. Regarding the distance issue, we have obtained in some cases that the order of the approximation of $d$ (the original distance) by $\tilde{d}$ is better than the order of the approximation of $d$ by $\widehat{d}$ (Theorem 2.3.1). Assuming that for some pairs ( $q, q^{\prime}$ ) of points translated from the singular locus the distance is decreasing (Theorem 2.3.2) and considering the order of approximation by $\widetilde{d}$, we prove that the difference $\left|d\left(q, q^{\prime}\right)-\widetilde{d}\left(q, q^{\prime}\right)\right|$ is strictly smaller than $\left|d\left(q, q^{\prime}\right)-\widehat{d}(q, q)\right|$. Finally, we provide the Hamiltonian associated to the flow defined by the solvable approximation in the 3D generic case and we compute the geodesics with initial condition $x(0)=y(0)=z(0)=0$.

Chapter 4 is devoted to general nilpotent and solvable approximations of almost-Riemannian structures. Firstly, we show that it is always possible to define the ARS locally, around the point $p=0$ in local privileged coordinates, by a set of $n$ orthonormal vector fields $X_{1}, \ldots, X_{n}$ such that the solvable approximation

$$
\widehat{X}_{1}, \ldots, \widehat{X}_{k}, \widehat{X}_{k+1}, \ldots, \widehat{X}_{m}, \widetilde{X}_{m+1}, \ldots, \widetilde{X}_{n},
$$

satisfies

- $\widehat{X}_{i}(0) \neq 0$ for $i=1, \ldots, k$;
- $\widehat{X}_{i} \neq 0$ but $\widehat{X}_{i}(0)=0$ for $i=k+1, \ldots, m$;
- $\widehat{X}_{i}=0$ for $i=m+1, \ldots, n$.

Secondly, we prove that the nilpotent or solvable approximation of an ARS at a singular point is a linear almost-Riemannian structure on a Lie group or a homogeneous space (Theorem 1.4.2 and 1.4.4 respectively), except in some very degenerated cases where neither the nilpotent approximation nor the solvable one defines an ARS.

When we work with $n$-ARS, we can find different and complex structures. For this reason, thirdly we determine the generic structures. It is shown in particular that generically: (1) the singular set $\mathcal{Z}$ is a union of submanifolds $\mathcal{Z}_{r}$ of codimension $r^{2}$ where the rank is $n-r ;$ (2) the rank of $\Delta+[\Delta, \Delta]$ is everywhere full ( $\Delta$ stands for the distribution) (Theorem 2.4.1). The structure of the points of $\mathcal{Z}_{r}$ where $\operatorname{dim}\left(T_{p} \mathcal{Z}_{r}\right)+\operatorname{dim}\left(\Delta_{p}\right)$ is not maximal is described in Theorem 2.4.2. For example in $\mathcal{Z}_{1}$ these points are the so-called tangency points (see [Boscain et al., 2015]), i.e. the points where $T_{p} \mathcal{Z}_{1}=\Delta_{p}$. They are generically isolated in $\mathcal{Z}_{1}$.

Thanks to these genericity results and with the help of local normal forms (see Section 2.2.4) it is finally shown that generically, there are only two possibilities for the nilpotent/solvable approximation at a point $p \in \mathcal{Z}$ (Theorem 2.4.3):

1. At a tangency point $p$ in $\mathcal{Z}_{1}$ one vector field of the nilpotent approximation vanishes, but the solvable approximation is not degenerated and defines a linear ARS.
2. At all other points, that is to say, at non-tangency points of $\mathcal{Z}_{1}$ and at all points in $\mathcal{Z}_{r}$ with $r \geq 2$, the nilpotent approximation is not degenerated.

In conclusion, the only generic points where the solvable approximation is useful are tangency points in $\mathcal{Z}_{1}$.

In the last section, we deal with the distance induced by the solvable approximation, at tangency points, of an $n$-dimensional ARS considering generic assumptions and the normal form when the point belonging to the singular set is a tangency point.

Let $d, \widetilde{d}$, and $\widehat{d}$ be the distances induced by the original structure, the solvable approximation, and the nilpotent one, respectively. The main result is Theorem 3.4.4 which states that generically, the distance $\widetilde{d}$ is closer to $d$ than $\widehat{d}$ for pairs of points translated in an appropriate direction (Section 4.3.4). This translation condition is significant because the distance $\widetilde{d}$ is not closer to $d$ than the distance induced by the nilpotent approximation for any pair of points. Then to prove the main result, it is essential to determine two facts. First, to state the order of approximation of $d$ by $\widetilde{d}$ (Theorem 3.4.1), and second, to find translation directions such that the distance $\widetilde{d}$ of a pair of translated points is decreasing.

To state the order of approximation of $d$ by $\widetilde{d}$, we analyze the divergence of curves admissible for $d$ and $\widetilde{d}$, defined by the same control functions and starting at the same point (see Proposition 3.4.1). By using this fact, we obtain that the distance induced by the solvable approximation improves the order of approximation of $d$ given by $\widehat{d}$ (see Theorem 3.4.1).

To find translation directions, we consider a vector field $Y \in \mathfrak{g}^{1}$ and then $Y \in \mathfrak{g}^{2}$, where $\mathfrak{g}=\mathfrak{g}^{1} \oplus \mathfrak{g}^{2}$ is the ideal generated in $\mathcal{L}$ by $\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}$, with $\mathfrak{g}^{s}$ the set of homogeneous vector fields of order $-s$, and $\mathcal{L}$ the Lie algebra generated by the solvable approximation.

## CHAPTER 2

## Mathematical prerequisites

This chapter aims to present some fundamental notions briefly for the comprehension of this thesis. They are the almost-Riemannian structure, linear almost-Riemannian structure, nonholonomic order, privileged coordinates, and nilpotent approximation.

For everything related to general sub-Riemannian geometry, including almost-Riemannian one, the reader is referred to [Agrachev et al., 2019], except for the linear almost-Riemannian structure; in this case, the reader is referred to [Ayala and Jouan, 2016]. Regarding the rest of the topics, the reader is referred to [Jean, 2014] or [Bellaïche, 1996].

## 1.2 . Almost Riemannian structures

Let $M$ be a $n$-dimensional, connected, $C^{\infty}$ manifold. The $C^{\infty}$-module of $C^{\infty}$ vector fields on $M$ is denoted by $\Gamma(M)$. Let $\Delta$ be a sub-module of $\Gamma(M)$. The flag of submodules

$$
\begin{equation*}
\Delta=\Delta^{1} \subseteq \Delta^{2} \subseteq \cdots \subseteq \Delta^{k-1} \subseteq \Delta^{k} \subseteq \ldots \tag{2.1}
\end{equation*}
$$

is defined by induction: $\Delta^{2}=\Delta+[\Delta, \Delta]$ is the sub-module of $\Gamma(M)$ generated by $\Delta$ and the Lie algebra of its elements, and $\Delta^{k+1}=\Delta^{k}+\left[\Delta, \Delta^{k}\right]$. The Lie algebra generated by $\Delta$ is $\mathcal{L}(\Delta)=\bigcup_{k \geq 1} \Delta^{k}$. The sub-module $\Delta$ satisfies the rank condition if the evaluaction of $\mathcal{L}(\Delta)$ at each point $p$ is equal to $T_{p} M$.

Definition 1.2.1 An almost-Riemannian structure (resp. distribution) on a smooth $n$-dimensional manifold $M$ is a triple $(E, f,\langle\cdot, \cdot\rangle)$ (resp. a pair $(E, f)$ ) where $E$ is a rank $n$ vector bundle over $M, f: E \rightarrow T M$ is a morphism of vector bundles, and $(E,\langle\cdot, \cdot\rangle)$ is an Euclidean bundle, that is $\langle\cdot, \cdot\rangle_{p}$ is an inner product on the fiber $E_{p}$ of $E$, smoothly varying w.r.t. p, assumed to satisfy the following properties:

1. The set of points $p \in M$ such that the restriction of $f$ to $E_{p}$ is onto is a proper open and dense subset of $M$;
2. the module $\Delta$ of vector fields on $M$ defined as the image by $f$ of the module of smooth sections of $E$ satisfies the rank condition.

The set of points of $M$ where the rank of $f\left(E_{p}\right)=\Delta_{p}$ is less than $n$ is called the singular locus (or singular set) of the almost-Riemannian structure and denoted by $\mathcal{Z}$.

The inner product on $E$ induces a bilinear symmetric and positive definite mapping, also denoted by $\langle\cdot, \cdot\rangle$, from $\Delta \times \Delta$ to $C^{\infty}(M)$. Indeed, and element $X \in \Delta($ resp. $Y \in \Delta)$ is the image by $f$ of a unique section $\sigma$ (resp. $\eta$ ) of $E$ and we can set $\langle X, Y\rangle_{p}=\langle\sigma, \eta\rangle_{p}$. Consequently andalmost-Riemannian structure (ARS in short) can be alternately defined as follows.

Definition 1.2.2 An almost-Riemannian structure on a smooth $n$-dimensional manifold $M$ is a pair $(\Delta,\langle\cdot, \cdot\rangle)$ where $\Delta$ is a sub-module of $\Gamma(M)$ that can be locally defined by $n$ vector fields and satisfies the rank condition, and $\langle\cdot, \cdot\rangle$ is a bilinear symmetric and positive definite mapping from $\Delta \times \Delta$ to $C^{\infty}(M)$, such that the set $\mathcal{Z}$ of points $p$ where the dimension of $\Delta_{p}$ is less than $n$ is nonempty but with empty interior.

Around any point $p \in M$ the sub-module $\Delta$ can be locally defined by an orthonormal frame $\left(X_{1}, X_{2}, \ldots, X_{n}\right)$. It is enough to select a set of $n$ sections $\left(e_{1}, e_{2}, \ldots, e_{n}\right)$ of $E$ orthonormal in a neighborhood of $p$ and define $X_{i}=f_{*} e_{i}$, where $f_{*} e_{i}=f \circ e_{i}$.

The almost-Riemannian norm on $\Delta_{p}$ is defined by

$$
\|v\|=\min \left\{\|u\|^{2}: u \in E_{p}, f(u)=v\right\} .
$$

An absolutely continuous curve $\gamma:[0, T] \longrightarrow M$ is admissible for $E$ if there exists a measurable essentially bounded function $t \mapsto u(t)$ from $[0, T]$ into $E_{\gamma(t)}$ called control function such that $\dot{\gamma}(t)=f(u(t))$ for almost every $t \in[0, T]$. Locally this means that $\dot{\gamma}(t)=u_{1}(t) X_{1}(\gamma(t))+$ $u_{2}(t) X_{2}(\gamma(t))+\ldots+u_{n}(t) X_{n}(\gamma(t))$ for almost every $t \in[0, T]$, where $X_{1}, X_{2}, \ldots, X_{n} \in \Delta$.

Given an admissible curve $\gamma:[0, T] \longrightarrow M$, the length of $\gamma$ is defined by

$$
l(\gamma)=\int_{0}^{T}\|\dot{\gamma}(t)\| d t
$$

The almost-Riemannian distance (or Carnot-Caratheodory distance) on $M$ associated with the $n$-ARS is defined as

$$
d\left(p_{0}, p_{1}\right)=\inf \left\{l(\gamma): \gamma(0)=p_{0}, \gamma(T)=p_{1}, \gamma \text { admissible }\right\}
$$

It induces the usual topology on $M$.

Remark 1.2.1 The structure is Riemannian out of $\mathcal{Z}$.

### 2.2 ARS on Lie groups and homogeneous spaces

### 1.2.2 . Linear vector fields

Let $G$ be a connected Lie group and $\mathfrak{g}$ its Lie algebra (the set of left-invariant vector fields, identified with the tangent space at the identity). The set of analytic vector fields on $G$ is denoted by $V^{\omega}(G)$, and the normalizer of $\mathfrak{g}$ in $V^{\omega}(G)$ is by definition

$$
\mathcal{N}=\operatorname{norm}_{V^{\omega}(G)} \mathfrak{g}=\left\{X \in V^{\omega}(G): \forall Y \in \mathfrak{g} \quad[X, Y] \in \mathfrak{g}\right\}
$$

Definition 2.2.1 A vector field $\mathcal{X}$ on $G$ is said to be linear or to be infinitesimal automorphism (see [Bourbaki, 2007]), if $\mathcal{X}$ belongs to $\mathcal{N}$ and $\mathcal{X}(e)=0$, where $e$ is the identity of $G$.

We can see in [Jouan, 2010] that a vector field $\mathcal{X}$ on $G$ if and only its flow $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ is a one-parameter group of automorphisms of $G$ and a linear vector field is consequently analytic and complete.

### 2.2.2 . Linear ARS's on Lie groups

Linear and invariant vector fields make it possible to define almost-Riemannian structures on Lie groups.

Definition 2.2.2 A linear ARS on $G$ is an almost-Riemannian structure defined by a set of

- $n-s$ left-invariant vector fields $Y_{1}, Y_{2}, \ldots, Y_{n-s}$.
- $s>0$ linear vector fields $\mathcal{X}_{n-s+1}, \ldots, \mathcal{X}_{n}$,
assumed to satisfy the rank condition and to have full rank on a proper open and dense subset of $G$.

A linear ARS is said to be simple if $s=1$.
For instance, the famous Grushin plane on the Abelian Lie group $\mathbb{R}^{2}$ is a simple ARS. This structure was introduced in [Ayala and Jouan, 2016] and its isometries have been studied in [Jouan et al., 2018].

### 3.2.2 . Linear ARS's on homogeneous spaces

Consider a homogeneous space $G / H$ of a connected and simply connected Lie group $G$ by a closed subgroup $H$ (the elements of $G / H$ are right cosets of $H$ because we deal with leftinvariant vector fields). Since we are interested in simply connected quotients, we assume $H$ to be connected. The projection of a left-invariant vector field $Y$ onto $G / H$ is well-defined, is referred to as a left-invariant vector field, and we can assume that it vanishes identically only if $Y$ is the zero-field (see details in [Jouan, 2010]). On the other hand, the projection of a linear field $\mathcal{X}$ of $G$ does exist on $G / H$ if and only if $H$ is invariant under its flow, or equivalently,
because $H$ is connected if the Lie algebra of $H$ is $\operatorname{ad}(\mathcal{X})$-invariant. This allows to define linear vector fields and linear ARS on $G / H$ in the same way as on Lie groups.

In the sequel, we will need a version of the equivalence Theorem (see [Jouan, 2010] and [Ayala and Jouan, 2016]).

Theorem 2.2.1 (Equivalence Theorem) Let $f_{1}, \ldots, f_{n}$ be a set of $n$ complete vector fields on a manifold $M$ and let us assume:

1. $f_{1}, \ldots, f_{n}$ define an almost-Riemannian structure on $M$;
2. The Lie algebra $\mathcal{L}$ generated by $f_{1}, \ldots, f_{n}$ is finite dimensional;
3. The ideal $\mathfrak{g}$ generated in $\mathcal{L}$ by $f_{1}, \ldots, f_{n-s}$ is nilpotent;

Then $M$ is diffeomorphic to a homogeneous space $G / H$ of the nilpotent simply connected group $G$ generated by $\mathfrak{g}$ and $f_{1}, \ldots, f_{n}$ defines a linear ARS on $G / H$. More accurately the vector fields $f_{1}, \ldots, f_{n-s}$ are left-invariant and $f_{n-s+1}, \ldots, f_{n}$ are linear on this homogeneous space.

## 3.2 . Nilpotent approximation

### 1.3.2 . Nonholonomic orders

Definition 3.2.1 Let $f: M \rightarrow \mathbb{R}$ be a continuous function. The nonholonomic order of $f$ at $p$, denoted $\operatorname{ord}_{p}(f)$, is the real number defined by

$$
\operatorname{ord}_{p}(f)=\sup \left\{s \in \mathbb{R}: f(q)=O\left(d(p, q)^{s}\right)\right\} .
$$

This order is always nonnegative.
Let $C^{\infty}(\mathrm{p})$ denote the set of germs of smooth functions at $p$. For $f \in C^{\infty}(p)$, we call nonholonomic derivative of order 1 of $f$ the Lie derivatives $X_{1} f, \cdots, X_{n} f$. We call further $X_{i} X_{j} f, X_{i} X_{j} X_{k} f, \ldots$, the nonholonomic derivatives of $f$ of order $2,3, \ldots$ of $f$. The nonholonomic derivative of order 0 of $f$ at $p$ is $f(p)$.

As a consequence, the nonholonomic order of a smooth (germ of) function is given by the formula

$$
\operatorname{ord}_{p}(f)=\min \left\{s \in \mathbb{N}: \exists i_{1}, \ldots, i_{s} \in\{1, \ldots, n\} \text { s.t. }\left(X_{i_{1}} \ldots X_{i_{s}} f\right)(p) \neq 0\right\},
$$

where as usual we adopt the convention that $\min \emptyset=+\infty$.
Let $V F(p)$ denote the set of germs of smooth vector fields at $p$.
Definition 3.2.2 Let $X \in V F(p)$. The nonholonomic order of $X$ at $p$, denoted by $\operatorname{ord}_{p}(X)$, is the real number defined by:

$$
\operatorname{ord}_{p}(X)=\sup \left\{\sigma \in \mathbb{R}: \operatorname{ord}_{p}(X f) \geq \sigma+\operatorname{ord}_{p}(f), \quad \forall f \in C^{\infty}(p)\right\} .
$$

### 2.3.2 . Privileged coordinates

Let $p$ be a point of $M$ and let $\Delta_{p}^{k}, k \geq 1$ be the evaluation of the submodule $\Delta^{k}$ at $p$. Thanks to the rank condition these submodules at $p$ forms a flag of subspaces of $T_{p} M$,

$$
\begin{equation*}
\Delta_{p}^{1} \subset \Delta_{p}^{2} \subset \cdots \subset \Delta_{p}^{r-1} \varsubsetneqq \Delta_{p}^{r}=T_{p} M \tag{2.2}
\end{equation*}
$$

where $r=r(p)$ is called the degree of nonholonomy at $p$. Let $n_{i}(p)=\operatorname{dim} \Delta_{p}^{i}$. The $r$-tuple of integers $\left(n_{1}(p), \ldots, n_{r}(p)\right)$ is called the growth vector at $p$. The first integer in the growth vector is the rank $n_{1}(p) \leq n$ of the family $X_{1}(p), \ldots, X_{n}(p)$, and the last one $n_{r}(p)=n$ is the dimension of $M$. The structure of the flag (2.2) may also be described by another sequence of integers. We define the weights at $p, w_{i}=w_{i}(p), i=1, \ldots, n$, by setting $w_{j}=s$ if $n_{s-1}(p)<j \leq n_{s}(p)$, where $n_{0}=0$. In other words, we have

$$
w_{1}=\cdots=w_{n_{1}}=1, w_{n_{1}+1}=\cdots=w_{n_{2}}=2, \ldots, w_{n_{r-1}+1}=\cdots=w_{n_{r}}=r .
$$

Definition 3.2.3 A system of privileged coordinates at $p$ is a system of local coordinates $\left(x_{1}, \ldots, x_{n}\right)$ such that $\operatorname{ord}_{p}\left(x_{j}\right)=w_{j}$, for $j=1, \ldots, n$.

Constructions of privileged coordinates can be found in [Bellaïche, 1996] and [Jean, 2014].

On the other hand, systems of privileged coordinates always exist (under the rank condition) and in such a system, given a sequence of integers $\alpha=\left(\alpha_{1}, \ldots, \alpha_{n}\right)$, we define the weight of the monomial $x^{\alpha}=x_{1}^{\alpha_{1}} \cdots x_{n}^{\alpha_{n}}$ to be $w(\alpha)=w_{1} \alpha_{1}+\cdots+w_{n} \alpha_{n}$ and the weighted degree of the monomial vector field $x^{\alpha} \frac{\partial}{\partial x_{j}}$ to be $w(\alpha)-w_{j}$.

More generally, the nonholonomic order at $p$ of a function $f$ (resp. a vector field $X$ ) is the minimum of the homogeneous nonholonomic orders of the monomials of its Taylor series. It will be denoted by $\operatorname{ord}_{p}(f)\left(\operatorname{resp} . \operatorname{ord}_{p}(X)\right)($ see [Jean, 2014], Proposition 2.2).

Proposition 3.2.1 ([Bellaïche, 1996], Proposition 5.16) Let $X$ and $Y$ be vector fields on $M$. If $X$ and $Y$ are homogeneous of degree $k$ and $l$ respectively (in the chosen system of privileged coordinates), then $[X, Y]$ is homogeneous of degree $k+l$ or vanishes.

Definition 3.2.4 The function defined in privileged coordinates, by $x \mapsto\|x\|_{p}=\sum_{i=1}^{n}\left|x_{i}\right|^{\frac{1}{w_{i}}}$ is the so-called pseudo-norm at $p$.

Remark 3.2.1 Let $x=\left(x_{1}, \ldots, x_{n}\right)$ be a system of privileged coordinates defined on an open neighborhood $U$ of the point $p$. When composed with the coordinate functions, the pseudonorm at $p$ is (non smooth) homogeneous of order 1, that is, $\|x(q)\|_{p}=O(d(p, q))$, where $x(q)$ are the coordinates of $q \in U$.

### 3.3.2 . Nilpotent approximation

Fix a system of privileged coordinates $\left(x_{1}, \ldots, x_{n}\right)$ at $p$. Every vector field $X_{i}$ is of order $\geq-1$, hence it has, in $x$ coordinates, a Taylor expansion

$$
X_{i}(x) \sim \sum_{\alpha, j} a_{\alpha, j} x^{\alpha} \frac{\partial}{\partial x_{j}},
$$

where $w(\alpha) \geq w_{j}-1$ if $a_{\alpha, j} \neq 0$. Grouping together the monomial vector fields of same weighted degree we express $X_{i}$ as a series of homogeneous vector fields of the form

$$
\begin{equation*}
X_{i}=X_{i}^{(-1)}+X_{i}^{(0)}+X_{i}^{(1)}+X_{i}^{(2)}+\cdots, \tag{2.3}
\end{equation*}
$$

where $X_{i}^{(s)}$ has degree $s$. We set

$$
\widehat{X}_{i}=X_{i}^{(-1)}, \quad i=1, \ldots, n
$$

Definition 3.2.5 The family of vector fields $\left(\widehat{X}_{1}, \ldots, \widehat{X}_{n}\right)$ is called the nilpotent approximation of the system $\left(X_{1}, \ldots, X_{n}\right)$ at $p$.

Proposition 3.2.2 ([Bellaïche, 1996], Proposition 5.17) The vector fields $\widehat{X}_{i}, i=1, \ldots, n$, generate a nilpotent Lie algebra Lie $\left(\widehat{X}_{1}, \ldots, \widehat{X}_{n}\right)$ of step $r=w_{n}$. They satisfy Larc at every point $y \in \mathbb{R}^{n}$, and the distance $\widehat{d}$ is finite for every $x, y \in \mathbb{R}^{n}$.

To finish, we recall the very important Theorem 7.32 of [Bellaïche, 1996] stated here with a slight modification.

Theorem 3.2.1 (Theorem 7.32 in [Bellaïche, 1996]) There exist constants $\varepsilon>0$ and $C>0$ such that for any $q, q^{\prime} \in B(p, \varepsilon)$, we have

$$
-C \tau d\left(q, q^{\prime}\right)^{\frac{1}{r}} \leq d\left(q, q^{\prime}\right)-\widehat{d}\left(q, q^{\prime}\right) \leq C \widehat{\tau} \widehat{d}\left(q, q^{\prime}\right)^{\frac{1}{r}},
$$

where $\tau=\max \left(\|q\|_{p}, d\left(q, q^{\prime}\right)\right)$ and $\widehat{\tau}=\max \left(\|q\|_{p}, \widehat{d}\left(q, q^{\prime}\right)\right)$.
Remark 3.2.2 The original result stated in [Bellaïche, 1996] is as it follows

$$
-C \widehat{d}(p, q) d\left(p, q^{\prime}\right)^{\frac{1}{r}} \leq d\left(q, q^{\prime}\right)-\widehat{d}\left(q, q^{\prime}\right) \leq C \widehat{d}(p, q) \widehat{d}\left(q, q^{\prime}\right)^{\frac{1}{r}} .
$$

Notice that if we consider $q=p$, we obtain that $d\left(p, q^{\prime}\right)=\widehat{d}\left(p, q^{\prime}\right)$. Of course, this is not always true.

## CHAPTER 3

## Solvable approximation and the 3-dimensional generic case

The content of this chapter is the result of an article in collaboration with Philippe Jouan, which was published in the journal Mathematical Control \& Related Fields, volume 12 (2022), Issue 2 (June), pp.303-326. ([Jouan and Manríquez, 2022]).

### 1.3. Solvable approximation

In this section, we introduce the solvable approximation of an ARS, we analyze its algebraic structure and the distance induced by it.

Let $\left\{X_{1}, \ldots, X_{n}\right\}$ be a set of vector fields defining an almost-Riemannian structure on an open neighborhood of $0 \in \mathbb{R}^{n}$. The point $p=0$ is assumed to belong to the singular locus, the natural coordinates of $\mathbb{R}^{n}$ to be privileged and we consider the nilpotent approximation $\left\{\widehat{X}_{1}, \ldots, \widehat{X}_{n}\right\}$ of $\left\{X_{1}, \ldots, X_{n}\right\}$ at $p=0$.

It may happen that some of the vector fields $\widehat{X}_{i}$ vanish, possibly changing the almostRiemannian structure defined by $X_{1}, \ldots, X_{n}$ into a constant rank sub-Riemannian one. It is what happens in some cases of generic 3-dimensional ARSs (see [Boscain et al., 2015]) that are described in detail in Section 2.3. In what follows, we are interested in the case where only one of the $\widehat{X}_{i}$ 's vanishes, say $\widehat{X}_{n}=0$. The other ones are independent and define a left-invariant sub-Riemannian structure on a Lie group, or a homogeneous space, the underlying manifold of which is $\mathbb{R}^{n}$.

Recall that each $X_{i}$ can be expanded in a series of homogeneous vector fields in the system of privileged coordinates at $p=0$, this is

$$
X_{i}=X_{i}^{(-1)}+X_{i}^{(0)}+X_{i}^{(1)}+\ldots, \quad \forall i \in\{1, \ldots, n\},
$$

where $X_{i}^{(k)}$ is the homogeneous component of degree $k$ (see [Bellaïche, 1996]). Denoting $\widetilde{X}_{n}=X_{n}^{(0)} \neq 0$, we introduce the following definition:
Definition 1.3.1 (Solvable approximation) The family $\left\{\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}, \widetilde{X}_{n}\right\}$ is the solvable approximation of $\left\{X_{1}, \ldots, X_{n}\right\}$.

The Lie algebra generated by this new family of vector fields is finite dimensional and solvable. Formally,
Proposition 1.3.1 $\mathcal{L}=\operatorname{Lie}\left(\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}, \widetilde{X}_{n}\right)$ is a finite dimensional solvable Lie algebra. Its step of solvability is less than or equal to $\log _{2}(r)+1$, where $r$ is the degree of nonholonomy at $p=0$.

### 1.1.3 . Structure of the approximating system

Despite the previous result, we are not interested in the solvable Lie group associated to the Lie algebra Lie $\left\{\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}, \widetilde{X}_{n}\right\}$ but in some nilpotent Lie group on which $\widetilde{X}_{n}$ acts as a linear vector field. For this reason, we denote by $\mathfrak{h}$ the Lie algebra generated by $\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}$ and by $\mathfrak{g}$ the ideal generated by $\mathfrak{h}$ in $\mathcal{L}=\operatorname{Lie}\left\{\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}, \widetilde{X}_{n}\right\}$.
Proposition 1.3.2 The ideal $\mathfrak{g}$ is the space of vector fields of $\mathcal{L}$ whose nonholonomic order is negative. It is a nilpotent Lie algebra and

$$
\mathcal{L}=\mathfrak{g} \oplus \mathbb{R} \widetilde{X}_{n}
$$

Moreover $D=-\operatorname{ad}\left(\tilde{X}_{n}\right)$ is a derivation of $\mathfrak{g}$.
Let $G$ be the simply connected Lie group whose Lie algebra is isomorphic to $\mathfrak{g}$. According to [Jouan, 2010] there exists a linear vector field on $G$ associated to the derivation $D=-\operatorname{ad}\left(\widetilde{X}_{n}\right)$. With a clear abuse of notation we will denote it by $\widetilde{X}_{n}$.

Thanks to the Theorem 2.2.1 (equivalence theorem) of [Jouan, 2010], we know that the space $\mathbb{R}^{n}$ is diffeomorphic to a homogeneous space. The solvable approximation is equivalent to a simple ARS on a homogeneous space or a Lie group through this diffeomorphism.
Theorem 1.3.1 The space $\mathbb{R}^{n}$ is diffeomorphic to a homogeneous space $G / G_{0}$ of $G$. Through this diffeomorphism $\left\{\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}, \widetilde{X}_{n}\right\}$ is equivalent to a simple $A R S$ on $G / G_{0}$, and the Lie algebra $\mathfrak{g}_{0}$ of $G_{0}$ is isomorphic to the set of vector fields of $\mathfrak{g}$ that vanish at 0 .
We are also interested in conditions for which $G=\mathbb{R}^{n}$.
Theorem 1.3.2 With the previous notations, the following assertions are equivalent:
(i) $\operatorname{ad}\left(\widetilde{X}_{n}\right) \cdot \widehat{X}_{i}$ belongs to Span $\left\{\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}\right\}$ for $i=1, \ldots, n-1$;
(ii) $\mathfrak{h}$ is ad $\left(\widetilde{X}_{n}\right)$-invariant;
(iii) $\mathfrak{h}=\mathfrak{g}$.

Under these conditions $\widetilde{X}_{n}$ is a linear vector field on $\exp (\mathfrak{h})$.

### 2.1.3 . Distance

On the other hand, the solvable approximation gives rise to a distance denoted by $\widetilde{d}$. This distance has the advantage to be really almost-Riemannian unlike the distance $\widehat{d}$ associated to the nilpotent approximation in the degenerated cases. Moreover, the distance $\widetilde{d}$ always satisfies $\widetilde{d} \leq \widehat{d}$. This is to say,

Proposition 1.3.3 For all $x, y \in \mathbb{R}^{n}, \widetilde{d}(x, y) \leq \widehat{d}(x, y)$.
Remark 1.3.1 It is not possible to define the solvable approximation in a 2-dimensional ARS because if a vector field vanishes in the nilpotent approximation, then it does not satisfy LARC, which is a contradiction according to Proposition 3.2.2.

## 2.3 . The 3-dimensional generic case

In this section, we describe the solvable approximation in the 3-dimensional generic case, this is to say, the algebraic structure, the distance issue, and we compute the geodesic with initial condition $x(0)=y(0)=z(0)=0$ and covector $\lambda=(p, q, r) \in T^{*} \mathbb{R}^{3}$.

Regarding the results about the algebraic structure, we show that the solvable approximation is a simple ARS on $\mathbb{R}^{3}$ diffeomorphic to a quotient of the 5 -dimensional Heisenberg group $\mathbb{H}^{2}$. Concerning the distance issue, we have obtained that the solvable approximation improves the order of approximation of $d$ given by the nilpotent approximation (Theorem 2.3.1). Moreover, we set directions of translation for a pair of points in which $\widetilde{d}$ is decreasing (Theorem 2.3.2 and 2.3.3). The latter is fundamental because it allows us to show that the solvable distance is closer than the nilpotent one to the original distance for a pair of points translated in such directions.

Recall that $\Delta_{p}=\operatorname{span}\left\{X_{1}(p), \ldots, X_{n}(p)\right\}$ and the singular locus $\mathcal{Z}$ is the set of points of $\mathbb{R}^{n}$ where the rank of the linear span of the vector fields is less than $n$. From [Boscain et al., 2015] we take the following.

Proposition 2.3.1 Consider a 3-ARS. The following conditions are generic for 3-ARSs
(G1) $\operatorname{dim}\left(\Delta_{p}\right) \geq 2$ and $\Delta_{p}+[\Delta, \Delta]_{p}=T_{p} M$ for every $p \in M$;
(G2) $\mathcal{Z}$ is an embedded (possibly empty) two-dimensional submanifold of $M$;
(G3) the points where $\Delta_{p}=T_{p} \mathcal{Z}$ are isolated.
Proposition 2.3.2 Under the previous conditions there are three types of points:

1. Riemannian points where $\Delta_{p}=T_{p} M$.
2. type-1 points where $\Delta_{p}$ has dimension 2 and is transversal to $\mathcal{Z}$.
3. type-2 points where $\Delta_{p}$ has dimension 2 and is tangent to $\mathcal{Z}$.

Moreover, type-2 points are isolated, type-1 points form a 2-dimensional manifold, and all other points are Riemannian.

The local representation of the 3-dimensional ARS at type-2 points (see more details in [Boscain et al., 2015]) is given by the vector fields:

$$
X_{1}=\left(\begin{array}{l}
1 \\
0 \\
0
\end{array}\right), \quad X_{2}=\left(\begin{array}{c}
0 \\
1+\delta(x, y, z) \\
x(1+\theta(x, y, z))
\end{array}\right), \quad X_{3}=\left(\begin{array}{c}
0 \\
0 \\
a z+b x^{2}+c y^{2}+o\left(x^{2}+y^{2}+|z|\right)
\end{array}\right),
$$

where $\delta$ and $\theta$ are smooth functions of order greater than or equal to 1 and $a, b, c$ are not all zero. Furthermore, its nilpotent approximation in privileged coordinates is:

$$
\widehat{X}_{1}=X_{1}=\left(\begin{array}{l}
1 \\
0 \\
0
\end{array}\right), \quad \widehat{X}_{2}=\left(\begin{array}{l}
0 \\
1 \\
x
\end{array}\right), \quad \widehat{X}_{3}=0 .
$$

and

$$
\widetilde{X}_{3}=\left(\begin{array}{c}
0 \\
0 \\
a z+b x^{2}+c y^{2}
\end{array}\right) .
$$

Then $\left(\widehat{X}_{1}, \widehat{X}_{2}, \widetilde{X}_{3}\right)$ is the solvable approximation at $p=0$ in case when 0 is a tangential (type-2) point.

Notice that the Lie algebra generated by $\widehat{X}_{1}$ and $\widehat{X}_{2}$ is:

$$
\mathfrak{h}=\operatorname{Span}\left\{\widehat{X}_{1}, \widehat{X}_{2}, Z=\left[\widehat{X}_{1}, \widehat{X}_{2}\right]\right\} \quad \text { where } \quad Z=\left(\begin{array}{l}
0 \\
0 \\
1
\end{array}\right)=\frac{\partial}{\partial z},
$$

that is the Heisenberg algebra. On the other hand the algebra generated by $\widehat{X}_{1}, \widehat{X}_{2}$ and $\widetilde{X}_{3}$ is Span $\left\{\widehat{X}_{1}, \widehat{X}_{2}, Z,\left[\widehat{X}_{1}, \widetilde{X}_{3}\right],\left[\widehat{X}_{2}, \widetilde{X}_{3}\right], \widetilde{X}_{3}\right\}$, where:

$$
\left[\widehat{X}_{1}, \widetilde{X}_{3}\right]=\left(\begin{array}{c}
0 \\
0 \\
2 b x
\end{array}\right)=2 b x \frac{\partial}{\partial z} \text { and }\left[\widehat{X}_{2}, \widetilde{X}_{3}\right]=\left(\begin{array}{c}
0 \\
0 \\
2 c y+a x
\end{array}\right)=(2 c y+a x) \frac{\partial}{\partial z}
$$

and the ideal generated by $\widehat{X}_{1}$ and $\widehat{X}_{2}$ is:

$$
\mathfrak{g}=\operatorname{Span}\left\{\frac{\partial}{\partial x}, \frac{\partial}{\partial y}+x \frac{\partial}{\partial z}, 2 b x \frac{\partial}{\partial z}, 2 c y \frac{\partial}{\partial z}+a x \frac{\partial}{\partial z}, \frac{\partial}{\partial z}\right\} .
$$

A straightforward computation shows that $\widetilde{X}_{3}$ acts as a derivation on $\mathfrak{g}$. If we assume $b \neq 0$ and $c \neq 0$ then we have also:

$$
\mathfrak{g}=\operatorname{Span}\left\{\frac{\partial}{\partial x}, x \frac{\partial}{\partial z}, \frac{\partial}{\partial y}, y \frac{\partial}{\partial z}, \frac{\partial}{\partial z}\right\} .
$$

This is the 5-dimensional Heisenberg Lie algebra $\mathfrak{h}^{2}$, and in this basis, the derivation $D=$ $-\operatorname{ad}\left(\widetilde{X}_{3}\right)$ is given by the following matrix:

$$
D=\left(\begin{array}{ccccc}
0 & 0 & & & \\
2 b & a & & & \\
& & 0 & 0 & \\
& & 2 c & a & \\
& & & & a
\end{array}\right)
$$

Finally, the solvable approximation $\left(\widehat{X}_{1}, \widehat{X}_{2}, \widetilde{X}_{3}\right)$ is a simple ARS on $\mathbb{R}^{3}$ diffeomorphic to a quotient of the 5 -dimensional group Heisenberg $\mathbb{H}^{2}$.

### 1.2.3. Distance in the 3D generic case

We denote by $d$ the distance associated with the 3 -dimensional almost-Riemannian structure defined by the vector fields $X_{1}, X_{2}$ and $X_{3}$ at type- 2 points. The following Theorem state the order of approximation of $d$ by $\widetilde{d}$.

Theorem 2.3.1 If ord $(\delta) \geq 2$ and $\operatorname{ord}_{p}(\theta) \geq 2$, then there exist constants $C$ and $\epsilon>0$, such that, for all $q, q^{\prime} \in B(p, \epsilon)$, we have

$$
\begin{equation*}
-C \tau^{\frac{3}{2}} d\left(q, q^{\prime}\right)^{\frac{1}{2}} \leq d\left(q, q^{\prime}\right)-\widetilde{d}\left(q, q^{\prime}\right) \leq C \cdot \widetilde{\tau}^{\frac{3}{2}} \widetilde{d}\left(q, q^{\prime}\right)^{\frac{1}{2}} \tag{3.1}
\end{equation*}
$$

where

$$
\begin{aligned}
\tau & =\max \left(\|q\|_{p}, d\left(q, q^{\prime}\right)\right) \\
\widetilde{\tau} & =\max \left(\|q\|_{p}, \widetilde{d}\left(q, q^{\prime}\right)\right) .
\end{aligned}
$$

Thanks to formula (3.1) of Theorem 2.3.1, we know that, at least in some 3D generic cases, the order of the approximation of $d$ by $\tilde{d}$ is better than the one of the approximation of $d$ by $\widehat{d}$ (see Theorem 3.2.1). Indeed, this order is $d^{2}$ in the first case and $d^{\frac{3}{2}}$ in the second one. However, this does not prove that the solvable approximation is really better than the nilpotent one, and anyway, it is certainly not valid for any pair of points.

Since the nilpotent distance $\widehat{d}$ is left-invariant while $d$ and $\widetilde{d}$ are not (may be decreasing), we are interested in conditions under which $\widetilde{d}(g, g \cdot q) \leq \widetilde{d}(0, q)$, where the product is the Heisenberg one, $q$ is a point in a neighborhood of 0 , and $g \in \mathbb{R}^{3}$.

Let $\gamma(t)=(x(t), y(t), z(t))$ be a geodesic of $\widetilde{d}$ such that $\gamma(0)=0$ with control functions $u_{1}, u_{2}$ and $u_{3}$. We consider $g=\left(g_{1}, g_{2}, g_{3}\right) \in \mathbb{R}^{3}$. Let $\gamma_{g}(t)=L_{g}(\gamma(t))$. The goal is to find conditions for $g$ such that $\gamma_{g}$ has a length less than $\gamma$.

Since Lie $\left\{\widehat{X}_{1}, \widehat{X}_{2}\right\}$ is the Heisenberg algebra, then

$$
L_{g}(\gamma(t))=\left(x(t)+g_{1}, y(t)+g_{2}, z(t)+g_{1} y(t)+g_{3}\right) .
$$

We set $h(x, y, z)=a z+b x^{2}+c y^{2}$. Then $h\left(\gamma_{g}\right)=h(\gamma)+h(g)+f(g, \gamma)$, where $f(g, \gamma)=$ $(2 b x+a y) g_{1}+2 c y g_{2}$.

We assume that $h\left(\gamma_{g}\right)$ does not vanish, this is to say $\gamma_{g}$ is not on $\mathcal{Z}$. In particular for $t=0$, $h\left(\gamma_{g}\right)=h(g)$ then $h(g) \neq 0$ this is equivalent to $g \notin \mathcal{Z}$.

We have the following results.
Theorem 2.3.2 Let $\gamma:[0, T] \longrightarrow \mathbb{R}^{n}$ be a length minimizer of $\tilde{d}$ with control functions $u_{1}(t)$, $u_{2}(t), u_{3}(t)$ with $u_{3}(t) \neq 0$ a.e, and $h\left(\gamma_{g}\right) \neq 0$. If $|h(\gamma)| \leq\left|h\left(\gamma_{g}\right)\right|$ then $\widetilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right) \leq$ $\widetilde{d}(\gamma(0), \gamma(T))$.

Theorem 2.3.3 With the same conditions of the above. If $\frac{\frac{\partial}{\partial g_{i}}(h(g)+f(g, \gamma))}{h\left(\gamma_{g}\right)}>0$ then $\widetilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right)<\widetilde{d}(\gamma(0), \gamma(T))$.

Finally, assuming that for some pairs $\left(q, q^{\prime}\right)$ of points translated from the singular locus the distance $\widetilde{d}$ is decreasing and by considering the order of approximation by $\widetilde{d}$, we prove that the difference $\left|d\left(q, q^{\prime}\right)-\widetilde{d}\left(q, q^{\prime}\right)\right|$ is strictly smaller than $\left|d\left(q, q^{\prime}\right)-\widehat{d}\left(q, q^{\prime}\right)\right|$. We can find this proof in Chapter 4, Section 3.4, where this issue is adressed in the $n$-dimensional generic case.

### 2.2.3. Geodesics

In this section the Hamiltonian for the normal flow defined by the solvable approximation in the 3D generic case is given. We compute the geodesic with initial condition $x(0)=y(0)=$ $z(0)=0$ and covector $\lambda=(p, q, r) \in T^{*} \mathbb{R}^{3}$ with $p(0)=\cos (\theta), q(0)=\sin (\theta), r(0)=r$.

From the above sections, the solvable approximation is defined by

$$
\begin{equation*}
\widehat{X}_{1}=X_{1}, \quad \widehat{X}_{2}=\frac{\partial}{\partial y}+x \frac{\partial}{\partial z}, \quad \widetilde{X}_{3}=\left(a z+b x^{2}+c y^{2}\right) \frac{\partial}{\partial z} \tag{3.2}
\end{equation*}
$$

From (3.2), the Hamiltonian for the normal flow is given by

$$
\begin{aligned}
& H(\lambda)=\frac{1}{2}\left(\left\langle\lambda, \widehat{X}_{1}(x, y, z)\right\rangle^{2}+\left\langle\lambda, \widehat{X}_{2}(x, y, z)\right\rangle^{2}+\left\langle\lambda, \widetilde{X}_{3}(x, y, z)\right\rangle^{2}\right) \\
& H(\lambda)=\frac{1}{2}\left(p^{2}+(q+r x)^{2}+r^{2}\left(a z+b x^{2}+c y^{2}\right)^{2}\right)
\end{aligned}
$$

where $\lambda=(p, q, r) \in T^{*} \mathbb{R}^{3}$. Hence

$$
\begin{array}{ll}
\dot{x}(t)=p & \dot{p}(t)=-(q+r x) r-2 b x r^{2}\left(a z+b x^{2}+c y^{2}\right) \\
\dot{y}(t)=q+r x & \dot{q}(t)=-2 c y r^{2}\left(a z+b x^{2}+c y^{2}\right) \\
\dot{z}(t)=(q+r x) x+r\left(a z+b x^{2}+c y^{2}\right)^{2} & \dot{r}(t)=-a r^{2}\left(a z+b x^{2}+c y^{2}\right)
\end{array}
$$

are the associated Hamiltonian equations to the solvable approximation.

The geodesic with initial condition $x(0)=y(0)=z(0)=0$ and $p(0)=\cos (\theta), q(0)=$ $\sin (\theta)$ and $r(0)=r=0$ is given by

$$
\begin{align*}
& x(t)=t \cos (\theta) \\
& y(t)=t \sin (\theta)  \tag{3.3}\\
& z(t)=\frac{1}{4} t^{2} \sin (2 \theta)
\end{align*}
$$

because $p(t)=\cos (\theta)$ and $q(t)=\sin (\theta)$, this is to say $p$ and $q$ are constants.
Notice that the above geodesic for $\widetilde{d}$ is the same as the geodesic for $\widehat{d}$. The above implies that this geodesic is optimal for any time and has no conjugate time (see Theorem 5.1 and 5.2 in [Boscain et al., 2015]). We can see some geodesics in Figure 3.1 when $r=0$.


Figure 3.1: Geodesics for $\theta \in\left\{0, \frac{\pi}{3}, \frac{2 \pi}{3}, \pi, \frac{4 \pi}{6}, \frac{5 \pi}{6}\right\}$ when $r=0$.
Due to the complexity of the Hamiltonian system of equations, we compute the geodesics considering $a=c=0$ and $b=1$. Thus the Hamiltonian is

$$
H(\lambda)=\frac{1}{2}\left(p^{2}+(q+r x)^{2}+r^{2} x^{4}\right),
$$

hence

$$
\begin{align*}
\dot{x}(t) & =p & & \dot{p}(t)=-(q+r x) r-2 r^{2} x^{3} \\
\dot{y}(t) & =q+r x & & \dot{q}(t)=0  \tag{3.4}\\
\dot{z}(t) & =x q+r x^{2}+r x^{4} & & \dot{r}(t)=0
\end{align*}
$$

Considering the initial condition $x(0)=0$ then $p(0)=\cos (\theta), q(0)=\sin (\theta)$ and $r(0)=r$. If $r=0$ then the solution to the differential systems (3.4) is given by (3.3). If $r(0)=r \neq 0$, since $\dot{x}(t)=p$, we get

$$
\ddot{x}+r^{2} x+2 r^{2} x^{3}=-r q .
$$

Since $q(0)=\sin (\theta)$ and $\dot{q}=0$, then $q=\sin (\theta)$. Hence $\ddot{x}+r^{2} x+2 r^{2} x^{3}=-r \sin (\theta)$. The above equation is equivalent to

$$
\begin{equation*}
\ddot{x}+r^{2} x+2 r^{2} x^{3}=-r \sin (\theta) \operatorname{cn}\left(0, k^{2}\right), \tag{3.5}
\end{equation*}
$$

where $\mathrm{cn}\left(0 \cdot t, k^{2}\right)$ is the Jacobian elliptic function that has a period in $0 \cdot t$ equal to $4 K\left(k^{2}\right)$ and $K\left(k^{2}\right)$ is the complete elliptic integral of the first kind for the modulus $k$ (see more in [Byrd and Friedman, 1954]). This equivalence is due to the fact that $\mathrm{cn}\left(0, k^{2}\right)=1$.

In [Zúniiga, 2006] a general solution to equation (3.5) is given, and hence by using the initials conditions $x(0)=y(0)=z(0)=0$, and straightforward computation (see details in [Jouan and Manríquez, 2022]), we obain

$$
\begin{aligned}
x(t) & =a_{1}\left(\mathrm{cn}\left(w_{1} t+\phi, k_{1}^{2}\right)-\mathrm{cn}\left(\phi, k_{1}^{2}\right)\right) \\
y(t) & =\left(\sin (\theta)-r a_{1} \mathrm{cn}\left(\phi, k_{1}^{2}\right)\right) t+\frac{r a_{1}}{k_{1}^{2} w_{1}} \cdot y_{1} \\
z(t) & =-\left(r a_{1}^{4} \mathrm{cn}\left(\phi, k_{1}^{2}\right)^{4}+r a_{1}^{2} \mathrm{cn}\left(\phi, k_{1}^{2}\right)^{2}+\sin (\theta)\right) a_{1} \mathrm{cn}\left(\phi, k_{1}^{2}\right) t+\frac{r a_{1}^{4}}{3 k_{1}^{8} w_{1}} z_{1}(t) \\
& +\frac{4 r a_{1}^{4} \mathrm{cn}\left(\phi, k_{1}^{2}\right)}{2 k_{1}^{6} w_{1}} z_{2}(t)+\frac{6 r a_{1}^{4} \mathrm{cn}\left(\phi, k_{1}^{2}\right)^{2}+r a_{1}^{2}}{k_{1}^{4} w_{1}} z_{3}(t) \\
& +\frac{4 r a_{1}^{4} \mathrm{cn}\left(\phi, k_{1}^{2}\right)^{3}+2 r a_{1}^{2} \mathrm{cn}\left(\phi, k_{1}^{2}\right)+\sin (\theta)}{k_{1}^{2} w_{1}} z_{4}(t),
\end{aligned}
$$

where $k_{1}^{\prime 2}+k_{1}^{2}=1, E(\cdot)$ is the incomplete elliptic integral of the second kind and

$$
\begin{aligned}
y_{1}= & \left(\arccos \left(\operatorname{dn}\left(w_{1} t+\phi, k_{1}^{2}\right)\right)-\arccos \left(\operatorname{dn}\left(\phi, k_{1}^{2}\right)\right)\right) \\
z_{1}(t)= & \left(2-3 k_{1}^{4}\right) k_{1}^{4} w_{1} t+2\left(2 k_{1}^{4}-1\right)\left(E\left(w_{1} t+\phi\right)-E(\phi)\right) \\
& +k_{1}^{4}\left(\operatorname{sn}\left(w_{1} t+\phi, k_{1}^{2}\right) \operatorname{cn}\left(w_{1} t+\phi, k_{1}^{2}\right) \operatorname{dn}\left(w_{1} t+\phi, k_{1}^{2}\right)\right. \\
& \left.-\operatorname{sn}\left(\phi, k_{1}^{2}\right) \operatorname{cn}\left(\phi, k_{1}^{2}\right) \operatorname{dn}\left(\phi, k_{1}^{2}\right)\right) \\
z_{2}(t)= & \left(2 k_{1}^{4}-1\right)\left(\arcsin \left(k_{1}^{2} \operatorname{sn}\left(w_{1} t+\phi, k_{1}^{2}\right)\right)-\arcsin \left(k_{1}^{2} \operatorname{sn}\left(\phi, k_{1}^{2}\right)\right)\right) \\
& +k_{1}^{2}\left(\operatorname{sn}\left(w_{1} t+\phi, k_{1}^{2}\right) \operatorname{dn}\left(w_{1} t+\phi, k_{1}^{2}\right)-\operatorname{sn}\left(\phi, k_{1}^{2}\right) \operatorname{dn}\left(w_{1} t+\phi, k_{1}^{2}\right)\right) \\
z_{3}(t)= & E\left(w_{1} t+\phi\right)-E(\phi)-k_{1}^{\prime 4} w_{1} t \\
z_{4}(t)= & \left.\arccos \left(\operatorname{dn}\left(w_{1} t+\phi, k_{1}^{2}\right)\right)-\arccos \left(\operatorname{dn}\left(\phi, k_{1}^{2}\right)\right)\right) .
\end{aligned}
$$



Figure 3.2: Ball in 3-D generic case.

## CHAPTER 4

## General Nilpotent and Solvable Approximations of Almost-Riemannian Structures

The content of this chapter is the result of two articles in collaboration with Philippe Jouan, and Yacine Chitour, one of which was submitted to the journal Discrete and Continuous Dynamical Systems ([Manríquez et al., 2022]) and the other is in progress. ([Manríquez, 2022]).

### 1.4. Nilpotent and Solvable approximations are linear

This section shows that the nilpotent or the solvable approximation of an almost-Riemannian structure at a singular point is always a linear almost-Riemannian structure on a Lie group or a homogeneous space.

The following definition is an extension of the model shown in Chapter 3.
Consider a set $X_{1}, X_{2}, \ldots, X_{n}$ of vector fields that generates $\Delta$ around $p$. In privileged coordinates each $X_{j}$ can be decomposed into

$$
X_{j}=X_{j}^{(-1)}+X_{j}^{(0)}+X_{j}^{(1)}+\cdots+X_{j}^{(s)}+\ldots
$$

where $X_{j}^{(s)}$ is the component of $X_{j}$ of homogeneous order $s$.
It may happen that some of the vector fields $\widehat{X}_{j}$ globally vanish. In that case they can be replaced by $\widetilde{X}_{j}=X_{j}^{(0)}$. Let us assume that only $m$ elements of $\widehat{X}_{1}, \widehat{X}_{2}, \ldots, \widehat{X}_{n}$ are linearly independent (as vector fields). As explained below, we can assume without lost of generality that these vector fields are $\widehat{X}_{1}, \widehat{X}_{2}, \ldots, \widehat{X}_{m}$ and that $\widehat{X}_{m+1}, \ldots, \widehat{X}_{n}$ vanish. The set of vector fields

$$
\widehat{X}_{1}, \ldots, \widehat{X}_{k}, \ldots, \widehat{X}_{m}, \widetilde{X}_{m+1}, \ldots, \widetilde{X}_{n},
$$

is called the solvable approximation of $X_{1}, X_{2}, \ldots, X_{n}$.

Though this section deals with local questions, around a point $p$ belonging to the singular locus, it will be more convenient to assume the ARS defined by a bundle $E$ and a morphism $f$ from $E$ to $T M$ as in Definition 1.2.1.

Firstly it is necessary to show that it is always possible to define the ARS locally, around the point $p=0$ in local privileged coordinates, by a set of $n$ orthonormal vector fields $X_{1}, \ldots, X_{n}$ such that the solvable approximation

$$
\widehat{X}_{1}, \ldots, \widehat{X}_{k}, \widehat{X}_{k+1}, \ldots, \widehat{X}_{m}, \widetilde{X}_{m+1}, \ldots, \widetilde{X}_{n},
$$

satisfies

- $\widehat{X}_{i}(0) \neq 0$ for $i=1, \ldots, k$;
- $\widehat{X}_{i} \neq 0$ but $\widehat{X}_{i}(0)=0$ for $i=k+1, \ldots, m$;
- $\widehat{X}_{i}=0$ for $i=m+1, \ldots, n$.

Let $K_{p}$ be the kernel of the restriction of $f$ to $E_{p}$, and let $V_{p}$ be an orthogonal complement to $K_{p}$ in $E_{p}$, that is $K_{p} \perp V_{p}$ and $K_{p} \oplus V_{p}=E_{p}$.

Let $e_{1}, \ldots, e_{n}$ be a set of $n$ sections of $E$, orthonormal in a neighborhood of $p$, such that $e_{j}(p) \in V_{p}$ for $j=1, \ldots, k$ and $e_{j}(p) \in K_{p}$ for $j=k+1, \ldots, n$.

The vector fields $X_{j}=f_{*}\left(e_{j}\right), j=1, \ldots, n$ define the ARS around $p$. Let $\left(x_{1}, \ldots, x_{n}\right)$ be a set of privileged coordinates and $\widehat{X}_{1}, \ldots, \widehat{X}_{n}$ be the related nilpotent approximation. Let $\mathcal{L}$ be the submodule of $\Gamma(E)$ generated by $e_{j}$ for $j=k+1, \ldots, n$. Consider now the mapping $e \in \mathcal{L} \longmapsto \widehat{f_{*} e}$.

Its rank is $m-k$ with $k \leq m \leq n$ and we can assume without lost of generality that $e_{m+1}, \ldots, e_{n}$ belong to the kernel of that linear map, and that $e_{k+1}, \ldots, e_{m}$ are orthogonal to that kernel.

The vector fields $X_{1}, \ldots, X_{n}$ satisfy the above conditions.
It may happen that the vector fields $\widetilde{X}_{m+1}, \ldots, \widetilde{X}_{n}$ fail to be linearly independent. In that case, neither the nilpotent approximation nor the solvable one defines an almost-Riemannian structure.

For that reason, we will always assume in what follows that the vector fields

$$
\widehat{X}_{1}, \ldots, \widehat{X}_{k}, \widehat{X}_{k+1}, \ldots, \widehat{X}_{m}, \widetilde{X}_{m+1}, \ldots, \widetilde{X}_{n}
$$

are linearly independent. Denote by $\mathcal{Z}_{a}$ ( $a$ for approximation) the set of points where their rank is not full. It is not empty because at least one vector field vanishes at $p=0$. On the other hand, the approximating vector fields are polynomial, and the interior of $\mathcal{Z}_{a}$ is empty. Consequently, $\mathcal{Z}_{a}$ is a proper subset of $\mathbb{R}^{n}$ with empty interior, and the set of approximating vector fields defines an ARS.

Remark 1.4.1 In the case where some of the $\widetilde{X}_{j}$ vanish or are linearly dependent, it seems difficult, if not impossible, to go one step further by considering homogeneous approximations of nonholonomic order $s>0$ because, as explained in the sequel, two important properties
could be lost. First, a homogeneous vector field of degree $s>0$ need not be complete. Second, the Lie algebra generated by the approximating vector fields would not generally be finite-dimensional. These two drawbacks are related, see [Palais, 1957] or [Jouan, 2010].

### 1.1.4 . The generated Lie algebra

In view of the next sections, it is important to notice that all involved vector fields are complete because of their triangular form. This fact is well-known for the $\widehat{X}_{i}$ (see [Bellaïche, 1996] or [Jean, 2014]).

Proposition 1.4.1 The vector fields $\widehat{X}_{j}$ and $\widetilde{X}_{j}$ defined above are complete.
This important property does not hold for homogeneous vector fields of positive degree, for example the first coordinate such a vector field could be $x_{1}^{2} \frac{\partial}{\partial x_{1}}$, but $\dot{x}_{1}=x_{1}^{2}$ is not complete.

The second feature we will use in the following subsections is the finiteness of the generated Lie algebra;

Proposition 1.4.2 The Lie algebra $\mathcal{L}$ generated by

$$
\widehat{X}_{1}, \ldots, \widehat{X}_{k}, \widehat{X}_{k+1}, \ldots, \widehat{X}_{m}, \widetilde{X}_{m+1}, \ldots, \widetilde{X}_{n}
$$

is finite dimensional.

### 2.1.4 . The Nilpotent case

It is the case where the vector fields $\widehat{X}_{1}, \ldots, \widehat{X}_{n}$ are linearly independent and the vectors $\widehat{X}_{1}(0), \ldots, \widehat{X}_{k}(0)$ are independent in $\mathbb{R}^{n}$. In particular no vector field $\widehat{X}_{i}$ vanishes, and $m=n$.

For $j=k+1, \ldots, n$ let $D_{j}$ stand for $\operatorname{ad}\left(\widehat{X_{j}}\right)$ and for any multi-index $J=\left(j_{1}, \ldots, j_{s}\right)$ let $D_{J}=D_{j_{s}} \circ \cdots \circ D_{j_{1}}$ (here $k+1 \leq j_{u} \leq n$ and $s \geq 0$ ). Let

$$
\mathcal{D}=\operatorname{Span}\left\{D_{J}\left(\widehat{X}_{i}\right) / i=1, \ldots, k ; J \text { as above }\right\} .
$$

Lemma 1.4.1 The Lie algebra $\mathfrak{g}$ generated by $\mathcal{D}$ is $D_{j}$-invariant for $j=k+1, \ldots, n$.
Let $\mathcal{L}$ stand for the Lie algebra generated by $\widehat{X}_{1}, \ldots, \widehat{X}_{n}$. It is a well-known fact that this Lie algebra is nilpotent and finite-dimensional (see [Bellaïche, 1996] or [Jean, 2014]).

Theorem 1.4.1 1. The ideal generated in $\mathcal{L}$ by $\widehat{X}_{1}, \ldots, \widehat{X}_{k}$ is $\mathfrak{g}$. It is a nilpotent Lie algebra.
2. The vector fields $\widehat{X}_{k+1}, \ldots, \widehat{X}_{n}$ do not belong to $\mathfrak{g}$ and act on $\mathfrak{g}$ as derivations.
3. The rank at $p=0$ of the elements of $\mathfrak{g}$ is full.

After this analysis at the algebra level, we can turn our attention to the Lie group level.
Let $G$ be the simply connected Lie group whose Lie algebra is $\mathfrak{g}$. Since $\mathfrak{g}$ is nilpotent the underlying manifold of $G$ is $\mathbb{R}^{N}, N=\operatorname{dim}(\mathfrak{g})$. The first task is to show that $\mathbb{R}^{n}$ is a homogeneous space of $G$. This is mainly because homogeneous vector fields generate $\mathfrak{g}$.

Lemma 1.4.2 The set $\mathbb{R}^{n}$ is a homogeneous space of $G$. More accurately, if $H$ stands for the connected subgroup of $G$ whose Lie algebra is the set of elements of $\mathfrak{g}$ that vanish at 0 , then $\mathbb{R}^{n}$ is diffeomorphic to the quotient $G / H$.

To complete the construction we associate to the derivation $D_{j}=\operatorname{ad}\left(\widehat{X}_{j}\right)$ of $\mathfrak{g}$ a linear vector field $\mathcal{X}_{j}$ on $G$ for $j>k$ ( $\mathcal{X}_{j}$ does exist because $G$ is simply connected). It is clear that the projection of $\mathcal{X}_{j}$ on $\mathbb{R}^{n}$ is $\widehat{X}_{j}$ (see [Jouan, 2010] for details). Finally the vector fields $\widehat{X}_{1}, \ldots, \widehat{X}_{k}$ are invariant, and $\widehat{X}_{k+1}, \ldots, \widehat{X}_{n}$ are linear vector fields on the homogeneous space $\mathbb{R}^{n}=G / H$.

We can state:
Theorem 1.4.2 The space $\mathbb{R}^{n}$ is a homogeneous space of the nilpotent Lie group $G$ whose Lie algebra is $\mathfrak{g}$.

The vector fields $\widehat{X}_{1}, \ldots, \widehat{X}_{k}$ are projections of invariant vector fields of $G$ and $\widehat{X}_{k+1}, \ldots, \widehat{X}_{n}$ are projections of linear vector fields of $G$.

Consequently the set $\widehat{X}_{1}, \ldots, \widehat{X}_{n}$ defines a linear ARS on the homogeneous space $\mathbb{R}^{n}$.

### 3.1.4 . The non-nilpotent case

We set $D_{j}=\operatorname{ad}\left(\widehat{X_{j}}\right)$ for $j=k+1, \ldots, m$ and $D_{j}=\operatorname{ad}\left(\widetilde{X_{j}}\right)$ for $j=m+1, \ldots, n$. As well as in the nilpotent case we set $D_{J}=D_{j_{s}} \circ \cdots \circ D_{j_{1}}$ for any multi-index $J=\left(j_{1}, \ldots, j_{s}\right)$ where $s \geq 0$ and $k+1 \leq j_{u} \leq n$, and

$$
\mathcal{D}=\operatorname{Span}\left\{D_{J}\left(X_{i}\right) / i=1, \ldots k ; J \text { as above }\right\} .
$$

The Lie algebra $\mathfrak{g}$ generated by $\mathcal{D}$ is again $D_{j}$-invariant for $j=k+1, \ldots, n$, which shows that $\mathfrak{g}$ is the ideal generated in $\mathcal{L}$ by $\widehat{X}_{1}, \ldots, \widehat{X}_{k}$.

Theorem 1.4.3 1. The ideal generated in $\mathcal{L}$ by $\widehat{X}_{1}, \ldots, \widehat{X}_{k}$ is $\mathfrak{g}$. It is a nilpotent Lie algebra.
2. The vector fields $\widetilde{X}_{m+1}, \ldots, \widetilde{X}_{n}$ do not belong to $\mathfrak{g}$ and act on $\mathfrak{g}$ as derivations.
3. The vector fields $\widehat{X}_{j}$, with $k+1 \leq j \leq m$ that do not belong to $\mathfrak{g}$ act on $\mathfrak{g}$ as derivations.
4. The rank at $p=0$ of the elements of $\mathfrak{g}$ is full.

Opposite to the nilpotent case we cannot assert that the vector fields $\widehat{X}_{k+1}, \ldots, \widehat{X}_{m}$ do not belong to $\mathfrak{g}$. Because of this phenomenon, illustrated by Example 3 in Section 4.1.4, we are lead to introduce one more index. Up to a reordering we can assume that $\widehat{X}_{k+1}, \ldots, \widehat{X}_{l}$ belong to $\mathfrak{g}$ and that $\widehat{X}_{l+1}, \ldots, \widehat{X}_{m}$ do not belong to $\mathfrak{g}$, where $k+1 \leq l \leq m$.

Theorem 1.4.4 The space $\mathbb{R}^{n}$ is a homogeneous space of the nilpotent Lie group $G$ whose Lie algebra is $\mathfrak{g}$.

The vector fields $\widehat{X}_{1}, \ldots, \widehat{X}_{l}$ are projections of invariant vector fields of $G$. The vector fields $\widehat{X}_{l+1}, \ldots, \widehat{X}_{m}$ are projections of linear or affine vector fields of $G$ and $\widetilde{X}_{m+1}, \ldots, \widetilde{X}_{n}$ are projections of linear ones.

Consequently the set of vector fields $\widehat{X}_{1}, \ldots, \widehat{X}_{m}, \widetilde{X}_{m+1}, \ldots, \widetilde{X}_{n}$ defines a linear ARS on the homogeneous space $\mathbb{R}^{n}$.

Remark 1.4.2 As shown by Example 4 in Section 4.1.4, the Lie algebra $\mathcal{L}$ generated by the vecotr fields $\widehat{X}_{1}, \ldots, \widehat{X}_{m}, \widetilde{X}_{m+1}, \ldots, \widetilde{X}_{n}$ need not be solvable when $m \leq n-2$. However, it is solvable if $m=n-1$ (see Propositon 1.3.1), and it will be proven in the next section that generically $m=n-1$ or $m=n$.

It is why we call solvable the approximations of the previous kind.

### 4.1.4 . Examples

The first one is a standard example of a solvable approximation on the group Heisenberg. Example 2 shows that the elements of the Lie algebra $\mathcal{L}$ of nonholonomic order smaller than -1 are not necessarily in $\mathfrak{g}$. We exhibit a vector field of the distribution that vanishes at $p$ but belongs to the ideal $\mathfrak{g}$ in Example 3 (which implies that the rank of the nilpotent approximation is not full).

To finish, the Lie algebra $\mathcal{L}$ of Example 4 is not solvable, it contains a semi-simple subalgebra. Recall that this is not generic.

In these four examples, the vector fields are equal to their nilpotent or solvable approximations at 0 . It is of course possible to add terms of higher nonholonomic order without modifying the conclusions.

## Example 1

Consider in $\mathbb{R}^{3}$ the almost-Riemannian structure defined by the vector fields:

$$
X_{1}=\partial_{x}, \quad X_{2}=\partial_{y}+x \partial_{z}, \quad X_{3}=x \partial_{y}+\frac{1}{2} x^{2} \partial_{z}
$$

At $p=(0,0,0)$ the coordinates $(x, y, z)$ are privileged with weights $(1,1,2)$, the vector fields $X_{1}$ and $X_{2}$ are homogeneous of order -1 and $X_{3}$ is homogeneous of order 0 , so that $\widehat{X}_{1}=X_{1}$, $\widehat{X}_{2}=X_{2}, \widehat{X}_{3}=0$ and $\widetilde{X}_{3}=X_{3}$. The nilpotent approximation at $p$ is not an almostRiemannian structure, it is the constant rank 2 sub-Riemannian structure defined by $X_{1}$ and $X_{2}$.

The Lie algebra generated by $\widehat{X}_{1}, \widehat{X}_{2}, \widetilde{X}_{3}$ is $\mathcal{L}=\operatorname{Span}\left\{\widehat{X}_{1}, \widehat{X}_{2}, \widetilde{X}_{3}, \partial_{z}\right\}$, the ideal $\mathfrak{g}=$ Span $\left\{\widehat{X}_{1}, \widehat{X}_{2}, \partial_{z}\right\}$ is here the Heisenberg Lie algebra, and $\widetilde{X}_{3}$ is a linear vector field on $\mathfrak{g}$. Finally $\widehat{X}_{1}, \widehat{X}_{2}, \widehat{X}_{3}$ is a linear ARS on the Heisenberg group.

## Example 2

The almost-Riemannian structure is here defined in $\mathbb{R}^{3}$ by:

$$
X_{1}=\partial_{x}, \quad X_{2}=x \partial_{y}, \quad X_{3}=y^{2} \partial_{z} .
$$

The Lie algebra $\mathcal{L}$ contains $X_{1}, X_{2}, X_{3}$ and

$$
\begin{aligned}
& X_{4}=\left[X_{1}, X_{2}\right]=\partial_{y}, X_{5}=\frac{1}{2}\left[X_{4}, X_{3}\right]=y \partial_{z}, \\
& X_{6}=\left[X_{2}, X_{5}\right]=x \partial_{z}, X_{7}=\left[X_{1}, X_{6}\right]=\left[X_{4}, X_{5}\right]=\partial_{z} \\
& X_{8}=\frac{1}{2}\left[X_{2}, X_{3}\right]=x y \partial_{z}, X_{9}=\left[X_{2}, X_{8}\right]=x^{2} \partial_{z} .
\end{aligned}
$$

Therefore:

$$
\begin{gathered}
\Delta^{1}=\Delta=\left\{X_{1}, X_{2}, X_{3}\right\}, \quad \Delta^{2}=\Delta^{1}+\left\{X_{4}, X_{8}\right\} \\
\Delta^{3}=\Delta^{2}+\left\{X_{5}, X_{9}\right\}, \quad \Delta^{4}=\Delta^{3}+\left\{X_{6}\right\}, \quad \Delta^{5}=\Delta^{4}+\left\{X_{7}\right\} .
\end{gathered}
$$

The canonical coordinates are privileged with weights $(1,2,5)$ and the vector fields $X_{1}, X_{2}$, and $X_{3}$ are homogeneous of order -1 hence equal to their nilpotent approximations.

The algebra $\mathfrak{g}$ is here the ideal of $\mathcal{L}$ generated by $X_{1}$ that is

$$
\mathfrak{g}=\operatorname{Span}\left\{X_{1}, X_{4}, X_{5}, X_{6}, X_{7}\right\} .
$$

The vector fields $X_{2}$ and $X_{3}$ are linear, as well as $X_{8}=\frac{1}{2}\left[X_{2}, X_{3}\right]$ and $X_{9}=\left[X_{2}, X_{8}\right]$.
The orders of $X_{8}$ and $X_{9}$ are respectively -2 and -3 , which shows that the vector fields of order smaller than -1 are not necessarily in $\mathfrak{g}$.

Notice that the singular locus is here $\mathcal{Z}=\{x y=0\}$.

## Example 3

Consider in $\mathbb{R}^{4}$ the almost-Riemannian structure defined by the vector fields:

$$
X_{1}=\partial_{x}, \quad X_{2}=\partial_{y}+x \partial_{z}, \quad X_{3}=y \partial_{w}, \quad X_{4}=\frac{1}{2} x^{2} \partial_{z}+\frac{1}{2} y^{2} \partial_{w}
$$

Since $\left[X_{1}, X_{2}\right]=\partial_{z}$ and $\left[X_{2}, X_{3}\right]=\partial_{w}$, the coordinates $(x, y, z, w)$ are privileged with weights $(1,1,2,2)$ at $p=(0,0,0,0)$. The vector fields $X_{1}, X_{2}$ are homogeneous of order -1 and independent at 0 , and the vector field $X_{3}$ is homogeneous of order -1 but vanishes at 0 . The last field $X_{4}$ is homogeneous of order 0 . Consequently the first three are equal to their nilpotent approximation and $X_{4}=\widetilde{X}_{4}$. According to the notations of Section 1.4 we have $k=2$ and $m=3$.

The Lie algebra $\mathcal{L}$ is spanned by $X_{1}, X_{2}, X_{3}, X_{4}$ and

$$
X_{5}=\left[X_{1}, X_{2}\right]=\partial_{z}, X_{6}=\left[X_{2}, X_{3}\right]=\partial_{w}, X_{7}=\left[X_{1}, X_{4}\right]=x \partial_{z} .
$$

Despite the fact that $\widehat{X}_{3}(0)=X_{3}(0)=0$ we cannot assert as in the nilpotent case that $X_{3}$ does not belong to $\mathfrak{g}$ (see subsection 3.1.4 after Theorem 1.4.3). Indeed the ideal generated in $\mathcal{L}$ by $X_{1}$ and $X_{2}$ is here

$$
\mathfrak{g}=\operatorname{Span}\left\{X_{1}, X_{2}, X_{3}, X_{5}, X_{6}, X_{7}\right\}
$$

because $X_{3}=\left[X_{2}, X_{4}\right]$.
As explained in subsection 3.1.4 this may happen when $k<m<n$.
Notice that the determinant of $X_{1}, X_{2}, X_{3}, X_{4}$ is $-\frac{1}{2} x^{2} y$. Therefore the singular locus is $\mathcal{Z}=\{x y=0\}$ which shows that the structure is not generic.

In the general case, the Lie algebra $\mathcal{L}$ need not be solvable. Indeed it is a subalgebra of the semi-direct product of $\mathfrak{g}$ by its algebra of derivations. But the algebra of derivations of a nilpotent Lie algebra is not solvable in general. For instance, the derivations of the Heisenberg algebra is the set of endomorphisms, the matrix of which writes in the canonical basis:

$$
D=\left(\begin{array}{ccc}
a & b & 0 \\
c & d & 0 \\
e & f & a+d
\end{array}\right)
$$

The subalgebra of such derivations that moreover satisfy $e=f=a+d=0$ is equal to $\mathfrak{s l}_{2}$ hence semisimple.

Example 4 illustrates that possibility.

## Example 4

Consider in $\mathbb{R}^{5}$, with coordinates $(x, y, z, w, t)$ :

$$
\begin{aligned}
& X_{1}=\partial_{x}, \quad X_{2}=\partial_{y}+x \partial_{w}+z \partial_{t}, \quad X_{3}=\partial_{z}+x \partial_{t}, \\
& X_{4}=x \partial_{y}+\frac{1}{2} x^{2} \partial_{w}, \quad X_{5}=y \partial_{x}+\frac{1}{2} y^{2} \partial_{w}
\end{aligned}
$$

The Lie algebra $\mathcal{L}$ contains also

$$
\begin{aligned}
& X_{6}=\left[X_{1}, X_{2}\right]=\partial_{w}, \quad X_{7}=\left[X_{1}, X_{3}\right]=\left[X_{3}, X_{2}\right]=\partial_{t}, \\
& X_{8}=\left[X_{3}, X_{5}\right]=-y \partial_{t}, \quad X_{9}=\left[X_{8}, X_{4}\right]=x \partial_{t}, \\
& X_{10}=\left[X_{4}, X_{5}\right]=x \partial_{x}-y \partial_{y} .
\end{aligned}
$$

The coordinates $(x, y, z, w, t)$ are privileged with weights $(1,1,1,2,2)$ at the origin. At this point the vector fields $X_{1}, X_{2}, X_{3}$ (resp. $X_{4}, X_{5}$ ) are homogeneous of order -1 (resp. 0), hence equal to their nilpotent approximations (resp. $\widetilde{X}_{4}=X_{4}$ and $\widetilde{X}_{5}=X_{5}$ ).

Since $X_{10}=\left[X_{4}, X_{5}\right],\left[X_{10}, X_{4}\right]=2 X_{4}$ and $\left[X_{10}, X_{5}\right]=-2 X_{5}$ the vector fields $X_{4}$, $X_{5}$ and $X_{10}$, that do not belong to $\mathfrak{g}$, generate a semi-simple Lie algebra isomorphic to $\mathfrak{s l}_{2}$. Consequently the algebra $\mathcal{L}$ is not solvable.

The singular locus is here $\mathcal{Z}=\{x y z=0\}$, and again the structure is not generic.

### 2.4. Genericity

In this section, we investigate the generic properties of almost-Riemannian structures in all dimensions, and we identify the generic nilpotent and solvable approximations.

The examples of Section 4.1.4 show that many different, complicated structures may arise, and the aim of this section is to determine the generic ones.

In what follows, we will say that a property of almost-Riemannian distributions (resp. structures) on a manifold $M$ is generic if for any rank $n$ vector bundle $E$ (resp. Euclidean vector bundle ( $E,\langle. .\rangle$,$) ) over M$ the set of smooth morphisms of vector bundles from $E$ to $T M$ for
which this property is satisfied is open and dense in the $\mathcal{C}^{2}$ Whitney topology.
Let $U$ be an open subset of $M$ on which $E$ and $T M$ are trivializable, and let $\Pi$ be the projection from $E$ onto $M$. Then the restriction to $\Pi^{-1}(U)$ of a vector bundle morphism $f$ is equivalent to a smooth mapping $X$ from $U$ to the set $\mathcal{M}_{n}(\mathbb{R})$ of $n \times n$ square matrices.

Alternately $X$ can be viewed as a mapping $\left(X_{1}, X_{2}, \ldots, X_{n}\right)$ from $U$ to the set $\Gamma(U)^{n}$ of $n$ vector fields on $U$.

It is not helpful to assume that $f$ satisfies the properties of almost-Riemannian distributions because these conditions will turn out to be generic.

The first two theorems deal with distributions only and do not require neither metric, nor normal forms.

In what follows, we denote by $\mathcal{M}_{n \times m}(\mathbb{R})$ the set of real $n \times m$ matrices (simply $\mathcal{M}_{n}(\mathbb{R})$ ) if $m=n$ ) and by $L^{r}$ the set of elements of $\mathcal{M}_{n \times m}(\mathbb{R})$ of corank $r$. It is a submanifold of $\mathcal{M}_{n \times m}(\mathbb{R})$ of codimension $(n-q+r)(m-q+r)$ where $q=\min \{n, m\}$ (see [Golubitsky and Guillemin, 1986]).

Recall from Chapter 2, Section 1.2 that $f$ being given, $\Delta$ stands for the submodule of $\Gamma(M)$ it defines.

### 1.2.4 . Generic distribution

Theorem 2.4.1 The following properties are generic:

1. Let $R$ be the largest integer such that $R^{2} \leq n$. For $1 \leq r \leq R$ let $\mathcal{Z}_{r}$ be the set of points where the rank of $f_{p}$, or locally the rank of $\left\{X_{1}, X_{2}, \ldots, X_{n}\right\}$, is $n-r$. Each $\mathcal{Z}_{r}$ is a codimension $r^{2}$ submanifold and the singular locus $\mathcal{Z}$ is the union of these disjoint submanifolds.
2. The submanifold $\mathcal{Z}_{r+1}$ is included in the closure $\overline{\mathcal{Z}}_{r}$ of $\mathcal{Z}_{r}$ for $r=1, \ldots, R-1$.
3. For any local representation $X$ of the distribution the mapping $x \longmapsto \operatorname{det}(X(x))$ is a submersion at all points $x \in \mathcal{Z}_{1}$.
4. For $n \geq 3$ the rank of $\Delta+[\Delta, \Delta]$ is full at all points.

Two subspaces of $T_{p} M$ are attached to a point $p$ belonging to the strate $\mathcal{Z}_{r}$ of the singular locus, namely $\Delta_{p}$, the distribution at $p$, and $T_{p} \mathcal{Z}_{r}$, the tangent subspace to $\mathcal{Z}_{r}$ at $p$. Their dimensions being respectively $n-r$ and $n-r^{2}$, the dimension of $T_{p} \mathcal{Z}_{r}+\Delta_{p}$ is at most equal to $\min \left(n, 2 n-r^{2}-r\right)$. We are interested in the cases where the actual dimension of $T_{p} \mathcal{Z}_{r}+\Delta_{p}$ is less than $\min \left(n, 2 n-r^{2}-r\right)$. For example in $\mathcal{Z}_{1}$ this means that $\Delta_{p}=T_{p} \mathcal{Z}_{r}$ (tangency points).

In what follows, we note $s=\min \left(n, 2 n-r^{2}-r\right)-\operatorname{dim}\left(T_{p} \mathcal{Z}_{r}+\Delta_{p}\right)$, and $\lfloor\alpha\rfloor$ stands for the integer part of the real number $\alpha$.

Theorem 2.4.2 The following properties are generic:

- $r=1$. The points $p \in \mathcal{Z}_{1}$ where $T_{p} \mathcal{Z}_{1}=\Delta_{p}$ are isolated in $\mathcal{Z}_{1}$.
- $r \geq 2$. Let $m(n, r)$ be the largest dimension that $T_{p} \mathcal{Z}_{r}+\Delta_{p}$ may reach, that is $m(n, r)=$ $\min \left(n, 2 n-r^{2}-r\right)$, and let $s=m(n, r)-\operatorname{dim}\left(T_{p} \mathcal{Z}_{r}+\Delta_{p}\right)$. Then

1. The set of points $p \in \mathcal{Z}_{r}$ where $s=1$ is a submanifold of $\mathcal{Z}_{r}$ for $n \geq r^{2}+r-$ $\left\lfloor\frac{r-1}{2}\right\rfloor$. It is empty if $n<r^{2}+r-\left\lfloor\frac{r-1}{2}\right\rfloor$.
2. The set of points $p \in \mathcal{Z}_{r}$ where $s \geq 2$ and $s^{2} \leq r$ is a submanifold of $\mathcal{Z}_{r}$ for $r^{2}+r-\left\lfloor\frac{r-s^{2}}{s-1}\right\rfloor \leq n \leq r^{2}+r+\left\lfloor\frac{r-s^{2}}{s-1}\right\rfloor$. It is empty if $n$ is not in this interval.
3. The set of points $p \in \mathcal{Z}_{r}$ where $s \geq 2$ and $s^{2}>r$ is empty.

### 2.2.4 . Normal Form

We consider an almost-Riemannian structure defined by a Euclidean vector bundle ( $E,\langle\rangle$, and a vector bundle morphism $f$. We are interested in local normal forms of orthonormal vector fields defining the structure in a neighborhood of a point $p$ that we can assume to be $p=0$ in local coordinates.

These normal forms turn out to be the key to the next section.
First we follow the lines of [Agrachev et al., 2008] (also used in [Boscain et al., 2013a] and [Boscain et al., 2015]).

Let $W$ be a codimension 1 submanifold transversal to the distribution. We can define a coordinate system $y=\left(x_{2}, \ldots, x_{n}\right)$ in $W$, and choose an orientation transversal to $W$. Let $\gamma_{y}$ be the family of normal geodesics parametrized by arclength, transversal to $W$ at $y$, and positively oriented. The mapping $\left(x_{1}, y\right) \mapsto \gamma_{y}\left(x_{1}\right)$ is a local diffeomorphism and the geodesics $x_{1} \mapsto \gamma_{y}\left(x_{1}\right)$ realize the minimal distance between $W=\left\{x_{1}=0\right\}$ and the surfaces $\left\{x_{1}=c\right\}$ for $c$ small enough. The transversality conditions of the PMP are consequently satisfied along all these surfaces: if $\lambda\left(x_{1}\right)$ is a covector associated to one of these geodesics then the tangent space to $\left\{x_{1}=c\right\}$ at $\gamma_{y}\left(x_{1}\right)$ is $\operatorname{ker}\left(\lambda\left(x_{1}\right)\right)$.

Now let $X_{1}=\partial_{x_{1}}$ be the vector field defined by $X_{1}(q)=\frac{d}{d x_{1}} \gamma_{y}\left(x_{1}\right)$ at the point $q=$ $\gamma_{y}\left(x_{1}\right)$. It is a unitary vector field belonging to $\Delta$. Let $X_{2}, \ldots, X_{n}$ be $n-1$ vector fields such that $\left\{X_{1}, X_{2}, \ldots, X_{n}\right\}$ be an orthonormal frame of $\Delta$. For geodesics the control functions $\left(u_{1}, u_{2}, \ldots, u_{n}\right)$ from the PMP satisfy $u_{j}=\left\langle\lambda, X_{j}\right\rangle$. But here $\left(u_{1}, u_{2}, \ldots, u_{n}\right)=(1,0 \ldots, 0)$ so that $\left\langle\lambda, X_{j}\right\rangle=0$ for $j=2, \ldots, n$ and the vector fields $X_{2}, \ldots, X_{n}$ are tangent to the surfaces $\left\{x_{1}=c\right\}$. Consequently the vector fields have the following form:

$$
X_{1}=\left(\begin{array}{c}
1 \\
0 \\
\cdot \\
\cdot \\
0
\end{array}\right), \quad X_{j}=\left(\begin{array}{c}
0 \\
a_{2, j} \\
\cdot \\
\cdot \\
a_{n, j}
\end{array}\right) \quad \text { for } 1<j<n
$$

for any choice of the coordinates in $W$ and any choice of $X_{2}, \ldots, X_{n}$, under the condition that they provide an orthonormal frame related to the sub-Riemannian metric.

Notice that thanks to the transversality conditions, the vectors fields $X_{2}, \ldots, X_{n}$ are not only orthogonal to $X_{1}$ for the sub-Riemannian metric but also for the canonical inner product of $\mathbb{R}^{n}$ for the chosen coordinates.

Let us assume now that $p=0$ belongs to $\mathcal{Z}_{r}$ with $r \geq 1$ and $r^{2} \leq n$. We want to show that the coordinates and $X_{2}, \ldots, X_{n}$ can be chosen in such a way that they write:

$$
\begin{aligned}
& X_{1}=\left(\begin{array}{c}
1 \\
0 \\
\cdot \\
\cdot \\
\cdot \\
\cdot \\
0
\end{array}\right) X_{2}=\left(\begin{array}{c}
0 \\
1+b_{2}(x) \\
a_{3,2}(x) \\
\cdot \\
\cdot \\
\cdot \\
a_{n, 2}(x)
\end{array}\right) \quad \ldots \quad X_{n-r}=\left(\begin{array}{c}
0 \\
\cdot \\
\cdot \\
1+b_{n-r}(x) \\
a_{n-r+1, n-r} \\
\cdot \\
a_{n, n-r}(x)
\end{array}\right) \text { and } \\
& \left.\left(\begin{array}{l}
X_{n-r+1}
\end{array}\right) \quad . \quad X_{n}\right)=\left(\begin{array}{c}
0 \\
\cdot \\
\cdot \\
0 \\
----- \\
D(x)
\end{array}\right) \quad \text { where } D(x) \in \mathcal{M}_{r}(\mathbb{R}),
\end{aligned}
$$

## $b_{j}(0)=a_{i, j}(0)=0$, and $D(0)=0$.

Firstly we can assume that $X_{j}(0)=\partial_{x_{j}}$ for $j=2, \ldots, n-r$ and $X_{j}(0)=0$ for $j=$ $n-r+1, \ldots, n$. Indeed $X_{2}, \ldots, X_{n}$ are tangent to $W$ where we can choose freely the cooordinates $\left(x_{2}, \ldots, x_{n}\right)$.

If $r=n-1$, which is generically possible only if $n=2$ and $r=1$, it is finished. Otherwise we can first replace $X_{2}$ by the normalization of $\sum_{j=2}^{n} a_{2 j} X_{j}$. Then we replace $X_{j}$ by $X_{j}-\frac{a_{2 j}}{a_{22}} X_{2}$ for $j>2$. These vector fields belong to $\Delta$, their first two coordinates vanish, and they are orthonormal to (the new) $X_{2}$. It remains to orthonormalize these $n-2$ vector fields. This cannot be done directly because some of them vanish in the singular locus. However, they are images by the vector bundle morphism $f$ of locally nonvanishing smooth sections of $E$ that can be orthonormalized.

The desired form of the vector fields is obtained by induction.

### 3.2.4 . Nilpotent and solvable approximation of generic distribtion

Thanks to the genericity results of the previous sections and with the help of local normal forms, it is finally shown that generically there are only two possibilities for the nilpotent/solvable approximation at a point $p \in \mathcal{Z}$ :

Theorem 2.4.3 For a generic distribution holds:
(i) Let $p$ be a tangency point in $\mathcal{Z}_{1}$, that is a point where $T_{p} \mathcal{Z}_{1}=\Delta_{p}$. Then $\widehat{X}_{n}=0$ but $\widetilde{X}_{n} \neq 0$, in normal form.
(ii) At all other points, including all points in $\mathcal{Z}_{r}$ with $r \geq 2$, the nilpotent approximation $\widehat{X}_{1}, \widehat{X}_{2}, \ldots, \widehat{X}_{n}$ is a set of $n$ linearly independant vector fields.

In conclusion, the only generic points where the solvable approximation is useful are tangency points in $\mathcal{Z}_{1}$.

### 3.4. Distance Induced by the Solvable Approximation of $n$-dimensional AlmostRiemannian Structures

The original system, the nilpotent and solvable approximations give rise to three different distances: $d, \widehat{d}$ and $\widetilde{d}$ respectively. This section deals with the almost-Riemannian distance defined by the solvable approximation at tangency points, of an $n$-dimensional ARS considering generic assumptions. The main result is Theorem 3.4 .4 which states that generically, the distance $\widetilde{d}$ is closer to $d$ than $\widehat{d}$ for pairs of points translated in an appropriate direction (Section 4.3.4). This translation condition is significant because the distance $\tilde{d}$ is not closer to $d$ than the distance induced by the nilpotent approximation for any pair of points. Then to prove the main result, it is essential to determine two facts. First, to state the order of approximation of $d$ by $\widetilde{d}$ (Theorem 3.4.1), and second, to find translation directions such that the distance $\tilde{d}$ of a pair of translated points is decreasing (see Section 3.3.4).

### 1.3.4 . Solvable approximation

By Theorem 2.4.1, generically there are points where the rank is $n-r$, as long as $r^{2}<n$. In this section, we consider only points belonging to $\mathcal{Z}_{1}$ because the only generic points where the solvable approximation is useful are tangency points in $\mathcal{Z}_{1}$ (see Theorem 2.4.3). Hence the distribution at $p=0$ has always dimension $n-1$, then assume that $X_{j}(0)=\frac{\partial}{\partial x_{j}}$ for $j=2, \ldots, n-1$. In Section 2.2.4, it is shown that the coordinates and $X_{1}, \ldots, X_{n}$ can be chosen in such a way that:

$$
X_{1}=\left(\begin{array}{c}
1 \\
0 \\
\vdots \\
\vdots \\
\vdots \\
0
\end{array}\right), X_{2}=\left(\begin{array}{c}
0 \\
1+\beta_{2}(x) \\
\alpha_{3,2}(x) \\
\vdots \\
\vdots \\
\alpha_{n, 2}(x)
\end{array}\right), \ldots, X_{n-1}=\left(\begin{array}{c}
0 \\
\vdots \\
\vdots \\
0 \\
1+\beta_{n-1}(x) \\
\alpha_{n, j}(x)
\end{array}\right), X_{n}=\left(\begin{array}{c}
0 \\
\vdots \\
\vdots \\
\vdots \\
0 \\
\alpha_{n}(x)
\end{array}\right)
$$

where $\beta_{j}(0)=\alpha_{i, j}(0)=0 i, j=2, \ldots, n$.
By the normal form of the vector fields, the singular locus is locally $\mathcal{Z}=\mathcal{Z}_{1}=\left\{\alpha_{n}(x)=0\right\}$.
Let $\alpha_{n}(x)=\widehat{\alpha}_{n}(x)+\widetilde{\alpha}_{n}(x)+\bar{\alpha}_{n}(x)$, that is $\alpha_{n}$ decomposed into its components of nonholonomic order 1,2 and greater than or equal to 3 respectively. Let $p$ be a tangency point in $\mathcal{Z}_{1}$, then by Theorem 2.4.3, $\widehat{X}_{n}=0$ and $\widetilde{X}_{n} \neq 0$, hence $\alpha_{n}(x)=\widetilde{\alpha}_{n}(x)+\bar{\alpha}_{n}(x)$.

Moreover, since $\operatorname{ord}_{p}\left(\widetilde{\alpha}_{n}\right)=2$ we get $\alpha_{n}(x)=a x_{n}+Q\left(x_{1}, x_{2}, \ldots, x_{n-1}\right)+\bar{\alpha}_{n}(x)$, where $Q\left(x_{1}, x_{2}, \ldots, x_{n-1}\right)$ is quadratic. Notice that, since the determinant is a submersion at $p$ (Theorem 2.4.1 item 3) $a \neq 0$.

We can find the nilpotent and solvable approximation in the coordinate system constructed in the normal form. Since $p=0$ is a tangency point and the weights of the coordinates are $w_{1}=\cdots=w_{n-1}=1$, and $w_{n}=2$, the nilpotent approximation is defined by

$$
\begin{aligned}
& \widehat{X}_{1}=X_{1}=\frac{\partial}{\partial x_{1}}, \\
& \widehat{X}_{j}=\frac{\partial}{\partial x_{j}}+\widehat{\alpha}_{n, j}(x) \frac{\partial}{\partial x_{n}} \text { for } j=2, \ldots, n-1, \text { and } \\
& \widehat{X}_{n}=0
\end{aligned}
$$

where $\widehat{\alpha}_{n, j}$ is the component of $\alpha_{n, j}$ of nonholonomic order 1 , hence it is linear in $x_{1}, x_{2}, \ldots, x_{n-1}$.
Therefore, the solvable approximation in the tangent case is defined by $\left\{\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}, \widetilde{X}_{n}\right\}$ where

$$
\widetilde{X}_{n}=\left(a x_{n}+Q\left(x_{1}, x_{2}, \ldots, x_{n-1}\right)\right) \frac{\partial}{\partial x_{n}}=\widetilde{a}_{n}(x) \frac{\partial}{\partial x_{n}},
$$

### 2.3.4 Comparison of distances

In this subsection, we state the order of approximation of the original distance by $\tilde{d}$. An important conclusion is that $\widetilde{d}$ improves the order of approximation of $d$ given by the nilpotent approximation.

Let $q$ and $q^{\prime}$ belong to the ball centered at $p$ and radius $\epsilon$, denoted by $B(p, \epsilon)$. We start by analyzing the divergence of curves respectively admissible for $d$ and $\widetilde{d}$, defined by the same control functions and starting at the same point $q$. Let us consider the vector fields $X_{1}, \ldots, X_{n}$ in normal form as in the above subsection. Hence we can express each vector field $X_{j}$ for $j=2, \ldots, n-1$, as

$$
X_{j}=\left(\begin{array}{c}
0 \\
\vdots \\
1+\beta_{j}(x) \\
\alpha_{j+1, j}(x) \\
\vdots \\
\alpha_{n, j}(x)
\end{array}\right) .
$$

Notice that $\alpha_{n, j}(x)$ can be split into components of order 1 and the remainder i.e $\alpha_{n, j}(x)=$ $\widehat{\alpha}_{n, j}(x)+\alpha_{n, j}^{+}(x)$. We denote by $\rho_{n, j}^{+}$the order of $\alpha_{n, j}^{+}$for $j=2, \ldots, n-1$.
Proposition 3.4.1 Let $\gamma$ be the geodesic for $d$ such that $\gamma(0)=q, \gamma(T)=q^{\prime}$ and associated with the control function $u(\cdot)$ satisfying $\|u\|=1$. Let $\widetilde{\gamma}$ be the admissible curve for $\widetilde{d}$ defined by the same control functions as $\gamma$ and $\widetilde{\gamma}(0)=q$. If $\rho_{n, j}^{+} \geq 3$ for $j=2, \ldots, n-1$, then

$$
\begin{equation*}
\|\gamma(t)-\widetilde{\gamma}(t)\|_{p} \leq C s t \cdot \tau^{\frac{3}{2}} \cdot t^{\frac{1}{2}} \tag{4.1}
\end{equation*}
$$

where $\tau=\max \left(\|q\|_{p}, t\right)$.
In order to state the result related to the comparison of distances, we need upper bounds for the distances $d$ and $\widetilde{d}$. So, from Theorems 7.31 and 7.26 of [Bellaïche, 1996] we get

$$
\begin{aligned}
& d\left(q, q^{\prime}\right) \leq \mathrm{Cst} \sum_{k, j \mid w_{k} \leq w_{j}}\|q\|_{p}^{1-\frac{w_{k}}{w_{j}}}\left|q_{k}-q_{k}^{\prime}\right|^{\frac{1}{w_{j}}}, \text { and } \\
& \widehat{d}\left(q, q^{\prime}\right) \leq \mathrm{Cst} \sum_{k, j \mid w_{k} \leq w_{j}}\|q\|_{p}^{1-\frac{w_{k}}{w_{j}}}\left|q_{k}-q_{k}^{\prime}\right|^{\frac{1}{w_{j}}} .
\end{aligned}
$$

Since $\widetilde{d}\left(q, q^{\prime}\right) \leq \widehat{d}\left(q, q^{\prime}\right)$, we get

$$
\widetilde{d}\left(q, q^{\prime}\right) \leq \operatorname{Cst} \sum_{k, j \mid w_{k} \leq w_{j}}\|q\|^{1-\frac{w_{k}}{w_{j}}}\left|q_{k}-q_{k}^{\prime}\right|^{\frac{1}{w_{j}}} .
$$

Since the weights of the coordinates are $w_{1}=\cdots=w_{n-1}=1$, and $w_{n}=2$, we obtain

$$
\begin{align*}
& d\left(q, q^{\prime}\right) \leq \operatorname{Cst}\left(\left\|q-q^{\prime}\right\|_{p}+\|q\|_{p}^{\frac{1}{2}}\left(\sum_{k=1}^{n-1}\left|q_{k}-q_{k}^{\prime}\right|^{\frac{1}{2}}\right)\right)  \tag{4.2}\\
& \widetilde{d}\left(q, q^{\prime}\right) \leq \operatorname{Cst}\left(\left\|q-q^{\prime}\right\|_{p}+\|q\|_{p}^{\frac{1}{2}}\left(\sum_{k=1}^{n-1}\left|q_{k}-q_{k}^{\prime}\right|^{\frac{1}{2}}\right)\right) . \tag{4.3}
\end{align*}
$$

The following notation and proposition are required for the comparison of distances result.
We denote by $\rho_{i, j}$ the order of $\alpha_{i, j}$ for $j=2, \ldots, n-1$ and $i=j, \ldots, n-1$ with the convention that $\alpha_{i, j}=\beta_{j}$ if $i=j$.

Proposition 3.4.2 Let $\gamma$ be the geodesic for $d$ such that $\gamma(0)=q, \gamma(T)=q^{\prime}$ and associated with the control function $u(\cdot)$ satisfying $\|u\|=1$. Let $\widetilde{\gamma}$ be the admissible curve for $\widetilde{d}$ defined by the same control functions as $\gamma$ and $\widetilde{\gamma}(0)=q$. If $\rho_{i, j} \geq 2$ and $\rho_{n, j}^{+} \geq 3$ for $j=2, \ldots, n-1$ and $i=j, \ldots, n-1$, then

1. $\widetilde{d}(\gamma(t), \widetilde{\gamma}(t)) \leq$ Cst $\cdot t^{\frac{1}{2}} \cdot \tau^{\frac{3}{2}}$.
2. $d(\gamma(t), \widetilde{\gamma}(t)) \leq C s t \cdot t^{\frac{1}{2}} \cdot \tau^{\frac{3}{2}}$,
where $\tau=\max \left(\|q\|_{p}, t\right)$.
Theorem 3.4.1 (Comparison of distances) If $\rho_{i, j} \geq 2$ and $\rho_{n, j}^{+} \geq 3$ for $j=2, \ldots, n-1$ and $i=j, \ldots, n-1$, then there exist constants $C$ and $\epsilon>0$, such that, for all $q, q^{\prime} \in B(p, \epsilon)$, we have

$$
\begin{equation*}
-C \tau^{\frac{3}{2}} d\left(q, q^{\prime}\right)^{\frac{1}{2}} \leq d\left(q, q^{\prime}\right)-\widetilde{d}\left(q, q^{\prime}\right) \leq C \cdot \widetilde{\tau}^{\frac{3}{2}} \widetilde{d}\left(q, q^{\prime}\right)^{\frac{1}{2}} \tag{4.4}
\end{equation*}
$$

where $\tau=\max \left(\|q\|_{p}, d\left(q, q^{\prime}\right)\right), \widetilde{\tau}=\max \left(\|q\|_{p}, \widetilde{d}\left(q, q^{\prime}\right)\right)$.
Notice that, if $d\left(q, q^{\prime}\right) \geq d(p, q)$ we get $\left|d\left(q, q^{\prime}\right)-\widetilde{d}\left(q, q^{\prime}\right)\right| \leq \mathrm{C} d\left(q, q^{\prime}\right)^{2}$. The similar inequality for $\widehat{d}$ is $\left|d\left(q, q^{\prime}\right)-\widehat{d}\left(q, q^{\prime}\right)\right| \leq \mathrm{C} d\left(q, q^{\prime}\right)^{\frac{3}{2}}$. This show that the order of bound of $\left|d\left(q, q^{\prime}\right)-\widetilde{d}\left(q, q^{\prime}\right)\right|$ is stricty better than the one of $\left|d\left(q, q^{\prime}\right)-\widehat{d}\left(q, q^{\prime}\right)\right|$.

### 3.3.4 . Translation

In this subsection, we address the second fact needed to prove the main result of this paper, that is, to find directions where the distance of a pair of translated points is decreasing. These directions are the appropriate ones where $\widetilde{d}$ is closer to $d$ than $\widehat{d}$.

It is well known that the distance defined by the nilpotent approximation is left-invariant (cf. [Jean, 2014]) while $\widetilde{d}$ is not. Let $p_{2}$ be a point in a neighborhood of $p=0$ and $g \in \mathbb{R}^{n}$. We are interested in conditions under which $\widetilde{d}\left(g, g \cdot p_{2}\right) \leq \widetilde{d}\left(0, p_{2}\right)$ (this means decreasing), where the product is the Lie group one. For this, some elements are required.
Let $\mathcal{L}=\operatorname{Lie}\left(\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}, \widetilde{X}_{n}\right), \mathfrak{g}$ the ideal generated in $\mathcal{L}$ by $\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}$ and $G$ the simply connected Lie group whose Lie algebra is $\mathfrak{g}$ that is, the set of left-invariant vector fields on $G$. We know that $\mathfrak{g}$ is a nilpotent Lie algebra. This Lie algebra $\mathfrak{g}$ can be split into homogeneous components

$$
\mathfrak{g}=\mathfrak{g}^{1} \oplus \mathfrak{g}^{2}
$$

where $\mathfrak{g}^{s}$ is the set of homogeneous vector fields of order $-s$.
The translation will be dealt with by considering a vector field $Y \in \mathfrak{g}^{1}$ and then $Y \in \mathfrak{g}^{2}$. We start by considering $Y \in \mathfrak{g}^{1}$. For the above, Definition 3.4.1 and Theorem 3.4.2 are necessary and they come from [Lee, 2006] and [Do Carmo, 1992].

Definition 3.4.1 Let $\gamma:[0, T] \longrightarrow \mathbb{R}^{n}$ be a smooth curve, and $\epsilon>0$. A variation of $\gamma$ is a smooth map $F:[0, T] \times(-\epsilon, \epsilon) \longrightarrow \mathbb{R}^{n}$ such that

$$
F(t, 0)=\gamma(t)
$$

for all $t \in[0, T]$.
For each $s \in(-\epsilon, \epsilon)$, the curve $\gamma_{s}:[0, T] \rightarrow \mathbb{R}^{n}$ given by $\gamma_{s}(t)=F(s, t)$, is called a curve of the variation $F$.

A variation $F$ of $\gamma$ determines a differentiable vector field $V(t)$ along $\gamma$ by $V(t)=\frac{\partial F}{\partial s}(t, 0)$.
We denote by $l\left(\gamma_{s}\right)$ the length of the curve $\gamma_{s}$.
Theorem 3.4.2 (First variation of length, [Lee, 2006]) Let $\gamma:[0, T] \longrightarrow \mathbb{R}^{n}$ be any unit speed admissible curve and $F(t, s)$ a smooth variation of $\gamma$. Then

$$
\frac{d}{d s} l\left(\gamma_{s}\right)(0)=-\int_{0}^{T}\left\langle V(t), \nabla_{\dot{\gamma}} \dot{\gamma}\right\rangle d t+\langle V(T), \dot{\gamma}(T)\rangle-\langle V(0), \dot{\gamma}(0)\rangle
$$

where $\nabla$ is the Levi-Civita connection.
From [Manríquez et al., 2022], we have that $\mathbb{R}^{n}$ is diffeomorphic to the quotient $G / H$ where $H$ stands for the connected subgroup of $G$ whose Lie algebra is the set of elements of $\mathfrak{g}$ that vanish at 0 . Moreover, the homogeneous space $G / H$ is the manifold of the right cosets
of $H$. We denote by $\Pi$ the canonical projection of $G$ onto $G / H$.
Let $\gamma:[0, T] \longrightarrow \mathbb{R}^{n}$ be a geodesic of $\widetilde{d}$ such that $\gamma(t) \notin \mathcal{Z}$ for $\left.\left.t \in\right] 0, T\right]$.
Let $Y \in \mathfrak{g}^{1}$, and $F:[0, T] \times(-\epsilon, \epsilon) \longrightarrow \mathbb{R}^{n}$ a variation of $\gamma$ such that

$$
F(t, s)=\gamma_{s}(t)=\gamma(t) \Pi(\exp (s Y))=L_{\gamma(t)}(\Pi(\exp (s Y)))
$$

where $L_{\gamma(t)}$ is the left translation by $\gamma(t)$.
For each $Y \in \mathfrak{g}^{1}$, the projection of $Y$ onto $G / H$ is denoted by $\Pi_{*} Y$, the latter is an invariant vector field on $G / H$ (cf. [Jouan, 2010])
Denoting by $V(t)$ the variation field of $F$, we get

$$
V(t)=\frac{\partial F}{\partial s}(t, 0)=\left.\frac{\partial}{\partial s}\right|_{s=0} \gamma_{s}(t)=T L_{\gamma(t)} \cdot \Pi_{*} Y=\left(\Pi_{*} Y\right)(\gamma(t)) .
$$

Then by Theorem 3.4.2,

$$
\frac{d}{d s} l\left(\gamma_{s}\right)(0)=\langle V(T), \dot{\gamma}(T)\rangle-\langle V(0), \dot{\gamma}(0)\rangle=\left\langle\left(\Pi_{*} Y\right)(\gamma(T)), \dot{\gamma}(T)\right\rangle-\left\langle\left(\Pi_{*} Y\right)(0), \dot{\gamma}(0)\right\rangle .
$$

Notice that the integral vanishes because $\gamma$ is a geodesic for $\widetilde{d}$ and hence $\nabla_{\dot{\gamma}} \dot{\gamma}=0$.
Since we want to study the translation of the curve $\gamma$ then $Y(0)$ must not be zero. Indeed, if $Y(0)=0$ then $\gamma(t) \Pi(\exp (0))=\gamma(t)$. So, we must look for a vector field $Y \in \mathfrak{g}^{1}$ such that $Y(0) \neq 0$.

So, we have the following proposition.
Proposition 3.4.3 Let $\gamma:[0, T] \longrightarrow \mathbb{R}^{n}$ be a length minimizer of $\tilde{d}$ with control functions $u_{1}, \ldots, u_{n}$, and $Y \in \mathfrak{g}^{1}$ such that $Y(\gamma(t))=\sum_{i=1}^{n-1} \alpha_{i} \widehat{X}_{i}(\gamma(t))$.

$$
\text { If } \sum_{i=1}^{n-1} \alpha_{i}\left(u_{i}(T)-u_{i}(0)\right)<0 \text { then } \frac{d}{d s} l\left(\gamma_{s}\right)(0)<0 .
$$

To deal with the case where the translation is in direction of a vector field $Y \in \mathfrak{g}^{2}$, i.e., a vector field $Y$ such that its evaluation does not belong to the tangent space at 0 , we must change the above strategy since it depends on $\langle Y(0), \dot{\gamma}(0)\rangle$, so if $Y(0)$ does not belong to the tangent space at $0,\langle Y(0), \dot{\gamma}(0)\rangle$ does not make sense.

The below (Proposition 3.4.4) is necessary and comes from [Bonfiglioli, 2007].
Let $G$ be a connected, simply connected Lie group of dimension $n$ such that $G$ is a Carnot (or stratified) group of step $r$ (see more details in [Bonfiglioli, 2007]) and $\mathfrak{g}$ its Lie algebra. After the choice of a basis $X_{1}, \ldots, X_{n}$ for $\mathfrak{g}$, the group $G$ is identified with $\mathbb{R}^{n}$ via the exponential mapping; this means that a point $x=\left(x_{1}, \ldots, x_{n}\right) \in \mathbb{R}^{n}$ is identified with the point $\exp \left(x_{1} X_{1}+\right.$ $\ldots+x_{n} X_{n}$ ) of the group. Hence we have the following result obtained from [Bonfiglioli, 2007] and [Serra Cassano, 2016].

Proposition 3.4.4 The group product has the form

$$
\begin{equation*}
x \cdot y=x+y+Q(x, y) \quad \forall x, y \in \mathbb{R}^{n}, \tag{4.5}
\end{equation*}
$$

where $Q=\left(Q_{1}, \ldots, Q_{n}\right): \mathbb{R}^{n} \times \mathbb{R}^{n} \rightarrow \mathbb{R}^{n}$ and $Q_{j}$ is a homogeneous polynomial of degree $w_{j}$. Moreover, for all $x, y \in \mathbb{R}^{n} Q_{1}(x, y)=\cdots=Q_{n_{1}}(x, y)=0$, where $n_{1}$ is such that $w_{1}=\cdots=w_{n_{1}}=1$, and

$$
Q_{i}(x, y)=\sum_{h, k} \mathcal{R}_{k, h}^{i}(x, y)\left(x_{k} y_{h}-x_{h} y_{k}\right),
$$

where the functions $\mathcal{R}_{k, h}^{i}$ are polynomials, homogenous of degree $w_{i}-w_{h}-w_{k}$ with respect to group dilations, and the sum is extended to all $h, k$ such that $w_{k}+w_{h} \leq w_{i}$.

Remark 3.4.1 In the context of the generic case, which is the case that interests us in this section, $n_{1}=n-1$, and since rank $\left(\Delta_{p}+[\Delta, \Delta]_{p}\right)=n$ at all points then $\mathcal{R}_{k, h}^{n}$ is a constant, and $w_{k}=w_{h}=1$.

Let $Y \in \mathfrak{g}^{2}$. In local coordinates $Y(x)=\sum_{i=1}^{n} f_{i}(x) \frac{\partial}{\partial x_{i}}$. Since $w_{1}=\ldots=w_{n-1}=1$ and $w_{n}=2$ then $f_{i}(x) \equiv 0$ for $i=1, \ldots, n-1$, and $f_{n}(x)$ is a constant different from 0 , hence $Y(x)=\nu \frac{\partial}{\partial x_{n}}$, with $\nu \neq 0$ (see Figure 4.1).


Figure 4.1: Singular set and its tangent space at $p$ with the translation vector $Y(x)$.
Let $\mu_{Y}$ be the integral curve of $Y$ passing through the identity of $G$ when $t=0$, then

$$
\dot{\mu}_{Y}(t)=Y\left(\mu_{Y}(t)\right)=\nu \frac{\partial}{\partial x_{n}},
$$

hence

$$
\mu_{Y}(t)=\left(\begin{array}{c}
0 \\
\vdots \\
0 \\
t \nu
\end{array}\right) .
$$

Since $\exp (Y)=\mu_{Y}(1)$ (see more details in [Bonfiglioli, 2007]), then $\exp (Y)=\left(\begin{array}{c}0 \\ \vdots \\ 0 \\ \nu\end{array}\right)$.
Let $\gamma:[0, T] \longrightarrow \mathbb{R}^{n}$ be a geodesic of $\widetilde{d}$ such that $\gamma(t)=\left(\gamma_{1}(t), \ldots, \gamma_{n}(t)\right), \gamma(t) \notin \mathcal{Z}$ for $t \in] 0, T]$, and $\gamma(0)=0$, with control functions $u_{1}, \ldots, u_{n}$. We consider $Y \in \mathfrak{g}^{2}$ such that, $Y(x)=\nu \frac{\partial}{\partial x_{n}}$. Let $\gamma_{L Y}(t)=L_{\exp (Y)}(\gamma(t))=\left(\gamma_{1 Y}(t), \ldots, \gamma_{n Y}(t)\right)$ and $\bar{u}_{1}, \ldots, \bar{u}_{n}$ its control functions.

Recall that $\widetilde{a}_{n}(x)=a x_{n}+Q\left(x_{1}, x_{2}, \ldots, x_{n-1}\right)$, where $Q\left(x_{1}, x_{2}, \ldots, x_{n-1}\right)$ is quadratic. We set $\widetilde{a}_{n}(\gamma)=\widetilde{a}_{n}(\gamma(t))$.

The following result provides conditions on $Y$ such that $\gamma_{L Y}$ has a length less than $\gamma$.
Theorem 3.4.3 Let $\gamma:[0, T] \longrightarrow \mathbb{R}^{n}$ be a length minimizer of $\widetilde{d}$ with control functions $u_{1}(t), \ldots, u_{n}(t)$ with $u_{n}(t) \neq 0$ almost everywhere, and $\widetilde{a}_{n}(\gamma)+a \nu \neq 0$. If $\left|\widetilde{a}_{n}(\gamma)\right|<$ $\left|\widetilde{a}_{n}(\gamma)+a \nu\right|$ then $\widetilde{d}\left(\gamma_{L Y}(0), \gamma_{L Y}(T)\right)<\widetilde{d}(\gamma(0), \gamma(T))$.

### 4.3.4 . Solvable distance better than the nilpotent one

It is known that the almost-Riemannian distance $d$ of the original system, close to $p=0$, behaves at the first-order as the distance defined by the nilpotent approximation at $p=0$. However, thanks to Theorem 3.4.1 we know that the solvable approximation improves the order of approximation of $d$ given by the nilpotent approximation. Despite the above, we can not state that the solvable distance is closer than the nilpotent one to the original distance for all pairs of points.

In this section, we prove that the approximation by $\widetilde{d}$ is better than the one by $\widehat{d}$ for a pair of points translated in a direction where the distance $\widetilde{d}$ is decreasing.

Before that, notice that we know by Proposition 1.3.3 that $\widetilde{d}\left(q, q^{\prime}\right) \leq \widehat{d}\left(q, q^{\prime}\right)$ then we can conclude that for $\widetilde{d}$ to be better than $\widehat{d}$ it must be satisfied that $\frac{\widetilde{d}\left(q, q^{\prime}\right)+\widehat{d}\left(q, q^{\prime}\right)}{2}>d\left(q, q^{\prime}\right)$.

Let $\gamma(t)=\left(\gamma_{1}(t), \ldots, \gamma_{n}(t)\right)$ be a (normal) geodesic of $\widetilde{d}$ such that $\gamma(0)=0$ with control functions $u_{1}, u_{2}, \ldots, u_{n}$ parametrized by arc length on $[0, T]$, and the length of the curve $\gamma$ is denoted by $l(\gamma)$. We consider $g \in \mathbb{R}^{n}$ such that $g=\exp (Y)$ with $Y \in \mathfrak{g}$ and $Y$ satisfiying Theorem 3.4.3. Let $\gamma_{g}(t)=L_{g}(\gamma(t))$ and $\bar{u}_{1}, \bar{u}_{2}, \ldots, \bar{u}_{n}$ its control functions. Note that $\gamma_{g}$ is admissible for $\widetilde{d}$ as long as it does not meet $\mathcal{Z}$. Indeed, all absolutely continuous curves are admissible out of the singular locus since the metric is Riemannian.

Let $\epsilon>0$ such that $C_{2} \cdot \widetilde{d}(0, \gamma(T))=C_{2} \cdot T \leq \epsilon$, where $C_{2}$ is the constant of inequality 4.4 of Theorem 3.4.1.

On the other hand, by Proposition 3.4.4 we have that $u_{i}(t)=\bar{u}_{i}$ for $i=1, \ldots, n-1$ an since $l\left(\gamma_{g}\right)<l(\gamma)$, we can assume that there exists $C:[0, T] \rightarrow[0,1[$ such that

$$
\begin{equation*}
\left|u_{n}(t)\right| C(t)=\left|\bar{u}_{n}(t)\right| . \tag{4.6}
\end{equation*}
$$

Moreover, from Pontryagin's maximum principle (more details see [Agrachev et al., 2019]) we know that

$$
\begin{equation*}
u_{n}(t)=\left\langle\lambda(t), \widetilde{X}_{i}\right\rangle=\lambda_{n}(t) \cdot\left(a x_{n}+Q\left(x_{1}, x_{2}, \ldots, x_{n-1}\right)\right) \tag{4.7}
\end{equation*}
$$

where $\lambda(t) \in T_{\gamma(t)}^{*} \mathbb{R}^{n}$.
Let $b=\max (C(t), t \in[0, T])$ and $S=\frac{\left(1-b^{2}\right) \lambda_{n}(0)^{2}(n-2)^{2} a^{2}}{40(n-1)^{2}} T^{4}$.
The following Theorem is the main result in this section.
Theorem 3.4.4 With the previous notations. If $\frac{2 \epsilon}{1+2 \epsilon}<S$ then

$$
\left|\widehat{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right)-d\left(\gamma_{g}(0), \gamma_{g}(T)\right)\right|>\left|d\left(\gamma_{g}(0), \gamma_{g}(T)\right)-\widetilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right)\right|
$$

and in consequence $\widetilde{d}$ is closer than $\widehat{d}$ to $d$.

## Part II

## Continuation Method in the rolling problem with obstacles

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A nonholonomic system on a $n$-dimensional manifold $M$ is a control system which is of the form

$$
\begin{equation*}
\dot{x}=\sum_{i=1}^{m} u_{i} X_{i}(x), \quad x \in M, \tag{1.1}
\end{equation*}
$$

where $m>1$ is an integer and $X_{1}, X_{2}, \ldots, X_{m}$ are $C^{\infty}$ vector field on $M$ (cf. [Jean, 2014]). From the control viewpoint, nonholonomic systems are nonlinear systems, and whose tangent linearization does not preserve controllability. These systems provide an important class of mechanical control systems such as rolling contact or sliding contact (cf. [Bloch, 2003]). With the great development of robotics, these systems have become very important; in particular, nonholonomy of rolling is relevant to robotic manipulation (see for instance [Chung, 2004], [Kolmanovsky and McClamroch, 1995], [Marigo and Bicchi, 2000], [Murray et al., 2017]).

The rolling-body problem (without slipping nor spinning) is a control system $\Sigma$, which models the rolling of an embedded connected surface $A_{2}$ in $\mathbb{R}^{3}$ on another $A_{1}$ one. A state of the system is an orientation-preserving isometry. The state space of the system $\Sigma$, denoted by $\mathcal{Q}\left(A_{1}, A_{2}\right)$, is a 5 -dimensional connected manifold (details in [Agrachev and Sachkov, 2004]). As a consequence of the rolling constraints (no slipping and no spinning), and given an absolutely continuous (a.c.) curve $\gamma_{1}$ on $A_{1}$, there exists a unique a.c. curve $\Gamma$ in $\mathcal{Q}$, which describes the rolling of the surface $A_{2}$ onto the surface $A_{1}$ along the curve $\gamma_{1}$. Thus, the admissible controls of $\Sigma$ correspond to the a.c. curves $\gamma_{1}$ of $A_{1}$ by their derivatives $\dot{\gamma}_{1}$. Then, the system of control can be written, in local coordinates, as a nonholonomic system

$$
\Sigma: \quad \dot{x}=u_{1} X_{1}+u_{2} X_{2},
$$

where $\left(u_{1}, u_{2}\right) \in \mathbb{R}^{2}$ is the control, and $X_{1}$ and $X_{2}$ are vector fields defined in the domain of the chart.

We can identify three aspects that arise in the rolling bodies problem: modeling of rolling surfaces, motion planning for rolling systems, and controllability. Regarding the modeling of rolling surfaces, several papers can be found where first and second-order kinematic contact equations are derived, see for instance the papers [Cai and Roth, 1987], [Montana, 1988], [Sarkar et al., 1997] and [Woodruff and Lynch, 2019]. A historical account of the problem of rolling two Riemannian manifolds one on the other, ranging from classical to modern results, can be found in [Chitour et al., 2014]. Concerning the motion planning problem, the simplest
model is a sphere rolling on the plane. The first time that this problem was considered as worthy of study was in the papers of Chaplygin ([Chaplygin, 2002] and [Chaplygin, 2012]) afterwards, it has been dealt in many papers, for instance [Li and Canny, 1990], [Oriolo and Vendittelli, 2005], and [Jurdjevic, 1993]. Considering a more general convex body in the motion planning problem, Marigo and Bicchi proposed an approximate motion-planning algorithm for general convex body $A_{2}$ [Marigo and Bicchi, 2000]. Alouges et al. in [Alouges et al., 2010] provide an algorithm, which is based on the Continuation Method for a convex surface rolling on a plane (more detail about the Continuation Method, see Section 3.2). In [Woodruff et al., 2020], the authors presented a method to generate motion plans and stabilize feedback controllers for general, smooth, 3 -dimensional objects in rolling contact. Chelouah and Chitour give an approach resting on the Liouvillian character of $\Sigma$. More precisely, if just one of the manifolds has a symmetry of revolution (i.e., $A_{1}$ or $A_{2}$ ), then $\Sigma$ is shown to be a Liouvillian system. If, in addition, that manifold is convex and the other one is a plane, then a maximal linearizing output is explicitly computed [Chelouah and Chitour, 2003]. In the framework of the motion planning problem with obstacles, Bicchi and Marigo used a lattice structure on the state space and translated Li-Canny's (see [Li and Canny, 1990]) global and exact computation into a series of local and approximate ones with good topological properties so that it could be incorporated into a more-general motion-planning algorithm dealing with obstacles in the plane [Bicchi and Marigo, 2002]. Grushkovskaya and Zuyev present a theoretical analysis of the obstacle avoidance problem and prove asymptotic stability. The time-varying control strategy is defined explicitly in terms of the gradient of a potential function [Grushkovskaya and Zuyev, 2018]. Divelbiss and Wen consider the path planning problem without obstacles by transforming it into a nonlinear least-squares problem in an augmented space which is then iteratively solved. Obstacle avoidance is included as inequality constraints, and exterior penalty functions are used to convert the inequality constraints into equality restrictions [Divelbiss and Wen, 1997]. Concerning the controllability issues, Agrachev and Sachkov [Agrachev and Sachkov, 2004] show that the control system $\Sigma$ is locally controllable at a point $q \in \mathcal{Q}$ if the local Gaussian curvatures of $A_{1}$ and $A_{2}$ are not equal. Furthermore, they proved that $\Sigma$ is completely controllable if and only if $A_{1}$ and $A_{2}$ are not isometric. In [Marigo and Bicchi, 2000], necessary conditions for the reachability of rolling contacts are defined.

In this part of the thesis, we address the motion planning problem of a strictly convex body $A_{2}$ rolling (without slipping nor spinning) on the Euclidean plane $\mathbb{R}^{2}$ with obstacles, with a numerical implementation of the Continuation Method.

The Continuation Method, also called homotopy method or continuous Newton's algorithm [Allgower and Georg, 1993], was introduced in [Sussmann, 1992] and [Sussmann, 1993] and widely developed in [Chitour and Sussmann, 1998], [Chitour, 2002], and [Chitour, 2006]. It is used to solve nonlinear equations of the form $\mathcal{L}(x)=y$, where $x$ is the unknown, and $\mathcal{L}: X \rightarrow Y$ is surjective. In the context of the motion planning, $\mathcal{L}$ is the endpoint map (associated with some fixed point $p$ ) from the space of admissible inputs to the state space, that is $E_{p}: \mathcal{H} \rightarrow M$, and it proceeds by starting from a value $u_{0} \in \mathcal{H}$ and $y_{0}=E_{p}\left(u_{0}\right)$, then by joining $y_{0}$ to the given $y$ by a continuous path $\pi$ and by trying to lift $\pi$ to the path
$\Pi$ so that $E_{p} \circ \Pi=\pi$. To construct such a path $\Pi$, we may differentiate $E_{p}(\Pi(s))=\pi(s)$ to get $D E_{p}(\Pi(s)) \dot{\Pi}(s)=\dot{\pi}(s)$. The latter is satisfied if we can solve $\dot{\Pi}(s)=F(\Pi(s)) \dot{\pi}(s)$ where $F(x)$ is a right inverse of $D E_{p}(u)$. Therefore, solving $E_{p}(u)=y$ is equivalent to firstly show that $F(\Pi(s))$ exists (for instance if $D E_{p}(\Pi(s))$ is surjective) and secondly, proving that the ODE in $\mathcal{H}, \dot{\Pi}(s)=F(\Pi(s)) \dot{\pi}(s)$, which is also called the path-lifting equation (PLE) or Wazewski equation [Ważewski, 1947], admits a global solution. The singularities of the $E_{p}$ are exactly the abnormal extremals of the sub-Riemannian metric induced by the dynamics of the system, which are usually a significant obstacle to efficiently apply the Continuation Method to the motion planning problem. In the case of $\Sigma$, nontrivial abnormal extremals and their trajectories were determined in [Chelouah and Chitour, 2003], and they are precisely the horizontal geodesics of $\Sigma$. However, by assuming that the surface $A_{2}$ is strictly convex and possesses a stable periodic geodesic, and $A_{1}$ is a plane, it was shown in [Chelouah and Chitour, 2003] that the PLE $\Pi(s)=F(\Pi(s)) \dot{\pi}(s)$ has a global solution, that is, the Continuation Method offers a total solution to the motion planning problem.

The convex body $A_{2}$, which can be embedded as a convex surface in $\mathbb{R}^{3}$, is assumed to have a stable periodic geodesic. We denote the state space $\mathcal{Q}\left(\mathbb{R}^{2}, A_{2}\right)$ simply by $M$. An obstacle $\mathcal{W}$ in $\mathbb{R}^{2}$ is a nonempty compact subset of $\mathbb{R}^{2} . \mathcal{W}$ maps in $M$ a region $C$. Thus, an obstacle in $M$ is a nonempty closed subset of $M$ such that $\widehat{M}=M \backslash C$ is also nonempty. Let us consider the control system on $M$ defined by

$$
\dot{y}=v_{1} \bar{X}_{1}(y)+v_{2} \bar{X}_{2}(y)
$$

where $\bar{X}_{i}=\zeta X_{i}, i=1,2$ and $\zeta: M \rightarrow \mathbb{R}$ such that $\zeta>0$ on $\widehat{M}, \zeta=0$ on $C$. Then $\widehat{M}$ is invariant under the above control system. Thus, the motion planning problem with obstacles is reduced to a motion planning problem for each connected component of $\widehat{M}$.

Therefore, we provide a complete numerical implementation of the Continuation Method presented above to solve the motion planning problem with forbidden regions on the plane.

The second part of this thesis is organized as follows.
Chapter 2 contains generalities about differential geometry (geodesic coordinates, orthonormal moving frame, Christoffel symbols, connection form, Gaussian curvature, and geodesics), rolling body problem, and Continuation Method.

In Chapter 3, we fix the definition of an obstacle, and we provide the dynamics of the control system on the state space $\mathcal{Q}\left(A_{1}, A_{2}\right)$ such that $A_{1}=\mathbb{R}^{2}$ and $A_{2}$ is an oriented surface of $\mathbb{R}^{3}$ with a geometric condition. Moreover, Section 2.3 provides a detailed description of the fundamental points for the numerical implementation: discretizing the control space, computing the differential of endpoint map, and the lift of the curve on the convex body $A_{2}$. Section 3.3 provides three examples of the bodies rolling on the Euclidean plane: the sphere, the flattened ball, and an egg. Finally, Section 4.3 contains conclusions and some comments.

## CHAPTER 2

## Mathematical prerequisites

This section is dedicated to present some definitions and results used in this part of the thesis related to differential geometry, rolling body problem, and the Continuation Method.

### 1.2. Differential geometry

The following definitions and results come from [Agrachev et al., 2019], [Klingenberg, 1982], and [Berger and Gostiaux, 2012].

Let $(A,\langle\cdot, \cdot\rangle)$ be a 2-dimensional, connected, oriented, smooth, complete Riemannian manifold for the Riemannian metric $\langle\cdot, \cdot\rangle$. We use TS to denote the tangent bundle over $A$ and $U A$ the unit tangent bundle, i.e., the subset of TA of points $(x, v)$ such that $x \in A$ and $v \in T_{x} A$, $\langle v, v\rangle=1$.

Let $\left\{U_{\alpha}, \alpha\right\}_{\alpha \in \mathcal{A}}$ be an atlas on $A$. For $\alpha_{1}, \alpha_{2} \in \mathcal{A}$ such that $U_{\alpha_{1}} \cap U_{\alpha_{2}}$ is not empty, we denote by $J_{\alpha_{2} \alpha_{1}}$ the Jacobian matrix of $\varphi^{\alpha_{2}} \circ\left(\varphi^{\alpha_{1}}\right)^{-1}$ the coordinate transformation on $\varphi^{\alpha_{1}}\left(U_{\alpha_{1}} \cap U_{\alpha_{2}}\right)$. For $\alpha_{1} \in \mathcal{A}$, the Riemannian metric is represented by the symmetric positivedefinite matrix $\mathcal{I}^{\alpha_{1}}$ and set $M^{\alpha_{1}}=\sqrt{\mathcal{I}^{\alpha_{1}}}$.

### 1.1.2 Geodesic coordinates

The geodesic coordinates on $A$ are charts $(v, w)$ defined such that $G^{\alpha_{1}}$ is diagonal with $g_{11}=1$ and $g_{22}=\vartheta^{2}(v, w)$. The function $\vartheta$ is defined in an open neighborhood of $(0,0)$ (the domain of the chart) and satisfies $\vartheta(0, w)=1, \vartheta_{v}(0, w)=0$, and $\vartheta_{v v}+K \vartheta=0$, where $K$ denotes the Gaussian curvature of $A$ at $(v, w)$; and $\vartheta_{v}\left(\vartheta_{v v}\right)$ is the (double) partial derivative of $\vartheta$ with respect to $v$.

### 2.1.2 . Orthonormal moving frame (OMF)

Definition 1.2.1 For $x \in A$, a frame $\Phi$ at $x$ is an ordered basis for $T_{x} A$ and, for $\alpha_{1}, \alpha_{2} \in \mathcal{A}$, we have $\Phi^{\alpha_{2}}=J_{\alpha_{2} \alpha_{1}} \Phi^{\alpha_{1}}$. The frame $\Phi$ is orthonormal if, in addition, $M^{\alpha_{1}} \Phi^{\alpha_{1}}$ is an orthogonal
matrix.
An orthonormal moving frame (OMF), defined on an open subset $U$ of $A$, is a smooth map assigning to each $x \in U$ a positively oriented orthonormal frame $\Phi(x)$ of $T_{x} A$.

### 3.1.2 . Christoffel symbols and connection form

Let $\nabla$ be the Riemannain connection on $A$.
For a given OMF $\Phi$ defined on $U \subset A$, the Christoffel symbols associated with $\Phi=\left(\Phi_{1}, \Phi_{2}\right)$ are defined by $\nabla_{\Phi_{i}} \Phi_{j}=\sum_{k} \Gamma_{i j}^{k} \Phi_{k}$ where $1 \leq i, j, k \leq 2$.

The connection form $\omega$ is the mapping defined on $U$ such that, for every $x \in U, \omega_{x}$ is the linear application from $T_{x} A$ to the set of $2 \times 2$ skew-symmetric matrices given as follows. For $i, j, k=1$ and 2 , the $(i, j)$ th coefficient of $\omega_{x}\left(\Phi_{k}\right)$ is equal to $\Gamma_{i j}^{k}$.

### 4.1.2 Parallel vector field

Let $\gamma: I \rightarrow A$ be an a.c curve in $A$ with $I$ compact interval of $\mathbb{R}$. Set $X(t):=\dot{\gamma}(t)$ in $I$, which defines a vector field along $\gamma$. Let $Y: I \rightarrow T A$ be an a.c. assignment such that, for every $t \in I, Y(t) \in T_{\gamma(t)} A$.
Definition 1.2.2 We say that $Y$ is parallel along $\gamma$ if $\nabla_{X} Y=0$ for almost all $t \in I$. Moreover, in the domain of an OMF $\Phi$, that equation can be written as follows: $\dot{Y}^{k}=-\sum_{1 \leq i, j, k \leq 2} \Gamma_{i j}^{k} X^{i} Y^{j}$ or equivalently $\dot{Y}=-\omega(X) Y$.

### 5.1.2 Gaussian curvature and geodesics

Let $A$ be an oriented surface of $\mathbb{R}^{3}$ with metrics induced by the Euclidean metric of $\mathbb{R}^{3}$. We assume that $A$ is defined as one bounded connected component of the zero-level set of a smooth real-valued function $a: \mathbb{R}^{3} \rightarrow \mathbb{R}$, and let $K$ be the Gaussian curvature of $A$.

Proposition 1.2.1 ([Berger and Gostiaux, 2012]) With the previous notation, we have

$$
K=-\frac{\operatorname{det}\left(\begin{array}{cc}
\nabla^{2} a & \nabla a  \tag{2.1}\\
(\nabla a)^{T} & 0
\end{array}\right)}{\|\nabla a\|^{4}},
$$

where $\nabla^{2} a$ is the Hessian matrix of $a$, and $(\nabla a)^{T}$ denotes the transpose matrix of $\nabla a$.
Definition 1.2.3 A smooth curve $\gamma:[0,1] \rightarrow A$ parametrized with constant speed is called geodesic if it satisfies

$$
\ddot{\gamma}(t) \perp T_{\gamma(t)} A
$$

for almost $t$ in $[0,1]$.

Since $A$ is defined as (a bounded connected component of) the zero-level set of a real-valued function $a$, we have the following Proposition

Proposition 1.2.2 (see [Agrachev et al., 2019]) A smooth curve $\gamma:[0,1] \rightarrow A$ is a geodesic if and only if it satisfies

$$
\begin{equation*}
\ddot{\gamma}=-\frac{\dot{\gamma}^{T} \nabla^{2} a(\gamma) \dot{\gamma}}{\|\nabla a(x)\|^{2}} \nabla a(\gamma) . \tag{2.2}
\end{equation*}
$$

## 2.2 . Rolling body problem

Let us consider the rolling body problem with no slipping or spinning of $A_{2}$ on top of $A_{1}$. We adopt the viewpoint presented in [Agrachev and Sachkov, 2004].

Let $A_{1}$ and $A_{2}$ be two-dimensional connected, oriented Riemannian manifold (surfaces of the rolling bodies). At the contact points of the bodies $x_{1} \in A_{1}$ and $x_{2} \in A_{2}$, their tangent spaces are identified by an orientation-preserving isometry $q: T_{x_{2}} A_{2} \rightarrow T_{x_{1}} A_{1}$ (see Figure 2.1). Such an isometry $q$ is a state of the system, and the state space is given by

$$
\mathcal{Q}\left(A_{1}, A_{2}\right)=\left\{q: T_{x_{2}} A_{2} \rightarrow T_{x_{1}} A_{1} \mid x_{1} \in A_{1}, x_{2} \in A_{2}, q \text { an isometry }\right\} .
$$



Figure 2.1: Identification of tangent spaces at a contact point.
As the set of all orientation-preserving isometries in $\mathbb{R}^{2}$ is $S O(2)$, which can be identified with the unit circle $S^{1}$ in $\mathbb{R}^{2}, \mathcal{Q}\left(A_{1}, A_{2}\right)$ is a 5 -dimensional connected manifold. A point $q \in \mathcal{Q}\left(A_{1}, A_{2}\right)$ is locally parametrized by $\left(x_{1}, x_{2}, R\right)$ with $x_{1} \in A_{1}, x_{2} \in A_{2}$, and $R \in S O(2)$.

Let $\Phi_{1}$ and $\Phi_{2}$ be two OMFs defined on the chart domains of $\alpha_{1}, \alpha_{2}$. For $i=1,2$, consider a curve $\gamma_{i}^{\alpha_{i}}$ defined inside the chart domian $\alpha_{i}$ on the body $A_{i}$. Let $b_{i}(t)=\Phi_{i}\left(\gamma_{i}(t)\right) R_{i}(t)$ parallel along $\gamma_{i}^{\alpha_{1}}, i=1,2$, and $R:=R_{2}(t) R_{1}(t)^{-1} \in S O(2)$, which measures, by definition, the relative position of $\Phi_{2}$ with respect to $\Phi_{1}$ along $\left(\gamma_{1}^{\alpha_{1}}, \gamma_{2}^{\alpha_{2}}\right)$. The variation of $R_{i}$ along $\gamma_{i}^{\alpha_{i}}$, for $i=1,2$, is given by $\dot{R}_{i}=-\omega_{i}\left(\dot{\gamma}_{i}^{\alpha_{i}}\right) R_{i}$.

Given an a.c. curve $\gamma_{1}:[0, T] \rightarrow A_{1}$, the rolling of $A_{2}$ on $A_{1}$ without slipping or spinning along $\gamma_{1}$ is characterized by a curve $\Gamma=\left(\gamma_{1}, \gamma_{2}, R\right):[0, T] \rightarrow \mathcal{Q}\left(A_{1}, A_{2}\right)$, which is defined by
the following two equations:

$$
\begin{array}{rlrl}
M^{\alpha_{2}} \dot{\gamma}_{2}^{\alpha_{2}}(t) & =R M^{\alpha_{1}} \dot{1}_{1}^{\alpha_{1}}(t) & & \text { (no-slipping condition) } \\
\dot{R} R^{-1} & =\omega_{1}\left(\dot{\gamma}_{1}^{\alpha_{1}}\right)-\omega_{2}\left(\dot{\gamma}_{2}^{\alpha_{2}}\right) & \text { (no-spinning condition) } \tag{2.4}
\end{array}
$$

Definition 2.2.1 The surface $A_{2}$ rolls on $A_{1}$ without slipping nor spinning if, for every $x=$ $\left(x_{1}, x_{2}, R_{0}\right) \in \mathcal{Q}\left(A_{1}, A_{2}\right)$ and a.c curve $\gamma_{1}:[0, T] \rightarrow A_{1}$ starting at $x_{1}$, there exists an a.c curve $\Gamma:[0, T] \rightarrow \mathcal{Q}\left(A_{1}, A_{2}\right)$, with $\Gamma(t)=\left(\gamma_{1}(t), \gamma_{2}(t), R(t)\right), \Gamma(0)=x$ and for every $t \in[0, T]$, such that, on appropriate charts, equations (2.3) and (2.4) are satisfied. We call the curve $\Gamma(t)$ an admissible trajectory.

Equations (2.3) and (2.4) can be rewritten in local coordinates by considering $f_{1}$ and $f_{2}$ two OMFs and if the state $x=\left(\gamma_{1}, \gamma_{2}, R\right)$ then for almost all $t$ such that $x(t)$ remains in the domain of an appropriate chart, there exists a measurable function $u:[b, d] \rightarrow \mathbb{R}^{2}$ (which is called control) such that

$$
\begin{aligned}
\dot{\gamma}_{1}(t) & =u_{1}(t) \Phi_{1}^{1}\left(\gamma_{1}(t)\right)+u_{2}(t) \Phi_{2}^{1}\left(\gamma_{1}(t)\right) \\
\dot{\gamma}_{2}(t) & =u_{1}(t)\left(\Phi^{2}\left(\gamma_{2}(t)\right) R(t)\right)_{1}+u_{2}(t)\left(\Phi^{2}\left(\gamma_{2}(t)\right) R(t)\right)_{2} \\
\dot{R}(t) R^{-1}(t) & =\sum_{i=1}^{2} u_{i}(t)\left(\omega_{1}\left(\Phi_{i}^{1}\left(\gamma_{1}(t)\right)\right)-\omega_{2}\left(\Phi^{2}\left(\gamma_{2}(t)\right) R(t)\right)_{i}\right) .
\end{aligned}
$$

Then, we can rewrite the above system of equations as follows:

$$
\begin{equation*}
\dot{x}=u_{1} X_{1}(x)+u_{2} X_{2}(x), \tag{2.5}
\end{equation*}
$$

where $X_{i}=\left(\Phi_{i}^{1},\left(\Phi^{2} R\right)_{i},\left(\omega_{1}\left(\Phi_{i}^{1}\right)-\omega_{2}\left(\Phi^{2} R(t)\right)_{i}\right)\right)^{T}, i=1,2$.
The following proposition describes a fundamental property of the rolling-body problem (see [Chelouah and Chitour, 2003] for more details).

Proposition 2.2.1 Let $u \in \mathcal{H}$ be an admissible control that gives rise to the admissible trajectory $\Gamma=\left(\gamma_{1}, \gamma_{2}, R\right):[0,1] \rightarrow \mathcal{Q}\left(A_{1}, A_{2}\right)$. Then, the following statements are equivalent:

1. The curve $\gamma_{1}:[0,1] \rightarrow A_{1}$ is a geodesic;
2. The curve $\gamma_{2}:[0,1] \rightarrow A_{2}$ is a geodesic;
3. The curve $\Gamma:[0,1] \rightarrow \mathcal{Q}\left(A_{1}, A_{2}\right)$ is a horizontal geodesic.

### 3.2. Continuation Method

This section provides a complete and general description of the Continuation Method and comes from [Chitour, 2006].

We fix some notations. The admissible inputs $u$ are elements of $\mathcal{H}=L^{2}\left([0,1], \mathbb{R}^{2}\right)$. To denote $\left(u_{1}(t)^{2}+u_{2}(t)^{2}\right)^{\frac{1}{2}}$ and $\left(\int_{0}^{1}\|u(t)\|^{2} d t\right)^{\frac{1}{2}}$ we use $\|u(t)\|$ and $\|u\|_{\mathcal{H}}$ respectively. If $u, v \in \mathcal{H}$, then $(u, v)_{\mathcal{H}}=\int_{0}^{1} u(t)^{T} v(t) d t$.

As mentioned in the Introduction Section, the Continuation Method is used to solve nonlinear equations of the form $\mathcal{L}(x)=y$, where $x$ is the unknown, and $\mathcal{L}: X \rightarrow Y$ is surjective.

In the context of the motion planning problem, the map $\mathcal{L}$ is the end-point $E_{p}: \mathcal{H} \rightarrow M$ associated with some fixed $p \in M$. For $u \in \mathcal{H}$ and $p \in M$, let $\gamma_{p, u}$ be the trajectory of $\Sigma$ starting at $p$ for $t=0$ and corresponding to $u$. Then for $v \in \mathcal{H}, E_{p}(v)$ is given by $E_{p}(v):=\gamma_{p, v}(1)$.

Recall that $E_{p}(v)$ is defined for every $v \in \mathcal{H}$. The motion planning problem can be reformulated as follows:
For every $p, q \in M$, exhibit a control $u_{p, q} \in \mathcal{H}$ such that $E_{p}\left(u_{p, q}\right)=q$. Meaning, we must find the right inverse of the endpoint map $E_{p}$. This right inverse exists in a neighborhood of any point $u \in \mathcal{H}$ such that $D E_{p}$ is surjective (by the contrability assumption). More accurately, we begin with an arbitrary control $u_{0}$. Set $q_{0}:=E_{p}\left(u_{0}\right)$, and choose a path $\pi:[0,1] \rightarrow M$ such that $\pi(0):=q_{0}$ and $\pi(1):=q$. We now look for a path $\Pi:[0,1] \rightarrow \mathcal{H}$ such that, for every $s \in[0,1]$

$$
\begin{equation*}
E_{p}(\Pi(s))=\pi(s) \tag{2.6}
\end{equation*}
$$

Differentiating equation (2.6) yields

$$
\begin{equation*}
D E_{p}(\Pi(s)) \cdot \frac{d \Pi}{d s}(s)=\frac{d \pi}{d t}(s) . \tag{2.7}
\end{equation*}
$$

If $D E_{p}(\Pi(s))$ has full rank, then equation (2.7) can be solved for $\Pi(s)$ by taking $\Pi$ such that

$$
\begin{equation*}
\frac{d \Pi}{d s}(s)=F(\Pi(s)) \cdot \frac{d \pi}{d t}(s) . \tag{2.8}
\end{equation*}
$$

where $F(\cdot)$ is a right inverse of $D E_{p}(\cdot)$. For instance, we can choose $F(\cdot)$ to be the MoorePenrose pseudoinverse of $D E_{p}(\cdot)$. Equation (2.8) is called the Path Lifting Equation (PLE) which is an Ordinary Differential Equation (ODE) on $\mathcal{H}$ (cf. [Ważewski, 1947]).
Notice that, by construction, the control defined by $u^{\text {final }}:=\Pi(1)$ steers the system from $p$ to $q$. In order to get the value of $\Pi(1)$, it suffices, to solve the initial value problem defined in the control space $\mathcal{H}$

$$
\begin{align*}
\frac{d \Pi}{d s}(s) & =F(\Pi(s)) \cdot \frac{d \pi}{d t}(s)  \tag{2.9}\\
\Pi(0) & =u_{0} .
\end{align*}
$$

To successfully apply the Continuation Method to the motion planning problem, we must guarantee the existence of $F(\Pi(s))$ for every $s \in[0,1]$ in such a way equation (2.8) is always well defined, this is to say $D E_{p}(\Pi(s))$ must always have rank equal to 5 . Moreover, since we
need to evaluate $\Pi(1)$ to get a control steering the system from $p$ to $q$, the PLE defined in (2.8) must have a global solution on $[0,1]$.

In conclusion, the application of the Continuation Method to the motion planning problem is decomposed into two steps: to characterize (when possible) $S_{p}=\left\{u \in \mathcal{H}: \operatorname{rank}\left(D E_{p}(u)\right)<\right.$ $5\}$ and its image under $E_{p}$, and lifting the paths $\pi:[0,1] \rightarrow M$ avoiding $E_{p}\left(S_{p}\right)$ to paths $\Pi:[0,1] \rightarrow \mathcal{H}$ globally defined on $[0,1]$ by (2.8).

In section 1.3, we address the above two issues.

## CHAPTER 3

## Numerical implementation

### 1.3. Rolling surface and obstacles

In this section, we set the definition of an obstacle, and we provide the dynamics of the control system on $\mathcal{Q}\left(A_{1}, A_{2}\right)$ such that $A_{1}=\mathbb{R}^{2}$ and $A_{2}$ is an oriented surface of $\mathbb{R}^{3}$.

Since $A_{2}$ is a two-dimensional connected oriented Riemannian manifold, then $A_{2}$ can be embedded as a convex surface in $\mathbb{R}^{3}$ (see [Berger and Gostiaux, 2012]). Hence, hereafter, we assume that $A_{2}$ is an oriented surface of $\mathbb{R}^{3}$ with metrics induced by the Euclidean metric of $\mathbb{R}^{3}$. Moreover, we assume that $A_{2}$ is defined as one bounded connected component of the zero-level set of a smooth real-valued function $a: \mathbb{R}^{3} \rightarrow \mathbb{R}$. The normal vector field to $A_{2}$ is given by $\frac{\nabla a}{\|a\|}$, where $\nabla a=\left(a_{x}, a_{y}, a_{z}\right)$ denotes the gradient vector of $a$. The Gaussian curvature of $A_{2}$ is denoted by $K$.

Recall that the rolling-body problem assumes that the tangent spaces at the contact points are identified. In $\mathbb{R}^{3}$, this is equivalent to identify the normal vectors. Then at contact points, we assign $\frac{\nabla a}{\|a\|}$ to $-N_{\mathbb{R}^{2}}$, where $N_{\mathbb{R}^{2}}$ is the normal vector to the plane $\mathbb{R}^{2}$.

Using the fact that $\mathcal{Q}\left(\mathbb{R}^{2}, A_{2}\right)$ is a circle bundle when $A_{2}$ is a 2-dimensional manifold, and taking geodesic coordinates $\vartheta$ for $A_{2}$ at contact point $x_{2}$, and consider coordinates $x=$ ( $v_{1}, w_{1}, v_{2}, w_{2}, \psi$ ) in some neighborhood of $\left(0, \psi_{0}\right)$ in $\mathbb{R}^{4} \times S^{1}$, the control system (2.5) can be locally written as

$$
\begin{equation*}
\dot{x}=u_{1} X_{1}(x)+u_{2} X_{2}(x), \tag{3.1}
\end{equation*}
$$

with

$$
\begin{align*}
& X_{1}(x)=\left(1,0, \cos (\psi),-\frac{\sin (\psi)}{\vartheta},-\frac{\vartheta_{v_{2}}}{\vartheta} \sin (\psi)\right)^{T}  \tag{3.2}\\
& X_{2}(x)=\left(0,1,-\sin (\psi),-\frac{\cos (\psi)}{\vartheta},-\frac{\vartheta_{v_{2}}}{\vartheta} \cos (\psi)\right)^{T} . \tag{3.3}
\end{align*}
$$

For a detailed development of how obtaining the system (3.1) with the vector fields $X_{1}$ and $X_{2}$ defined by (3.2) and (3.3) respectively, see [Chelouah and Chitour, 2003].

### 1.1.3. Obstacles and dynamics of the control system

Hereafter, the estate space $\mathcal{Q}\left(\mathbb{R}^{2}, A_{2}\right)$ is simply denoted by $M$.
To deal with the problem of motion planning with obstacles, we follow a strategy suggested by E. Sontag in [Sontag, 1995].

Definition 1.3.1 We say that $\mathcal{W}$ is an obstacle of $\mathbb{R}^{2}$ if $\mathcal{W}$ is a nonempty compact subset of $\mathbb{R}^{2}$.

From the above definition, we can assume that there exists a function $\xi: \mathbb{R}^{2} \rightarrow \mathbb{R}$ such that $\mathcal{W}=\left\{w \in \mathbb{R}^{2} \mid \xi(w) \leq 0\right\}$.

An obstacle $\mathcal{W}$ maps in $M$ a region (see Figure 3.1), denoted by $C$, defined by

$$
C=\left\{x=\left(\gamma_{1}, \gamma_{2}, R\right) \in M \mid \gamma_{1} \in \mathcal{W}\right\}=\left\{x \in M \mid \xi\left(\gamma_{1}\right) \leq 0\right\} .
$$



Figure 3.1: Obstacle maps in $M$.
Hence, $C$ is an obstacle in the estate space $M$. This is to say, an obstacle in $M$ is a nonempty closed subset of $M$ such that $\widehat{M}=M \backslash C$ is also nonempty.

Consider an increasing sequence $\left(W_{i}\right)_{i \geq 0}$ of the compact subsets of $M$ such that $C=$ $\cup_{i \geq 0} W_{i}$. For every $i \geq 0$, we can construct a smooth function $\zeta_{i}: M \rightarrow \mathbb{R}$ such that $\zeta_{i} \equiv 0$ on $W_{i}$ and $0<\zeta_{i}(x) \leq \inf \left(1, d\left(\gamma_{1}, \mathcal{W}_{i}\right)\right)$ for $x \in M \backslash W_{i}$. Let

$$
M_{i}:=\sup _{x \in M}\left(\left\|D \zeta_{i}(x)\right\|+\left\|D^{2} \zeta_{i}(x)\right\|\right)
$$

Thus, we define

$$
\begin{equation*}
\zeta(x)=\sum_{i \geq 0} \frac{1}{2^{i} M_{i}} \zeta_{i}(x) . \tag{3.4}
\end{equation*}
$$

Then, there exists a smooth bounded function $\zeta: M \rightarrow \mathbb{R}$ such that $\zeta>0$ on $\widehat{M}, \zeta=0$ on $C$ and $\left(\widehat{M},\langle\langle\cdot, \cdot\rangle\rangle_{x}\right)$ is a complete Riemaniann manifold, where

$$
\begin{equation*}
\langle\langle\cdot, \cdot\rangle\rangle_{x}=\frac{\langle\cdot, \cdot\rangle}{\zeta^{2}} \tag{3.5}
\end{equation*}
$$

Let us consider the control system on $M$ defined by

$$
\begin{equation*}
\dot{y}=v_{1} \bar{X}_{1}(y)+v_{2} \bar{X}_{2}(y) \tag{3.6}
\end{equation*}
$$

where $\bar{X}_{i}=\zeta X_{i}, i=1,2$. Then $\widehat{M}$ is invariant under the control system (3.6), more accurately,
Lemma 1.3.1 $\forall v \in \mathcal{H}, \forall p \in \widehat{M}$, if $\bar{\gamma}_{v}$ is the solution of (3.6) with $\bar{\gamma}(0)=p$ then $\bar{\gamma}([0,1]) \subset$ $\widehat{M}$.
In other words, for every $t \in[0,1], \bar{\gamma}(t)$ belongs to the connected component of $\widehat{M}$ containing $p$. Thus, the motion planning problem with obstacles is reduced to a motion planning problem for each connected component of $\widehat{M}$. (For more details about $\zeta$ and the proof of Lemma 1.3.1 see [Chitour, 2006]).

We now describe how to apply the Continuation Method to solve the motion planning problem for the control system (3.6). First of all, it must be taken into account that $\mathcal{Q}\left(\mathbb{R}^{2}, A_{2}\right)=$ $M$ is simply equal to $\mathbb{R}^{2} \times T_{1} A_{2}$, where $T_{1} A_{2}$ is the unit tangent bundle of $A_{2}$ (more details see [Chelouah and Chitour, 2003] and [Bryant and Hsu, 1993]). Then, we can write for $z=(x, y) \in \mathbb{R}^{2} \times T_{1} A_{2}$, the vector fields $X_{1}$ and $X_{2}$ defined in (3.2) and (3.3) respectively, as

$$
\begin{aligned}
& X_{1}(z)=(1,0, f(y))^{T} \\
& X_{2}(z)=(0,1, h(y))^{T}
\end{aligned}
$$

where $f$ is the infinitesimal generator of the geodesic flow on $T_{1} A_{2}$ and $h$ is a vector field on $T_{1} A_{2}$ whose integral curves are also geodesics (cf. [Chelouah and Chitour, 2003]). Hence, the control system (3.6) can now be written as:

$$
\begin{align*}
\dot{v}_{1} & =\zeta(x) u_{1} \\
\dot{w}_{1} & =\zeta(x) u_{2}  \tag{3.7}\\
\dot{y} & =\zeta(x)\left(u_{1} f(y)+u_{2} h(y)\right)
\end{align*}
$$

with $x=\left(v_{1}, w_{1}\right) \in \mathbb{R}^{2} \backslash \bigcup_{i} \mathcal{W}_{i}$ and $y \in T_{1} A_{2}$.

### 2.1.3 . Conjeture

Since we want to implement the Continuation Method applied to the motion planning problem with obstacles, we must ensure that its application is successful, that is, $D E_{p}(u)$ must always have an equal rank to 5 , and the PLE defined in (2.8) must have a global solution on $[0,1]$.

A sufficient condition resolving the above is as follows (cf. [Chitour and Sussmann, 1998]). Let $S_{p}=\left\{u \in \mathcal{H}: \operatorname{rank}\left(D E_{p}(u)\right)<5\right\}$.

Condition 1.3.1 We say that a closed subset $\mathcal{K}$ of $\widehat{M}$ verifies this condition if

1. $\mathcal{K}$ is disjoint from $\overline{E_{p}\left(S_{p}\right)}$, where $\overline{E_{p}\left(S_{p}\right)}$ is the closure of $E_{p}\left(S_{p}\right)$;
2. there exists $c_{\mathcal{K}}>0$ such that, for every $u \in \mathcal{H}$ with $E_{p}(u) \in \mathcal{K}$, we have

$$
\begin{gathered}
\|F(u)\| \leq c_{\mathcal{K}}\|u\| \\
\text { where }\|F(u)\|=\left(\inf _{\|z\|=1} z^{T} D E_{p}(u) D E_{p}(u)^{T} z\right)^{-\frac{1}{2}}, \text { with } z \in T_{E_{p}(u)}^{*} \widehat{M} .
\end{gathered}
$$

In the case of a convex surface $A_{2}$ rolling on a plane without forbidden regions, it is shown in [Chelouah and Chitour, 2003] that if $A_{2}$ verifies Condition 1.3.2 (below), then there exists a compact subset $\mathcal{K}$ in the state space $M$ verifying the Condition 1.3.1, which is large enough to completely resolve the motion planning problem.

Let $d_{2}$ be the distance function associated to the Riemannian metric of $A_{2}$. The curve $\gamma: \mathbb{R}^{+} \rightarrow T_{1} A_{2}$ is a periodic geodesic of $T_{1} A_{2}$ if there exists $L \geq \frac{2 \pi}{\sqrt{K_{\max }}}$ such that $\gamma(t+L)=$ $\gamma(t)$ for all $t \geq 0$. We use $G$ to denote the closed subset of $T_{1} A_{2}, \gamma([0, L])$. For $r>0$, let $N_{r}(G)$ be the open set of points $y \in T_{1} A_{2}$ such that $d_{2}(y, G)<r$. Let $\phi(y, t)$ be the geodesic flow of $T_{1} A_{2}$.

Condition 1.3.2 ([Chelouah and Chitour, 2003]) We say that a surface $A_{2}$ verifies Condition 1.3.2 if there exists a geodesic curve $\gamma: \mathbb{R}^{+} \rightarrow T_{1} A_{2}, L>0$ and $\rho_{0}>0$ such that

1. $\gamma(t+L)=\gamma(t)$ for all $t \geq 0$;
2. For all $r<r_{0}, \exists \mu>0, \forall y_{0} \in N_{r}(G), \forall t \geq 0$ we have

$$
\phi\left(y_{0}, t\right) \in N_{\mu}(G) .
$$

Remark 1.3.1 The geometric property of having a stable periodic geodesic (and then satisfying Condition 1.3.2) is true for any convex compact surface having symmetry of revolution. It is generic within the convex compact surfaces verifying $\frac{K_{\min }}{K_{\max }}>\frac{1}{4}$, where $K_{\text {min }}$ and $K_{\text {max }}$ denote the minimum and the maximum of the Gaussian curvature over the surface, respectively (cf. [Klingenberg, 1982]).

Let $\bar{M}_{i}$ be a connected component of $\widehat{M}$ and let $S_{i}$ be an open line segment contained in the connected component of $\mathbb{R}^{2}$ which defines $\bar{M}_{i}$ for $i=1, \ldots, N$.

For $r \in\left(0, r_{0}\right)$, we define $\mathcal{K}_{r}$ as the complement in $\bar{M}_{i}$ of $S_{i} \times N_{r}(G)$. Since $\gamma$ is periodic $N_{r}(G)$ is diffeomorphic to the product of a small two-dimensional ball and a closed path in $T_{1} A_{2}$ (cf. [Chelouah and Chitour, 2003]). Therefore $\mathcal{K}_{r}$ is closed and arcwise-connected.

We conjecture the following.
Conjecture 1 With the previous notations, there exists $\bar{r} \in\left(0, r_{0}\right)$ such that for every $r \in$ $(0, \bar{r}), \mathcal{K}_{r}$ satisfies Condition 1.3.1.

If Conjecture 1 is true then an application of Gronwall lemma to the PLE (2.8) yields that, for every path $\pi:[0,1] \rightarrow \mathcal{K}_{r}$ of class $C^{1}$ and every control $\bar{u} \in \mathcal{H}$ such that $E_{p}(\bar{u})=\pi(0)$, the solution of the PLE defined in equation (2.8) with initial condition $\bar{u}$ exists globally on $[0,1]$.

In other words, the veracity of Conjecture 1 implies that the Continuation Method can be successfully applied to solving the rolling problem with obstacles.

An important fact for our purpose is that the singular controls in the case of convex surfaces rolling on the plane are exactly straight lines on the plane (see Proposition 5 in [Chelouah and Chitour, 2003]). We guess that, according to the Conjecture, the singular controls in each connected component of $\widehat{M}$ must also be straight lines (or segments of them).

In conclusion, from now on, $A_{2}$ is a strictly convex surface with symmetry of revolution, and the initial input $u_{0}$ is not a line on the plane. Finally, to delimit our proposal, we will consider a finite number of obstacles in the plane.

### 2.3. Key points for numerical implementation

In [Chitour, 2006] is proved that, for every compact subinterval $J=[0, a]$ of the interval of existence of the maximal solution of the equation (2.8), denoted by $I$, the numerical approximation of equation (2.8) defined by $E_{p, j}$ and associated to $\pi$ has a global solution on $J$ for a large enough $j$. In particular, if $J=I=[0,1]$ (i.e. the PLE has a global solution) the above result can be seen as a theoretical justification of the use of Galerkin's procedure in numerical implementations of the Continuation Method.

In this Section, we want to show how the Continuation Method, presented in Section 3.2 can be numerically implemented to solve the motion planning problem with obstacles when a convex surface of $\mathbb{R}^{3}$ rolls on the plane with forbidden regions (or obstacles).

In the following three subsections, we give details about the fundamental points for the numerical implementation, which are the discretization of the control space $\mathcal{H}$, the computation of $D E_{p}(u)$, and the lift of the curve $\widetilde{\gamma}_{1}$ on the convex body $A_{2}$. The ideas come from [Alouges et al., 2010].

### 1.2.3 . Discretizing $\mathcal{H}$

Recall that the control space $\mathcal{H}$ is, in general, an infinite-dimensional vector space, then we start by discretizing $\mathcal{H}$ because we need to solve the initial-value problem defined in (2.9). In this case, the controls are plane curves $\gamma_{1}:[0,1] \rightarrow \mathbb{R}^{2}$ such that

$$
\dot{\gamma}_{1}(t)=\zeta(x(t))\left(u_{1}, u_{2}\right), \quad x(t)=\left(\gamma_{1}(t), \gamma_{2}(t), R(t)\right)
$$

We divide the interval $[0,1]$ into $n-1$ parts and approximate the control space $\mathcal{H}$ by the $2 n$-dimensional subspace $\widetilde{\mathcal{H}}$ of piecewise-linear functions (more details see [Gautschi, 2012]). Then, $\gamma_{1}$ can be approximated by $\widetilde{\gamma}_{1}$, i.e., the linear interpolation of $\left(\gamma_{1}^{1}, \ldots, \gamma_{1}^{n}\right)$, where $\gamma_{1}^{i+1}=\gamma_{1}\left(\frac{i}{n-1}\right)=\left(x_{i+1}, y_{i+1}\right)^{T}$ for $i=0, \ldots, n-1$.

On each segment $\left[t_{i+1}, t_{i+2}\right]=\left[\frac{i}{n-1}, \frac{i+1}{n-1}\right]$ the approximate control $\left(\widetilde{u}_{1}^{i+1}, \widetilde{u}_{2}^{i+1}\right)^{T}$ is proporcional to the vector $\frac{1}{\zeta(x(t))}\left(x_{i+2}-x_{i+1}, y_{i+2}-y_{i+1}\right)^{T}$ for $i=0, \ldots, n-2$.

Notice that, when elements in $\widetilde{\mathcal{H}}$ are piecewise-linear functions with more than one piece, they are not singular controls. Then, the corresponding trajectories on $A_{2}$ are also easy to obtain by integrating some geodesic equations by using Proposition 2.2.1. Finally, Euler's method is used to integrate (2.9).

### 2.2.3 . Computing $D E_{p}(u)$

To evaluate $D E_{p}(u)$, for $u \in \mathcal{H}$, the following is necessary (see [Chelouah and Chitour, 2003]).
For $z \in T_{E_{p}(u)}^{*} M$, let $\lambda_{z, u}:[0,1] \rightarrow T^{*} M$ be the field of covectors along $\gamma_{p, u}$ such that it satisfies, in coordinates, the adjoint equation along $\gamma_{p, u}$ with terminal condition $z$, that is, $\lambda_{z, u}(1)=z$ and for a.e $t \in[0,1]$ we have

$$
\dot{\lambda}_{z, u}(t)=-\lambda_{z, u}(t) \cdot\left(u_{1}(t) D X_{1}\left(\lambda_{z, u}(t)\right)+u_{2}(t) D X_{2}\left(\lambda_{z, u}(t)\right)\right) .
$$

If $V$ is a $C^{\infty}$ vector field on $M$, the switching function $\varphi_{V, z, u}(t)$ associated to $V$ is the evaluation of $\lambda \cdot V(x)$, the Hamiltonian function of $V$ along $\left(\gamma_{p, u}, \lambda_{p, u}\right)$, i.e., for $t \in[0,1]$

$$
\varphi_{V, z, u}(t):=\lambda_{z, u} \cdot V\left(\gamma_{p, u}(t)\right) .
$$

Then $D E_{p}(u)$ can be computed thanks to the following formula: for $z \in T_{E_{p}(u)}^{*} M$ and $u, v \in \mathcal{H}$

$$
\begin{equation*}
z \cdot D E_{p}(u)(v)=\left(v, \varphi_{z, u}\right)_{\mathcal{H}}, \tag{3.8}
\end{equation*}
$$

where the switching-function vector $\varphi_{z, u}$ is the solution of the Cauchy problem, which is defined, in coordinates, by

$$
\begin{align*}
& \dot{\varphi}_{1}=-u_{2} K \varphi_{3} \\
& \dot{\varphi}_{2}=u_{1} K \varphi_{3} \\
& \dot{\varphi}_{3}=-u_{2} \varphi_{4}+u_{1} \varphi_{5}  \tag{3.9}\\
& \dot{\varphi}_{4}=-u_{2} K \varphi_{3} \\
& \dot{\varphi}_{5}=u_{1} K \varphi_{3}
\end{align*}
$$

with terminal condition $\varphi_{z, u}(1)=z$.
In practice, since the discrete $D E_{p}(u)$ is a $5 \times 5$ matrix and its image is given by (3.8), it suffices to take the five vectors of the canonical basis of $\mathbb{R}^{5}$, as final conditions $z$, and integrate (3.9) in reverse time. In the simulations, a fourth-order Runge-Kutta numerical scheme is used for integration, the scalar product $(\cdot, \cdot)_{\mathcal{H}}$ in control space $\mathcal{H}$ is evaluated by Gaussian quadrature, and the Gaussian curvature $K$ is computed by using the Proposition 1.2.1.

### 3.2.3 . Lifting the plane curve $\widetilde{\gamma}_{1}$ on $A_{2}$

Note that the curvature $K$ appearing in (3.9) is taken at the final contact point on the surface $A_{2}$ after it has rolled along the piecewise-constant curve $\widetilde{\gamma}_{1}$. Thus, to locate the final point, we need to lift the plane curve $\widetilde{\gamma}_{1}$ on $A_{2}$, and the lifting dynamics are given by (3.7). However, since the geodesic coordinates involved in (3.7) are not given explicitly in practice, the numerical lifting method is based on Proposition 2.2.1.

On each interval $\left[t_{i+1}, t_{i+2}\right]$, the approximate control curve $\widetilde{\gamma}_{1}$ is a straight line (i.e., a geodesic in $\mathbb{R}^{2}$ ), and then, by Proposition 2.2.1, the lifting curve $\widetilde{\gamma}_{1}$ on $A_{2}$ is also a geodesic on each interval $\left[t_{i+1}, t_{i+2}\right]$ for all $i=0, \ldots, n-2$. Then, from the initial contact point $x_{0}$ on $A_{2}$, we can integrate successively (with fourth-order Runge-Kutta scheme) the geodesic equation, given by Proposition 1.2.2 (equation 2.2), on each $\left[t_{i+1}, t_{i+2}\right]$ with initial conditions equal to $\widetilde{\gamma}_{1}\left(t_{i+1}\right)$ and $\left(\widetilde{u}_{1}^{i+1}, \widetilde{u}_{2}^{i+1}\right)$, for $i=0, \ldots n-2$.

A very important difficulty is that the numerical integration is performed on the manifold $A_{2}, t>0$. Assuming that we are at point $x \in A_{2}$ at time $t$, then, at time $t+\delta t$, we move to $x_{\text {new }}=x+(\delta t) e$, with $e \in T_{x} A_{2}$, but $x_{\text {new }}$ does not belong to $A_{2}$ if $e$ is nonzero. Therefore, at each integration step, we have to project $x_{\text {new }}$ on $A_{2}$. In details, let us assume that the point $(0,0,0)$ is inside the convex body $A_{2}$. Since $A_{2}$ is defined as (a bounded connected component of) the zero-level set of a smooth function $a$, we assume that $\left|a\left(x_{\text {new }}\right)\right| \leq \varepsilon$ for some $\varepsilon \ll 1$, i.e., $x_{\text {new }}$ is close to $A_{2}$. Then, by the convexity of $A_{2}$, there exists a unique real number $\mu$ close to 1 such that $a\left(\mu x_{\text {new }}\right)=0$. The projection issue to be addressed is a local one, and therefore, Newton's method is efficient to find $\mu$. The derivative with respect to $\mu$ is also needed, which is evaluated by a finite-difference scheme.

### 3.3. Examples

This section provides three examples of the rolling bodies on the Euclidean plane: the sphere, the flattened ball, and an egg. The simulations were developed in Matlab software, and the code was run on a Macbook Pro Apple M1 8Go of RAM for 100 iterations with $n-1=100$ for the discretization of control space.

To illustrate how the convex body rolls along the curve obtained by the application of the Continuation Method, see Figure 3.2.


Figure 3.2: Convex body rolling along the curve obtained by the application of the Continuation Method (blue curve).

Each convex body $A_{2}$ (the sphere, the flattened ball, and an egg) is defined by a function $a(x, y, z)$ such that $A_{2}=a^{-1}(0)$. The initial points of contact are blue and the end contact points are magenta. The initial and end orientations have the same color as the initial and end points. Obstacles are circular regions in the plane, and they are colored red (see Figures 3.3, 3.5 and 3.7).

The initial chosen curve to apply the Continuation Method is the green curve, and the curve obtained by this process is colored blue. In each example, three different curves are chosen to evaluate the method. Finally, the red curve on the convex body is obtained by the lifting of the curve computed by the Continuation Method (blue curve) (see Figures 3.4, 3.6 and 3.8).

### 1.3.3. Sphere

The sphere is defined by the zero-level set of the function $a$, with $a(x, y, z)=x^{2}+y^{2}+z^{2}-1$.


Figure 3.3: Left side: sphere with inital and end contact points and its orientations. Right side: obstacles (in red) and inital and end contact points on the plane.


Figure 3.4: Left side: curve $\gamma_{2}$ on the sphere. Right side: Curve obtained by applying the Continuation Method (blue) to inital chosen curve (green).

### 2.3.3 . Egg

This egg is defined by the zero-level set of the function $a$, defined by $a(x, y, z)=\frac{x^{2}+y^{2}}{1-0.4 z}+\frac{1}{4} z^{2}-1$


Figure 3.5: Left side: egg with inital and end contact points and its orientations. Right side: obstacles (in red) and inital and end contact points on the plane.


Figure 3.6: Left side: curve $\gamma_{2}$ on the egg. Right side: Curve obtained by applying the Continuation Method (blue) to inital chosen curve (green).

### 3.3.3 . Flattened ball

The flattened ball $A_{2}$ is defined by the function $a$ such that $A_{2}=a^{-1}(0)$, with $a(x, y, z)=$ $x^{2}+y^{2}+5 z^{2}-1$.

(a)

(b)

Figure 3.7: Left side: flattened ball with inital and end contact points and its orientations. Right side: obstacles (in red) and inital and end contact points on the plane.


Figure 3.8: Left side: curve $\gamma_{2}$ on the flattened ball. Right side: curve obtained by applying the Continuation Method (blue) to inital chosen curve (green).

Finally, to show how well the implemented method works, we have considered obstacles with more complexity and the same convex bodies. The initial chosen curve is colored green, and the one obtained by the Continuation Method is colored blue. The obstacles are colored green. See Figure 3.9.


Figure 3.9: Left side: convex body with initial and end contact points with orientations and obstacle on the plane (in green). Right side: curve $\gamma_{2}$ on the convex body and curve obtained by applying the Continuation Method (blue) to the initial chosen curve (green).

### 4.3. Conclusions and comments

This part of the thesis presents a numerical implementation of the Continuation Method to solve the motion planning problem when a convex body rolls on the Euclidean plane with obstacles, detailing the discretization of the control space $\mathcal{H}$, the evaluation of $D E_{p}(u)$, and the lift of the curve $\widetilde{\gamma}_{1}$ on the convex body $A_{2}$. The convex body is considered as an oriented surface of $\mathbb{R}^{3}$ and hence it is defined by a function $a: \mathbb{R}^{3} \rightarrow \mathbb{R}$ such that $A_{2}=a^{-1}(0)$. In addition, we conjecture the existence of a closed subset $\mathcal{K}_{r}$ in the state space $\widetilde{M}$ that guarantees the successful application of the Continuation Method if we consider that $A_{2}$ has a stable periodic geodesic. This property is true for any convex compact surface having symmetry of revolution (cf. [Klingenberg, 1982]). For this reason, the sphere, the flattened ball, and an egg have been considered for simulations. However, when we relax this geometric property, the method implemented still works. The above can be seen as the robustness of the method. Therefore, there are two challenges: 1) To prove that Conjecture 1 is true and 2 ) to research
a theoretical convergence result that ensures the successful application of the Continuation Method when the convex body does not satisfy the geometrical condition of symmetry of revolution. We believe that the above is related to the fact that $\nabla a(x, y, z) \neq 0$ on the zerolevel set of $a$, because, according to Propositions 1.2.1 and 1.2.2, the Gaussian curvature of the surface $A_{2}$ and the lift of the plane curve on $A_{2}$ respectively, depend on $\nabla a(x, y, z)$, indeed, if $\nabla a(x, y, z) \neq 0$ the equations (2.1) and (2.2) that define the Gaussian curvature and the lift respectively, are always well defined. For instance, let us consider $A_{2}$ an oriented surface without symmetric axis, defined by $a(x, y, z)=\frac{x^{2}}{1-0.7 y}+\frac{3 y^{2}}{1-0.1 z}+\frac{0.3 z^{2}}{1-0.3 x-0.1 y}-1$ and hence $\nabla a(x, y, z)=\left(\begin{array}{c}\frac{9 z^{2}}{(3 x+y-10)^{2}}-\frac{20 x}{7 y-10} \\ \frac{3 z^{2}}{(3 x+y-10)^{2}}-\frac{60 y}{z-10}+\frac{70 x^{2}}{(7 y-10)^{2}} \\ \frac{3 y^{2}}{(z-0.1)^{2}}-\frac{3 z}{0.75 x+0.5 y-5}\end{array}\right) \neq 0$ on the zero-level set of $a$. In Figure 3.10, we can see how well the method works.


Figure 3.10: Left side: Non-symmetric convex body with initial and end contact points and curve $\gamma_{2}$. Right side: Curve obtained by applying the Continuation Method (blue) to the initial chosen curve (green).

Furthermore, if we consider a more complex obstacle, its performance is still good. Figure 3.11 illustrates its performance. Finally, we are interested in adapting this proposal in the motion planning problem in other situations of the rolling body problem, which are to consider a convex body $A_{2}$ rolling on a plane with forbidden regions in $A_{2}$, and a strictly convex body rolling on the other without or with forbidden regions.


Figure 3.11: Top: Non-symmetric convex body with initial and end contact points with orientations and obstacle on the plane (in green). Bottom: curve $\gamma_{2}$ on the convex body and curve obtained by applying the Continuation Method (blue) to the initial chosen curve (green).

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## APPENDIX A

## Synthèse en français

## Introduction générale

La géométrie sub-riemannienne est une généralisation de la géométrie riemannienne (géométrie qui est née avec les travaux de Bernhard Riemann dans le but de généraliser certains résultats de Gauss concernant la courbure). Son origine se trouve dans l'article [Strichartz, 1986]. Cependant, certains éléments, exemples ou applications de la géométrie sub-riemannienne peuvent être situés avant, par exemple, dans des travaux relatifs aux problèmes de contrôle optimal [Brockett, 1982], de thermodynamique [Carathéodory, 1909], ou encore dans des articles relatifs à la géométrie riemannienne [Hermann, 1973] et à l'étude de la géométrie du groupe d'Heisenberg, qui est un exemple célèbre de géométrie sub-riemannienne [Gaveau, 1977]. Cette géométrie existe grâce à l'idée de contraintes non-intégrables ; c'est à dire une contrainte sur la direction admissible des mouvements, et ainsi, elle a reçu une attention croissante dans de nombreuses disciplines telles que la théorie du contrôle, la robotique, la mécanique classique, l'analyse de l'opérateur hypoelliptique, le collecteur de diffusion, et même dans d'autres branches de la géométrie comme la géométrie de Cauchy-Riemann.

Une structure sous-riemannienne particulière est ce que l'on appelle les structures presque riemanniennes. Formellement, une structure presque riemannienne à $n$-dimensions (ARS en abrégé) est une structure sub-riemannienne à rang variable qui peut être définie localement par un ensemble de $n$ champs vectoriels lisses sur une varieté à $n$ dimensions, satisfaisant la condition de rang de l'algèbre de Lie. L'ensemble des points où la dimension de l'étendue linéaire des champs de vecteurs n'est pas complète est appelé le lieu singulier (ou l'ensemble singulier) et est noté $\mathcal{Z}$. Les modèles attractifs de ARS peuvent être décrits sur des groupes de Lie en utilisant des champs vectoriels invariants et linéaires. Ils sont appelés ARS linéaires (ou simples) sur les groupes de Lie (voir [Ayala and Jouan, 2016]).

L'objectif de cette thèse est d'étudier deux sujets en géométrie sub-riemannienne. D'une part, l'approximation locale d'une structure presque riemannienne aux points singuliers, où l'approximation nilpotente a perdu la structure originale, et d'autre part, le système cinéma-
tique d'un varieté roulant sur un autre varieté sans tourner ni glissement, en particulier une implémentation numérique de la Méthode de Continuation lorsqu'un varieté à 2-dimensions roule sur le plan euclidien avec des régions interdites.

## Approximation locale par systèmes linéaires et structures quasi-riemanniennes sur les groupes de Lie

Dans cette première partie, nous nous sommes consacrés à l'approximation locale des structures quasi-riemanniennes aux points singuliers par des ARS sur les groupes de Lie et à montrer que cette approximation est généralement meilleure que celle du nilpotent.

Soit $M$ une varieté différentiel de dimension $n$ et considérons $\left\{X_{1}, X_{2}, \ldots, X_{n}\right\}$ un ensemble de champs vectoriels lisses sur $M$. Localement, un ARS sur $M$ peut être défini par $\left\{X_{1}, X_{2}, \ldots, X_{n}\right\}$ satisfaisant la condition de rang de l'algèbre de Lie (Larc en abrégé). Cet ensemble de champs vectoriels est considéré comme un cadre orthonormé. Nous désignons par $\Delta_{p}$ l'étendue linéaire des champs de vecteurs $\left\{X_{1}, X_{2}, \ldots, X_{n}\right\}$ au point $p$. Rappelons que $\mathcal{Z}=\left\{p \in M: \operatorname{rank}\left(\Delta_{p}\right)<n\right\}$. Si $\mathcal{Z}$ est vide, alors la structure presque riemannienne est une structure riemannienne (plus de détails dans [Agrachev et al., 2019]). Dans le cas générique à 3 de dimension, l'ensemble singulier est formé par deux types de points (voir Figure A.1): les points de type- 1 où $\Delta_{p}$ a une dimension 2 et est transversal à $\mathcal{Z}$, et les points de type- 2 où $\Delta_{p}$ a une dimension 2 et est tangent à $\mathcal{Z}$. De plus, les points de type 2 sont isolés (pour plus de détails, voir [Boscain et al., 2015]). Dans le cas générique de dimension 2, il existe également des points où $\Delta_{p}$ est transversal à $\mathcal{Z}$ et isolé. De tels points sont appelés points de tangence dans [Agrachev et al., 2008].


Figure A.1: Deux points différents forment $\mathcal{Z}$ dans le cas générique 3D.

D'autre part, les approximations nilpotentes sont utilisées pour étudier localement le comportement des structures quasi riemanniennes en raison de leur grande similitude avec la dynamique originale. Cependant, il existe des cas où l'approximation nilpotente d'une ARS s'avère être une structure sous-riemannienne de rang constant. En d'autres termes, certains champs vectoriels peuvent disparaître. C'est précisément ce qui se passe dans le cas générique à 3 de dimension, traité dans [Boscain et al., 2015], où aux points de type 2 (ou points de tangence), l'approximation nilpotente est la structure sous-riemannienne de Heisenberg et n'est donc pas une 3 -ARS.

L'approximation résoluble a été introduite dans [Jouan and Manríquez, 2022] pour récupérer la structure presque riemannienne perdue dans l'approximation nilpotente. Dans cet article, on a considéré le cas où un seul des champs vectoriels disparaît et les autres sont indépendants. L'approximation solvable est alors une approximation locale d'un ARS aux points singuliers où l'approximation nilpotente n'est plus un ARS mais une structure sous-riemannienne de rang constant. Une généralisation de cette approximation est donnée dans [Manríquez et al., 2022], où une description complète des approximations nilpotentes et solvables est abordée, y compris les approximations génériques.

Nous pouvons identifier, principalement, deux axes de travail dans la recherche, liés essentiellement aux structures des systèmes d'approximation et à l'estimation des différentes distances définies par la structure originale, l'approximation nilpotente et celle résoluble. En ce qui concerne les structures, l'algèbre de Lie générée par cette nouvelle famille de champs vectoriels est de dimension finie et résoluble (dans le cas générique). Grâce au théorème d'équivalence de [Jouan, 2010] nous savons que l'espace $\mathbb{R}^{n}$ est difféomorphe à un certain espace homogène ou groupes de Lie. Grâce à ce difféomorphisme, l'approximation résoluble et l'approximation nilpotente sont équivalentes à une ARS linéaire sur un espace homogène ou un groupe de Lie. De plus, nous pouvons trouver des structures dures ou complexes et de nombreuses structures différentes. Pour cette raison, nous avons traité les ARS génériques en déterminant les distributions génériques sur une varieté connecté de dimension $n$ et les avons utilisées pour exposer les approximations nilpotentes et résoluble génériques. En ce qui concerne l'axe d'estimation de la distance, nous avons établi l'ordre des approximations de la distance originale par celle définie par l'approximation résoluble, en obtenant que l'approximation résoluble améliore l'ordre d'approximation de la distance originale par l'approximation nilpotente. Soit $d, \widetilde{d}, \widehat{d}$ les distances induites par la structure originale, l'approximation résoluble, et l'approximation nilpotente. Nous montrons que l'approximation par $\widetilde{d}$ est strictement meilleure que celle par $\widehat{d}$ pour une paire de points translatés dans une direction où la distance $\widetilde{d}$ est décroissante.

## Méthode de continuation dans un problème de roulement avec obstacles.

Dans la deuxième partie de cette thèse, nous abordons le problème de la planification du mouvement d'un corps strictement convexe $A_{2}$ roulant (sans glisser ni tourner) sur le plan euclidien $\mathbb{R}^{2}$ avec obstacles, avec une implémentation numérique de la Méthode des Continuités.

Un système nonholonomique sur une varieté $n$-dimensionnel $M$ est un système de contrôle qui est de la forme

$$
\dot{x}=\sum_{i=1}^{m} u_{i} X_{i}(x), \quad x \in M,
$$

où $m>1$ est un entier et $X_{1}, X_{2}, \ldots, X_{m}$ sont des champs de vecteurs $C^{\infty}$ sur $M$ (cf. [Jean, 2014]). Ces systèmes ont attiré l'attention de nombreux auteurs de différentes disciplines pour leurs applications variées, principalement en robotique (cf. [Murray et al., 2017] et ses références). Le problème du corps roulant (sans glissement ni rotation) d'une varieté riemannien bidimensionnel sur un autre (qui est un excellent exemple de la fusion entre la géométrie subriemannienne et la théorie du contrôle (géométrique)), peut être écrit comme un système
nonholonomique. Plus précisément, le problème du corps roulant (sans glissement ni rotation) est un système de contrôle $\Sigma$, qui modélise le roulement d'une surface connectée encastrée $A_{2}$ dans $\mathbb{R}^{3}$ sur une autre surface $A_{1}$. En conséquence des contraintes de roulement, et étant donné une courbe absolument continue (a.c.) $\gamma_{1}$ sur $A_{1}$, il existe une unique courbe a.c. $\Gamma$ dans l'espace d'état, qui décrit le roulement de la surface $A_{2}$ sur la surface $A_{1}$ le long de la courbe $\gamma_{1}$. Ainsi, les commandes admissibles de $\Sigma$ correspondent aux courbes a.c. $\gamma_{1}$ de $A_{1}$ par leurs dérivées $\dot{\gamma}_{1}$. Alors, le système de contrôle peut être écrit, en coordonnées locales, comme un système nonholonomique

$$
\Sigma: \quad \dot{x}=u_{1} X_{1}+u_{2} X_{2},
$$

où $\left(u_{1}, u_{2}\right) \in \mathbb{R}^{2}$ est la commande, et $X_{1}$ et $X_{2}$ sont des champs vectoriels.
La Méthode de Continuation, a été introduite dans [Sussmann, 1992] et [Sussmann, 1993] et largement développée dans [Chitour and Sussmann, 1998], [Chitour, 2002], et [Chitour, 2006]. Il est utilisé pour résoudre des équations non linéaires de la forme $\mathcal{L}(x)=y$, où $x$ est l'inconnue, et $\mathcal{L}: X \rightarrow Y$ est surjectif. Dans le contexte de la planification du mouvement, $\mathcal{L}$ est la carte du point final (associée à un certain point fixe $p$ ) de l'espace des entrées admissibles à l'espace d'état, c'est-à-dire $E_{p}: \mathcal{H} \rightarrow M$.

On considère le corps convexe $A_{2}$, qui peut être intégré comme une surface convexe dans $\mathbb{R}^{3}$, avec une géodésique périodique stable et il est défini par la fonction $a(x, y, z)$ telle que $A_{2}=a^{-1}(0)$. Nous désignons l'espace des domaines simplement par $M$. Un obstacle $\mathcal{W}$ dans $\mathbb{R}^{2}$ est un sous-ensemble compact non vide de $\mathbb{R}^{2}$. Un obstacle $\mathcal{W}$ fait correspondre dans $M$ une région $C$. Ainsi, un obstacle dans $M$ est un sous-ensemble fermé non vide de $M$ tel que $\widehat{M}=M \backslash C$ est également non vide. Considérons le système de contrôle sur $M$ défini par

$$
\dot{y}=v_{1} \bar{X}_{1}(y)+v_{2} \bar{X}_{2}(y)
$$

où $\bar{X}_{i}=\zeta X_{i}, i=1,2$ et $\zeta: M \rightarrow \mathbb{R}$ tel que $\zeta>0$ sur $\widehat{M}, \zeta=0$ sur $C$. Alors $\widehat{M}$ est invariant sous le système de contrôle ci-dessus. Ainsi, le problème de planification du mouvement avec obstacles est réduit à un problème de planification du mouvement pour chaque composant connecté de $\widehat{M}$.

Par conséquent, nous fournissons une implémentation numérique complète de la méthode de continuation présentée ci-dessus pour résoudre le problème de planification de mouvement avec des régions interdites sur le plan. Nous détaillons les points fondamentaux de l'implémentation numérique, qui sont la discrétisation de l'espace de contrôle $\mathcal{H}$, le calcul de $D E_{p}(u)$, et la levée de la courbe $\widetilde{\gamma}_{1}$ sur le corps convexe $A_{2}$. Cette partie fournit trois exemples de corps roulant sur le plan euclidien (avec obstacles) : la sphère, la boule aplatie, et un œuf, c'est-à-dire avec une géodésique périodique stable (voir Figure A.2). Cependant, lorsque nous relâchons cette propriété géométrique, la méthode mise en œuvre fonctionne toujours.

## 1.A Approximation locale par des systèmes linéaires et structures presqueRiemanniennes

Comme nous l'avons dit dans l'introduction précédente, une structure presque-Riemannienne (ARS en abrégé) sur une varieté différentiel de dimension $n$ est une structure sous-Riemannienne


Figure A.2: Différents exemples de corps roulant sur le plan avec des obstacles.
à rang variable qui peut être définie, au moins localement, par un ensemble de champs vectoriels $n$ satisfaisant la condition de rang de l'algèbre de Lie (Larc en abrégé). Nous désignons par $\Delta_{p}$ l'étendue linéaire des champs vectoriels au point $p$. L'ensemble des points où $\operatorname{dim}\left(\Delta_{p}\right)<n$ est appelé locus singulier ou ensemble singulier et noté $\mathcal{Z}$. De nombreux articles dédiés à l'étude des ARSs peuvent être trouvés dans la littérature, par exemple: [Agrachev et al., 2010], [Bonnard et al., 2009], [Bonnard et al., 2011], [Boscain et al., 2013a] [Boscain et al., 2013b].

Dans le cas générique tridimensionnel, qui nous intéresse particulièrement, l'ensemble singulier est un sous-varieté intégré de codimension un et les points où $\Delta_{p}=T_{p} \mathcal{Z}$ sont isolés. De tels points sont appelés points de tangence dans [Agrachev et al., 2008] et points de type 2 dans [Boscain et al., 2015].

Nous nous intéressons également aux ARS linéaires sur les groupes de Lie (ou les espaces homogènes) (voir [Ayala and Jouan, 2016]) car elles seront utilisées comme structures d'approximation pour les ARS générales.

D'autre part, les approximations nilpotentes sont utilisées pour étudier localement le comportement des structures presque riemanniennes en raison de leur grande similitude avec la dynamique originale. Cependant, dans certains cas, l'approximation nilpotente d'une ARS dégénère, car il ne s'agit plus d'une ARS mais d'une structure sous-riemannienne de rang constant. En d'autres termes, il peut arriver que certains des champs vectoriels de l'approximation nilpotente disparaissent, transformant la structure presque riemannienne en une structure sousriemannienne de rang constant. Par exemple, si

$$
X_{1}=\left(\begin{array}{l}
1 \\
0 \\
0
\end{array}\right), \quad X_{2}=\left(\begin{array}{l}
0 \\
1 \\
x
\end{array}\right) \quad \text { and } \quad X_{3}=\left(\begin{array}{c}
0 \\
0 \\
x^{2}
\end{array}\right)
$$

alors son approximation nilpotente est

$$
\widehat{X}_{1}=\left(\begin{array}{l}
1 \\
0 \\
0
\end{array}\right), \quad \widehat{X}_{2}=\left(\begin{array}{l}
0 \\
1 \\
x
\end{array}\right) \quad \text { and } \quad \widehat{X}_{3}=0 .
$$

c'est-à-dire la structure sub-riemannienne de Heisenberg. C'est exactement ce qui se passe dans le cas générique à 3 dimensions, traité dans [Boscain et al., 2015].

Nous supposons, sans perte de généralité, que les derniers champs de vecteurs $m \geq 1$ disparaissent dans l'approximation nilpotente. Notre objectif consiste à récupérer la structure quasi riemannienne perdue dans l'approximation nilpotente grâce aux champs de vecteurs, notés $\widetilde{X}_{i}$ pour $i=n-m, \ldots, n$, qui sont la composante homogène de degré 0 du développement de Taylor en coordonnées privilégiées des champs de vecteurs qui disparaissent. La nouvelle famille de champs de vecteurs composée par l'approximation nilpotente et $\widetilde{X}_{n-m}, \widetilde{X}_{n-m+1}, \ldots, \widetilde{X}_{n}$ est appelée approximation résoluble. L'algèbre de Lie générée par cette nouvelle famille de champs vectoriels est de dimension finie. De plus, lorsqu'un seul des champs vectoriels disparait et que les autres sont indépendants, cette algèbre de Lie est solvable (d'où son nom). Nous nous intéressons également à un groupe de Lie nilpotent sur lequel $\widetilde{X}_{n-m}, \widetilde{X}_{n-m+1}, \ldots, \widetilde{X}_{n}$ agissent comme des champs de vecteurs linéaires.

Cette première partie est organisée comme suit.
Le chapitre 2 contient des généralités sur les ARS, l'ordre nonholonomique, les coordonnées privilégiées, l'approximation nilpotente, les champs vectoriels linéaires et les ARS linéaires sur les groupes de Lie ou les espaces homogènes.

Dans le Chapitre 3, dans la première section, nous introduisons la définition d'une approximation résoluble en considérant $m=1$; c'est-à-dire qu'un seul des champs vectoriels disparaît et les autres sont indépendants; nous analysons ses structures algébriques et la distance définie par l'approximation résoluble ( $\widetilde{d})$, en concluant que l'algèbre de Lie générée par cette nouvelle famille de champs vectoriels est de dimension finie et résoluble (Proposition 1.3.1), I'approximation résoluble est équivalente à une ARS simple sur un espace homogène ou un groupe de Lie (Théorème 1.3.1), et la distance $\widetilde{d}$ satisfait toujours $\widetilde{d} \leq \widehat{d}$ (Proposition 1.3.3), où $\widehat{d}$ est la distance induite par l'approximation nilpotente.

Dans la deuxième section, nous avons abordé le cas générique tridimensionnel. Ici, nous montrons que l'approximation résoluble est un ARS simple sur $\mathbb{R}^{3}$ difféomorphe à un quotient du groupe de Heisenberg à 5 dimension $\mathbb{H}^{2}$. Concernant la question de la distance, nous avons obtenu dans certains cas que l'ordre de l'approximation de $d$ (la distance originale) par $\widetilde{d}$ est meilleur que l'ordre de l'approximation de $d$ par $\widehat{d}$ (Théorème 3.4.1). En utilisant le fait que pour certaines paires $\left(q, q^{\prime}\right)$ de points translatés du lieu singulier la distance est décroissante (Théorème 2.3.2) et l'ordre d'approximation par $\widetilde{d}$, nous prouvons que la différence $\left|d\left(q, q^{\prime}\right)-\widetilde{d}\left(q, q^{\prime}\right)\right|$ est strictement plus petite que $\left|d\left(q, q^{\prime}\right)-\widehat{d}(q, q)\right|$. Enfin, nous donnons l'hamiltonien associé au flux défini par l'approximation résoluble dans le cas générique 3D et nous calculons les géodésiques avec la condition initiale $x(0)=y(0)=z(0)=0$ and $p(0)=\cos (\theta), q(0)=\sin (\theta), r(0)=r$ dans un cas particulier.

Le chapitre 4 est consacré aux approximations générales nilpotentes et solvables des structures presque-Riemanniennes. Tout d'abord, nous montrons qu'il est toujours possible de définir l'ARS localement, autour du point $p=0$ en coordonnées locales privilégiées, par un ensemble de $n$ champs de vecteurs orthonormés $X_{1}, \ldots, X_{n}$ tel que l'approximation résoluble

$$
\widehat{X}_{1}, \ldots, \widehat{X}_{k}, \widehat{X}_{k+1}, \ldots, \widehat{X}_{m}, \widetilde{X}_{m+1}, \ldots, \widetilde{X}_{n}
$$

satisfait à

- $\widehat{X}_{i}(0) \neq 0$ pour $i=1, \ldots, k$;
- $\widehat{X}_{i} \neq 0$ mais $\widehat{X}_{i}(0)=0$ pour $i=k+1, \ldots, m$;
- $\widehat{X}_{i}=0$ pour $i=m+1, \ldots, n$.

Le premier objectif est ici de prouver que l'approximation nilpotente ou résoluble d'un ARS en un point singulier est une structure linéaire presque Riemannienne sur un groupe de Lie ou un espace homogène (Théorème 1.4.2 et 1.4.4 respectivement), sauf dans certains cas très dégénérés où ni l'approximation nilpotente ni l'approximation résoluble ne définissent un ARS.

Lorsque nous travaillons avec $n$-ARS, nous pouvons trouver des structures difficiles ou complexes et de nombreuses structures différentes. Pour cette raison, le deuxième objectif est de déterminer les structures génériques. On montre en particulier que génériquement: (1) l'ensemble singulier $\mathcal{Z}$ est une union de sous-variété $\mathcal{Z}_{r}$ de codimension $r^{2}$ où le rang est $n-r$; (2) le rang de $\Delta+[\Delta, \Delta]$ est partout plein ( $\Delta$ représente la distribution) (Théorème 2.4.1). La structure des points de $\mathcal{Z}_{r}$ où $\operatorname{dim}\left(T_{p} \mathcal{Z}_{r}\right)+\operatorname{dim}\left(\Delta_{p}\right)$ n'est pas maximale est décrite dans le Théorème 2.4.2. Par exemple, dans $\mathcal{Z}_{1}$, ces points sont les soi-disant points de tangence (voir [Boscain et al., 2015]), c'est-à-dire les points où $T_{p} \mathcal{Z}_{1}=\Delta_{p}$. Ils sont génériquement isolés dans $\mathcal{Z}_{1}$.

Grâce à ces résultats de généricité et à l'aide des formes normales locales (voir la section 2.2.4), on montre finalement que génériquement il n'y a que deux possibilités pour l'approximation nilpotente/résoluble en un point $p \in \mathcal{Z}$ (Théorème 2.4.3):

1. En un point de tangence $p$ dans $\mathcal{Z}_{1}$, un champ de vecteurs de l'approximation nilpotente disparaît, mais l'approximation résoluble n'est pas dégénérée et définit une ARS linéaire.
2. En tous les autres points, c'est-à-dire les points non tangents de $\mathcal{Z}_{1}$ et tous les points dans $\mathcal{Z}_{r}$ avec $r \geq 2$, l'approximation nilpotente n'est pas dégénérée.

En conclusion, les seuls points generic où l'approximation résoluble est utile sont les points de tangence dans $\mathcal{Z}_{1}$.

Dans la dernière section, nous traitons de la distance induite par l'approximation résoluble, aux points de tangence, d'un ARS $n$-dimensionnel en considérant des hypothèses génériques et la forme normale lorsque le point appartenant à l'ensemble singulier est un point tangent.

## 2.A . Méthode de continuation dans un problème de roulement avec obstacles.

Dans cette partie de la thèse, nous abordons le problème de la planification du mouvement d'un corps $A_{2}$ strictement convexe roulant (sans glisser ni tourner) sur le plan euclidien $R^{2}$ avec obstacles, avec une implémentation numérique de la Méthode de Continuation.

Nous commençons par décrire le problème du corps roulant, puis la méthode de continuation.

## Le problème du corps roulant

Considérons le problème du corps roulant sans glissement ni rotation de $A_{2}$ sur le dessus de $A_{1}$. Nous adoptons le point de vue présenté dans [Agrachev and Sachkov, 2004].

Soit $A_{1}$ et $A_{2}$ des variétés riemanniennes bidimensionnelles connectées et orientées (surfaces des corps roulants). Aux points de contact des corps $x_{1} \in A_{1}$ et $x_{2} \in A_{2}$, leurs espaces tangents sont identifíés par une isométrie préservant l'orientation $q: T_{x_{2}} A_{2} \rightarrow T_{x_{1}} A_{1}$ (voir figure A.3). Une telle isométrie $q$ est un état du système, et l'espace d'état est donné par

$$
\mathcal{Q}\left(A_{1}, A_{2}\right)=\left\{q: T_{x_{2}} A_{2} \rightarrow T_{x_{1}} A_{1} \mid x_{1} \in A_{1}, x_{2} \in A_{2}, q \text { un isométrie }\right\} .
$$



Figure A.3: Idétification des espaces tangents au point de contact.
Comme l'ensemble de toutes les isométries préservant l'orientation dans $\mathbb{R}^{2}$ est $S O(2)$, qui peut être identifié avec le cercle unitaire $S^{1}$ dans $\mathbb{R}^{2}, \mathcal{Q}\left(A_{1}, A_{2}\right)$ est une variété. Un point $q \in \mathcal{Q}\left(A_{1}, A_{2}\right)$ est localement paramétré par $\left(x_{1}, x_{2}, R\right)$ avec $x_{1} \in A_{1}, x_{2} \in A_{2}$, et $R \in S O(2)$.

Soit $\Phi_{1}$ et $\Phi_{2}$ deux OMFs définis sur les domaines de diagramme de $\alpha_{1}, \alpha_{2}$. Pour $i=1,2$, on considère une courbe $\gamma_{i}^{\alpha_{i}}$ définie à l'intérieur du domaine graphique $\alpha_{i}$ sur le corps $A_{i}$. Soit $b_{i}(t)=\Phi_{i}\left(\gamma_{i}(t)\right) R_{i}(t)$ parallèle le long de $\gamma_{i}^{\alpha_{1}}, i=1,2$, et $R:=R_{2}(t) R_{1}(t)^{-1} \in S O(2)$, qui, par définition, mesure la position relative de $\Phi_{2}$ par rapport à $\Phi_{1}$ le long de $\left(\gamma_{1}^{\alpha_{1}}, \gamma_{2}^{\alpha_{2}}\right)$. La variation de $R_{i}$ le long de $\gamma_{i}^{\alpha_{i}}$, pour $i=1,2$, est donnée par $\dot{R}_{i}=-\omega_{i}\left(\dot{\gamma}_{i}^{\alpha_{i}}\right) R_{i}$.

Étant donné une courbe a.c. $\gamma_{1}:[0, T] \rightarrow A_{1}$, le roulement de $A_{2}$ sur $A_{1}$ sans glisser ou tourner le long de $\gamma_{1}$ est caractérisé par une courbe $\Gamma=\left(\gamma_{1}, \gamma_{2}, R\right):[0, T] \rightarrow \mathcal{Q}\left(A_{1}, A_{2}\right)$, qui
est définie par les deux équations suivantes :

$$
\begin{array}{ll}
M^{\alpha_{2}} \dot{\gamma}_{2}^{\alpha_{2}}(t)=R M^{\alpha_{1}} \dot{\gamma}_{1}^{\alpha_{1}}(t) & \text { (condition de non-glissement) } \\
N B 1 . \dot{R} R^{-1}=\omega_{1}\left(\dot{\gamma}_{1}^{\alpha_{1}}\right)-\omega_{2}\left(\dot{\gamma}_{2}^{\alpha_{2}}\right) & \text { (condition sans rotation) } \tag{A.2}
\end{array}
$$

Definition 2.A. 1 La surface $A_{2}$ roule sur $A_{1}$ sans glisser ni tourner si, pour chaque $x=$ $\left(x_{1}, x_{2}, R_{0}\right) \in \mathcal{Q}\left(A_{1}, A_{2}\right)$ et courbe a.c. $\gamma_{1}:[0, T] \rightarrow A_{1}$ commençant à $x_{1}$, il existe une courbe a.c. $\Gamma:[0, T] \rightarrow \mathcal{Q}\left(A_{1}, A_{2}\right)$, avec $\Gamma(t)=\left(\gamma_{1}(t), \gamma_{2}(t), R(t)\right), \Gamma(0)=x$ et pour chaque $t \in[0, T]$, telle que, sur des graphiques appropriés, les équations (A.1) et (A.2) sont satisfaites. Nous appelons la courbe $\Gamma(t)$ une trajectoire admissible.

Les équations (A.1) et (A.2) peuvent être réécrites en coordonnées locales en considérant $f_{1}$ et $f_{2}$ deux OMFs et si l'état $x=\left(\gamma_{1}, \gamma_{2}, R\right)$ alors pour presque tous les $t$ tels que $x(t)$ reste dans le domaine d'un graphique approprié, il existe une fonction mesurable $u:[b, d] \rightarrow \mathbb{R}^{2}$ (que I'on appelle contrôle) telle que

$$
\begin{aligned}
\dot{\gamma}_{1}(t) & =u_{1}(t) \Phi_{1}^{1}\left(\gamma_{1}(t)\right)+u_{2}(t) \Phi_{2}^{1}\left(\gamma_{1}(t)\right) \\
\dot{\gamma}_{2}(t) & =u_{1}(t)\left(\Phi^{2}\left(\gamma_{2}(t)\right) R(t)\right)_{1}+u_{2}(t)\left(\Phi^{2}\left(\gamma_{2}(t)\right) R(t)\right)_{2} \\
\dot{R}(t) R^{-1}(t) & =\sum_{i=1}^{2} u_{i}(t)\left(\omega_{1}\left(\Phi_{i}^{1}\left(\gamma_{1}(t)\right)\right)-\omega_{2}\left(\Phi^{2}\left(\gamma_{2}(t)\right) R(t)\right)_{i}\right)
\end{aligned}
$$

Ensuite, nous pouvons réécrire le système d'équations ci-dessus comme suit :

$$
\begin{equation*}
\dot{x}=u_{1} X_{1}(x)+u_{2} X_{2}(x), \tag{A.3}
\end{equation*}
$$

où $X_{i}=\left(\Phi_{i}^{1},\left(\Phi^{2} R\right)_{i},\left(\omega_{1}\left(\Phi_{i}^{1}\right)-\omega_{2}\left(\Phi^{2} R(t)\right)_{i}\right)\right)^{T}, i=1,2$.
La proposition suivante décrit une propriété fondamentale du problème du corps roulant (voir [Chelouah and Chitour, 2003] pour plus de détails).

Proposition 2.A. 1 Soit $u \in \mathcal{H}$ une commande admissible qui donne lieu à la trajectoire admissible $\Gamma=\left(\gamma_{1}, \gamma_{2}, R\right):[0,1] \rightarrow \mathcal{Q}\left(A_{1}, A_{2}\right)$. Alors, les affirmations suivantes sont équivalentes

1. La courbe $\gamma_{1}:[0,1] \rightarrow A_{1}$ est une géodésique ;
2. La courbe $\gamma_{2}:[0,1] \rightarrow A_{2}$ est une géodésique ;
3. La courbe $\Gamma:[0,1] \rightarrow \mathcal{Q}\left(A_{1}, A_{2}\right)$ est une géodésique horizontale.

## Méthode de Continuation

Nous fixons quelques notations.

Les entrées admissibles $u$ sont des éléments de $\mathcal{H}=L^{2}\left([0,1], \mathbb{R}^{2}\right)$. Pour désigner
$\left(u_{1}(t)^{2}+u_{2}(t)^{2}\right)^{\frac{1}{2}}$ et $\left(\int_{0}^{1}\|u(t)\|^{2} d t\right)^{\frac{1}{2}}$ nous utilisons respectivement $\|u(t)\|$ et $\|u\|_{\mathcal{H}}$. Si $u, v \in \mathcal{H}$, alors $(u, v)_{\mathcal{H}}=\int_{0}^{1} u(t)^{T} v(t) d t$.

Comme mentionné dans la section Introduction, la méthode de continuité est utilisée pour résoudre des équations non linéaires de la forme $\mathcal{L}(x)=y$, où $x$ est l'inconnue, et $\mathcal{L}: X \rightarrow Y$ est surjectif.

Dans le contexte du problème de planification du mouvement, la fonction $\mathcal{L}$ est le point final $E_{p}: \mathcal{H} \rightarrow M$ associé à un certain $p \in M$ fixe. Pour $u \in \mathcal{H}$ et $p \in M$, soit $\gamma_{p, u}$ la trajectoire de $\Sigma$ commençant à $p$ pour $t=0$ et correspondant à $u$. Alors pour $v \in \mathcal{H}, E_{p}(v)$ est donné par $E_{p}(v):=\gamma_{p, v}(1)$.

Rappelons que $E_{p}(v)$ est défini pour chaque $v \in \mathcal{H}$. Le problème de planification du mouvement peut être reformulé comme suit :
Pour chaque $p, q \in M$, trouver une commande $u_{p, q} \in \mathcal{H}$ telle que $E_{p}\left(u_{p, q}\right)=q$. En d'autres termes, nous devons trouver l'inverse droit de la carte du point final $E_{p}$. Cet inverse droit existe dans un voisinage de tout point $u \in \mathcal{H}$ tel que $D E_{p}$ est surjectif (par l'hypothèse de contrabilité). Plus précisément, on commence avec un contrôle arbitraire $u_{0}$. On définit $q_{0}:=E_{p}\left(u_{0}\right)$, et on choisit un chemin $\pi:[0,1] \rightarrow M$ tel que $\pi(0):=q_{0}$ et $\pi(1):=q$. Nous cherchons maintenant un chemin $\Pi:[0,1] \rightarrow \mathcal{H}$ tel que, pour tout $s \in[0,1]$, il y a

$$
\begin{equation*}
E_{p}(\Pi(s))=\pi(s) \tag{A.4}
\end{equation*}
$$

En différentiant l'équation (A.4), on obtient

$$
\begin{equation*}
D E_{p}(\Pi(s)) \cdot \frac{d \Pi}{d s}(s)=\frac{d \pi}{d t}(s) . \tag{A.5}
\end{equation*}
$$

Si $D E_{p}(\Pi(s))$ a un rang complet, alors l'équation (A.5) peut être résolue pour $\Pi(s)$ en prenant $\Pi$ tel que

$$
\begin{equation*}
\frac{d \Pi}{d s}(s)=F(\Pi(s)) \cdot \frac{d \pi}{d t}(s) . \tag{A.6}
\end{equation*}
$$

où $F(\cdot)$ est un inverse droit de $D E_{p}(\cdot)$. Par exemple, nous pouvons choisir $F(\cdot)$ comme étant le pseudo-inverse de Moore-Penrose de $D E_{p}(\cdot)$. L'équation (A.6) est appelée l'équation de levage de chemin d'accès (PLE), qui est une équation différentielle ordinaire (ODE) sur $\mathcal{H}$ (cf. [Ważewski, 1947]).
Remarquez que, par construction, la commande définie par $u^{\text {final }}:=\Pi(1)$ dirige le système de $p$ à $q$. Afin d'obtenir la valeur de $\Pi(1)$, il suffit de résoudre le problème de la valeur initiale défini dans l'espace de contrôle $\mathcal{H}$.

$$
\begin{align*}
\frac{d \Pi}{d s}(s) & =F(\Pi(s)) \cdot \frac{d \pi}{d t}(s)  \tag{A.7}\\
\Pi(0) & =u_{0}
\end{align*}
$$

Pour appliquer avec succès la méthode de continuation au problème de planification du mouvement, nous devons garantir l'existence de $F(\Pi(s))$ pour chaque $s \in[0,1]$ de telle sorte que l'équation (A.6) soit toujours bien définie, c'est-à-dire que $D E_{p}(\Pi(s))$ doit toujours avoir un rang égal à 5 . De plus, puisque nous avons besoin d'évaluer $\Pi(1)$ pour obtenir une commande dirigeant le système de $p$ à $q$, la PLE définie dans (A.6) doit avoir une solution globale sur $[0,1]$.

En conclusion, l'application de la méthode de continuation au problème de planification du mouvement se décompose en deux étapes: caractériser (lorsque cela est possible) $S_{p}=\{u \in$ $\left.\mathcal{H}: \operatorname{rank}\left(D E_{p}(u)\right)<5\right\}$ et son image sous $E_{p}$, et lever les chemins $\pi:[0,1] \rightarrow M$ évitant $E_{p}\left(S_{p}\right)$ vers les chemins $\Pi:[0,1] \rightarrow \mathcal{H}$ globalement définis sur $[0,1]$ par (A.6).

## Surface convexe roulant sur le plan avec obstacles

Puisque $A_{2}$ est une variété riemannienne orientée, connectée et bidimensionnelle, alors $A_{2}$ peut être intégrée comme une surface convexe dans $\mathbb{R}^{3}$ (voir [Berger and Gostiaux, 2012]). Par conséquent, dans la suite, nous supposons que $A_{2}$ est une surface orientée de $\mathbb{R}^{3}$ avec une métrique induite par la métrique euclidienne de $\mathbb{R}^{3}$. De plus, nous supposons que $A_{2}$ est définie comme une composante connexe bornée de l'ensemble de niveau zéro d'une fonction lisse à valeur réelle $a: \mathbb{R}^{3} \rightarrow \mathbb{R}$. Le champ de vecteurs normaux à $A_{2}$ est donné par $\frac{\nabla a}{\|a\|}$, où $\nabla a=\left(a_{x}, a_{y}, a_{z}\right)$ désigne le vecteur gradient de $a$. La courbure gaussienne de $A_{2}$ est désignée par $K$.

Rappelons que le problème du corps roulant suppose que les espaces tangents aux points de contact sont identifiés. Dans $\mathbb{R}^{3}$, ceci est équivalent à l'identification des vecteurs normaux. Soit $n_{2}$ le vecteur normal de $A_{2}$, Puis, aux points de contact, nous attribuons $\frac{\nabla a}{\|a\|}$ à $-N_{\mathbb{R}^{2}}$, où $N_{\mathbb{R}^{2}}$ est le vecteur normal au plan $\mathbb{R}^{2}$.

En utilisant le fait que $\mathcal{Q}\left(\mathbb{R}^{2}, A_{2}\right)$ est un bundle de cercle quand $A_{2}$ est une variété à 2 dimensions, et en prenant les coordonnées géodésiques $\vartheta$ pour $A_{2}$ au point de contact $x_{2}$, et en considérant les coordonnées $x=\left(v_{1}, w_{1}, v_{2}, w_{2}, \psi\right)$ dans un certain voisinage de $\left(0, \psi_{0}\right)$ dans $\mathbb{R}^{4} \times S^{1}$, le système de contrôle (A.3) peut s'écrire localement comme suit

$$
\begin{equation*}
\dot{x}=u_{1} X_{1}(x)+u_{2} X_{2}(x), \tag{A.8}
\end{equation*}
$$

avec

$$
\begin{align*}
& X_{1}(x)=\left(1,0, \cos (\psi),-\frac{\sin (\psi)}{\vartheta},-\frac{\vartheta_{v_{2}}}{\vartheta} \sin (\psi)\right)^{T}  \tag{A.9}\\
& X_{2}(x)=\left(0,1,-\sin (\psi),-\frac{\cos (\psi)}{\vartheta},-\frac{\vartheta_{v_{2}}}{\vartheta} \cos (\psi)\right)^{T} . \tag{A.10}
\end{align*}
$$

Pour un développement détaillé de la manière d'obtenir le système (A.8) avec les champs de vecteurs $X_{1}$ et $X_{2}$ définis respectivement par (A.9) et (A.10), voir [Chelouah and Chitour, 2003].

## Obstacles et dynamique du système de contrôle

Par la suite, l'espace des domaines $\mathcal{Q}\left(\mathbb{R}^{2}, A_{2}\right)$ est simplement désigné par $M$.
Pour traiter le problème de planification de mouvement avec obstacles, nous suivons une stratégie suggérée par E. Sontag dans [Sontag, 1995].
Definition 2.A. 2 Nous disons que $\mathcal{W}$ est un obstacle de $\mathbb{R}^{2}$ si $\mathcal{W}$ est un sous-ensemble compact non vide de $\mathbb{R}^{2}$.

D'après la définition ci-dessus, on peut supposer qu'il existe une fonction $\xi: \mathbb{R}^{2} \rightarrow \mathbb{R}$ telle que $\mathcal{W}=\left\{w \in \mathbb{R}^{2} \mid ; \xi(w) \leq 0\right\}$.

Un obstacle $\mathcal{W}$ fait correspondre dans $M$ une région (voir Figure A.4), notée $C$, définie par

$$
C=\left\{x=\left(\gamma_{1}, \gamma_{2}, R\right) \in M \mid ; \gamma_{1} \in \mathcal{W}\right\}=\left\{x \in M \mid ; \xi\left(\gamma_{1}\right) \leq 0\right\} .
$$



Figure A.4: Application d'obstacles dans $M$.
Par conséquent, $C$ est un obstacle dans l'espace des domaines $M$. Autrement dit, un obstacle dans $M$ est un sous-ensemble fermé non vide de $M$ tel que $\widehat{M}=M \backslash C$ est également non vide.

Considérons une suite croissante $\left(W_{i}\right)_{i \geq 0}$ des sous-ensembles compacts de $M$ telle que $C=\cup_{i \geq 0} W_{i}$. Pour chaque $i \geq 0$, on peut construire une fonction lisse $\zeta_{i}: M \rightarrow \mathbb{R}$ telle que $\zeta_{i} \equiv 0$ sur $W_{i}$ et $0<\zeta_{i}(x) \leq \inf \left(1, d\left(\gamma_{1}, \mathcal{W}_{i}\right)\right)$ pour $x \in M \backslash W_{i}$. Soit

$$
M_{i}:=\sup _{x \in M}\left(\left\|D \zeta_{i}(x)\right\|+\left\|D^{2} \zeta_{i}(x)\right\|\right) .
$$

On définit alors

$$
\begin{equation*}
\zeta(x)=\sum_{i \geq 0} \frac{1}{2^{i} M_{i}} \zeta_{i}(x) . \tag{A.11}
\end{equation*}
$$

Alors, il existe une fonction lisse bornée $\zeta: M \rightarrow \mathbb{R}$ telle que $\zeta>0$ sur $\widehat{M}, \zeta=0$ sur $C$ et $\left(\widehat{M},\langle\langle\cdot, \cdot\rangle\rangle_{x}\right)$ est une variété riemanienne complète, où

$$
\begin{equation*}
\langle\cdot, \cdot\rangle\rangle_{x}=\frac{\langle\cdot, \cdot\rangle}{\zeta^{2}} . \tag{A.12}
\end{equation*}
$$

Considérons le système de contrôle sur $M$ défini par

$$
\begin{equation*}
\dot{y}=v_{1} \bar{X}_{1}(y)+v_{2} \bar{X}_{2}(y) \tag{A.13}
\end{equation*}
$$

où $\bar{X}_{i}=\zeta X_{i}, i=1,2$. Nous pouvons vérifier que le système (A.13), restreint à $\widehat{M}$, satisfait la condition LARC (voir [Chitour and Sussmann, 1998]).

Alors $\widehat{M}$ est invariant sous le système de contrôle (A.13), plus précisément,
Lemma 2.A.1 pour tout $v \in \mathcal{H}$, pour tout $p \in \widehat{M}$, si $\bar{\gamma}_{v}$ est la solution de (A.13) avec $\bar{\gamma}(0)=p$ alors $\bar{\gamma}([0,1]) \subset \widehat{M}$.

En d'autres termes, pour chaque $t \in[0,1], \bar{\gamma}(t)$ appartient à la composante connectée de $\widehat{M}$ contenant $p$. Ainsi, le problème de planification du mouvement avec obstacles se réduit à un problème de planification du mouvement pour chaque composante connectée de $\widehat{M}$. (Pour plus de détails sur $\zeta$ et la preuve du lemme 2.A.1, voir [Chitour, 2006]).

## Points clés pour la mise en œuvre numérique

Dans [Chitour, 2006] est prouvé que, pour tout sous-intervalle compact $J=[0, a]$ de l'intervalle d'existence de la solution maximale de l'équation (2.8), noté $I$, l'approximation numérique de l'équation (2.8) définie par $E_{p, j}$ et associée à $\pi$ a une solution globale sur $J$ pour $j$ suffisamment grand. En particulier, si $J=I=[0,1]$ (c'est-à-dire que la PLE a une solution globale), le résultat ci-dessus peut être considéré comme une justification théorique de l'utilisation de la procédure de Galerkin dans les implémentations numériques de la méthode de continuité.

Dans cette Section, nous voulons montrer comment la Méthode de Continuation, présentée dans la Section 3.2 peut être implémentée numériquement pour résoudre le problème de planification de mouvement avec obstacles lorsqu'une surface convexe de $\mathbb{R}^{3}$ roule sur le plan avec des régions interdites (ou obstacles).

Dans les trois sous-sections suivantes, nous détaillons les points fondamentaux de l'implémentation numérique, qui sont la discrétisation de l'espace de contrôle $\mathcal{H}$, le calcul de $D E_{p}(u)$, et le soulèvement de la courbe $\widetilde{\gamma}_{1}$ sur le corps convexe $A_{2}$. Les idées proviennent de [Alouges et al., 2010].

## Discrétisation de $\mathcal{H}$

Rappelons que l'espace de contrôle $\mathcal{H}$ est, en qénéral, un espace vectoriel de dimension infinie, alors nous commençons par discrétiser $\mathcal{H}$ car nous devons résoudre le problème de valeur initiale défini dans (2.9). Dans ce cas, les commandes sont des courbes planes $\gamma_{1}:[0,1] \rightarrow \mathbb{R}^{2}$ telles que

$$
\dot{\gamma}_{1}(t)=\zeta(x(t))\left(u_{1}, u_{2}\right), \quad x(t)=\left(\gamma_{1}(t), \gamma_{2}(t), R(t)\right) .
$$

Nous divisons l'intervalle $[0,1]$ en $n-1$ parties et approximons l'espace de contrôle $\mathcal{H}$ par le sous-espace $\widetilde{\mathcal{H}}$ à $2 n$ dimensions des fonctions linéaires par morceaux (pour plus de détails, voir [Gautschi, 2012]). Alors, $\gamma_{1}$ peut être approximé par $\widetilde{\gamma}_{1}$, c'est-à-dire, l'interpolation linéaire de
$\left(\gamma_{1}^{1}, \ldots, \gamma_{1}^{n}\right)$, où $\gamma_{1}^{i+1}=\gamma_{1}\left(\frac{i}{n-1}\right)=\left(x_{i+1}, y_{i+1}\right)^{T}$ pour $i=0, \ldots, n-1$.
Sur chaque segment $\left[t_{i+1}, t_{i+2}\right]=\left[\frac{i}{n-1}, \frac{i+1}{n-1}\right]$ la commande approximative $\left(\widetilde{u}_{1}^{i+1}, \widetilde{u}_{2}^{i+1}\right)^{T}$ est proportionnelle au vecteur $\frac{1}{\zeta(x(t))}\left(x_{i+2}-x_{i+1}, y_{i+2}-y_{i+1}\right)^{T}$ pour $i=0, \ldots, n-2$.

Remarquez que lorsque les éléments dans $\widetilde{\mathcal{H}}$ sont des fonctions linéaires par morceaux avec plus d'un morceau, alors ce ne sont pas des contrôles singuliers. Ensuite, les trajectoires correspondantes sur $A_{2}$ sont également faciles à obtenir en intégrant certaines équations géodésiques en utilisant la Proposition 2.2.1. Enfin, on utilise la méthode d'Euler pour intégrer (2.9).

Calculer $D E_{p}(u)$.
Pour évaluer $D E_{p}(u)$, pour $u \in \mathcal{H}$, ce qui suit est nécessaire.
Plus de détails voir [Chelouah and Chitour, 2003].
Pour $z \in T_{E_{p}(u)}^{*} M$, soit $\lambda_{z, u}:[0,1] \rightarrow T^{*} M$ le champ de covecteurs le long de $\gamma_{p, u}$ tel qu'il satisfasse, en coordonnées, l'équation adjointe le long de $\gamma_{p, u}$ avec la condition terminale $z$, c'est-à-dire $\lambda_{z, u}(1)=z$ et pour a. e $t \in[0,1]$ on a

$$
\dot{\lambda}_{z, u}(t)=-\lambda_{z, u}(t) \cdot\left(u_{1}(t) D X_{1}\left(\lambda_{z, u}(t)\right)+u_{2}(t) D X_{2}\left(\lambda_{z, u}(t)\right)\right) .
$$

Si $V$ est un champ de vecteurs $C^{\infty}$ sur $M$, la fonction de commutation $\varphi_{V, z, u}(t)$ associée à $V$ est l'évaluation de $\lambda \cdot V(x)$, la fonction hamiltonienne de $V$ le long de $\left(\gamma_{p, u}, \lambda_{p, u}\right)$, c'est-à-dire que pour $t \in[0,1]$, on obtient

$$
\varphi_{V, z, u}(t):=\lambda_{z, u} \cdot V\left(\gamma_{p, u}(t)\right) .
$$

Ensuite, $D E_{p}(u)$ peut être calculé grâce à la formule suivante: pour $z \in T_{E_{p}(u)}^{*} M$ et $u, v \in \mathcal{H}$

$$
\begin{equation*}
z \cdot D E_{p}(u)(v)=\left(v, \varphi_{z, u}\right)_{\mathcal{H}}, \tag{A.14}
\end{equation*}
$$

où le vecteur de la fonction de commutation $\varphi_{z, u}$ est la solution du problème de Cauchy, qui est défini, en coordonnées, par

$$
\begin{align*}
& \dot{\varphi}_{1}=-u_{2} K \varphi_{3} \\
& \dot{\varphi}_{2}=u_{1} K \varphi_{3} \\
& \dot{\varphi}_{3}=-u_{2} \varphi_{4}+u_{1} \varphi_{5}  \tag{A.15}\\
& \dot{\varphi}_{4}=-u_{2} K \varphi_{3} \\
& \dot{\varphi}_{5}=u_{1} K \varphi_{3}
\end{align*}
$$

avec la condition terminale $\varphi_{z, u}(1)=z$.
En pratique, puisque la matrice discrète $D E_{p}(u)$ est une matrice $5 \times 5$ et que son image est donnée par (A.14), il suffit de prendre les cinq vecteurs de la base canonique de $\mathbb{R}^{5}$, comme conditions finales $z$, et d'intégrer (A.15) en temps inverse. Dans les simulations, un schéma numérique Runge-Kutta d'ordre 4 est utilisé pour l'intégration, le produit scalaire $(\cdot, \cdot)_{\mathcal{H}}$ dans l'espace de contrôle $\mathcal{H}$ est évalué par quadrature gaussienne, et la courbure gaussienne $K$ est calculée en utilisant la proposition 1.2.1.

Lifting de la courbe plane $\widetilde{\gamma}_{1}$ sur $A_{2}$.
Notez que la courbure $K$ apparaissant dans (3.9) est prise au point de contact final sur la surface $A_{2}$ après qu'elle ait roulé le long de la courbe à constante par morceaux $\widetilde{\gamma}_{1}$. Ainsi, pour localiser le point final, nous devons soulève la courbe plane $\widetilde{\gamma}_{1}$ sur $A_{2}$, et la dynamique de soulèvement est donnée par (3.7). Cependant, puisque les coordonnées géodésiques impliquées dans (3.7) ne sont pas données explicitement en pratique, la méthode numérique de soulèvement est basée sur la proposition 2.2.1.

Sur chaque intervalle $\left[t_{i+1}, t_{i+2}\right]$, la courbe de contrôle approximative $\widetilde{\gamma}_{1}$ est une ligne droite (c'est-à-dire une géodésique dans $\mathbb{R}^{2}$ ), et alors, par la Proposition 2.2.1, la courbe de soulèvement $\widetilde{\gamma}_{1}$ sur $A_{2}$ est également une géodésique sur chaque intervalle $\left[t_{i+1}, t_{i+2}\right]$ pour tous les $i=0, \ldots, n-2$. Alors, à partir du point de contact initial $x_{0}$ sur $A_{2}$, on peut intégrer successivement (avec un schéma Runge-Kutta d'ordre 4) l'équation géodésique, donnée par la Proposition 1.2.2 (équation 2.2), sur chaque $\left[t_{i+1}, t_{i+2}\right]$ avec des conditions initiales égales à $\widetilde{\gamma}_{1}\left(t_{i+1}\right)$ et $\left(\widetilde{u}_{1}^{i+1}, \widetilde{u}_{2}^{i+1}\right)$, pour $i=0, \ldots n-2$.

Une difficulté très importante est que l'intégration numérique s'effectue sur le collecteur $A_{2}, t>0$. En supposant que nous soyons au point $x \in A_{2}$ au temps $t$, alors, au temps $t+\delta t$, nous nous déplaçons à $x_{\text {new }}=x+(\delta t) e$, avec $e \in T_{x} A_{2}$, mais $x_{\text {new }}$ n'appartient pas à $A_{2}$ si $e$ est non nul. Il faut donc, à chaque pas d'intégration, projeter $x_{\text {new }}$ sur $A_{2}$. En détail, supposons que le point $(0,0,0)$ est à l'intérieur du corps convexe $A_{2}$. Puisque $A_{2}$ est défini comme (une composante connexe bornée de) l'ensemble de niveau zéro d'une fonction lisse $a$, nous supposons que $\left|a\left(x_{n e w}\right)\right| \leq \varepsilon$ pour un certain $\varepsilon \ll 1$, c'est-à-dire que $x_{\text {new }}$ est proche de $A_{2}$. Alors, par la convexité de $A_{2}$, il existe un unique nombre réel $\mu$ proche de 1 tel que $a\left(\mu x_{\text {new }}\right)=0$. Le problème de la projection à traiter est local, et par conséquent, la méthode de Newton est efficace pour trouver $\mu$. La dérivée par rapport à $\mu$ est également nécessaire, et est évaluée par un schéma à différences finies.

## Exemples

Cette section présente trois exemples de corps roulant sur le plan euclidien: la sphère, la boule aplatie et un œuf. Les simulations ont été développées dans le logiciel Matlab, et le code a été exécuté sur un Macbook Pro Apple M1 8Go de RAM pour 100 itérations avec $n-1=100$ pour la discrétisation de l'espace de contrôle.

Chaque corps convexe $A_{2}$ est défini par la fonction $a(x, y, z)$ telle que $A_{2}=a^{-1}(0)$. Les points de contact initiaux sont bleu et les points de contact finaux sont magenta. Les orientations initiales et finales ont la même couleur que les points initiaux et finaux. Les obstacles sont des régions circulaires dans le plan, et ils sont colorés en rouge (voir les figures A.5, A. 7 et A.9).

La courbe initiale choisie pour appliquer la méthode de continuation est la courbe verte, et la courbe obtenue par ce processus est colorée en bleu. Dans chaque exemple, trois courbes différentes sont choisies pour évaluer la méthode. Enfin, la courbe rouge sur le corps convexe est obtenue à partir du soulèvement de la courbe calculée par la méthode de continuation (courbe bleue) (voir figures A.6, A. 8 et A.10).

## Sphère

Cette sphère est définie par l'ensemble des niveaux zéro de la fonction $a$, définie par

$$
a(x, y, z)=x^{2}+y^{2}+z^{2}-1
$$



Figure A.5: Côté gauche: sphère avec points de contact initiaux et finaux et ses orientations. Côté droit : obstacles (en rouge) et points de contact initiaux et finaux sur le plan.

## Oeuf

Cet oeuf est défini par l'ensemble de niveau zéro de la fonction $a$, définie par

$$
a(x, y, z)=\frac{x^{2}+y^{2}}{1-0.4 z}+\frac{1}{4} z^{2}-1
$$



Figure A.6: À gauche: courbe $\gamma_{2}$ sur la sphère. Côté droit : Courbe obtenue en appliquant la méthode de continuation (bleu) à la courbe initiale choisie (vert).

## Boule aplatie

Cette boule aplatie est définie par l'ensemble de niveau zéro de la fonction $a$, définie par

$$
a(x, y, z)=x^{2}+y^{2}+5 z^{2}-1
$$



Figure A.7: Côté gauche: oeuf avec points de contact initiaux et finaux et ses orientations. Côté droit: obstacles (en rouge) et points de contact initiaux et finaux sur le plan.


Figure A.8: À gauche: courbe $\gamma_{2}$ sur le oeuf. Côté droit: Courbe obtenue en appliquant la méthode de continuation (bleu) à la courbe initiale choisie (vert).

Enfin, pour montrer le bon fonctionnement de la méthode mise en œuvre, nous avons considéré des obstacles plus complexes avec les mêmes corps convexes. La courbe initiale choisie est colorée en vert, et celle obtenue par la méthode de continuation est colorée en bleu.


Figure A.9: Côté gauche: boule aplatie avec points de contact initiaux et finaux et ses orientations. Côté droit: obstacles (en rouge) et points de contact initiaux et finaux sur le plan.


Figure A.10: À gauche: courbe $\gamma_{2}$ sur la boule aplatie. Côté droit : Courbe obtenue en appliquant la méthode de continuation (bleu) à la courbe initiale choisie (vert).

Les obstacles sont colorés en vert. Voir la figure A.11.


Figure A.11: Côté gauche: corps convexe avec points de contact initial et final avec orientations et obstacle sur le plan (en vert). Côté droit: courbe $\gamma_{2}$ sur le corps convexe et courbe obtenue en appliquant la méthode de continuation (en bleu) à la courbe initiale choisie (en vert).

## Conclusions et commentaires

Cette partie de la thèse présente une implémentation numérique pour résoudre le problème de planification de mouvement lorsqu'un corps convexe roule sur le plan euclidien avec des obstacles, en détaillant la discrétisation de l'espace de contrôle $\mathcal{H}$, l'évaluation de $D E_{p}(u)$, et la levée de la courbe $\widetilde{\gamma}_{1}$ sur le corps convexe $A_{2}$. Le corps convexe est considéré comme une surface orienté de $\mathbb{R}^{3}$ et donc il est défini par une fonction $a: \mathbb{R}^{3} \rightarrow \mathbb{R}$ telle que $A_{2}=a^{-1}(0)$. De plus, pour appliquer avec succès la méthode de continuation, $A_{2}$ doit avoir une géodésique périodique stable. Cette propriété est vraie pour toute surface compacte convexe ayant une symétrie de révolution (cf. [Chelouah and Chitour, 2003]). Pour cette raison, la sphère, la balle aplatie et un œuf ont été considérés pour les simulations. Cependant, lorsque nous relaxons cette propriété géométrique, la méthode mise en œuvre fonctionne toujours. Ce qui précède peut être considéré comme la robustesse de la méthode. Il y a un défi à relever ici: rechercher un résultat théorique de convergence qui assure l'application réussie de la méthode de continuation lorsque le corps convexe ne satisfait pas la condition géométrique de symétrie de révolution. Nous pensons que ce résultat est lié au fait que $\nabla a(x, y, z) \neq 0$ sur l'ensemble de niveau zéro de $a$, car, selon les propositions 1.2.1 et 1.2.2, la courbure gaussienne de la surface $A_{2}$ et la levée de la courbe plane sur $A_{2}$ respectivement, dépendent de $\nabla a(x, y, z)$, en effet, si $\nabla a(x, y, z) \neq 0$ les équations (2.1) et (2.2) qui définissent respectivement la courbure gaussienne et le lift, sont toujours bien définies. Par exemple, considérons $A_{2}$ une surface
orientée sans axe de symétrie, définie par $a(x, y, z)=\frac{x^{2}}{1-0.7 y}+\frac{3 y^{2}}{1-0.1 z}+\frac{0.3 z^{2}}{1-0.3 x-0.1 y}-1$ et donc $\nabla a(x, y, z)=\left(\begin{array}{c}\frac{9 z^{2}}{(3 x+y-10)^{2}}-\frac{20 x}{7 y-10} \\ \frac{3 z^{2}}{(3 x+y-10)^{2}}-\frac{60 y}{z-10}+\frac{70 x^{2}}{(7 y-10)^{2}} \\ \frac{3 y^{2}}{(z-0.1)^{2}}-\frac{3 z}{0.75 x+0.5 y-5}\end{array}\right) \neq 0$ sur l'ensemble de niveau zéro de $a$. Dans la Figure A.12, nous pouvons voir comment la méthode fonctionne.


Figure A.12: Côté gauche: Corps convexe non symétrique avec points de contact initial et final et courbe $\gamma_{2}$. Côté droit: Courbe obtenue en appliquant la méthode de continuation (bleu) à la courbe initiale choisie (vert).

## APPENDIX B

## Paper 1

Solvable approximation of 3-dimensional almost-Riemannian structures, article in collaboration with Philippe Jouan. It was published in the journal Mathematical Control \& Related Fields, volume 12 (2022), Issue 2 (June), pp.303-326. (see online link).

# Solvable Approximations of 3-dimensional Almost-Riemannian structures 

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#### Abstract

In some cases, the nilpotent approximation of an almost-Riemannian structure can degenerate into a constant rank sub-Riemannian one. In those cases, the nilpotent approximation can be replaced by a solvable one that turns out to be a linear ARS on a nilpotent Lie group or a homogeneous space.

The distance defined by the solvable approximation is analyzed in the 3Dgeneric cases. It is shown that it is a better approximation of the original distance than the nilpotent one.


Key words: Almost-Riemannian geometry, Nilpotent approximation.
AMS subject classifications: 53C15, 53C17, 22E25, 53B99.

## 1 Introduction

The aim of this paper is to locally approximate almost-Riemannian structures (ARS in short) at singular points, by ARSs on Lie groups and to show that this approximation is generally better than the nilpotent one.

An ARS on an $n$-dimensional differential manifold is a rank-varying sub-Riemannian structure that can be defined, at least locally, by a set of $n$ vector fields satisfying the Lie algebra rank condition (Larc in short). We denote by $\Delta_{p}$ the linear span of the vector fields at the point $p$. The set of points where $\operatorname{dim}\left(\Delta_{p}\right)<n$ is called the singular locus or the singular set and denoted by $\mathcal{Z}$. Many papers dedicated to the study of ARSs can be found in the literature, for instance [2], [8], [9], [11] [12].

In the generic 3 -dimensional case, in which we are particularly interested, the singular set is a codimension one embedded submanifold and the points where $\Delta_{p}=$ $T_{p} \mathcal{Z}$ are isolated (see [3] and [10]).

We are likewise interested in the so-called simple ARSs on Lie groups (or homogeneous space) because they will be used as approximating structures for general

[^0]ARSs: a simple ARS on an $n$-dimensional Lie group is an almost-Riemannian structure defined by $n-1$ left-invariant vector fields and one vector field whose flow is a one-parameter group of automorphisms, called linear in the sequel. Under some conditions, the singular set of such structures is a subgroup or an analytic, embedded, codimension one submanifold (see [5] and [17]).

In some cases the nilpotent approximation of an ARS degenerates, because it is no longer an ARS but a constant rank sub-Riemannian structure. In other words, it may happen that some of the vector fields of the nilpotent approximation vanish, changing the almost-Riemannian structure into a constant rank sub-Riemannian one. For instance, if

$$
X_{1}=\left(\begin{array}{l}
1 \\
0 \\
0
\end{array}\right), \quad X_{2}=\left(\begin{array}{c}
0 \\
1 \\
x
\end{array}\right) \quad \text { and } \quad X_{3}=\left(\begin{array}{c}
0 \\
0 \\
x^{2}
\end{array}\right)
$$

then its nilpotent approximation is

$$
\widehat{X}_{1}=\left(\begin{array}{l}
1 \\
0 \\
0
\end{array}\right), \quad \widehat{X}_{2}=\left(\begin{array}{l}
0 \\
1 \\
x
\end{array}\right) \quad \text { and } \quad \widehat{X}_{3}=0 .
$$

It is what happens in some generic 3-dimensional cases (see for instance [10]). In this paper we are interested in the case where only one of the vector fields vanishes and the other ones are independent: then they define a left-invariant sub-Riemannian structure on a Lie group (or a homogeneous space).

Our aim consists in recovering the almost-Riemannian structure lost in the nilpotent approximation, thanks to a vector field, denoted $\widetilde{X}_{n}$ which is the homogeneous component of degree 0 of the Taylor expansion in privileged coordinates of the vector field that vanishes. The new family of vector fields composed by the nilpotent approximation and $\widetilde{X}_{n}$ is called the solvable approximation.

The Lie algebra generated by this new family of vector fields is finite dimensional and solvable. However, we are interested in some nilpotent Lie group on which $\widetilde{X}_{n}$ acts as a linear vector field (these vector fields were generalized in [6]). Thanks to the equivalence theorem of [16] we know that the space $\mathbb{R}^{n}$ is diffeomorphic to a homogeneous space. Through this diffeomorphism, the solvable approximation is equivalent to a simple ARS on a homogeneous space or a Lie group. It is important to notice that in the 3D-generic case the non-degenerated nilpotent approximation, and the solvable one in the degenerated case, are simple ARSs on Lie groups or homogeneous space (see Section 3.3).

On the other hand, the solvable approximation gives rise to a distance denoted by $\widetilde{d}$. This distance has the advantage to be really almost-Riemannian unlike the distance $\widehat{d}$ associated to the nilpotent approximation in the degenerated cases. The distance $\widetilde{d}$ is not homogeneous but always satisfies $\widetilde{d} \leq \widehat{d}$.

Denoting by $d$ the distance associated to the original structure we show that in some 3D-generic cases the order of $|d-\widetilde{d}|$ is strictly better than the one of $|d-\widehat{d \mid}|$. More accurately, the order of $|d-\widehat{d}|$ is $d^{\frac{3}{2}}$ and the one of $|d-\widetilde{d}|$ is $d^{2}$ in the cases we consider.

Moreover, the nilpotent distance $\widehat{d}$ is left-invariant while $d$ and $\widetilde{d}$ are not. Using this fact we prove that for some pairs $\left(q, q^{\prime}\right)$ of points translated from the singular locus the difference $\left|d\left(q, q^{\prime}\right)-\widetilde{d}\left(q, q^{\prime}\right)\right|$ is strictly smaller than $\left|d\left(q, q^{\prime}\right)-\widehat{d}\left(q, q^{\prime}\right)\right|$.

The paper is organized as follows. Section 2 contains generalities about ARSs, nonholonomic order, privileged coordinates, the nilpotent approximation, linear vector fields and simple ARS on Lie groups or homogeneous spaces.

In Section 3 we introduce the definition of a solvable approximation, we analyze its algebraic structure and an example is detailed.

Section 4 is divided in two parts. In the first one, we state two propositions about the almost-Riemannian distance $\widetilde{d}$ defined by the solvable approximation. The second part is devoted to analyze $\widetilde{d}$ in the 3 -dimensional generic case.

Finally, in Section 5 we provide the Hamiltonian associated to the flow defined by the solvable approximation in the 3D generic case and we compute the geodesics with initial condition $x(0)=y(0)=z(0)=0$ and $p(0)=\cos (\theta), q(0)=\sin (\theta)$, $r(0)=r$ in a particular case.

## 2 Preliminaries

In this section some definitions and results are reviewed and come from [4], [7] and [15].

### 2.1 Almost-Riemannian structures

An almost-Riemannian structure can always be locally defined by a set of $n$ vector fields, where $n$ is the dimension of the state space. Since we are interested in local questions, the following definition will be enough in this paper, and the reader is referred to [4], [5] and [15] for the global definition on manifolds.

We denote by $\operatorname{Lie}\left(X_{1}, \ldots, X_{n}\right)$ the Lie algebra generated by the vector fields $X_{1}, \ldots, X_{n}$ on $\mathbb{R}^{n}$.

Definition 1. We say that the vector fields $X_{1}, \ldots, X_{n}$ satisfy the Lie algebra rank condition (Larc in short) on an open set $\Omega$ of $\mathbb{R}^{n}$ if

$$
\operatorname{Lie}\left(X_{1}, \ldots, X_{n}\right)(p)=T_{p} \mathbb{R}^{n}
$$

for all $p \in \Omega$.
Definition 2. The set $\left\{X_{1}, \ldots, X_{n}\right\}$ defines an almost-Riemannian structure (ARS in short) on the open and connected subset $\Omega$ of $\mathbb{R}^{n}$ if:
(i) It satisfies Larc.
(ii) The singular locus, that is $\mathcal{Z}=\left\{p \in \Omega / \operatorname{rank}\left(X_{1}(p), X_{2}(p), \ldots X_{n}(p)\right)<n\right\}$ is non-empty, but with empty interior.

The metric is defined by declaring the frame to be orthonormal.

## Remark 1.

1. The structure is Riemannian out of $\mathcal{Z}$.
2. Let $v \in T_{p} \Omega$. If $p$ is a Riemannian point then $v=\sum_{i=1}^{n} u_{i} X_{i}(p)$ in a unique way and its length is $\|v\|=\left(\sum_{i=1}^{n} u_{i}{ }^{2}\right)^{\frac{1}{2}}$. If $p \in \mathcal{Z}$ then the decomposition of $v$, if it exists, is not unique and $\|v\|=\inf \left\{\left(\sum_{i=1}^{n} u_{i}{ }^{2}\right)^{\frac{1}{2}}: v=\sum_{i=1}^{n} u_{i} X_{i}(p)\right\}$.
An absolutely continuous curve $\gamma:[0, T] \longrightarrow \mathbb{R}^{n}$ is admissible if there exists a measurable essentially bounded function $t \mapsto u(t)$ from $[0, T]$ into $\mathbb{R}^{n}$ called control function such that $\dot{\gamma}(t)=u_{1}(t) X_{1}(\gamma(t))+u_{2}(t) X_{2}(\gamma(t))+\ldots+u_{n}(t) X_{n}(\gamma(t))$ for almost every $t \in[0, T]$. Given an admissible curve $\gamma:[0, T] \longrightarrow \mathbb{R}^{n}$, the length of $\gamma$ is

$$
l(\gamma)=\int_{0}^{T}\|\dot{\gamma}(t)\| d t
$$

The almost-Riemannian distance (or Carnot-Caratheodory distance) on $\Omega$ associated with the $n$-ARS is defined by

$$
d\left(p_{0}, p_{1}\right)=\inf \left\{l(\gamma): \gamma(0)=p_{0}, \gamma(T)=p_{1}, \gamma \text { admissible }\right\}
$$

It induces the usual topology on $\Omega$.

### 2.2 Nonholonomic orders

In what follows $\left\{X_{1}, X_{2}, \ldots, X_{n}\right\}$ defines an ARS on $\Omega \subset \mathbb{R}^{n}$.
Definition 3. Let $f: M \rightarrow \mathbb{R}$ be a continuous function. The nonholonomic order of $f$ at $p$, denoted $\operatorname{ord}_{p}(f)$, is the real number defined by

$$
\operatorname{ord}_{p}(f)=\sup \left\{s \in \mathbb{R}: f(q)=O\left(d(p, q)^{s}\right)\right\} .
$$

This order is always nonnegative.
Let $C^{\infty}(\mathrm{p})$ denote the set of germs of smooth functions at $p$. For $f \in C^{\infty}(p)$, we call nonholonomic derivative of order 1 of $f$ the Lie derivatives $X_{1} f, \cdots, X_{n} f$. We call further $X_{i} X_{j} f, X_{i} X_{j} X_{k} f, \ldots$, the nonholonomic derivatives of $f$ of order $2,3, \ldots$ of $f$. The nonholonomic derivative of order 0 of $f$ at $p$ is $f(p)$.

As a consequence, the nonholonomic order of a smooth (germ of) function is given by the formula

$$
\operatorname{ord}_{p}(f)=\min \left\{s \in \mathbb{N}: \exists i_{1}, \ldots, i_{s} \in\{1, \ldots, n\} \text { s.t. }\left(X_{i_{1}} \ldots X_{i_{s}} f\right)(p) \neq 0\right\}
$$

where as usual we adopt the convention that $\min \emptyset=+\infty$.
Let $V F(p)$ denote the set of germs of smooth vector fields at $p$.
Definition 4. Let $X \in \operatorname{VF}(p)$. The nonholonomic order of $X$ at $p$, denoted by $\operatorname{ord}_{p}(X)$, is the real number defined by:

$$
\operatorname{ord}_{p}(X)=\sup \left\{\sigma \in \mathbb{R}: \operatorname{ord}_{p}(X f) \geq \sigma+\operatorname{ord}_{p}(f), \quad \forall f \in C^{\infty}(p)\right\}
$$

### 2.3 Privileged coordinates

We adopt the notation of [15] to define privileged coordinates. Let $\operatorname{VF}(\Omega)$ denote the set of smooth vector fields on $\Omega$. We define $\Delta^{1}=\operatorname{span}\left\{X_{1}, \ldots, X_{n}\right\}$. For $s \geq 1$, define $\Delta^{s+1}=\Delta^{s}+\left[\Delta^{1}, \Delta^{s}\right]$, where

$$
\left[\Delta^{1}, \Delta^{s}\right]=\operatorname{span}\left\{[X, Y]: X \in \Delta^{1}, Y \in \Delta^{s}\right\}
$$

For $p \in \Omega$, we set for $s \geq 1, \Delta^{s}(p)=\left\{X(p): X \in \Delta^{s}\right\}$. By definition these sets are linear subspaces of $T_{p} \Omega$.

The evaluation of these sets at $p$ forms a flag of subspaces of $T_{p} \Omega$, and since $X_{1}, \ldots, X_{n}$ satisfy Larc, we get,

$$
\begin{equation*}
\Delta^{1}(p) \subset \Delta^{2}(p) \subset \cdots \subset \Delta^{r-1}(p) \varsubsetneqq \Delta^{r}(p)=T_{p} \Omega \tag{1}
\end{equation*}
$$

where $r=r(p)$ is called the degree of nonholonomy at $p$. Let $n_{i}(p)=\operatorname{dim} \Delta^{i}(p)$. The $r$-tuple of integers $\left(n_{1}(p), \ldots, n_{r}(p)\right)$ is called the growth vector at $p$. The first integer in the growth vector is the rank $n_{1}(p) \leq n$ of the family $X_{1}(p), \ldots, X_{n}(p)$, and the last one $n_{r}(p)=n$ is the dimension of $\mathbb{R}^{n}$.
Definition 5. The point $p$ is regular if the growth vector is constant in some neighborhood of $p$. Otherwise we say that $p$ is a singular point.

The structure of the flag (1) may also be described by another sequence of integers. We define the weights at $p, w_{i}=w_{i}(p), i=1, \ldots, n$, by setting $w_{j}=s$ if $n_{s-1}(p)<j \leq n_{s}(p)$, where $n_{0}=0$. In other words, we have

$$
w_{1}=\cdots=w_{n_{1}}=1, w_{n_{1}+1}=\cdots=w_{n_{2}}=2, \ldots, w_{n_{r-1}+1}=\cdots=w_{n_{r}}=r .
$$

Definition 6. A system of privileged coordinates at $p$ is a system of local coordinates $\left(x_{1}, \ldots, x_{n}\right)$ such that $\operatorname{ord}_{p}\left(x_{j}\right)=w_{j}$, for $j=1, \ldots, n$.

On the other hand, given a sequence of integers $\alpha=\left(\alpha_{1}, \ldots, \alpha_{n}\right)$ we define the weight of the monomial $x^{\alpha}=x_{1}^{\alpha_{1}} \cdots x_{n}^{\alpha_{n}}$ to be $w(\alpha)=w_{1} \alpha_{1}+\cdots+w_{n} \alpha_{n}$ and the weighted degree of the monomial vector field $x^{\alpha} \frac{\partial}{\partial x_{j}}$ to be $w(\alpha)-w_{j}$. The weighted degrees allow to compute the orders of functions and vector fields in a purely algebraic way.

Constructions of privileged coordinates can be found in [7] and [15].
Proposition 1 ([15], Proposition 2.2). For a smooth function $f$ with a Taylor expansion in privileged coordinates

$$
f(x) \sim \sum_{\alpha} c_{\alpha} x^{\alpha}
$$

the order of $f$ is the least weighted degree of monomials having a nonzero coefficient in the Taylor series.
For a vector field $X$ with a Taylor expansion in privileged coordinates

$$
X(x) \sim \sum_{\alpha, j} a_{\alpha, j} x^{\alpha} \frac{\partial}{\partial x_{j}}
$$

the order of $X$ is the least weighted degree of a monomial vector field having a nonzero coefficient in the Taylor series.

Remark 2. A vector field of degree $<-r$ vanishes.
The one-parameter family of dilations $\delta_{\lambda}: \mathbb{R}^{n} \longrightarrow \mathbb{R}^{n}$ is defined by $\delta_{\lambda}(x)=$ $\left(\lambda^{w_{1}} x_{1}, \lambda^{w_{2}} x_{2}, \ldots, \lambda^{w_{n}} x_{n}\right), \lambda \geq 0$. A dilation $\delta_{\lambda}$ acts also on functions and vector fields by pull-back: $\delta_{\lambda}^{*} f=f \circ \delta_{\lambda}$ and $\delta_{\lambda}^{*} X$ is the vector field such that $\left(\delta_{\lambda}^{*} X\right)\left(\delta_{\lambda}^{*} f\right)=$ $\delta_{\lambda}^{*}(X f)$. So we have the following definition.

Definition 7. A function $f$ is homogeneous of degree s if $\delta_{\lambda}^{*} f=\lambda^{s} f$. A vector field $X$ is homogeneous of degree $s$ if $\delta_{\lambda}^{*} X=\lambda^{s} X$.

Proposition 2 ([7], Proposition 5.16). Let $X$ and $Y$ be vector fields on $M$. If $X$ and $Y$ are homogeneous of degree $k$ and $l$ respectively (in the chosen system of privileged coordinates) then $[X, Y]$ is homogeneous of degree $k+l$ or vanishes.

Definition 8. The function defined by $x \mapsto\|x\|_{p}=\sum_{i=1}^{n}\left|x_{i}\right|^{\frac{1}{w_{i}}}$ is the so-called pseudo-norm at $p$.

Remark 3. Let $x=\left(x_{1}, \ldots, x_{n}\right)$ be a system of privileged coordinates defined on an open neighborhood $U$ of the point $p$. When composed with the coordinate functions, the pseudo-norm at $p$ is (non smooth) homogeneous of order 1, that is, $\|x(q)\|_{p}=O(d(p, q))$, where $x(q)$ are the coordinates of $q \in U$.

### 2.4 Nilpotent approximation

Fix a system of privileged coordinates $\left(x_{1}, \ldots, x_{n}\right)$ at $p$. Every vector field $X_{i}$ is of order $\geq-1$, hence it has, in $x$ coordinates, a Taylor expansion

$$
X_{i}(x) \sim \sum_{\alpha, j} a_{\alpha, j} x^{\alpha} \frac{\partial}{\partial x_{j}}
$$

where $w(\alpha) \geq w_{j}-1$ if $a_{\alpha, j} \neq 0$. Grouping together the monomial vector fields of same weighted degree we express $X_{i}$ as a series of homogeneous vector fields of the form

$$
\begin{equation*}
X_{i}=X_{i}^{(-1)}+X_{i}^{(0)}+X_{i}^{(1)}+X_{i}^{(2)}+\cdots, \tag{2}
\end{equation*}
$$

where $X_{i}^{(s)}$ has degree $s$. We set

$$
\widehat{X}_{i}=X_{i}^{(-1)}, \quad i=1, \ldots, n
$$

Definition 9. The family of vector fields $\left(\widehat{X}_{1}, \ldots, \widehat{X}_{n}\right)$ is called the nilpotent approximation of the system $\left(X_{1}, \ldots, X_{n}\right)$ at $p$.

Proposition 3 ([7], Proposition 5.17). The vector fields $\widehat{X}_{i}, i=1, \ldots, n$, generate a nilpotent Lie algebra Lie $\left(\widehat{X}_{1}, \ldots, \widehat{X}_{n}\right)$ of step $r=w_{n}$. They satisfy Larc at every point $y \in \mathbb{R}^{n}$, and the distance $\widehat{d}$ is finite for every $x, y \in \mathbb{R}^{n}$.

The following results are found in [7] and [15] and will be useful in this work.

Proposition 4 ([7] Proposition 7.25 and [15], Lemma 2.1). There exist positive constants $C, C^{\prime}$ such that for all $q \in \mathbb{R}^{n}$ we have

$$
C\|q\|_{p} \leq \widehat{d}_{p}(p, q) \leq C^{\prime}\|q\|_{p} .
$$

Lemma 1 ([15], Lemma 2.2). There exists constant $C$ and $\varepsilon>0$, such that, for any $z^{0} \in \mathbb{R}^{n}$ and any $t \in \mathbb{R}^{+}$satisfying $\tau=\max \left(\left\|z^{0}\right\|_{p}, t\right)<\varepsilon$, we have

$$
\|z(t)-\widehat{z}(t)\|_{p} \leq C \tau t^{\frac{1}{r}}
$$

where $z(\cdot)$ and $\widehat{z}(\cdot)$ are trajectories of the nonholonomic systems defined respectively by $X_{1}, \ldots, X_{n}$ and $\widehat{X}_{1}, \ldots, \widehat{X}_{n}$ starting at the same point $z^{0}$, associated with the same control function $u(\cdot)$, and satisfying $\|u(t)\|=1$ a.e.

To finish, we recall the very important Theorem 7.32 of [7] stated here with a slight modification.

Theorem 1 (Theorem 7.32 in [7]). There exist constants $\varepsilon>0$ and $C>0$ such that for any $q, q^{\prime} \in B(p, \varepsilon)$, we have

$$
-C \tau d\left(q, q^{\prime}\right)^{\frac{1}{r}} \leq d\left(q, q^{\prime}\right)-\widehat{d}\left(q, q^{\prime}\right) \leq C \widehat{\tau} \widehat{d}\left(q, q^{\prime}\right)^{\frac{1}{r}},
$$

where $\tau$ is as in Lemma 1 and $\widehat{\tau}$ is similarly defined, this is $\tau=\max \left(\|q\|_{p}, d\left(q, q^{\prime}\right)\right)$ and $\widehat{\tau}=\max \left(\|q\|_{p}, \widehat{d}\left(q, q^{\prime}\right)\right)$.

### 2.5 Linear vector fields

The definition of linear vector fields comes from [6] and [16].
Let $G$ be a connected Lie group and $\mathfrak{g}$ its Lie algebra (the set of left-invariant vector fields, identified with the tangent space at the identity). The set of analytic vector fields on $G$ is denoted by $V^{\omega}(G)$, and the normalizer of $\mathfrak{g}$ in $V^{\omega}(G)$ is by definition

$$
\mathcal{N}=\operatorname{norm}_{V^{\omega}(G)} \mathfrak{g}=\left\{X \in V^{\omega}(G): \forall Y \in \mathfrak{g} \quad[X, Y] \in \mathfrak{g}\right\}
$$

Definition 10. A vector field $\mathcal{X}$ on $G$ is said to be linear or to be infinitesimal automorphism (see [13]), if $\mathcal{X}$ belongs to $\mathcal{N}$ and $\mathcal{X}(e)=0$, where $e$ is the identity of $G$.

We can see in [16] that a vector field $\mathcal{X}$ on $G$ if and only its flow $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ is a one-parameter group of automorphisms of $G$ and a linear vector field is consequently analytic and complete.

### 2.5.1 Simple ARS's on Lie groups

Linear and invariant vector fields make it possible to define almost-Riemannian structures on Lie groups. The following definition is given in [5].

Definition 11. A simple $A R S$ is an $A R S$ defined on a connected Lie group $G$ by a set of $n$ vector fields $\left\{\mathcal{X}, Y_{1}, \ldots, Y_{n-1}\right\}$ where $\mathcal{X}$ is linear, $Y_{1}, \ldots, Y_{n-1}$ are leftinvariant, $\operatorname{dim} G=n$ and the rank of $\mathcal{X}, Y_{1}, \ldots, Y_{n-1}$ is full on a non empty subset of $G$ and the set $\left\{\mathcal{X}, Y_{1}, \ldots, Y_{n-1}\right\}$ satisfies Larc.

For instance, the famous Grushin plane on the Abelian Lie group $\mathbb{R}^{2}$ is a simple ARS. This structure was introduced in [5] and its isometries have been study in [17].

In Section 3.3 a 3-dimensional example will be provided.

### 2.5.2 Simple ARS's on homogeneous spaces

Consider a homogeneous space $G / H$ of a connected and simply connected Lie group $G$ by a closed subgroup $H$ (the elements of $G / H$ are right cosets of $H$ because we deal with left-invariant vector fields). Since we are interested in simply connected quotients we assume $H$ to be connected. Let $\mathfrak{g}$ be the Lie algebra of $G$, identified with the space of left-invariant vector fields. The projection of such a vector field $Y$ on $G / H$ is well-defined, is referred to as a left-invariant vector field, and we can assume that it vanishes identically only if $Y$ is the zero field (see details in [16]). On the other hand the projection of a linear field $\mathcal{X}$ of $G$ does exist on $G / H$ if and only if $H$ is invariant under its flow, or equivalently, because $H$ is connected, if the Lie algebra of $H$ is $\operatorname{ad}(\mathcal{X})$-invariant. This allows to define linear vector fields and simple ARS on $G / H$ :

Let $Y_{1}, \ldots, Y_{n-1}, \mathcal{X}$ be a set of $n=\operatorname{dim}(G / H)$ vectors fields on $G / H$, where $Y_{1}, \ldots, Y_{n-1}$ are invariant and $\mathcal{X}$ is linear. It defines a simple ARS if

1. They satisfy Larc.
2. The singular set $\mathcal{Z}$ where their rank is less than $n$ is proper with empty interior.

In the sequel, we will need (a simplified version of) the equivalence Theorem (see [16] and [5]).

Theorem 2 (Equivalence Theorem). Let $f_{1}, \ldots, f_{n}$ be a set of $n$ complete vector fields on $\mathbb{R}^{n}$ and let us assume:

1. $f_{1}, \ldots, f_{n}$ define an Almost-Riemannian Structure on $\mathbb{R}^{n}$;
2. The Lie algebra $\mathcal{L}$ generated by $f_{1}, \ldots, f_{n}$ is finite dimensional;
3. The ideal $\mathfrak{g}$ generated in $\mathcal{L}$ by $f_{1}, \ldots, f_{n-1}$ is nilpotent and of codimension 1 in $\mathcal{L}$.

Then $\mathbb{R}^{n}$ is diffeomorphic to a homogeneous space $G / H$ of the nilpotent simply connected group $G$ generated by $\mathfrak{g}$ and $f_{1}, \ldots, f_{n}$ defines a simple ARS on $G / H$. More accurately the vector fields $f_{1}, \ldots, f_{n-1}$ are left-invariant and $f_{n}$ is linear on this homogeneous space.

## 3 Solvable Approximation

In this section we introduce the solvable approximation of an ARS and we analyze its algebraic structure.

### 3.1 Definition

Let $\left\{X_{1}, \ldots, X_{n}\right\}$ be a set of vector fields defining an almost-Riemannian structure on an open neighborhood of $0 \in \mathbb{R}^{n}$. The point $p=0$ is assumed to belong to the singular locus, the natural coordinates of $\mathbb{R}^{n}$ to be privileged and we consider the nilpotent approximation $\left\{\widehat{X}_{1}, \ldots, \widehat{X}_{n}\right\}$ of $\left\{X_{1}, \ldots, X_{n}\right\}$ at $p=0$.

It may happen that some of the vector fields $\widehat{X}_{i}$ vanish, possibly changing the almost-Riemannian structure defined by $X_{1}, \ldots, X_{n}$ into a constant rank subRiemannian one. It is what happens in some cases of generic 3-dimensional ARSs that are described in detail in Section 3.3. In what follows we are interested in the case where only one of the $\widehat{X}_{i}$ 's vanishes, say $\widehat{X}_{n}=0$, and the other ones are independent and define a left-invariant sub-Riemannian structure on a Lie group, or a homogeneous space, the underlying manifold of which is $\mathbb{R}^{n}$. Recall that each $X_{i}$ can be expanded in a series of homogeneous vector fields in the system of privileged coordinates at $p=0$, this is

$$
X_{i}=X_{i}^{(-1)}+X_{i}^{(0)}+X_{i}^{(1)}+\ldots, \quad \forall i \in\{1, \ldots, n\}
$$

where $X_{i}^{(k)}$ is the homogeneous component of degree $k$. Denoting $\widetilde{X}_{n}=X_{n}^{(0)}$, we introduce the following definition:

Definition 12 (Solvable approximation). The family $\left\{\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}, \widetilde{X}_{n}\right\}$ is the solvable approximation of $\left\{X_{1}, \ldots, X_{n}\right\}$.

Proposition 5. $\mathcal{L}=\operatorname{Lie}\left(\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}, \widetilde{X}_{n}\right)$ is a finite dimensional solvable Lie algebra. Its step of solvability is less than or equal to $\log _{2}(r)+1$, where $r$ is the degree of nonholonomy at $p=0$.

Proof. Let $\mathcal{D}^{k} \mathcal{L}$ stand for the $k^{\text {th }}$ derived algebra of $\mathcal{L}$, with $\mathcal{L}=\mathcal{D}^{0} \mathcal{L}$. According to Proposition 2 and Remark 2 the algebra $\mathcal{L}$ is generated by homogeneous vector fields of degree $0,-1, \ldots,-r$ because the $\widehat{X}_{i}$ 's are homogeneous of degree -1 , for $i=1, \ldots, n-1$, and $\widetilde{X}_{n}$ is homogeneous of degree 0 . According to Proposition 2 again $\mathcal{D}^{1} \mathcal{L}$ is generated by homogeneous vector fields of degree $-1, \ldots,-r$. More generality $\mathcal{D}^{s} \mathcal{L}$ is generated by homogeneous vector fields of degree $-2^{s-1},-2^{s}, \ldots,-r$, so that $\mathcal{D}^{s} \mathcal{L}=0$ if $2^{s}>r$. Therefore $\mathcal{L}$ is solvable and the step of solvability $\sigma$ of $\mathcal{L}$ satisfies $\sigma \leq \log _{2}(r)+1$. On the other hand, the Lie algebra $\mathcal{L}$ splits into homogeneous components

$$
\mathcal{L}=\mathcal{L}^{0} \oplus \mathcal{L}^{-1} \oplus \mathcal{L}^{-2} \oplus \ldots \oplus \mathcal{L}^{-r}
$$

where $\mathcal{L}^{-s}$ is the set of homogeneous vector fields of degree $-s$ under the action of $\delta_{\lambda}$. A homogeneous vector field $X$ of degree $w \in\{0,-1,-2, \ldots,-r\}$ writes in coordinates $X(x)=\sum_{i=1}^{n} f_{i}(x) \frac{\partial}{\partial x_{i}}$, where $f_{i}(x)$ is a homogeneous polynomial function of degree $w+\operatorname{ord}\left(x_{i}\right)$. Since the space of polynomials of degree $w+\operatorname{ord}\left(x_{i}\right)$ is finite dimensional, $\mathcal{L}^{-w}$ is finite dimensional. Therefore $\mathcal{L}$ is finite dimensional.

Remark 4. It is clear that the families of vector fields $\left\{\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}, \widetilde{X}_{n}\right\}$ and $\left\{X_{1}, \ldots, X_{n}\right\}$ have the same nilpotent approximation. Consequently the family of vector fields $\left\{\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}, \widetilde{X}_{n}\right\}$ satisfies Larc on $\mathbb{R}^{n}$ and the growth vector at 0 of $\left\{\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}, \widetilde{X}_{n}\right\}$ is equal to the one of $\left\{X_{1}, \ldots, X_{n}\right\}$.

### 3.2 Structure of the approximating system

Fundamental remark. Despite the previous result we are not interested in the solvable Lie group associated to the Lie algebra Lie $\left\{\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}, \widetilde{X}_{n}\right\}$ but in some nilpotent Lie group on which $\widetilde{X}_{n}$ acts as a linear vector field.

For this reason we denote by $\mathfrak{h}$ the Lie algebra generated by $\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}$ and by $\mathfrak{g}$ the ideal generated by $\mathfrak{h}$ in $\mathcal{L}=\operatorname{Lie}\left\{\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}, \widetilde{X}_{n}\right\}$.

Proposition 6. The ideal $\mathfrak{g}$ is the space of vector fields of $\mathcal{L}$ whose nonholonomic order is negative. It is a nilpotent Lie algebra and

$$
\mathcal{L}=\mathfrak{g} \oplus \mathbb{R} \widetilde{X}_{n}
$$

Moreover $D=-\operatorname{ad}\left(\widetilde{X}_{n}\right)$ is a derivation of $\mathfrak{g}$.
Proof. Since $\mathfrak{g}$ is the ideal generated by $\mathfrak{h}$ in $\mathcal{L}$, we have $\mathcal{L}=\mathbb{R} \widetilde{X}_{n} \oplus \mathfrak{g}$ and ad $\left(\widetilde{X}_{n}\right)$ is a derivation of $\mathfrak{g}$.

Let $G$ be the simply connected Lie group whose Lie algebra is isomorphic to $\mathfrak{g}$. According to [16] there exists a linear vector field on $G$ associated to the derivation $D=-\operatorname{ad}\left(\widetilde{X}_{n}\right)$. With a clear abuse of notation we will denote it by $\widetilde{X}_{n}$. Thanks to the equivalence theorem we have the following:

Theorem 3. The space $\mathbb{R}^{n}$ is diffeomorphic to a homogeneous space $G / G_{0}$ of $G$. Through this diffeomorphism $\left\{\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}, \widetilde{X}_{n}\right\}$ is equivalent to a simple $A R S$ on $G / G_{0}$, and the Lie algebra $\mathfrak{g}_{0}$ of $G_{0}$ is isomorphic to the set of vector fields of $\mathfrak{g}$ that vanish at 0 .
Proof. First of all, notice that the vector fields $\left\{\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}, \widetilde{X}_{n}\right\}$ are complete from their triangular form, namely, in the equation $\dot{x}=\sum_{i=1}^{n} u_{i} \widehat{X}_{i}(x)+u_{n} \widetilde{X}_{n}(x), \dot{x}_{j}$ is linear with respect to the coordinates of weight $w_{j}$ and polynomial with respect to coordinates of weight $<w_{j}$ (see [7] or [15] for details). They define an ARS, hence in particular satisfy Larc and generate a finite dimensional Lie algebra (Proposition 5). According to Theorem $2 \mathbb{R}^{n}$ is diffeomorphic to a homogeneous space $G / G_{0}$ of $G$, where $G_{0}$ is the connected subgroup of $G$ whose Lie algebra is, after identification of $L(G)$ with $\mathfrak{g}$, the set of elements of $\mathfrak{g}$ that vanish at 0 . Thanks to the diffeomorphism between $\mathbb{R}^{n}$ and $G / G_{0}$ the system $\left\{\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}, \widetilde{X}_{n}\right\}$ can be identified to a simple ARS on $G / G_{0}$.

We are also interested in conditions for which $G=\mathbb{R}^{n}$.
Theorem 4. With the previous notations the following assertions are equivalent:
(i) ad $\left(\widetilde{X}_{n}\right) \cdot \widehat{X}_{i}$ belongs to Span $\left\{\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}\right\}$ for $i=1, \ldots, n-1$;
(ii) $\mathfrak{h}$ is ad $\left(\widetilde{X}_{n}\right)$-invariant;
(iii) $\mathfrak{h}=\mathfrak{g}$.

Under these conditions $\widetilde{X}_{n}$ is a linear vector field on $\exp (\mathfrak{h})$.
Proof.
(i) $\Rightarrow$ (ii) It is an immdiate consequence of the Jacobi identity.
(ii) $\Rightarrow$ (iii) Condition (ii) implies $\mathcal{L}=\mathfrak{h} \oplus \mathbb{R} \widetilde{X}_{n}$, hence $\mathfrak{h}=\mathfrak{g}$.
(iii) $\Rightarrow$ (i) Condition (iii) implies that $\mathfrak{h}$ is ad $\left(\widetilde{X}_{n}\right)$-invariant. According to Proposition 2 the set of elements of $\mathfrak{h}$ of order -1 is $\operatorname{Span}\left\{\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}\right\}$. For $i=$ $1, \ldots, n-1$ the bracket $\left[\widehat{X}_{i}, \widetilde{X}_{n}\right]$ belongs to $\mathfrak{h}$, since $\mathfrak{h}=\mathfrak{g}$, and is of order -1 or is equal to 0 . Therefore it belongs to $\operatorname{Span}\left\{\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}\right\}$.

### 3.3 Example. The 3D-generic case

The local representation of a generic ARS in dimension 3 is detailed in Section 4.2. Its nilpotent approximation at a point of the singular locus is the following:

$$
\widehat{X}_{1}=\left(\begin{array}{l}
1 \\
0 \\
0
\end{array}\right), \quad \widehat{X}_{2}=\left(\begin{array}{c}
0 \\
1 \\
x \cos \sigma
\end{array}\right), \quad \widehat{X}_{3}=\left(\begin{array}{c}
0 \\
0 \\
x \sin \sigma
\end{array}\right), \quad \text { where } \sigma \in\left[0, \frac{\pi}{2}\right]
$$

There are two particular cases, according to the value of the parameter $\sigma$. The first one for $\sigma=\frac{\pi}{2}$, because the bracket of $\widehat{X}_{1}$ and $\widehat{X}_{2}$ vanishes, and the second one for $\sigma=0$, because $\widehat{X}_{3}$ vanishes, it is the tangent case.

More accurately the Lie brackets are:

$$
\left[\widehat{X}_{1}, \widehat{X}_{2}\right]=\left(\begin{array}{c}
0 \\
0 \\
\cos \sigma
\end{array}\right), \quad\left[\widehat{X}_{1}, \widehat{X}_{3}\right]=\left(\begin{array}{c}
0 \\
0 \\
\sin \sigma
\end{array}\right), \quad\left[\widehat{X}_{2}, \widehat{X}_{3}\right]=0 .
$$

The analysis of the different cases is as follows:

1. General case $\sigma \in] 0, \frac{\pi}{2}[$.

We can set $Z=\left(\begin{array}{c}0 \\ 0 \\ \cos \sigma\end{array}\right)$ so that $\left[\widehat{X}_{1}, \widehat{X}_{2}\right]=Z$ and $\left(\widehat{X}_{1}, \widehat{X}_{2}, Z\right)$ is the Heisenberg Lie algebra. Since $\widehat{X}_{3}$ vanishes at ( $0,0,0$ ) and since its Lie brackets with $\widehat{X}_{1}, \widehat{X}_{2}$ and $Z$ are

$$
\left[\widehat{X}_{1}, \widehat{X}_{3}\right]=\tan (\sigma) Z \quad \text { and } \quad\left[\widehat{X}_{2}, \widehat{X}_{3}\right]=\left[Z, \widehat{X}_{3}\right]=0
$$

it is a linear vector field on the Heisenberg group, the associated derivation of which is

$$
D=\left(\begin{array}{ccc}
0 & 0 & 0 \\
0 & 0 & 0 \\
\tan \sigma & 0 & 0
\end{array}\right) .
$$

The conclusion is that $\left(\widehat{X}_{1}, \widehat{X}_{2}, \widehat{X}_{3}\right)$ defines a simple ARS on the 3D Heisenberg group.
2. Particular case $\sigma=\frac{\pi}{2}$. Then

$$
\widehat{X}_{1}=\left(\begin{array}{l}
1 \\
0 \\
0
\end{array}\right), \quad \widehat{X}_{2}=\left(\begin{array}{l}
0 \\
1 \\
0
\end{array}\right), \quad \widehat{X}_{3}=\left(\begin{array}{l}
0 \\
0 \\
x
\end{array}\right) .
$$

It is a simple ARS on the Abelian Lie group $\mathbb{R}^{3}$. Indeed $\widehat{X}_{1}$ and $\widehat{X}_{2}$ are (left and right) invariant and $\widehat{X}_{3}$ is linear.
3. Tangent case $\sigma=0$. Here the nilpotent approximation degenerates into the following sub-Riemannian structure on the Heisenberg group

$$
\widehat{X}_{1}=\left(\begin{array}{l}
1 \\
0 \\
0
\end{array}\right), \quad \widehat{X}_{2}=\left(\begin{array}{l}
0 \\
1 \\
x
\end{array}\right), \quad \widehat{X}_{3}=\left(\begin{array}{l}
0 \\
0 \\
0
\end{array}\right) .
$$

However we will see after the next remark that in case where the component $\widetilde{X}_{3}$ of order 0 of the vector field $X_{3}$ does not vanish then $\left(\widehat{X}_{1}, \widehat{X}_{2}, \widetilde{X}_{3}\right)$ defines a simple ARS on a homogeneous space.

Remark 5. In the cases $\sigma=\frac{\pi}{2}$ and $\left.\sigma \in\right] 0, \frac{\pi}{2}[$ the Lie algebra generated by $\widehat{X}_{1}, \widehat{X}_{2}, \widehat{X}_{3}$ is 4-dimensional and two points of view are possible. The usual one consists in considering $\mathbb{R}^{3}$ as a homogeneous space of a nilpotent 4-dimensional Lie group. Our point of view is to consider $\widehat{X}_{3}$ as a linear vector field on the 3dimensional Lie group $\mathbb{R}^{3}$ endowed with the Abelian structure if $\sigma=\frac{\pi}{2}$ and the Heisenberg one if $\sigma \in] 0, \frac{\pi}{2}[$.

Following this way we will consider the solvable approximation whenever $\sigma=0$, and finally all approximations of generic $3 D-A R S$ will appear as being simple $A R S$.

The solvable approximation of the tangent case $\sigma=0$.
The homogeneous component of nonholonomic order 0 of $X_{3}$ is

$$
\widetilde{X}_{3}=\left(\begin{array}{c}
0 \\
0 \\
a z+b x^{2}+c y^{2}
\end{array}\right)=\left(a z+b x^{2}+c y^{2}\right) \frac{\partial}{\partial z} \quad \text { (See Section } 4.2 \text { again). }
$$

As well as in the general case the Lie algebra generated by $\widehat{X}_{1}$ and $\widehat{X}_{2}$ is:

$$
\mathfrak{h}=\operatorname{Span}\left\{\widehat{X}_{1}, \widehat{X}_{2}, Z=\left[\widehat{X}_{1}, \widehat{X}_{2}\right]\right\} \quad \text { where } \quad Z=\left(\begin{array}{l}
0 \\
0 \\
1
\end{array}\right)=\frac{\partial}{\partial z},
$$

that is the Heisenberg algebra. On the other hand the algebra generated by $\widehat{X}_{1}, \widehat{X}_{2}$ and $\widetilde{X}_{3}$ is Span $\left\{\widehat{X}_{1}, \widehat{X}_{2}, Z,\left[\widehat{X}_{1}, \widetilde{X}_{3}\right],\left[\widehat{X}_{2}, \widetilde{X}_{3}\right], \widetilde{X}_{3}\right\}$, where:

$$
\left[\widehat{X}_{1}, \widetilde{X}_{3}\right]=\left(\begin{array}{c}
0 \\
0 \\
2 b x
\end{array}\right)=2 b x \frac{\partial}{\partial z} \text { and }\left[\widehat{X}_{2}, \widetilde{X}_{3}\right]=\left(\begin{array}{c}
0 \\
0 \\
2 c y+a x
\end{array}\right)=(2 c y+a x) \frac{\partial}{\partial z},
$$

and the ideal generated by $\widehat{X}_{1}$ and $\widehat{X}_{2}$ is:

$$
\mathfrak{g}=\operatorname{Span}\left\{\frac{\partial}{\partial x}, \frac{\partial}{\partial y}+x \frac{\partial}{\partial z}, 2 b x \frac{\partial}{\partial z}, 2 c y \frac{\partial}{\partial z}+a x \frac{\partial}{\partial z}, \frac{\partial}{\partial z}\right\} .
$$

A straightforward computation shows that $\widetilde{X}_{3}$ acts as a derivation on $\mathfrak{g}$. If we assume $b \neq 0$ and $c \neq 0$ then we have also:

$$
\mathfrak{g}=\operatorname{Span}\left\{\frac{\partial}{\partial x}, x \frac{\partial}{\partial z}, \frac{\partial}{\partial y}, y \frac{\partial}{\partial z}, \frac{\partial}{\partial z}\right\} .
$$

This is the 5 -dimensional Heisenberg Lie algebra $\mathfrak{h}^{2}$ and in this basis the derivation $D=-\operatorname{ad}\left(\widetilde{X}_{3}\right)$ is given by the following matrix:

$$
D=\left(\begin{array}{ccccc}
0 & 0 & & & \\
2 b & a & & & \\
& & 0 & 0 & \\
& & 2 c & a & \\
& & & & a
\end{array}\right)
$$

Finally the solvable approximation $\left(\widehat{X}_{1}, \widehat{X}_{2}, \widetilde{X}_{3}\right)$ is a simple ARS on $\mathbb{R}^{3}$ diffeomorphic to a quotient of the 5 -dimensional group Heisenberg $\mathbb{H}^{2}$.

## 4 Distances

We can distinguish three different families of vector fields from the above section: $\left\{X_{1}, X_{2}, \ldots, X_{n}\right\},\left\{\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}\right\}$ and $\left\{\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}, \widetilde{X}_{n}\right\}$ which satisfy Larc. Assuming orthonormality, they define three different distances: $d, \widehat{d}$ and $\widetilde{d}$ respectively,
where $\widetilde{d}$ and $\widehat{d}$ are defined on $\mathbb{R}^{n}$. This section is divided in two parts. In the first one, we give two propositions about the almost-Riemannian distance $\widetilde{d}$ defined by the solvable approximation. The second part is devoted to analyze $\widetilde{d}$ in the 3 -dimensional generic case.

### 4.1 The almost-Riemannian distance $\widetilde{d}$

The following proposition establishes a relation between $\widetilde{d}$ and $\widehat{d}$. It is important because it allows us to find an upper bound for $\widetilde{d}$ (see Section 4.2.2) and to compare the distances of the solvable and nilpotent approximation.
Proposition 7. For all $x, y \in \mathbb{R}^{n}, \widetilde{d}(x, y) \leq \widehat{d}(x, y)$.
Proof. Let $x, y \in \mathbb{R}^{n}$ and let $\gamma$ be a minimizing geodesic for $\dot{x}=\sum_{i=1}^{n-1} u_{i} \widehat{X}_{i}$, such that $\gamma(0)=x, \gamma(T)=y$. Setting $u_{n}=0$ the curve $\gamma$ is admissible for $\dot{x}=$ $\sum_{i=1}^{n-1} u_{i} \widehat{X}_{i}+u_{n} \widetilde{X}_{n}$. Since $u_{n}=0$ the length of $\gamma$ is the same for both metrics, hence $\widetilde{d}(x, y) \leq l(\gamma)=\widehat{d}(x, y)$.

Let $\delta_{\lambda}$ be the dilation related to the privileged coordinates and the weights at $p=0$. We know that the distance $\widehat{d}$ is homogeneous of degree 1 with respect to $\delta_{\lambda}$ (see [15]). However $\widetilde{d}$ does not possess this property. This is due to the fact that $\widetilde{X}_{n}$ and the $\widehat{X}_{i}$ 's do not have the same degree of homogeneity.
Proposition 8. The almost-Riemannian distance $\widetilde{d}$ is not homogeneous.
Proof. Let $\gamma$ be an admissible curve for $\widetilde{d}$, that is

$$
\begin{equation*}
\dot{\gamma}(t)=\sum_{i=1}^{n-1} u_{i} \widehat{X}_{i}(\gamma(t))+u_{n} \widetilde{X}_{n}(\gamma(t)) . \tag{3}
\end{equation*}
$$

Since $\widehat{X}_{i}$ and $\widetilde{X}_{n}$ are homogeneous of degree -1 and 0 respectively and the pullback by $\delta_{\lambda}$ of a vector field $X$ (see [1]) is defined by

$$
\begin{equation*}
d \delta_{\lambda}(q)\left(\delta_{\lambda}^{*} X(q)\right)=X\left(\delta_{\lambda}(q)\right), \tag{4}
\end{equation*}
$$

we get

$$
\begin{align*}
d \delta_{\lambda}(q) \cdot \widehat{X}_{i}(q)=d \delta_{\lambda}(q) \cdot \lambda \cdot \delta_{\lambda}^{*} \widehat{X}_{i}(q) & =\lambda \widehat{X}_{i}\left(\delta_{\lambda}(q)\right) \quad \text { and } \\
d \delta_{\lambda}(q) \cdot \widetilde{X}_{n}(q)=d \delta_{\lambda}(q) \cdot \delta_{\lambda}^{*} \widetilde{X}_{n}(q) & =\widetilde{X}_{n}\left(\delta_{\lambda}(q)\right) . \tag{5}
\end{align*}
$$

Therefore

$$
\begin{aligned}
\frac{d}{d t}\left(\delta_{\lambda} \circ \gamma\right)(t) & =d \delta_{\lambda}(\gamma(t)) \cdot \dot{\gamma}(t) \\
& =\sum_{i=1}^{n-1} u_{i} d \delta_{\lambda}(\gamma(t)) \cdot \widehat{X}_{i}(\gamma(t))+u_{n} d \delta_{\lambda}(\gamma(t)) \cdot \widetilde{X}_{n}(\gamma(t)) \\
& =\sum_{i=1}^{n-1} \lambda u_{i} \widehat{X}_{i}\left(\delta_{\lambda}(\gamma(t))\right)+u_{n} \widetilde{X}_{n}\left(\delta_{\lambda}(\gamma(t))\right) .
\end{aligned}
$$

This implies that $l\left(\delta_{\lambda} \gamma\right) \neq \lambda l(\gamma)$, except if $u_{n}(t)$ vanishes a.e. This proves the non homogeneity of $\widetilde{d}$.

### 4.2 The 3D-tangential case

In Section 3, we have established a model to locally approximate an $n$-ARS whose nilpotent approximation is a constant rank sub-Riemannian structure, by a solvable approximation. In this context we want to determine conditions for $|d-\widetilde{d}|$ to be smaller than $\mid d-\widehat{d \mid}$.

Recall that $\Delta(p)=\operatorname{span}\left\{X_{1}(p), \ldots, X_{n}(p)\right\}$ and the singular locus $\mathcal{Z}$ is the set of points of $\mathbb{R}^{n}$ where the rank of the linear span of the vector fields is less than $n$. From [10] we take the following.
Proposition 9. Consider a 3-ARS. The following conditions are generic for 3-ARSs
(G1) $\operatorname{dim}(\Delta(p)) \geq 2$ and $\Delta(p)+[\Delta(p), \Delta(p)]=T_{p} M$ for every $p \in M$;
(G2) $\mathcal{Z}$ is an embedded (possibly empty) two-dimensional submanifold of $M$;
(G3) the points where $\Delta(p)=T_{p} \mathcal{Z}$ are isolated.
Proposition 10. Under the previous conditions there are three types of points:

1. Riemannian points where $\Delta(p)=T_{p} M$.
2. type-1 points where $\Delta(p)$ has dimension 2 and is transversal to $\mathcal{Z}$.
3. type-2 points where $\Delta(p)$ has dimension 2 and is tangent to $\mathcal{Z}$.

Moreover type-2 points are isolated, type-1 points form a 2 dimensional manifold and all other points are Riemannian.

The local representation of the 3-dimensional ARS at type-2 points is given by the vector fields
$X_{1}=\left(\begin{array}{l}1 \\ 0 \\ 0\end{array}\right), \quad X_{2}=\left(\begin{array}{c}0 \\ 1+\delta(x, y, z) \\ x(1+\theta(x, y, z))\end{array}\right), \quad X_{3}=\left(\begin{array}{c}0 \\ 0 \\ a z+b x^{2}+c y^{2}+o\left(x^{2}+y^{2}+|z|\right)\end{array}\right)$,
where $\delta$ and $\theta$ are smooth functions of order greater than or equal to 1 and $a, b, c$ are not all zero. Furthermore, from Subsection 3.3, the nilpotent approximation in privileged coordinates is

$$
\widehat{X}_{1}=X_{1}=\left(\begin{array}{l}
1 \\
0 \\
0
\end{array}\right), \quad \widehat{X}_{2}=\left(\begin{array}{l}
0 \\
1 \\
x
\end{array}\right), \quad \widehat{X}_{3}=0 .
$$

and

$$
\widetilde{X}_{3}=\left(\begin{array}{c}
0 \\
0 \\
a z+b x^{2}+c y^{2}
\end{array}\right) .
$$

$\left(\widehat{X}_{1}, \widehat{X}_{2}, \widetilde{X}_{3}\right)$ is the solvable approximation at $p=0$ in case when 0 is a tangential (type-2) point.

### 4.2.1 Divergence of curves

Let $p=0$ be a type-2 point such that the coordinates centered at $p$ are privileged and $q, q^{\prime}$ belonging to the ball centered at $p$ and radius $\epsilon$, denoted by $B(p, \epsilon)$.

In this subsection we analyze the divergence of curves respectively admissible for $d$ and $\widetilde{d}$, defined by the same control functions and starting at the same point $q$. More accurately: let $\gamma$ be the geodesic for $d$ such that $\gamma(0)=q, \gamma(T)=q^{\prime}$ with $u_{1}^{2}+u_{2}^{2}+u_{3}^{2}=1$ and let $\widetilde{\gamma}$ be the admissible curve for $\widetilde{d}$ defined by the same control functions as $\gamma$ with $\widetilde{\gamma}(0)=q$. We have the following:

$$
\begin{aligned}
\dot{\gamma}(t)-\dot{\tilde{\gamma}}(t) & =u_{1}\left(X_{1}(\gamma)-\widehat{X}_{1}(\widetilde{\gamma})\right)+u_{2}\left(X_{2}(\gamma)-\widehat{X}_{2}(\widetilde{\gamma})\right)+u_{3}\left(X_{3}(\gamma)-\widetilde{X}_{3}(\widetilde{\gamma})\right) \\
\left(\begin{array}{c}
\dot{x}(t)-\dot{\widetilde{x}}(t) \\
\dot{y}(t)-\tilde{\tilde{y}}(t) \\
\dot{z}(t)-\dot{\tilde{z}}(t)
\end{array}\right) & =\left(\begin{array}{c}
0 \\
u_{2} \delta(x, y, z) \\
u_{2} x \theta(x, y, z)+u_{3}\left(a(z-\widetilde{z})+b\left(x^{2}-\widetilde{x}^{2}\right)+c\left(y^{2}-\widetilde{y}^{2}\right)+o\left(x^{2}+y^{2}+|z|\right)\right)
\end{array}\right)
\end{aligned}
$$

We have successively:

- $\dot{x}(t)=\dot{\widetilde{x}}(t)$, hence $x(t)=\widetilde{x}(t)$.
- $\dot{y}(t)-\dot{\widetilde{y}}(t)=u_{2} \delta(x, y, z)$, hence $y(t)-\widetilde{y}(t)=\int_{0}^{t} u_{2}(s) \delta(x, y, z) d s$.

We denote by $\rho \geq 1$ the order of $\delta$. Then $|\delta(x, y, z)| \leq$ Cst $\cdot\|\gamma(s)\|_{p}^{\rho} \leq \mathrm{Cst} \cdot \tau^{\rho}$ because $\|\gamma(s)\|_{p} \leq$ Cst $\cdot \tau$, where $\tau=\max \left(\|q\|_{p}, t\right)$ (the proof of the above inequality is given in the proof of Lemma 1 of [15]), hence

$$
\begin{equation*}
|y(t)-\widetilde{y}(t)| \leq \int_{0}^{t} \mathrm{Cst} \cdot\|\gamma(s)\|_{p}^{\rho} d s \leq \int_{0}^{t} \mathrm{Cst} \cdot \tau^{\rho} d s=\mathrm{Cst} \cdot \tau^{\rho} \cdot t \tag{6}
\end{equation*}
$$

- $\dot{z}(t)-\dot{\widetilde{z}}(t)=u_{2} x \theta(x, y, z)+u_{3}\left(a(z-\widetilde{z})+c\left(y^{2}-\widetilde{y}^{2}\right)+o\left(x^{2}+y^{2}+|z|\right)\right)$, hence
$z(t)-\widetilde{z}(t)=\int_{0}^{t} u_{2}(s) x \theta(x, y, z) d s+\int_{0}^{t} u_{3}(s) a(z-\widetilde{z}) d s+\int_{0}^{t} u_{3}(s) c\left(y^{2}-\widetilde{y}^{2}\right) d s+$ $\int_{0}^{t} u_{3}(s) o\left(x^{2}+y^{2}+|z|\right) d s$.
Since $(x, y, z)$ are privileged coordinates at 0 , then $x^{2}+y^{2}+|z| \leq C \cdot d(0,(x, y, z))^{2}$. Moreover, if $f((x, y, z))=o\left(x^{2}+y^{2}+|z|\right)$, then $f((x, y, z))=o\left(d(0,(x, y, z))^{2}\right)$. This implies that $\operatorname{ord}_{p}(f)>2$, hence $f(\gamma(t))=O\left(d(0, \gamma(t))^{3}\right)$. Therefore $|f(\gamma(t))| \leq$ Cst $\cdot \tau^{3}$. On the other hand, let us denote by $m \geq 1$ the order of $\theta$. Then $|x \cdot \theta(x, y, z)| \leq$ Cst $\cdot \tau \cdot\|\gamma(s)\|_{p}^{m} \leq$ Cst $\cdot \tau^{m+1}$ because $\dot{x}=u_{1}$, hence $|x| \leq$ Cst $\cdot t+\|q\|_{p} \leq$ Cst $\cdot \tau$. Also notice

$$
\left|y^{2}-\widetilde{y}^{2}\right|=|y+\widetilde{y}||y-\widetilde{y}| \leq \mathrm{Cst} \cdot \tau^{\rho} \cdot t \cdot \tau=\mathrm{Cst} \cdot t \cdot \tau^{\rho+1}
$$

because $\dot{y}=u_{2}(1+\delta(x, y, z))$, hence $|y| \leq$ Cst $\cdot t+\|q\|_{p} \leq$ Cst $\cdot \tau$. Similarly for $\widetilde{y}$.
Then

$$
\begin{align*}
& |z(t)-\widetilde{z}(t)| \leq \operatorname{Cst} \cdot t \cdot \tau^{m+1}+\operatorname{Cst} \cdot t^{2} \cdot \tau^{\rho+1}+\operatorname{Cst} \cdot t \cdot \tau^{3}+\int_{0}^{t}|a||z-\widetilde{z}| d s \\
& |z(t)-\widetilde{z}(t)| \leq \mathrm{Cst} \cdot\left(t \cdot \tau^{m+1}+t^{2} \cdot \tau^{\rho+1}+t \cdot \tau^{3}\right) e^{|a| t} \\
& |z(t)-\widetilde{z}(t)| \leq \operatorname{Cst} \cdot t \cdot \tau^{\min (m+1, \rho+1,3)} . \tag{7}
\end{align*}
$$

Finally,

$$
\begin{align*}
\|\gamma(t)-\widetilde{\gamma}(t)\|_{p} & =|x(t)-\widetilde{x}(t)|+|y(t)-\widetilde{y}(t)|+|z(t)-\widetilde{z}(t)|^{\frac{1}{2}} \\
& \leq \operatorname{Cst}\left(t \cdot \tau^{\rho}+t^{\frac{1}{2}} \cdot \tau^{\frac{\min (m+1, p+1,3)}{2}}\right) \\
& \leq \operatorname{Cst} \cdot t^{\frac{1}{2}} \cdot \tau^{\frac{\min (m+1, p+1,3)}{2}} . \tag{8}
\end{align*}
$$

Remark 6. The order of $\delta(x, y, z)$ does not change the inequality

$$
\|\gamma(t)-\widehat{\gamma}(t)\|_{p} \leq C s t \cdot \tau \cdot t^{\frac{1}{2}}
$$

that comes from Lemma 1. Indeed,

$$
\|\gamma(t)-\widehat{\gamma}(t)\|_{p} \leq \operatorname{Cst}\left(t \cdot \tau^{\rho}+t^{\frac{1}{2}} \cdot \tau\right) \leq C s t \cdot t^{\frac{1}{2}} \cdot \tau
$$

### 4.2.2 Upper bounds

In order to state our main result in the next section, we need upper bounds for the distances $d$ and $\widetilde{d}$.

We know that the distance $\widehat{d}$ is left-invariant, this is to say, $\widehat{d}\left(q, q^{\prime}\right)=\widehat{d}\left(a \cdot q, a \cdot q^{\prime}\right)$, for all $a \in \mathbb{R}^{3}$. Here $\cdot$ stands for the Heisenberg product defined by

$$
\begin{equation*}
(x, y, z) \cdot\left(x^{\prime}, y^{\prime}, z^{\prime}\right)=\left(x+x^{\prime}, y+y^{\prime}, z+z^{\prime}+x y^{\prime}\right) \tag{9}
\end{equation*}
$$

Recall also that $\widetilde{d}\left(q, q^{\prime}\right) \leq \widehat{d}\left(q, q^{\prime}\right)$. Considering the above, we have

$$
\begin{equation*}
\widetilde{d}\left(q, q^{\prime}\right) \leq \widehat{d}\left(q, q^{\prime}\right)=\widehat{d}\left(0, q^{-1} q^{\prime}\right) \leq C\left\|q^{-1} q^{\prime}\right\|_{p} \tag{10}
\end{equation*}
$$

Considering $q=(x, y, z)$ and $q^{\prime}=\left(x^{\prime}, y^{\prime}, z^{\prime}\right)$ and from (10) and (9) we have

$$
\begin{align*}
\widetilde{d}\left(q, q^{\prime}\right) \leq C\left\|q^{-1} \cdot q^{\prime}\right\|_{p} & =C\left(\left|x^{\prime}-x\right|+\left|y^{\prime}-y\right|+\left|z^{\prime}-z+x\left(y-y^{\prime}\right)\right|^{\frac{1}{2}}\right) \\
& \leq C\left(\left|x-x^{\prime}\right|+\left|y-y^{\prime}\right|+\left|z-z^{\prime}\right|^{\frac{1}{2}}+\left|x\left(y-y^{\prime}\right)\right|^{\frac{1}{2}}\right) \\
& \leq C\left(\left\|q-q^{\prime}\right\|_{p}+\|q\|_{p}^{\frac{1}{2}}\left|y-y^{\prime}\right|^{\frac{1}{2}}\right) \tag{11}
\end{align*}
$$

### 4.2.3 Upper bound for $d$

From Theorems 7.31 and 7.26 of [7] we get

$$
\begin{equation*}
d\left(q, q^{\prime}\right) \leq \mathrm{Cst} \sum_{k, j \mid w_{k} \leq w_{j}}\|q\|^{1-\frac{w_{k}}{w_{j}}}\left|q_{k}-q_{k}^{\prime}\right|^{\frac{1}{w_{j}}} \tag{12}
\end{equation*}
$$

Since $w_{1}=w_{2}=1$ and $w_{3}=2$ we obtain that

$$
\begin{equation*}
d\left(q, q^{\prime}\right) \leq \operatorname{Cst}\left(\left\|q-q^{\prime}\right\|_{p}+\|q\|_{p}^{\frac{1}{2}}\left(\left|y-y_{2}\right|^{\frac{1}{2}}+\left|x-x_{2}\right|^{\frac{1}{2}}\right)\right) . \tag{13}
\end{equation*}
$$

### 4.2.4 Main result

We have seen that the order of $\delta$ does not change the estimation of $\|\gamma(t)-\widehat{\gamma}(t)\|_{p}$ according to Lemma 1, moreover, it does not change the estimates of $\widehat{d}$ as well, this is to say, $\widehat{d}(\gamma(t), \widehat{\gamma}(t)) \leq$ Cst $\cdot t^{\frac{1}{2}} \cdot \tau$. Indeed, from inequality (10) and Remark 6 , we get

$$
\widehat{d}(\gamma(t), \widehat{\gamma}(t)) \leq \operatorname{Cst}\left(t^{\frac{1}{2}} \cdot \tau+t \cdot \tau^{\frac{1}{2}+\rho}\right)
$$

then $\widehat{d}(\gamma(t), \widehat{\gamma}(t)) \leq$ Cst $\cdot t^{\frac{1}{2}} \cdot \tau$. However, the above is not true when we want to estimate $\widetilde{d}$, because the estimates depend of $\rho$ and $m$. Indeed, from inequalities (8) and (11) we get

$$
\widetilde{d}(\gamma(t), \widetilde{\gamma}(t)) \leq \operatorname{Cst} \cdot\left(t^{\frac{1}{2}} \cdot \tau^{\frac{\min (m+1, \rho+1,3)}{2}}+t^{\frac{1}{2}} \cdot t^{\frac{1}{2}} \cdot \tau^{\frac{\rho}{2}}\right) \leq \operatorname{Cst} \cdot t^{\frac{1}{2}} \cdot \tau^{\frac{\min (m+1, \rho+1,3)}{2}}
$$

Therefore, we can conclude that if $\rho \geq 2$ and $m \geq 2$ then

$$
\widetilde{d}(\gamma(t), \widetilde{\gamma}(t)) \leq \operatorname{Cst} \cdot t^{\frac{1}{2}} \cdot \tau^{\frac{3}{2}}
$$

Finally, we obtain:

Proposition 11. If $\operatorname{ord}_{p}(\delta) \geq 2$ and $\operatorname{ord}_{p}(\theta) \geq 2$, then

1. $\widetilde{d}(\gamma(t), \widetilde{\gamma}(t)) \leq C s t \cdot t^{\frac{1}{2}} \cdot \tau^{\frac{3}{2}}$.
2. $d(\gamma(t), \widetilde{\gamma}(t)) \leq C s t \cdot t^{\frac{1}{2}} \cdot \tau^{\frac{3}{2}}$.

Proof. The proof of item 2 follows from inequalities (8) and (13).
Theorem 5. If $m \geq 2$ and $\rho \geq 2$, then there exists constants $C$ and $\epsilon>0$, such that, for all $q, q^{\prime} \in B(p, \epsilon)$, we have

$$
\begin{equation*}
-C \tau^{\frac{3}{2}} d\left(q, q^{\prime}\right)^{\frac{1}{2}} \leq d\left(q, q^{\prime}\right)-\widetilde{d}\left(q, q^{\prime}\right) \leq C \cdot \widetilde{\tau}^{\frac{3}{2}} \widetilde{d}\left(q, q^{\prime}\right)^{\frac{1}{2}} \tag{14}
\end{equation*}
$$

where

$$
\begin{aligned}
\tau & =\max \left(\|q\|_{p}, d\left(q, q^{\prime}\right)\right) \\
\widetilde{\tau} & =\max \left(\|q\|_{p}, \widetilde{d}\left(q, q^{\prime}\right)\right)
\end{aligned}
$$

Proof. Let $q$ belonging to $B(p, \epsilon)$. Let us consider the geodesics $\gamma:[0, T] \rightarrow M$ for the distance $d$ such that $\gamma(0)=q, \gamma(T)=q^{\prime}$ and associated with the control function $u(\cdot)$ satisfying $\|u(t)\|=1$ and $\widetilde{\gamma}$ the admissible curve for $\widetilde{d}$ defined by the same control functions that $\gamma$ with $\widetilde{\gamma}(0)=q$. By Proposition 11 item 1

$$
\begin{equation*}
\widetilde{d}(\gamma(T), \widetilde{\gamma}(T)) \leq \mathrm{Cst} \cdot T^{\frac{1}{2}} \cdot \tau^{\frac{3}{2}} \tag{15}
\end{equation*}
$$

On the other hand, note that

$$
d\left(q, q^{\prime}\right)=l(\gamma)=l(\widetilde{\gamma}) \geq \widetilde{d}(q, \widetilde{\gamma}(T))
$$

Moreover, by triangle inequality, we have

$$
\widetilde{d}(q, \widetilde{\gamma}(T)) \geq \widetilde{d}\left(q, q^{\prime}\right)-\widetilde{d}\left(q^{\prime}, \widetilde{\gamma}(T)\right),
$$

Then, from (15), transitivity and since $\gamma(T)=q^{\prime}$, we get

$$
\begin{align*}
d\left(q, q^{\prime}\right) & \geq \widetilde{d}\left(q, q^{\prime}\right)-\mathrm{Cst} \cdot T^{\frac{1}{2}} \cdot \tau^{\frac{3}{2}} \\
d\left(q, q^{\prime}\right)-\widetilde{d}\left(q, q^{\prime}\right) & \geq-\mathrm{Cst} \cdot T^{\frac{1}{2}} \cdot \tau^{\frac{3}{2}} \tag{16}
\end{align*}
$$

Now, we change the roles of $d$ and $\widetilde{d}$ and by Proposition 11 item 2, we obtain

$$
\begin{equation*}
d\left(q, q^{\prime}\right)-\widetilde{d}\left(q, q^{\prime}\right) \leq \mathrm{Cst} \cdot \widetilde{T}^{\frac{1}{2}} \cdot \widetilde{\tau}^{\frac{3}{2}} \tag{17}
\end{equation*}
$$

where $\widetilde{T}=\widetilde{d}\left(q, q^{\prime}\right)$.
Therefore from (16) and (17)

$$
-C \tau^{\frac{3}{2}} d\left(q, q^{\prime}\right)^{\frac{1}{2}} \leq d\left(q, q^{\prime}\right)-\widetilde{d}\left(q, q^{\prime}\right) \leq C \cdot \widetilde{\tau}^{\frac{3}{2}} \widetilde{d}\left(q, q^{\prime}\right)^{\frac{1}{2}} .
$$

The proof is complete.

### 4.2.5 Translations

We mentioned in Section 4.2.2 that the distance $\widehat{d}$ is left-invariant. It is not the case of $\tilde{d}$. Let $q$ be a point in a neighborhood of 0 and $g \in \mathbb{R}^{3}$. We are interested in conditions under which $\widetilde{d}(g, g \cdot q) \leq \widetilde{d}(0, q)$ (the product is the Heisenberg one).

Let $\gamma(t)=(x(t), y(t), z(t))$ be a geodesic of $\tilde{d}$ such that $\gamma(0)=0$ with control functions $u_{1}, u_{2}$ and $u_{3}$. We consider $g=\left(g_{1}, g_{2}, g_{3}\right) \in \mathbb{R}^{3}$. Let $\gamma_{g}(t)=L_{g}(\gamma(t))=$ $\left(x_{g}(t), y_{g}(t), z_{g}(t)\right)$ and $\bar{u}_{1}, \bar{u}_{2}, \bar{u}_{3}$ its control functions. Note that $\gamma_{g}$ is admissible for $\widetilde{d}$ as long as it does not meet $\mathcal{Z}$. Indeed, all absolutely continuous curves are admissible out of the singular locus since the metric is Riemannian. The goal is to find conditions for $g$ such that $\gamma_{g}$ has a length less than $\gamma$. Since Lie $\left\{\widehat{X}_{1}, \widehat{X}_{2}\right\}$ is the Heisenberg algebra, then

$$
L_{g}(\gamma(t))=\left(x(t)+g_{1}, y(t)+g_{2}, z(t)+g_{1} y(t)+g_{3}\right) .
$$

We set $h(x, y, z)=a z+b x^{2}+c y^{2}$. Then

$$
\begin{aligned}
h\left(\gamma_{g}\right) & =a\left(z+g_{1} y+g_{3}\right)+b\left(x+g_{1}\right)^{2}+c\left(y+g_{2}\right)^{2} \\
& =h(\gamma)+h(g)+(2 b x+a y) g_{1}+2 c y g_{2}=h(\gamma)+h(g)+f(g, \gamma),
\end{aligned}
$$

where $f(g, \gamma)=(2 b x+a y) g_{1}+2 c y g_{2}$.
We assume that $h\left(\gamma_{g}\right)$ does not vanish, this is to say $\gamma_{g}$ is not on $\mathcal{Z}$. In particular for $t=0, h\left(\gamma_{g}\right)=h(g)$ then $h(g) \neq 0$ this is equivalent to $g \notin \mathcal{Z}$.

We have the following result.
Theorem 6. Let $\gamma:[0, T] \longrightarrow \mathbb{R}^{n}$ be a length minimizer of $\widetilde{d}$ with control functions $u_{1}(t), u_{2}(t), u_{3}(t)$ with $u_{3}(t) \neq 0$ a.e, and $h\left(\gamma_{g}\right) \neq 0$. If $|h(\gamma)| \leq\left|h\left(\gamma_{g}\right)\right|$ then $\widetilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right) \leq \widetilde{d}(\gamma(0), \gamma(T))$.

Proof. Since $x_{g}(t)=x(t)+g_{1}$ then $\dot{x}_{g}(t)=\dot{x}(t)=u_{1}(t)$. This implies that $u_{1}(t)=\bar{u}_{1}(t)$. In the same way, we obtain $u_{2}(t)=\bar{u}_{2}(t)$ because $y_{g}(t)=y(t)+g_{2}$. Furthermore, $z_{g}(t)=z(t)+g_{1} y(t)+g_{3}$ then

$$
\begin{equation*}
\dot{z}_{g}(t)=\dot{z}(t)+g_{1} \dot{y}(t)=u_{2}(t) x_{g}(t)+u_{3}(t) h(\gamma(t)) . \tag{18}
\end{equation*}
$$

Besides the above equation, $z_{g}(t)$ satisfies the equation

$$
\begin{equation*}
\dot{z}_{g}(t)=\bar{u}_{2}(t) x_{g}(t)+\bar{u}_{3}(t) h\left(\gamma_{g}(t)\right), \tag{19}
\end{equation*}
$$

because $\gamma_{g}$ is an admissible curve for $\widetilde{d}$. Finally, from the equations (19) and (18) and as $u_{2}(t)=\bar{u}_{2}(t)$ we get

$$
\bar{u}_{3}(t)=\frac{u_{3}(t) h(\gamma(t))}{h\left(\gamma_{g}(t)\right)} .
$$

The condition $|h(\gamma)| \leq\left|h\left(\gamma_{g}\right)\right|$ implies that $\left|\bar{u}_{3}(t)\right| \leq\left|u_{3}(t)\right|$, hence $\bar{u}_{3}\left(t_{\sim}^{2} \leq u_{3}(t)^{2}\right.$. Therefore the length of $\gamma_{g}$ decreases and consequently $\widetilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right) \leq \widetilde{d}(\gamma(0), \gamma(T))$.

In the same sense of the above theorem, the following result gives us a sufficient condition to determine when the distance of the points translated by $g$ is less than the distance from the origin to $\gamma(T)$.
Theorem 7. With the same conditions of the above. If $\frac{\frac{\partial}{\partial g_{i}}(h(g)+f(g, \gamma))}{h\left(\gamma_{g}\right)}>0$ then $\widetilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right)<\widetilde{d}(\gamma(0), \gamma(T))$.

Proof. From Theorem 6, the control functions of $\gamma_{g}$ are $u_{1}(t), u_{2}(t)$ and $\bar{u}_{3}(t)$, hence

$$
\begin{aligned}
l\left(\gamma_{g}\right) & =\int_{0}^{T}\left(u_{1}(t)^{2}+u_{2}(t)^{2}+\bar{u}_{3}(t)^{2}\right)^{\frac{1}{2}} d t \\
& =\int_{0}^{T}\left(u_{1}(t)^{2}+u_{2}(t)^{2}+\frac{u_{3}(t)^{2} h(\gamma(t))^{2}}{h\left(\gamma_{g}(t)\right)^{2}}\right)^{\frac{1}{2}} d t \\
\frac{\partial}{\partial g_{i}} l\left(\gamma_{g}\right) & =\int_{0}^{T} \frac{1}{2}\left(\frac{\left(u_{1}(t)^{2}+u_{2}(t)^{2}\right) h\left(\gamma_{g}(t)\right)^{2}+u_{3}(t)^{2} h(\gamma(t))^{2}}{h\left(\gamma_{g}(t)\right)^{2}}\right)^{-\frac{1}{2}} \cdot \frac{\partial}{\partial g_{i}}\left(\frac{u_{3}(t)^{2} h(\gamma(t))^{2}}{h\left(\gamma_{g}(t)\right)^{2}}\right) d t \\
& =\int_{0}^{T} \frac{-\left|h\left(\gamma_{g}(t)\right)\right| u_{3}(t)^{2} h(\gamma(t))^{2}}{\left(\left(u_{1}(t)^{2}+u_{2}(t)^{2}\right) h\left(\gamma_{g}(t)\right)^{2}+u_{3}(t)^{2} h(\gamma(t))^{2}\right)^{\frac{1}{2}} h\left(\gamma_{g}(t)\right)^{3}} \cdot \frac{\partial}{\partial g_{i}}(h(g)+f(g, \gamma)) d t \\
& =-\int_{0}^{T} \frac{u_{3}(t)^{2} h(\gamma(t))^{2} \frac{\partial}{\partial g_{i}}(h(g)+f(g, \gamma))}{\left(\left(u_{1}(t)^{2}+u_{2}(t)^{2}\right) h\left(\gamma_{g}(t)\right)^{2}+u_{3}(t)^{2} h(\gamma(t))^{2}\right)^{\frac{1}{2}}\left|h\left(\gamma_{g}(t)\right)\right| h\left(\gamma_{g}(t)\right)} d t .
\end{aligned}
$$

Note that the function $S$ defined by

$$
S(t)=\frac{u_{3}(t)^{2} h(\gamma(t))^{2} \frac{\partial}{\partial g_{i}}(h(g)+f(g, \gamma))}{\left(u_{1}(t)^{2}+u_{2}(t)^{2} h\left(\gamma_{g}(t)\right)^{2}+u_{3}(t)^{2} h(\gamma(t))^{2}\right)^{\frac{1}{2}}\left|h\left(\gamma_{g}(t)\right)\right| h\left(\gamma_{g}(t)\right)}
$$

is positive if and only if

$$
\frac{\frac{\partial}{\partial g_{i}}(h(g)+f(g, \gamma))}{h\left(\gamma_{g}(t)\right)}>0 .
$$

In this case $\frac{\partial l\left(\gamma_{g}\right)}{\partial g_{i}}<0$ and $\widetilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right)<\widetilde{d}(\gamma(0), \gamma(T))$.
In particular at $g=(0,0,0)$,

$$
\left.\frac{\partial}{\partial g_{i}} l\left(\gamma_{g}\right)\right|_{\left(g_{1}, g_{2}, g_{3}\right)=(0,0,0)}=-\int_{0}^{T} \frac{u_{3}(t)^{2} \frac{\partial}{\partial g_{i}}(h(g)+f(g, \gamma))}{h(\gamma(t))} d t .
$$

### 4.2.6 Conclusion

In Section 3, we have shown that in case where the nilpotent approximation of an ARS degenerates, that is when it is no longer an ARS but a sub-Riemannian structure, we can replace it by a simple ARS on a Lie group or a homogeneous space. Thanks to formula (14) of Theorem 5 we know that, at least in some 3D generic cases, the order of the approximation of $d$ by $\tilde{d}$ is better than the one of the approximation of $d$ by $\widehat{d}$. Indeed, this order is $d^{2}$ in the first case and $d^{\frac{3}{2}}$ in the second one. However, this does not prove that the solvable approximation is really better than the nilpotent one, and anyway it is certainly not true for any pair of points.

Since under left translations the nilpotent distance $\widehat{d}$ is invariant while the solvable distance $\widetilde{d}$ may be decreasing, we can expect to prove that the approximation by $\widetilde{d}$ is strictly better than the one by $\widehat{d}$ for pairs of points translated in a suitable direction.

For this purpose, we consider here the 3D-generic case of Section 4.2 with the particular values $a=1, b=c=0$, that is $\widetilde{X}_{3}=z \frac{\partial}{\partial z}$. The singular locus of the solvable approximation is then the plane $\{z=0\}$.

In what follows we consider a (normal) geodesic $\gamma$ for $\widetilde{d}$, originated at $(0,0,0)$ and parametrized by arc length on $[0, T]$. Denoting $\gamma(t)=(x(t), y(t), z(t))$ it is moreover assumed that $z(t)>0$ on $] 0, T]$.

This geodesic is translated by $\mathbf{g}=(0,0, g)$, with $g \geq 0$, into $\gamma_{g}=L_{\mathbf{g}} \gamma$. Since g belongs to the center of the Heisenberg group the curve $\gamma_{g}$ is simply $\gamma_{g}(t)=$ $(x(t), y(t), z(t)+g)$.

The different distances between $g$ and $\gamma_{g}(T)$ are analyzed in several steps.

1. Since the controls associated to $\gamma_{g}$ are $u_{1}, u_{2}$, and $\frac{z(t)}{z(t)+g} u_{3}$ the length of $\gamma_{g}$ related to $\widetilde{d}$ is

$$
\tilde{l}\left(\gamma_{g}\right)=\int_{0}^{T}\left(u_{1}^{2}+u_{2}^{2}+\left(\frac{z(t)}{z(t)+g}\right)^{2} u_{3}^{2}\right)^{\frac{1}{2}} d t
$$

Since $z \longmapsto \frac{z}{z+g}$ is increasing we have $\widetilde{l}\left(\gamma_{g}\right) \leq I_{g}$, where $I_{g}$ stands for

$$
I_{g}=\int_{0}^{T}\left(u_{1}^{2}+u_{2}^{2}+\left(\frac{z_{m}}{z_{m}+g}\right)^{2} u_{3}^{2}\right)^{\frac{1}{2}} d t
$$

with $z_{m}=\max \{z(t) ; \quad t \in[0, T]\}$.
2. We apply now formula (14), which writes here:

$$
d\left(\mathbf{g}, \gamma_{g}(T)\right) \leq \widetilde{d}\left(\mathbf{g}, \gamma_{g}(T)\right)+C \cdot \widetilde{\tau}^{\frac{3}{2}} \widetilde{d}\left(\mathbf{g}, \gamma_{g}(T)\right)^{\frac{1}{2}}
$$

where $\widetilde{\tau}=\max \left\{\|\mathbf{g}\|_{p}, \widetilde{d}\left(\mathbf{g}, \gamma_{g}(T)\right)\right\}$. It will always be assumed that $\|\mathbf{g}\|_{p} \leq$ $\widetilde{d}\left(\mathbf{g}, \gamma_{g}(T)\right)$, that is $g^{\frac{1}{2}} \leq \widetilde{d}\left(\mathbf{g}, \gamma_{g}(T)\right)$. Taking into account $\widetilde{d}\left(\mathbf{g}, \gamma_{g}(T)\right) \leq$ $\widehat{d}\left(\mathbf{g}, \gamma_{g}(T)\right)=\widehat{d}(0, \gamma(T))$, we get

$$
\begin{aligned}
d\left(\mathbf{g}, \gamma_{g}(T)\right) & \leq \widetilde{d}\left(\mathbf{g}, \gamma_{g}(T)\right)+C \cdot \widetilde{d}\left(\mathbf{g}, \gamma_{g}(T)\right)^{2} \\
& \leq \frac{I_{g}}{T} \widetilde{d}(0, \gamma(T))+C \cdot\left(\frac{I_{g}}{T}\right)^{2} \widetilde{d}(0, \gamma(T))^{2} \\
& \leq \frac{I_{g}}{T} \widehat{d}(0, \gamma(T))\left(1+\frac{I_{g}}{T} C \cdot \widetilde{d}(0, \gamma(T))\right)
\end{aligned}
$$

3. In order to approximate $z(t)$ and $u_{3}(t)$, we consider the Hamiltonian equations (see details in the next section), for the values $a=1, b=c=0$. They are:

$$
\left\{\begin{array} { l } 
{ \dot { x } = p } \\
{ \dot { y } = q + r x } \\
{ \dot { z } = ( q + r x ) x + r z ^ { 2 } }
\end{array} \quad \left\{\begin{array}{l}
\dot{p}=-(q+r x) r \\
\dot{q}=0 \\
\dot{r}=-r^{2} z
\end{array}\right.\right.
$$

It is important to notice that $r_{0}$ can be chosen arbitrarily large because $H(t=$ $0)=\frac{1}{2}\left(p_{0}^{2}+q_{0}^{2}\right)$. We make the choice $p_{0}=q_{0}$ and the following approximations hold:

$$
x(t) \approx p_{0} t, \quad \dot{z} \approx q_{0} p_{0} t=\frac{1}{2} t, \quad z(t) \approx \frac{1}{4} t^{2}, \quad u_{3}(t)=r(t) z(t) \approx \frac{1}{4} r_{0} t^{2} .
$$

In order to compute $\widetilde{l}\left(\gamma_{g}\right)$, we need to apply the condition $g^{\frac{1}{2}} \leq \widetilde{d}\left(\mathbf{g}, \gamma_{g}(T)\right)$ of point 2 . We do not know $\widetilde{d}\left(\mathbf{g}, \gamma_{g}(T)\right)$ but we can set $\widetilde{d}\left(\mathbf{g}, \gamma_{g}(T)\right)=\beta T$ with $0<\beta<1$ (this estimation "a priori" will be justified later), and set $g=\beta^{2} T^{2}$. Then we get

$$
\frac{z}{z+g} \leq \frac{z_{m}}{z_{m}+g} \approx \frac{\frac{1}{4} T^{2}}{\frac{1}{4} T^{2}+\beta^{2} T^{2}}=\frac{1}{1+4 \beta^{2}}
$$

Therefore

$$
\begin{aligned}
\widetilde{l}\left(\gamma_{g}\right) & \leq I_{g}=\int_{0}^{T}\left(u_{1}^{2}+u_{2}^{2}+\left(\frac{1}{1+4 \beta^{2}}\right)^{2} u_{3}^{2}\right)^{\frac{1}{2}} d t \\
& =\int_{0}^{T}\left(1-u_{3}^{2}+\left(\frac{1}{1+4 \beta^{2}}\right)^{2} u_{3}^{2}\right)^{\frac{1}{2}} d t=\int_{0}^{T}\left(1-\frac{16 \beta^{4}+8 \beta^{2}}{\left(1+4 \beta^{2}\right)^{2}} u_{3}^{2}\right)^{\frac{1}{2}} d t \\
& \approx \int_{0}^{T}\left(1-\frac{16 \beta^{4}+8 \beta^{2}}{\left(1+4 \beta^{2}\right)^{2}} \frac{1}{16} r_{0}^{2} t^{4}\right)^{\frac{1}{2}} d t=\int_{0}^{T}\left(1-\frac{2 \beta^{4}+\beta^{2}}{2\left(1+4 \beta^{2}\right)^{2}} r_{0}^{2} t^{4}\right)^{\frac{1}{2}} d t .
\end{aligned}
$$

We write $\delta=\frac{2 \beta^{4}+\beta^{2}}{2\left(1+4 \beta^{2}\right)^{2}}$ and we set $\delta r_{0}^{2} T^{4}=\frac{1}{2}$. Notice that this is possible, even if $T$ is small, by increasing $r_{0}$. Thanks to $(1-c)^{\frac{1}{2}} \leq 1-0.5 c$ whenever $0 \leq c \leq 1$ we get:

$$
\widetilde{l}\left(\gamma_{g}\right) \leq I_{g} \leq \int_{0}^{T}\left(1-0.5 \delta r_{0}^{2} t^{4}\right) d t=T-0.5 \delta r_{0}^{2} \frac{T^{5}}{5}=T\left(1-0.1 \delta r_{0}^{2} T^{4}\right) \approx 0.95 T
$$

4. Assuming $C \cdot \widetilde{d}(0, \gamma(T))=C \cdot T \leq 0.01$ we get on one hand:

$$
\begin{aligned}
d\left(\mathbf{g}, \gamma_{g}(T)\right) & \leq \widetilde{d}\left(\mathbf{g}, \gamma_{g}(T)\right)+C \cdot \widetilde{d}\left(\mathbf{g}, \gamma_{g}(T)\right)^{2} \\
& =\widetilde{d}\left(\mathbf{g}, \gamma_{g}(T)\right)\left(1+C \cdot \widetilde{d}\left(\mathbf{g}, \gamma_{g}(T)\right)\right) \\
& \leq 1.01 \widetilde{d}\left(\mathbf{g}, \gamma_{g}(T)\right)
\end{aligned}
$$

On the other hand:

$$
\begin{aligned}
\widehat{d}\left(\mathbf{g}, \gamma_{g}(T)\right)-d\left(\mathbf{g}, \gamma_{g}(T)\right) & \geq \widehat{d}\left(\mathbf{g}, \gamma_{g}(T)\left(1-1.01 \frac{I_{g}}{T}\right)\right. \\
& =\widetilde{d}\left(\mathbf{g}, \gamma_{g}(T)\right) \frac{\widehat{d}\left(\mathbf{g}, \gamma_{g}(T)\right)}{\widetilde{d}\left(\mathbf{g}, \gamma_{g}(T)\right)}\left(1-1.01 \frac{I_{g}}{T}\right) \\
& \geq \widetilde{d}\left(\mathbf{g}, \gamma_{g}(T)\right) \frac{T}{I_{g}}\left(1-1.01 \frac{I_{g}}{T}\right) \\
& =\widetilde{d}\left(\mathbf{g}, \gamma_{g}(T)\right)\left(\frac{T}{I_{g}}-1.01\right)
\end{aligned}
$$

Therefore $\widehat{d}\left(\mathbf{g}, \gamma_{g}(T)\right)-d\left(\mathbf{g}, \gamma_{g}(T)\right)>d\left(\mathbf{g}, \gamma_{g}(T)\right)-\widetilde{d}\left(\mathbf{g}, \gamma_{g}(T)\right)$ as soon as $\frac{T}{I_{g}}-1.01>0.01$ hence as soon as $\frac{I_{g}}{T}<0.98$.
According to Point 3, we can obtain $I_{g} \leq 0.95 T$ and in that case the solvable distance between $\mathbf{g}$ and $\gamma_{g}(T)$ is strictly closer to the original distance between these points than the nilpotent one.

## 5 Geodesics

In this section the Hamiltonian for the normal flow defined by the solvable approximation in the 3D generic case is given. We compute the geodesic with initial condition $x(0)=y(0)=z(0)=0$ and covector $\lambda=(p, q, r) \in T^{*} \mathbb{R}^{3}$ with $p(0)=\cos (\theta)$,
$q(0)=\sin (\theta), r(0)=r$.
From the above sections, the solvable approximation is defined by

$$
\begin{equation*}
\widehat{X}_{1}=X_{1}, \quad \widehat{X}_{2}=\frac{\partial}{\partial y}+x \frac{\partial}{\partial z}, \quad \widetilde{X}_{3}=\left(a z+b x^{2}+c y^{2}\right) \frac{\partial}{\partial z} . \tag{20}
\end{equation*}
$$

From (20), the Hamiltonian for the normal flow is given by

$$
\begin{aligned}
& H(\lambda)=\frac{1}{2}\left(\left\langle\lambda, \widehat{X}_{1}(x, y, z)\right\rangle^{2}+\left\langle\lambda, \widehat{X}_{2}(x, y, z)\right\rangle^{2}+\left\langle\lambda, \widetilde{X}_{3}(x, y, z)\right\rangle^{2}\right) \\
& H(\lambda)=\frac{1}{2}\left(p^{2}+(q+r x)^{2}+r^{2}\left(a z+b x^{2}+c y^{2}\right)^{2}\right)
\end{aligned}
$$

where $\lambda=(p, q, r) \in T^{*} \mathbb{R}^{3}$. Hence

$$
\begin{array}{ll}
\dot{x}(t)=p & \dot{p}(t)=-(q+r x) r-2 b x r^{2}\left(a z+b x^{2}+c y^{2}\right) \\
\dot{y}(t)=q+r x & \dot{q}(t)=-2 c y r^{2}\left(a z+b x^{2}+c y^{2}\right) \\
\dot{z}(t)=(q+r x) x+r\left(a z+b x^{2}+c y^{2}\right)^{2} & \\
\dot{r}(t)=-a r^{2}\left(a z+b x^{2}+c y^{2}\right)
\end{array}
$$

are the associated Hamiltonian equations to the solvable approximation.
The geodesic with initial condition $x(0)=y(0)=z(0)=0$ and $p(0)=\cos (\theta)$, $q(0)=\sin (\theta)$ and $r(0)=r=0$ is given by

$$
\begin{align*}
& x(t)=t \cos (\theta) \\
& y(t)=t \sin (\theta)  \tag{21}\\
& z(t)=\frac{1}{4} t^{2} \sin (2 \theta),
\end{align*}
$$

because $p(t)=\cos (\theta)$ and $q(t)=\sin (\theta)$, this is to say $p$ and $q$ are constants.
Notice that the above geodesic for $\widetilde{d}$ is the same as the geodesic for $\widehat{d}$. The above implies that this geodesic is optimal for any time and has no conjugate time (see Theorem 5.1 and 5.2 in [10]). We can see some geodesics in Figure 1 when $r=0$.

Due to the complexity of the Hamiltonian system of equations, we compute the geodesics considering $a=c=0$ and $b=1$. Thus the Hamiltonian is

$$
H(\lambda)=\frac{1}{2}\left(p^{2}+(q+r x)^{2}+r^{2} x^{4}\right),
$$

hence

$$
\begin{array}{ll}
\dot{x}(t)=p & \dot{p}(t)=-(q+r x) r-2 r^{2} x^{3} \\
\dot{y}(t)=q+r x & \dot{q}(t)=0  \tag{22}\\
\dot{z}(t)=x q+r x^{2}+r x^{4} & \dot{r}(t)=0
\end{array}
$$

Considering the initial condition $x(0)=0$ then $p(0)=\cos (\theta), q(0)=\sin (\theta)$ and $r(0)=r$. If $r=0$ then the solution to the differential systems (22) is given by (21). If $r(0)=r \neq 0$, since $\dot{x}(t)=p$, we get

$$
\ddot{x}=-r(q+r x)-2 r^{2} x^{3}
$$

$$
\ddot{x}+r^{2} x+2 r^{2} x^{3}=-r q .
$$



Figure 1: Geodesics for $\theta \in\left\{0, \frac{\pi}{3}, \frac{2 \pi}{3}, \pi, \frac{4 \pi}{6}, \frac{5 \pi}{6}\right\}$ when $r=0$.

Since $q(0)=\sin (\theta)$ and $\dot{q}=0$, then $q=\sin (\theta)$. Hence

$$
\begin{equation*}
\ddot{x}+r^{2} x+2 r^{2} x^{3}=-r \sin (\theta) . \tag{23}
\end{equation*}
$$

The equation (23) is equivalent to

$$
\begin{equation*}
\ddot{x}+r^{2} x+2 r^{2} x^{3}=-r \sin (\theta) \operatorname{cn}\left(0, k^{2}\right), \tag{24}
\end{equation*}
$$

where $\mathrm{cn}\left(0 \cdot t, k^{2}\right)$ is the Jacobian elliptic function that has a period in $0 \cdot t$ equal to $4 K\left(k^{2}\right)$ and $K\left(k^{2}\right)$ is the complete elliptic integral of the first kind for the modulus $k$ (see more in [14]). This equivalence is due to the fact that $\mathrm{cn}\left(0, k^{2}\right)=1$.

In [18] a general solution to

$$
\ddot{x}+c_{n} \dot{x}+w_{n} x+\epsilon x^{3}=F \operatorname{cn}\left(w t, k^{2}\right),
$$

is given by

$$
x(t)=a_{1}(t) \operatorname{cn}\left(w_{1} t+\phi, k_{1}^{2}\right)+A_{1}(t) \operatorname{cn}\left(w t, k^{2}\right)+B_{1}(t) \cdot \operatorname{sn}\left(w t, k^{2}\right) .
$$

Therefore, the solution for the equation (24) is given by

$$
x(t)=a_{1}(t) \operatorname{cn}\left(w_{1} t+\phi, k_{1}^{2}\right)+A_{1}(t),
$$

where $a_{1}(t), A_{1}(t), w_{1}, \phi$ and $k_{1}$ need to be determined. Notice that $B_{1}(t) \cdot \operatorname{sn}\left(w t, k^{2}\right)$ vanishes because $\operatorname{sn}\left(0, k^{2}\right)=0$.

From [18] is possible to obtain that $a_{1}(t)$ and $A_{1}(t)$ are constants. Then

$$
\begin{equation*}
x(t)=a_{1} \operatorname{cn}\left(w_{1} t+\phi, k_{1}^{2}\right)+A_{1} . \tag{25}
\end{equation*}
$$

Moreover, since $x(0)=0$

$$
\begin{equation*}
-A_{1}=a_{1} \cdot \operatorname{cn}\left(\phi, k_{1}^{2}\right) \tag{26}
\end{equation*}
$$

Furthermore, differentiating in (25) and since $\dot{x}(0)=p(0)=\cos (\theta)$, we have

$$
a_{1}=\frac{\cos (\theta) \operatorname{ns}\left(\phi, k_{1}^{2}\right) \operatorname{nd}\left(\phi, k_{1}^{2}\right)}{w_{1}}
$$

Finally, since $y(0)=z(0)=0$,

$$
\begin{aligned}
x(t) & =a_{1}\left(\operatorname{cn}\left(w_{1} t+\phi, k_{1}^{2}\right)-\operatorname{cn}\left(\phi, k_{1}^{2}\right)\right) \\
y(t) & =\left(\sin (\theta)-r a_{1} \operatorname{cn}\left(\phi, k_{1}^{2}\right)\right) t+\frac{r a_{1}}{k_{1}^{2} w_{1}}\left(\arccos \left(\operatorname{dn}\left(w_{1} t+\phi, k_{1}^{2}\right)\right)-\arccos \left(\operatorname{dn}\left(\phi, k_{1}^{2}\right)\right)\right) \\
z(t) & =-\left(r a_{1}^{4} \operatorname{cn}\left(\phi, k_{1}^{2}\right)^{4}+r a_{1}^{2} \operatorname{cn}\left(\phi, k_{1}^{2}\right)^{2}+\sin (\theta)\right) a_{1} \operatorname{cn}\left(\phi, k_{1}^{2}\right) t+\frac{r a_{1}^{4}}{3 k_{1}^{8} w_{1}} z_{1}(t) \\
& +\frac{4 r a_{1}^{4} \operatorname{cn}\left(\phi, k_{1}^{2}\right)}{2 k_{1}^{6} w_{1}} z_{2}(t)+\frac{6 r a_{1}^{4} \operatorname{cn}\left(\phi, k_{1}^{2}\right)^{2}+r a_{1}^{2}}{k_{1}^{4} w_{1}} z_{3}(t) \\
& +\frac{4 r a_{1}^{4} \operatorname{cn}\left(\phi, k_{1}^{2}\right)^{3}+2 r a_{1}^{2} \operatorname{cn}\left(\phi, k_{1}^{2}\right)+\sin (\theta)}{k_{1}^{2} w_{1}} z_{4}(t),
\end{aligned}
$$

where $k_{1}^{\prime 2}+k_{1}^{2}=1, E(\cdot)$ is the incomplete elliptic integral of the second kind and

$$
\begin{aligned}
z_{1}(t)= & \left(2-3 k_{1}^{4}\right) k_{1}^{\prime 4} w_{1} t+2\left(2 k_{1}^{4}-1\right)\left(E\left(w_{1} t+\phi\right)-E(\phi)\right) \\
& +k_{1}^{4}\left(\operatorname{sn}\left(w_{1} t+\phi, k_{1}^{2}\right) \operatorname{cn}\left(w_{1} t+\phi, k_{1}^{2}\right) \operatorname{dn}\left(w_{1} t+\phi, k_{1}^{2}\right)-\operatorname{sn}\left(\phi, k_{1}^{2}\right) \operatorname{cn}\left(\phi, k_{1}^{2}\right) \operatorname{dn}\left(\phi, k_{1}^{2}\right)\right) \\
z_{2}(t)= & \left(2 k_{1}^{4}-1\right)\left(\arcsin \left(k_{1}^{2} \operatorname{sn}\left(w_{1} t+\phi, k_{1}^{2}\right)\right)-\arcsin \left(k_{1}^{2} \operatorname{sn}\left(\phi, k_{1}^{2}\right)\right)\right) \\
& +k_{1}^{2}\left(\operatorname{sn}\left(w_{1} t+\phi, k_{1}^{2}\right) \operatorname{dn}\left(w_{1} t+\phi, k_{1}^{2}\right)-\operatorname{sn}\left(\phi, k_{1}^{2}\right) \operatorname{dn}\left(w_{1} t+\phi, k_{1}^{2}\right)\right) \\
z_{3}(t)= & E\left(w_{1} t+\phi\right)-E(\phi)-k_{1}^{4} w_{1} t \\
z_{4}(t)= & \left.\arccos \left(\operatorname{dn}\left(w_{1} t+\phi, k_{1}^{2}\right)\right)-\arccos \left(\operatorname{dn}\left(\phi, k_{1}^{2}\right)\right)\right) .
\end{aligned}
$$

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Figure 2: Ball in 3-D generic case.
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## APPENDIX C

## Paper 2

General Nilpotent and Solvable Approximations of Almost-Riemannian Structures, article in collaboration with Yacine Chitour and Philippe Jouan. Submitted to the journal Discrete and Continuous Dynamical Systems (DCDS) (see online link).

# General Nilpotent and Solvable Approximations of Almost-Riemannian Structures 

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#### Abstract

It is first shown that the nilpotent or the solvable approximation of an almost-Riemannian structure at a singular point is always a linear almost-Riemannian structure on a Lie group or a homogeneous space.

The generic properties of almost-Riemannian structures are then investigated in all dimensions and the generic nilpotent and solvable approximations are identified.


Keywords: Almost-Riemannian geometry; Nilpotent approximation; Solvable approximation; Genericity.

## 1 Introduction

The aim of this paper is twofold. On the first hand it is shown that the nilpotent or the solvable approximation of an almost-Riemannian structure at a singular point is always a linear almost-Riemannian structure on a Lie group or a homogeneous space. On the other hand the generic almost-Riemannian structures are described and used to exhibit the generic nilpotent and solvable approximations.

More precisely an almost-Riemannian structure, ARS in short, on an $n$ dimensional differential manifold is a rank-varying sub-Riemannian structure that can be locally defined by a set of $n$ vector fields satisfying the Lie algebra rank condition and such that the singular locus $\mathcal{Z}$, that is the set of points where their rank is not full, is a proper but with empty interior subset (see [2], [3], [4], [6], [7], [12]).

In particular almost-Riemannian structures on Lie groups or homogeneous spaces that are defined by invariant and linear vector fields only are referred to as linear ARSs (a vector field on a Lie group is linear if its flow is a one-parameter group of automorphisms)(see [4]).

Following [5] and [10] the nilpotent approximation of such a structure at a point $p$ of the singular locus has been used in [7] and [6] to investigate the behaviour around $p$ in dimension respectively 2 and 3 .

However it may happen that some vector fields of the nilpotent approximation vanish, changing the almost-Riemannian structure into a subRiemannian one. In that case the nilpotent approximation can be replaced by a solvable one, the definition of which was stated in [13] and is recalled in Section 2.2. In [13] the solvable approximation is used to study the distance around some singular points in the generic $3 D$-case.

The first aim is herein to prove that the nilpotent or solvable approximation of an ARS at a singular point is a linear almost-Riemannian structure on a Lie group or a homogeneous space, excepted in some very degenerated cases where nor the nilpotent approximation neither the solvable one defines an ARS. It is the purpose of Section 3, and the statements are illustrated by examples postponed to Section 5.

The generic properties of almost-Riemannian structures are then examined in Section 4. It is in particular shown that generically: (1) the singular set $\mathcal{Z}$ is a union of submanifolds $\mathcal{Z}_{r}$ of codimension $r^{2}$ where the rank is $n-r ;(2)$ the rank of $\Delta+[\Delta, \Delta]$ is everywhere full ( $\Delta$ stands for the distribution). The structure of the points of $\mathcal{Z}_{r}$ where $\operatorname{dim}\left(T_{p} \mathcal{Z}_{r}\right)+\operatorname{dim}\left(\Delta_{p}\right)$ is not maximal is described in Theorem 6. For example in $\mathcal{Z}_{1}$ these points are the so-called tangency points (see [6]), i.e. the points where $T_{p} \mathcal{Z}_{1}=\Delta_{p}$. They are generically isolated in $\mathcal{Z}_{1}$.

Thanks to these genericity results and with the help of local normal forms (see Section 4.2) it it finally shown that generically there are only two possibilities for the nilpotent/solvable approximation at a point $p \in \mathcal{Z}$ :

1. At a tangency point $p$ in $\mathcal{Z}_{1}$ one vector field of the nilpotent approximation vanish, but the solvable approximation is not degenerated and defines a linear ARS.
2. At all other points, that is nontangent points of $\mathcal{Z}_{1}$ and all points in $\mathcal{Z}_{r}$ with $r \geq 2$, the nilpotent approximation is not degenerated.

In conclusion the only generic points where the solvable approximation is useful are tangency points in $\mathcal{Z}_{1}$.

As shown by the examples the picture is very different for nongeneric ARSs.

## 2 Basic definitions

### 2.1 Almost-Riemannian structures

For all that concern general sub-Riemannian geometry, including almostRiemannian one, the reader is referred to [1].

Let $M$ be a $n$-dimensional, connected, $\mathcal{C}^{\infty}$ manifold. The $\mathcal{C}^{\infty}$-module of $\mathcal{C}^{\infty}$ vector fields on $M$ is denoted by $\Gamma(M)$. Let $\Delta$ be a submodule of $\Gamma(M)$. The flag of submodules

$$
\Delta=\Delta^{1} \subseteq \Delta^{2} \subseteq \Delta^{3} \subseteq \cdots \subseteq \Delta^{k} \subseteq \ldots
$$

is defined by induction: $\Delta^{2}=\Delta+[\Delta, \Delta]$ is the submodule of $\Gamma(M)$ generated by $\Delta$ and the Lie brackets of its elements, and $\Delta^{k+1}=\Delta^{k}+\left[\Delta, \Delta^{k}\right]$. The Lie algebra generated by $\Delta$ is $\mathcal{L}(\Delta)=\bigcup_{k \geq 1} \Delta^{k}$. The submodule $\Delta$ satisfies the rank condition if the evaluation of $\mathcal{L}(\bar{\Delta})$ at each point $q$ is equal to $T_{q} M$.

Definition 1 An almost-Riemannian structure (resp. distribution) on a smooth $n$-dimensional manifold $M$ is a triple $(E, f,\langle.,\rangle$.$) (resp. a pair (E, f)$ ) where $E$ is a rank $n$ vector bundle over $M, f: E \longmapsto T M$ is a morphism of vector bundles, and $(E,\langle.,\rangle$.$) is an Euclidean bundle, that is \langle., .\rangle_{q}$ is an inner product on the fiber $E_{q}$ of $E$, smoothly varying w.r.t. $q$, assumed to satisfy the following properties:
(i) The set of points $q \in M$ such that the restriction of $f$ to $E_{q}$ is onto is a proper open and dense subset of $M$;
(ii) The module $\Delta$ of vector fields of $M$, defined as the image by $f$ of the module of smooth sections of $E$, satisfies the rank condition.

The set of points of $M$ where the rank of $f\left(E_{q}\right)=\Delta_{q}$ is less than $n$ is called the singular locus of the $A R S$ and denoted by $\mathcal{Z}$.

The inner product on $E$ induces a bilinear symmetric and positive definite mapping, also denoted by $\langle.,$.$\rangle , from \Delta \times \Delta$ to $\mathcal{C}^{\infty}(M)^{1}$. Indeed an element $X \in \Delta$ (resp. $Y \in \Delta$ ) is the image by $f$ of an unique section $\sigma$ (resp. $\eta$ ) of $E$ and we can set $\langle X, Y\rangle_{q}=\langle\sigma, \eta\rangle_{q}$.

Consequently an ARS can be alternately defined as follows.

[^1]Definition 2 An almost-Riemannian structure on a smooth n-dimensional manifold $M$ is a pair $(\Delta,\langle.,\rangle$.$) where \Delta$ is a submodule of $\Gamma(M)$ that can be locally defined by $n$ vector fields and satisfies the rank condition, and $\langle.,$.$\rangle is$ a bilinear symmetric and positive definite mapping from $\Delta \times \Delta$ to $\mathcal{C}^{\infty}(M)$, such that the set $\mathcal{Z}$ of points $q$ where the dimension of $\Delta_{q}$ is less than $n$ is nonempty but with empty interior.

Around any point $p \in M$ the submodule $\Delta$ can be locally defined by an orthonormal frame $\left(X_{1}, X_{2}, \ldots, X_{n}\right)$. It is enough to select a set of $n$ sections ( $e_{1}, e_{2}, \ldots, e_{n}$ ) of $E$ orthonormal in a neighborhood of $p$ and define $X_{i}=f_{*} e_{i}$ where by definition $f_{*} e_{i}=f \circ e_{i}$.
Norm
The almost-Riemannian norm on $\Delta_{q}$ is defined by

$$
\|v\|=\min \left\{\|u\| ; u \in E_{q} \text { and } f(u)=v\right\}
$$

### 2.2 Privileged coordinates and approximations

All the material of this section comes from [5] and [10], excepted the solvable approximations that are defined in [13].

Let $p$ be a point of $M$ and let $\Delta^{k}(p), k \geq 1$ be the evaluation of the submodule $\Delta^{k}$ at $p$. Thanks to the rank condition these submodules verify $\Delta^{1}(p) \subset \Delta^{2}(p) \subset \cdots \subset \Delta^{r}(p)=T_{p} M$ for some integer $r$ referred to as the degree of nonholonomy at $p$. Let $n_{j}$ stand for the dimension of $\Delta^{j}(p)$. The nonholonomic weights $w_{1}, w_{2}, \ldots, w_{n}$ at $p$ are defined by $w_{i}=j \Longleftrightarrow n_{j-1}<$ $i \leq n_{j}$.

Let $\left(x_{1}, x_{2}, \ldots, x_{n}\right)$ be a system of coordinates centered at $p$. These coordinates are privileged if for each $i$ there exist $w_{i}$ vector fields in $\Delta$ such that the Lie derivative $X_{j_{1}} X_{j_{2}} \ldots X_{j_{i}} x_{i}$ does not vanish at $p=0$ but that any such Lie derivative of length smaller than $w_{i}$ vanishes at 0 (see [5] or [10]).

Systems of privileged coordinates always exist (under the rank condition) and in such a system the weighted degree (homogeneous nonholonomic order) of the monomial $x_{1}^{\alpha_{1}} x_{2}^{\alpha_{1}} \ldots x_{n}^{\alpha_{n}}$ is $\alpha_{1} w_{1}+\alpha_{2} w_{2}+\cdots+\alpha_{n} w_{n}$, the weighted degree of the vector field $\frac{\partial}{\partial x_{j}}$ is $-w_{j}$, and the weighted degree of the vector field $x_{1}^{\alpha_{1}} x_{2}^{\alpha_{1}} \ldots x_{n}^{\alpha_{n}} \frac{\partial}{\partial x_{j}}$ is $\alpha_{1} w_{1}+\alpha_{2} w_{2}+\cdots+\alpha_{n} w_{n}-w_{j}$.

More generally the nonholonomic order at $p$ of a function $f$ (resp. a vector field $X$ ) is the minimum of the homogeneous nonholonomic orders of the monomials of its Taylor series.

It is important to notice that the nonholonomic degree $\operatorname{ord}_{p} X$ of a vector field $X$ at $p$ cannot be less than $-r$, and that $\operatorname{ord}_{p}[X, Y]=\operatorname{ord}_{p} X+\operatorname{ord}_{p} Y$ if $X$ and $Y$ are homogeneous and $[X, Y] \neq 0$ (see [5] or [10]).

The nonholonomic order of a vector field $X$ belonging to $\Delta$ is at least equal to -1 . Consider a set $X_{1}, X_{2}, \ldots, X_{n}$ of vector fields that generates $\Delta$ around $p$. In privileged coordinates each $X_{j}$ can be decomposed into

$$
X_{j}=X_{j}^{(-1)}+X_{j}^{(0)}+X_{j}^{(1)}+\cdots+X_{j}^{(s)}+\ldots
$$

where $X_{j}^{(s)}$ is the component of $X_{j}$ of homogeneous order $s$.
The nilpotent approximation of $X_{j} \in \Delta$ is $\widehat{X}_{j}=X_{j}^{(-1)}$.
The Lie algebra generated by $\widehat{X}_{1}, \widehat{X}_{2}, \ldots, \widehat{X}_{n}$ is nilpotent and finite dimensional. The rank condition is preserved: it is satisfied by $\widehat{X}_{1}, \widehat{X}_{2}, \ldots, \widehat{X}_{n}$ as soon it is satisfied by $X_{1}, X_{2}, \ldots, X_{n}$.

It may happen that some of the vector fields $\widehat{X}_{j}$ globally vanish. In that case they can be replaced by $\widetilde{X}_{j}=X_{j}^{(0)}$. Let us assume that only $m$ elements of $\widehat{X}_{1}, \widehat{X}_{2}, \ldots, \widehat{X}_{n}$ are linearly independant (as vector fields). As explained in the next section we can assume without lost of generality that these vector fields are $\widehat{X}_{1}, \widehat{X}_{2}, \ldots, \widehat{X}_{m}$ and that $\widehat{X}_{m+1}, \ldots, \widehat{X}_{n}$ vanish. The set of vector fields

$$
\widehat{X}_{1}, \ldots, \widehat{X}_{k}, \ldots, \widehat{X}_{m}, \widetilde{X}_{m+1}, \ldots, \widetilde{X}_{n}
$$

is called the solvable approximation of $X_{1}, X_{2}, \ldots, X_{n}$.

### 2.3 ARSs on Lie groups and homogeneous spaces

The purpose of this section is to recall the definition of linear ARSs on Lie groups and homogeneous spaces (see [4]).

Let $G$ be a connected Lie group and $\mathfrak{g}$ its Lie algebra (the set of leftinvariant vector fields, identified with the tangent space at the identity). A vector field $\mathcal{X}$ on $G$ is linear if its flow is a one-parameter group of automorphisms or equivalently if $\mathcal{X}(e)=0$ and for any left-invariant field $Y$ the Lie bracket $[\mathcal{X}, Y]$ is also left-invariant.

A linear ARS on $G$ is an almost-Riemannian structure defined by a set of:

- $n-s$ left-invariant vector fields $Y_{1}, \ldots, Y_{n-s}$.
- $s>0$ linear vector fields $\mathcal{X}_{n-s+1}, \ldots, \mathcal{X}_{n}$.
assumed to satisfy the rank condition and to have full rank on a proper open and dense subset of $G$. The almost-Riemannian metric is defined by the orthonormality of this set of vector fields.

The singular set $\mathcal{Z}$ where their rank is less than $n$ cannot be empty because at least $\mathcal{X}_{n}$ vanishes at the identity. On the other hand its interior is empty by analyticity as soon as the rank is full at one point.

A linear ARS is said to be simple if $s=1$. For instance, the famous Grushin plane on the Abelian Lie group $\mathbb{R}^{2}$ is a simple ARS.

Consider a homogeneous space $G / H$ of $G$ by a closed and connected subgroup $H$ (the elements of $G / H$ are the right cosets of $H$ ). The projection of a left-invariant vector field $Y$ onto $G / H$ is well-defined and will be referred to as a left-invariant vector field (see details in [11]). On the other hand the projection of a linear field $\mathcal{X}$ of $G$ does exist on $G / H$ if and only if $H$ is invariant under its flow, or equivalently, because $H$ is connected, if the Lie algebra of $H$ is $\operatorname{ad}(\mathcal{X})$-invariant. This allows to define linear vector fields and linear ARSs on homogeneous spaces in exactly the same way than on Lie groups.

## 3 Nilpotent and solvable approximations are linear

Though this section deals with local questions, around a point $p$ belonging to the singular locus, it will be more convenient to assume the ARS defined by a bundle $E$ and a morphism $f$ from $E$ to $T M$ as in Definition 1.

Firstly it is necessary to show that it is always possible to define the ARS locally, around the point $p=0$ in local privileged coordinates, by a set a $n$ orthonormal vector fields $X_{1}, \ldots, X_{n}$ such that the solvable approximation

$$
\widehat{X}_{1}, \ldots, \widehat{X}_{k}, \widehat{X}_{k+1}, \ldots, \widehat{X}_{m}, \widetilde{X}_{m+1}, \ldots, \widetilde{X}_{n}
$$

satisfies

- $\widehat{X}_{i}(0) \neq 0$ for $i=1, \ldots, k$;
- $\widehat{X}_{i} \neq 0$ but $\widehat{X}_{i}(0)=0$ for $i=k+1, \ldots, m$;
- $\widehat{X}_{i}=0$ for $i=m+1, \ldots, n$.

Let $K_{p}$ be the kernel of the restriction of $f$ to $E_{p}$, and let $V_{p}$ be an orthogonal complement to $K_{p}$ in $E_{p}$, that is $K_{p} \perp V_{p}$ and $K_{p} \oplus V_{p}=E_{p}$.

Let $e_{1}, \ldots, e_{n}$ be a set of $n$ sections of $E$, orthonormal in a neighborhood of $p$, such that $e_{j}(p) \in V_{p}$ for $j=1, \ldots, k$ and $e_{j}(p) \in K_{p}$ for $j=k+1, \ldots, n$.

The vector fields $X_{j}=f_{*}\left(e_{j}\right), j=1, \ldots, n$ define the ARS around $p$. Let $\left(x_{1}, \ldots, x_{n}\right)$ be a set of privileged coordinates and $\widehat{X}_{1}, \ldots, \widehat{X}_{n}$ be the related nilpotent approximation. Let $\mathcal{L}$ be the submodule of $\Gamma(E)$ generated by $e_{j}$ for $j=k+1, \ldots, n$. Consider now the mapping $e \in \mathcal{L} \longmapsto \widehat{f_{*} e}$.

Its rank is $m-k$ with $k \leq m \leq n$ and we can assume without lost of generality that $e_{m+1}, \ldots, e_{n}$ belong to the kernel of that linear map, and that $e_{k+1}, \ldots, e_{m}$ are orthogonal to that kernel.

The vector fields $X_{1}, \ldots, X_{n}$ satisfy the above conditions.
It may happen that the vector fields $\widetilde{X}_{m+1}, \ldots, \widetilde{X}_{n}$ fail to be linearly independant. In that case neither the nilpotent approximation nor the solvable one define an almost-Riemannian structure.

For that reason we will always assume in what follows that the vector fields

$$
\widehat{X}_{1}, \ldots, \widehat{X}_{k}, \widehat{X}_{k+1}, \ldots, \widehat{X}_{m}, \widetilde{X}_{m+1}, \ldots, \widetilde{X}_{n}
$$

are linearly independant. Denote by $\mathcal{Z}_{a}$ ( $a$ for approximation) the set of points where their rank is not full. It is not empty because at least one vector field vanishes at $p=0$. On the other hand the approximating vector fields are polynomial and the interior of $\mathcal{Z}_{a}$ is empty. Consequently $\mathcal{Z}_{a}$ is a proper with empty interior subset of $\mathbb{R}^{n}$ and the set of approximating vector fields defines an ARS.

## Remarks.

1. It will be shown in the second part of the paper that generically $m=n$ or $m=n-1$ and that in the second case $\widetilde{X}_{n} \neq 0$.
2. In the case where some of the $\widetilde{X}_{j}$ vanish or are linearly dependant it seems difficult, if not impossible, to go one step further by considering homogeneous approximations of nonholonomic order $s>0$ because, as explained in the sequel, two important properties could be lost. First a homogeneous vector field of degree $s>0$ need not be complete. Second the Lie algebra generated by the approximating vector fields would not be finite dimensional in general. These two drawbacks are related, see [14] or [11].

### 3.1 The generated Lie algebra

In view of the next sections it is very important to notice that all involved vector fields are complete, because of their triangular form.

This fact is well-known for the $\widehat{X}_{i}$ (see [5] or [10]).

It is as well true for the order zero homogeneous fields $\tilde{X}_{i}$. Indeed such a field writes:

$$
\widetilde{X}_{i}=\sum_{j=1}^{n} P_{i j}(x) \frac{\partial}{\partial x_{j}}
$$

where $P_{i j}$ is polynomial of nonholonomic degree $w_{j}$. In particular $P_{i j}$ depends linearly on the coordinates $\left(x_{j_{1}}, \ldots, x_{j_{s}}\right)$ of weight $w_{j}$, is polynomial in the coordinates of weight smaller than $w_{j}$, and does not depend on the coordinates of weight greater than $w_{j}$. The associated differential equation is consequently triangular: the coordinates of weight 1 satisfy a linear homogeneous equation, the coordinates of weight 2 satisfy a linear equation with a second member that depends on the the coordinates of weight 1 only, and so on. All solutions are therefore defined on $\mathbb{R}$ and $\widetilde{X}_{i}$ is complete. We can state:

Proposition 1 The vector fields $\widehat{X}_{j}$ and $\widetilde{X}_{j}$ defined above are complete.
This important property does not hold for homogeneous vector fields of positive degree, for example the first coordinate such a vector field could be $x_{1}^{2} \frac{\partial}{\partial x_{1}}$, but $\dot{x}_{1}=x_{1}^{2}$ is not complete.

The second feature we will use in the next subsections is the finiteness of the generated Lie algebra;

Proposition 2 The Lie algebra $\mathcal{L}$ generated by

$$
\widehat{X}_{1}, \ldots, \widehat{X}_{k}, \widehat{X}_{k+1}, \ldots, \widehat{X}_{m}, \tilde{X}_{m+1}, \ldots, \tilde{X}_{n}
$$

is finite dimensional.
Proof. The nonholonomic order of these vector fields is 0 or -1 , and the nonholonomic order of their brackets is in the range $-r, \ldots, 0$ where $r$ is the nonholonomic degree of the set of vector fields. Consequently, their components are polynomials of degrees less than or equal to $r$. The Lie algebra $\mathcal{L}$ is thus a subspace of a finite-dimensional vector space of polynomials.

### 3.2 The nilpotent case

It is the case where the vector fields $\widehat{X}_{1}, \ldots, \widehat{X}_{n}$ are linearly independent and the vectors $\widehat{X}_{1}(0), \ldots, \widehat{X}_{k}(0)$ are independent in $\mathbb{R}^{n}$. In particular no vector field $\widehat{X}_{i}$ vanishes, and $m=n$.

For $j=k+1, \ldots, n$ let $D_{j}$ stand for $\operatorname{ad}\left(\widehat{X_{j}}\right)$ and for any multi-index $J=\left(j_{1}, \ldots, j_{s}\right)$ let $D_{J}=D_{j_{s}} \circ \cdots \circ D_{j_{1}}$ (here $k+1 \leq j_{u} \leq n$ and $\left.s \geq 0\right)$. Let

$$
\mathcal{D}=\operatorname{Span}\left\{D_{J}\left(\widehat{X}_{i}\right) / i=1, \ldots, k ; J \text { as above }\right\} .
$$

Lemma 1 The Lie algebra $\mathfrak{g}$ generated by $\mathcal{D}$ is $D_{j}$-invariant for $j=k+$ $1, \ldots, n$.

Proof. Let $D_{J_{1}}\left(\widehat{X}_{i_{1}}\right)$ and $D_{J_{2}}\left(\widehat{X}_{i_{2}}\right)$ in $\mathcal{D}$. Then
$D_{j}\left[D_{J_{1}}\left(\widehat{X}_{i_{1}}\right), D_{J_{2}}\left(\widehat{X}_{i_{2}}\right)\right]=\left[D_{j} \circ D_{J_{1}}\left(\widehat{X}_{i_{1}}\right), D_{J_{2}}\left(\widehat{X}_{i_{2}}\right)\right]+\left[D_{J_{1}}\left(\widehat{X}_{i_{1}}\right), D_{j} \circ D_{J_{2}}\left(\widehat{X}_{i_{2}}\right)\right]$
belongs to $\mathfrak{g}$.
Let $\mathcal{L}$ stand for the Lie algebra generated by $\widehat{X}_{1}, \ldots, \widehat{X}_{n}$. It is a wellknown fact that this Lie algebra is nilpotent and finite dimensional (see [5] or [10]).

Theorem 1 1. The ideal generated in $\mathcal{L}$ by $\widehat{X}_{1}, \ldots, \widehat{X}_{k}$ is $\mathfrak{g}$. It is a nilpotent Lie algebra.
2. The vector fields $\widehat{X}_{k+1}, \ldots, \widehat{X}_{n}$ do not belong to $\mathfrak{g}$ and act on $\mathfrak{g}$ as derivations.
3. The rank at $p=0$ of the elements of $\mathfrak{g}$ is full.

Proof.

1. Since $\mathfrak{g}$ is $D_{j}$-invariant for $j>k$ it is clear that it is an ideal of $\mathcal{L}$ that contains $\widehat{X}_{1}, \ldots, \widehat{X}_{k}$. The ideal generated by these vector fields contains all the $D_{J}\left(X_{i}\right)$ hence is equal to $\mathfrak{g}$.
2. Let us assume that $\widehat{X}_{j}$ belongs to $\mathfrak{g}$ for some $j>k$. Because of the rule about the nonholonomic order of brackets of homogeneous vector fields, all the elements of $\mathfrak{g}$ of order -1 are linear combinations of $\widehat{X}_{1}, \ldots, \widehat{X}_{k}$. The vector field $\widehat{X}_{j}$ is homogeneous of order -1 and can consequently be written as:

$$
\widehat{X}_{j}=\sum_{i=1}^{k} \lambda_{i} \widehat{X}_{i} .
$$

But $\widehat{X}_{j}$ vanishes at 0 and the vectors $\widehat{X}_{i}(0)$ are independent by assumption, so that the $\lambda_{i}$ 's are all equal to 0 , a contradiction.
3. Let $Y \in \mathcal{L} \backslash \mathfrak{g}$. It can be obtained only as brackets of $\widehat{X}_{k+1}, \ldots, \widehat{X}_{n}$, hence $Y(0)=0$ because all these fields vanish at 0 .
If $\mathfrak{g}$ did not satisfy the rank condition then the Lie algebra $\mathcal{L}$ neither would satisfy it. This proves item 3 .

After this analysis at the algebra level we can turn our attention to the Lie group level.

Let $G$ be the simply connected Lie group whose Lie algebra is $\mathfrak{g}$. Since $\mathfrak{g}$ is nilpotent the underlying manifold of $G$ is $\mathbb{R}^{N}, N=\operatorname{dim}(\mathfrak{g})$. The first task is to show that $\mathbb{R}^{n}$ is a homogeneous space of $G$. This is mainly due to the fact that $\mathfrak{g}$ is generated by homogeneous vector fields.
Lemma 2 The set $\mathbb{R}^{n}$ is a homogeneous space of $G$. More accurately if $H$ stands for the connected subgroup of $G$ whose Lie algebra is the set of elements of $\mathfrak{g}$ that vanish at 0 , then $\mathbb{R}^{n}$ is diffeomorphic to the quotient $G / H$.

## Proof.

Since the elements of $\mathfrak{g}$ are complete vector fields of $\mathbb{R}^{n}$ the group $G$ acts naturally on $\mathbb{R}^{n}$ as a group of diffeomorphisms, and it is enough to show that this action is transitive. More accurately $G$ is the set of $\exp \left(t_{s} Y_{s}\right) \ldots \exp \left(t_{1} Y_{1}\right)$ where $Y_{i} \in \mathfrak{g}$ is a vector field on $\mathbb{R}^{n}$ and $t_{i} \in \mathbb{R}$, i.e. a group of diffeomorphisms of $\mathbb{R}^{n}$.

Let us begin by a simple remark. Let $Y=\sum_{j=1}^{n} y_{j}(x) \frac{\partial}{\partial x_{j}}$ be an element of $\mathfrak{g}$, and assume that $Y$ is homogeneous and that $Y(0) \neq 0$. There exists an index $i$ such that $y_{i}(0) \neq 0$, and by homogeneity $y_{i}$ is constant: $\forall x \in \mathbb{R}^{n}$, $y_{i}(x)=a_{i} \neq 0$. By homogeneity again the polynomials $y_{j}$ are homogeneous of order $w_{j}-w_{i}$. Consequently

$$
Y=\sum_{w_{j}=w_{i}} a_{j} \frac{\partial}{\partial x_{j}}+\sum_{w_{j}>w_{i}} y_{j}(x) \frac{\partial}{\partial x_{j}} \quad \text { where } \quad a_{i} \neq 0
$$

Since $\mathfrak{g}$ is generated by homogeneous elements and its rank is full at 0 , we can choose in $\mathfrak{g}$ a set $n$ homogeneous elements $Y_{1}, Y_{2}, \ldots, Y_{n}$ linearly independant at 0 . We can also assume that the nonholonomic order at $p$ of $Y_{1}, \ldots, Y_{n_{1}}$ is -1 , the one of $Y_{n_{1}+1}, \ldots, Y_{n_{2}}$ is -2 and so on. Up to linear combinations we can assume that $Y_{1}, \ldots, Y_{n_{1}}$ have the following form:

$$
Y_{i}=\frac{\partial}{\partial x_{i}}+\sum_{w_{j}>1} y_{j}(x) \frac{\partial}{\partial x_{j}}
$$

More generally we can assume that if the order of $Y_{i}$ is $w_{i}$ then

$$
Y_{i}=\frac{\partial}{\partial x_{i}}+\sum_{w_{j}>w_{i}} y_{j}(x) \frac{\partial}{\partial x_{j}}
$$

This way it is clear that the rank of the set of vector fields $Y_{1}, Y_{2}, \ldots, Y_{n}$ is full everywhere in $\mathbb{R}^{n}$. This implies that the action of $G$ on $\mathbb{R}^{n}$ is transitive.

To complete the construction we associate to the derivation $D_{j}=\operatorname{ad}\left(\widehat{X}_{j}\right)$ of $\mathfrak{g}$ a linear vector field $\mathcal{X}_{j}$ on $G$ for $j>k\left(\mathcal{X}_{j}\right.$ does exist because $G$ is simply connected). It is clear that the projection of $\mathcal{X}_{j}$ on $\mathbb{R}^{n}$ is $\widehat{X}_{j}$ (see [11] for details). Finally the vector fields $\widehat{X}_{1}, \ldots, \widehat{X}_{k}$ are invariant, and $\widehat{X}_{k+1}, \ldots, \widehat{X}_{n}$ are linear vector fields on the homogeneous space $\mathbb{R}^{n}=G / H$.

We can state:
Theorem 2 The space $\mathbb{R}^{n}$ is a homogeneous space of the nilpotent Lie group $G$ whose Lie algebra is $\mathfrak{g}$.

The vector fields $\widehat{X}_{1}, \ldots, \widehat{X}_{k}$ are projections of invariant vector fields of $G$ and $\widehat{X}_{k+1}, \ldots, \widehat{X}_{n}$ are projections of linear vector fields of $G$.

Consequently the set $\widehat{X}_{1}, \ldots, \widehat{X}_{n}$ defines a linear $A R S$ on the homogeneous space $\mathbb{R}^{n}$.

### 3.3 The non-nilpotent case

We set $D_{j}=\operatorname{ad}\left(\widehat{X_{j}}\right)$ for $j=k+1, \ldots, m$ and $D_{j}=\operatorname{ad}\left(\widetilde{X}_{j}\right)$ for $j=m+$ $1, \ldots, n$. As well as in the nilpotent case we set $D_{J}=D_{j_{s}} \circ \cdots \circ D_{j_{1}}$ for any multi-index $J=\left(j_{1}, \ldots, j_{s}\right)$ where $s \geq 0$ and $k+1 \leq j_{u} \leq n$, and

$$
\mathcal{D}=\operatorname{Span}\left\{D_{J}\left(X_{i}\right) / i=1, \ldots k ; J \text { as above }\right\}
$$

The Lie algebra $\mathfrak{g}$ generated by $\mathcal{D}$ is again $D_{j}$-invariant for $j=k+1, \ldots, n$, which shows that $\mathfrak{g}$ is the ideal generated in $\mathcal{L}$ by $\widehat{X}_{1}, \ldots, \widehat{X}_{k}$.

Theorem 3 1. The ideal generated in $\mathcal{L}$ by $\widehat{X}_{1}, \ldots, \widehat{X}_{k}$ is $\mathfrak{g}$. It is a nilpotent Lie algebra.
2. The vector fields $\widetilde{X}_{m+1}, \ldots, \widetilde{X}_{n}$ do not belong to $\mathfrak{g}$ and act on $\mathfrak{g}$ as derivations.
3. The vector fields $\widehat{X}_{j}$, with $k+1 \leq j \leq m$ that do not belong to $\mathfrak{g}$ act on $\mathfrak{g}$ as derivations.

## 4. The rank at $p=0$ of the elements of $\mathfrak{g}$ is full.

Proof. The proofs that $\mathfrak{g}$ is the ideal generated in $\mathcal{L}$ by $\widehat{X}_{1}, \ldots, \widehat{X}_{k}$ and that its rank at $p=0$ is full are identical to the nilpotent case.

The Lie algebra $\mathfrak{g}$ is generated by homogeneous vector fields of order at most -1 . Since the order of a Lie bracket is the sum of the orders of the factors and a vector field of order less than $-r$ vanishes, all brackets of length larger than $r$ vanish, which shows that $\mathfrak{g}$ is nilpotent.

The points 2 . and 3 . are clear.

Opposite to the nilpotent case we cannot assert that the vector fields $\widehat{X}_{k+1}, \ldots, \widehat{X}_{m}$ do not belong to $\mathfrak{g}$. Because of this phenomenon, illustrated by Example 3 in Section 5, we are lead to introduce one more index. Up to a reordering we can assume that $\widehat{X}_{k+1}, \ldots, \widehat{X}_{l}$ belong to $\mathfrak{g}$ and that $\widehat{X}_{l+1}, \ldots, \widehat{X}_{m}$ do not belong to $\mathfrak{g}$, where $k+1 \leq l \leq m$.

Theorem 4 The space $\mathbb{R}^{n}$ is a homogeneous space of the nilpotent Lie group $G$ whose Lie algebra is $\mathfrak{g}$.

The vector fields $\widehat{X}_{1}, \ldots, \widehat{X}_{l}$ are projections of invariant vector fields of $G$. The vector fields $\widehat{X}_{l+1}, \ldots, \widehat{X}_{m}$ are projections of linear or affine vector fields of $G$ and $\widetilde{X}_{m+1}, \ldots, \widetilde{X}_{n}$ are projections of linear ones.

Consequently the set of vector fields $\widehat{X}_{1}, \ldots, \widehat{X}_{m}, \widetilde{X}_{m+1}, \ldots, \widetilde{X}_{n}$ defines a linear ARS on the homogeneous space $\mathbb{R}^{n}$.

Proof. Similar to the one of Theorem 2.

## Remark.

As shown by Example 4 in Section 5 the Lie algebra $\mathcal{L}$ generated by $\widehat{X}_{1}, \ldots, \widehat{X}_{m}, \widetilde{X}_{m+1}, \ldots, \widetilde{X}_{n}$ need not be solvable when $m \leq n-2$.

However it is solvable if $m=n-1$ and it will be proven in the next section that generically $m=n-1$ or $m=n$.

It is why we call solvable the approximations of the previous kind.

## 4 Genericity

The examples of Section 5 show that many different, complicated structures may arise and the aim of this section is to determine the generic ones.

In what follows we will say that a property of almost-Riemannian distributions (resp. structures) on a manifold $M$ is generic if for any rank $n$
vector bundle $E$ (resp. Euclidean vector bundle $(E,\langle.,\rangle$.$) ) over M$ the set of smooth morphisms of vector bundles from $E$ to $T M$ for which this property is satisfied is open and dense in the $\mathcal{C}^{2}$ Whitney topology.

Let $U$ be an open subset of $M$ on which $E$ and $T M$ are trivializable, and let $\Pi$ be the projection from $E$ onto $M$. Then the restriction to $\Pi^{-1}(U)$ of a vector bundle morphism $f$ is equivalent to a smooth mapping $X$ from $U$ to the set $\mathcal{M}_{n}(\mathbb{R})$ of $n \times n$ square matrices.

Alternately $X$ can be viewed as a mapping $\left(X_{1}, X_{2}, \ldots, X_{n}\right)$ from $U$ to the set $\Gamma(U)^{n}$ of $n$ vector fields on $U$.

It is not useful to assume that $f$ satisfies the properties of almostRiemannian distributions because these conditions will turn out to be generic.

The first two theorems deal with distributions only, and do not require neither metric, nor normal forms.

In what follows we denote by $\mathcal{M}_{n \times m}(\mathbb{R})$ the set of real $n \times m$ matrices (simply $\left.\mathcal{M}_{n}(\mathbb{R})\right)$ if $m=n$ ) and by $L^{r}$ the set of elements of $\mathcal{M}_{n \times m}(\mathbb{R})$ of corank $r$. It is a submanifold of $\mathcal{M}_{n \times m}(\mathbb{R})$ of codimension $(n-q+r)(m-q+r)$ where $q=\min \{n, m\}$ (see [8])

Recall from Section 2 that $f$ being given, $\Delta$ stands for the submodule of $\Gamma(M)$ it defines.

### 4.1 Generic distributions

Theorem 5 The following properties are generic:

1. Let $R$ be the largest integer such that $R^{2} \leq n$. For $1 \leq r \leq R$ let $\mathcal{Z}_{r}$ be the set of points where the rank of $f_{p}$, or locally the rank of $\left\{X_{1}, X_{2}, \ldots, X_{n}\right\}$, is $n-r$. Each $\mathcal{Z}_{r}$ is a codimension $r^{2}$ submanifold and the singular locus $\mathcal{Z}$ is the union of these disjoint submanifolds.
2. The submanifold $\mathcal{Z}_{r+1}$ is included in the closure $\overline{\mathcal{Z}}_{r}$ of $\mathcal{Z}_{r}$ for $r=$ $1, \ldots, R-1$.
3. For any local representation $X$ of the distribution the mapping $x \longmapsto$ $\operatorname{det}(X(x))$ is a submersion at all points $x \in \mathcal{Z}_{1}$.
4. For $n \geq 3$ the rank of $\Delta+[\Delta, \Delta]$ is full at all points.

Proof.

1. Let $E$ be a rank $n$ vector bundle over $M$ and let $\mathcal{H o m}$ be the vector bundle over $M$ whose fiber at $p \in M$ is the vector space $\mathcal{H o m}=\mathcal{H o m}\left(E_{p}, T_{p} M\right)$ of homomorphisms from $E_{p}$ to $T_{p} M$.

A smooth vector bundle morphism $f$ from $E$ to $T M$ can be viewed as a smooth section of $\mathcal{H o m}$.

In each fiber $\mathcal{H o m}_{p}$ the set of morphisms of corank $r>0$ is a submanifold of $\mathcal{H o m}_{p}$ of codimension $r^{2}$. Their union over $M$ defines a codimension $r^{2}$ submanifold $S^{r}$ of $\mathcal{H o m}$. Since the union of the $S^{r}$ for $r=1, \ldots, n$ is closed, the set of smooth sections of $\mathcal{H o m}$ that are tranversal to all the $S^{r}$ is open and dense in the set of smooth sections of $\mathcal{H o m}$ endowed with the $\mathcal{C}^{2}$ Whitney topology.

The codimension of $S^{r}$ is $r^{2}$ and transversality means nonintersection if $r^{2}>n$. If $f^{-1}\left(S^{r}\right)$ is not empty transversality implies that it is a submanifold of $M$ of codimension $r^{2}$. This proves 1 .
2. In what follows we will always assume that the distributions under consideration are tranversal to the manifolds $S^{r}$ for $r=1, \ldots, n$.

Let such a distribution be locally defined by $X$ around a point $p$ assumed to belong to $\mathcal{Z}_{r}$ with $r \geq 2$. The rank of $\Lambda=X(p)$ is $n-r$ and there exists an invertible matrix $P$ such that $P \Lambda P^{-1}=\left(\begin{array}{ll}A_{\Lambda} & B_{\Lambda} \\ C_{\Lambda} & D_{\Lambda}\end{array}\right)$ where $A_{\Lambda} \in \mathcal{M}_{n-r}(\mathbb{R})$ is invertible. As in the proof of the product of coranks Theorem we consider $T=\left(\begin{array}{cc}I & 0 \\ -C_{\Lambda} A_{\Lambda}^{-1} & I\end{array}\right)$ so that $T P \Lambda P^{-1}=\left(\begin{array}{cc}A_{\Lambda} & B_{\Lambda} \\ 0 & D_{\Lambda}-C_{\Lambda} A_{\Lambda}^{-1} B_{\Lambda}\end{array}\right)$.

Let the matrix $P$ be fixed. The set $\Omega$ of matrices $Q$ such that, with obvious notations, $A_{Q}$ is invertible, hence $\Phi(Q)=D_{Q}-C_{Q} A_{Q}^{-1} B_{Q}$ is well defined, is open and the mapping $\Phi$ is a submersion from this open subset of $\mathcal{M}_{n}(\mathbb{R})$ onto $\mathcal{M}_{r}(\mathbb{R})$. Clearly rank $(Q)=n-r+\operatorname{rank}(\Phi(Q))$ and for $0 \leq s \leq r$ we have $L_{\mathcal{M}_{n}(\mathbb{R})}^{s}=\Phi^{-1}\left(L_{\mathcal{M}_{r}(\mathbb{R})}^{s}\right)$ (with obvious notations again).

Since the distribution is generic we have $X \Pi_{p} L^{r}$ which implies that $\Phi \circ X$ is a submersion at $p$ (see [8], Lemma 4.3). Let $V$ be an open neighborhood of $p$ and $W=V \bigcap X^{-1}(\Omega)$. Then $\Phi \circ X(W)$ is a neighborhood of 0 in $\mathcal{M}_{r}(\mathbb{R})$ that encounters $L_{\mathcal{M}_{r}(\mathbb{R})}^{s}$ for $0 \leq s \leq r$. Consequently

$$
\forall s, \quad 0<s<r, \quad V \bigcap \mathcal{Z}_{s} \neq \emptyset
$$

3. Consider now the determinant mapping, denoted by det, from $\mathcal{M}_{n}(\mathbb{R})$ to $\mathbb{R}$. The differential of det at $\Lambda$ applied to $H$ is $d \operatorname{det}(\Lambda) \cdot H=\operatorname{Trace}(\widetilde{\Lambda} H)$ where $\widetilde{\Lambda}$ stands for the transpose of the matrix of cofactors of $\Lambda$. Consequently $d \operatorname{det}(\Lambda) \neq 0$ if and only if $\operatorname{rank}(\Lambda) \geq n-1$. This shows that det is a submersion on $\mathcal{M}_{n}(\mathbb{R}) \backslash\left(\bigcup_{r \geq 2} L^{r}\right)$. In particular det is a submersion in a neighborhood $W$ of $L^{1}$ small enough for $L^{1}=\{\operatorname{det}=0\} \bigcap W$.

Let us consider a distribution locally defined by $X$ around a point $p$ assumed to belong to $\mathcal{Z}_{1}$. Since $X \Pi_{p} L^{1}$ and according to [8] (Lemma 4.3) det $\circ X$ is a submersion at $p$. In other words $q \longmapsto \operatorname{det}(X(q))$ is a submersion at $p$.
4. It remains to show that the rank of $\Delta+[\Delta, \Delta]$ is generically full everywhere. Let $X=\left(X_{1}, \ldots, X_{n}\right)$ be a set of $n$ vector fields on an open subset $U$ of $M$ that define locally the distribution. Up to a system of coordinates $X$ is a mapping from $U$ to $\left(\mathbb{R}^{n}\right)^{n} \equiv \mathcal{M}_{n}(\mathbb{R})$. Let $J^{1}$ be a typical fiber of the space of 1-jets of sections of $T^{n} U$, identified with $\left(\mathbb{R}^{n}\right)^{n} \times\left(\mathcal{M}_{n}(\mathbb{R})\right)^{n}$. The set $F_{i}$ of elements $\bar{X}=(X, d X)$ of $J^{1}$ such that $X_{i}=0$ and $\operatorname{rank}\left(d X_{i}\right)<n$ is closed with empty interior: it is a finite union of submanifolds of codimension $n+\rho^{2}$, where $\rho$ is the corank of $d X_{i}$. Consequently the set $\mathcal{O}=J^{1} \backslash\left(\bigcup_{i=1}^{n} F_{i}\right)$ is the open and dense subset of $J^{1}$ of elements that verify $X_{i}=0 \Longrightarrow d X_{i}$ invertible.

Let $\Psi$ be the mapping from $J^{1}$ to $\mathcal{M}_{n \times \frac{n(n+1)}{2}}(\mathbb{R})$ defined by

$$
\Psi(\bar{X})=\left(X,\left(\left[X_{i}, X_{j}\right]=d X_{j} \cdot X_{i}-d X_{i} \cdot X_{j}\right)_{1 \leq i<j \leq n}\right)
$$

The mapping $\Psi$ is a submersion on $\mathcal{O}$.
On the other hand the set $L^{r}$ of elements of $\mathcal{M}_{n \times \frac{n(n+1)}{2}}(\mathbb{R})$ of corank $r$ is a submanifold of codimension $r\left(\frac{n(n-1)}{2}+r\right)$ and $\Psi^{-1}\left(L^{r}\right) \bigcap \mathcal{O}$ is a submanifold of $\mathcal{O}$ of the same codimension.

But $n<r\left(\frac{n(n-1)}{2}+r\right)$ excepted in the particular case $n=2$ and $r=1$ where there is equality (this case has been studied in [3]).

To finish the set of $X$ that are transversal to the $F_{i}$, the union of which is closed, is open and dense. Such distributions take their values in $\mathcal{O}$, and the set of $X$ that are moreover transversal to all the $\Psi^{-1}\left(L^{r}\right) \bigcap \mathcal{O}$ is open and dense.

Excepted in the case $n=2$ and $r=1$ transversality means nonintersection and implies that the rank of $\Delta+[\Delta, \Delta]$ is full at all points.

Two subspaces of $T_{p} M$ are attached to a point $p$ belonging to the strate $\mathcal{Z}_{r}$ of the singular locus, namely $\Delta_{p}$, the distribution at $p$, and $T_{p} \mathcal{Z}_{r}$, the tangent subspace to $\mathcal{Z}_{r}$ at $p$. Their dimensions being respectively $n-r$ and $n-r^{2}$, the dimension of $T_{p} \mathcal{Z}_{r}+\Delta_{p}$ is at most equal to $\min \left(n, 2 n-r^{2}-r\right)$. We are interested in the cases where the actual dimension of $T_{p} \mathcal{Z}_{r}+\Delta_{p}$ is less than $\min \left(n, 2 n-r^{2}-r\right)$. For example in $\mathcal{Z}_{1}$ this means that $\Delta_{p}=T_{p} \mathcal{Z}_{r}$ (tangency points).

In what follows we note $s=\min \left(n, 2 n-r^{2}-r\right)-\operatorname{dim}\left(T_{p} \mathcal{Z}_{r}+\Delta_{p}\right)$, and $\lfloor\alpha\rfloor$ stands for the integer part of the real number $\alpha$.

Theorem 6 The following properties are generic:

- $r=1$. The points $p \in \mathcal{Z}_{1}$ where $T_{p} \mathcal{Z}_{1}=\Delta_{p}$ are isolated in $\mathcal{Z}_{1}$.
- $r \geq 2$. Let $m(n, r)$ be the largest dimension that $T_{p} \mathcal{Z}_{r}+\Delta_{p}$ may reach, that is $m(n, r)=\min \left(n, 2 n-r^{2}-r\right)$, and let $s=m(n, r)-\operatorname{dim}\left(T_{p} \mathcal{Z}_{r}+\right.$ $\Delta_{p}$ ). Then

1. The set of points $p \in \mathcal{Z}_{r}$ where $s=1$ is a submanifold of $\mathcal{Z}_{r}$ for $n \geq r^{2}+r-\left\lfloor\frac{r-1}{2}\right\rfloor$. It is empty if $n<r^{2}+r-\left\lfloor\frac{r-1}{2}\right\rfloor$.
2. The set of points $p \in \mathcal{Z}_{r}$ where $s \geq 2$ and $s^{2} \leq r$ is a submanifold of $\mathcal{Z}_{r}$ for $r^{2}+r-\left\lfloor\frac{r-s^{2}}{s-1}\right\rfloor \leq n \leq r^{2}+r+\left\lfloor\frac{r-s^{2}}{s-1}\right\rfloor$. It is empty if $n$ is not in this interval.
3. The set of points $p \in \mathcal{Z}_{r}$ where $s \geq 2$ and $s^{2}>r$ is empty.

Proof. In this proof all distributions are assumed to satisfy the transversality conditions of Theorem 5 .

1. Let us consider such a distribution and a point $p$ in $\mathcal{Z}_{r}$ for some $r>0$, locally defined by $X$ on an open set $U$ containing $p$.
Moreover the set $L^{r}$ is a codimension $r^{2}$ submanifold of $\mathcal{M}_{n}(\mathbb{R})$ and it can be locally defined, in a neighborhood $V$ of the point $X(p) \in L^{r}$, by a submersion $\Phi$ from $V$ to $\mathbb{R}^{r^{2}}$, so that $L^{r} \bigcap V=\Phi^{-1}(0)$.
2. Let us denote by $d(p)$ the dimension of the quotient $\left(T_{p} \mathcal{Z}_{r}+\Delta_{p}\right) / T_{p} \mathcal{Z}_{r}$. Since the dimension of $\mathcal{Z}_{r}$ is $n-r^{2}$ and the one of $\Delta_{p}$ is $n-r$ we have $n-r \leq \operatorname{dim}\left(T_{p} \mathcal{Z}_{r}+\Delta_{p}\right) \leq \min \left(n, 2 n-r^{2}-r\right)$, and:

$$
r^{2}-r \leq d(p) \leq \min \left(n-r, r^{2}\right)
$$

The distribution being assumed to be transversal to $L^{r}$ at $p$ the restriction of $d X$ at any supplementary subspace to $T_{p} \mathcal{Z}_{r}$ in $T_{p} U$ is an isomorphism onto its image. On the other hand $d X(p)\left(T_{p} \mathcal{Z}_{r}\right) \subset \operatorname{ker}\left(d \Phi\left(X_{p}\right)\right)$ because $\mathcal{Z}_{r}=X^{-1}\left(L^{r}\right)$, and since $\Phi$ is a submersion at $p$ we obtain $d(p)=\operatorname{rank} d(\Phi \circ X)(p) \cdot X(p)$.
Moreover there are exactly $n-r$ indices such that the vector fields $X_{i_{1}}, X_{i_{2}}, \ldots, X_{i_{n-r}}$ are independant at $p$ so that

$$
d(p)=\operatorname{rank}\left(d(\Phi \circ X)(p) \cdot X_{i_{1}}(p), \ldots, d(\Phi \circ X)(p) \cdot X_{i_{n-r}}(p)\right)
$$

3. In a neighborhood of $p$ where the vector fields $X_{i_{1}}, X_{i_{2}}, \ldots, X_{i_{n-r}}$ are independant at all points we define:

$$
x \longmapsto \Theta(x)=\left(X(x) ; d(\Phi \circ X)(x) \cdot X_{i_{1}}(x), \ldots, d(\Phi \circ X)(x) \cdot X_{i_{n-r}}(x)\right)
$$

The mapping $\Theta$ takes its values in $\mathcal{M}_{n}(\mathbb{R}) \times \mathcal{M}_{r^{2} \times(n-r)}(\mathbb{R})$ and $\Theta(x)$ belongs to $L^{r} \times L^{s}$ if and only if $x \in \mathcal{Z}_{r}$ and $d(x)=\min \left(r^{2}, n-r\right)-s$.
This shows that $s$ measures the gap between $d(p)$ and the maximal value it can take. Consequently we are interested in the values of $s$ that verify:

$$
0 \leq s \leq \min \left(r, n-r^{2}\right)
$$

4. The sets $L^{r} \times L^{s}$ are submanifolds in the typical fibers of the space of 1 -jets of smooth sections of $\mathcal{H o m}$. Since their union is closed the set of smooth sections of $\mathcal{H}$ om that are transversal to all these submanifolds is open and dense in the set of smooth sections of $\mathcal{H o m}$ endowed with the Whitney $\mathcal{C}^{2}$ topology.
5. Let us first assume that $n-r \geq r^{2}$. The codimension of $L^{r} \times L^{s}$ in $\mathcal{M}_{n}(\mathbb{R}) \times \mathcal{M}_{r^{2} \times(n-r)}(\mathbb{R})$ is $r^{2}+s\left(n-r^{2}-r+s\right)$.
For $s=1$ it is $n-r+1$. Since $n \geq n-r+1$ the set of points of $\mathcal{Z}_{r}$ where $s=1$ is generically a submanifold of codimension $n-r+1$. In particular for $r=1$ the codimension of this submanifold is $n$ and it consists in isolated points.
For $s \geq 2$ transversality means nonintersection if $n<r^{2}+s\left(n-r^{2}-\right.$ $r+s)$. But the two conditions $n \geq r^{2}+s\left(n-r^{2}-r+s\right)$ and $n-r \geq r^{2}$ are equivalent to $0 \leq n-r^{2}-r \leq \frac{r-s^{2}}{s-1}$, hence to:

$$
\begin{equation*}
r^{2}+r \leq n \leq r^{2}+r+\left\lfloor\frac{r-s^{2}}{s-1}\right\rfloor \tag{1}
\end{equation*}
$$

Notice that $s^{2} \leq r$ is a necessary condition for this inequality to hold.
6. Let us now assume that $n-r \leq r^{2}$. The codimension of $L^{r} \times L^{s}$ is $r^{2}+s\left(r^{2}+r+s-n\right)$. Again transversality means nonintersection if $n<r^{2}+s\left(r^{2}+r+s-n\right)$. The two conditions $n-r \leq r^{2}$ and $n \geq r^{2}+s\left(r^{2}+r+s-n\right)$ are equivalent to:

$$
\begin{equation*}
r^{2}+r-\left\lfloor\frac{r-s^{2}}{s-1}\right\rfloor \leq n \leq r^{2}+r \tag{2}
\end{equation*}
$$

Again $s^{2} \leq r$ is necessary for this inequality to hold.
7. The results of items 4 and 5 provide all the statements of the theorem.

### 4.2 Normal forms of almost-Riemannian structures

We consider an almost-Riemannian structure defined by a Euclidean vector bundle $(E,\langle\rangle$,$) and a vector bundle morphism f$, and we are interested in local normal forms of orthonormal vector fields defining the structure in a neighborhood of a point $p$ that we can assume to be $p=0$ in local coordinates.

These normal forms turn out to be the key of the next section.
First we follow the lines of [2] (also used in [7] and [6]).
Let $W$ be a codimension 1 submanifold transversal to the distribution. We can define a coordinate system $y=\left(x_{2}, \ldots, x_{n}\right)$ in $W$, and choose an orientation transversal to $W$. Let $\gamma_{y}$ be the family of normal geodesics parametrized by arclength, transversal to $W$ at $y$, and positively oriented. The mapping $\left(x_{1}, y\right) \mapsto \gamma_{y}\left(x_{1}\right)$ is a local diffeomorphism and the geodesics $x_{1} \mapsto \gamma_{y}\left(x_{1}\right)$ realize the minimal distance between $W=\left\{x_{1}=0\right\}$ and the surfaces $\left\{x_{1}=c\right\}$ for $c$ small enough. The transversality conditions of the PMP are consequently satisfied along all these surfaces: if $\lambda\left(x_{1}\right)$ is a covector associated to one of these geodesics then the tangent space to $\left\{x_{1}=c\right\}$ at $\gamma_{y}\left(x_{1}\right)$ is $\operatorname{ker}\left(\lambda\left(x_{1}\right)\right)$.

Now let $X_{1}=\partial_{x_{1}}$ be the vector field defined by $X_{1}(q)=\frac{d}{d x_{1}} \gamma_{y}\left(x_{1}\right)$ at the point $q=\gamma_{y}\left(x_{1}\right)$. It is a unitary vector field belonging to $\Delta$. Let $X_{2}, \ldots, X_{n}$ be $n-1$ vector fields such that $\left\{X_{1}, X_{2}, \ldots, X_{n}\right\}$ be an orthonormal frame of $\Delta$. For geodesics the control functions $\left(u_{1}, u_{2}, \ldots, u_{n}\right)$ from the PMP satisfy $u_{j}=\left\langle\lambda, X_{j}\right\rangle$. But here $\left(u_{1}, u_{2}, \ldots, u_{n}\right)=(1,0 \ldots, 0)$ so that $\left\langle\lambda, X_{j}\right\rangle=0$ for $j=2, \ldots, n$ and the vector fields $X_{2}, \ldots, X_{n}$ are tangent to the surfaces $\left\{x_{1}=c\right\}$. Consequently the vector fields have the following form:

$$
X_{1}=\left(\begin{array}{c}
1 \\
0 \\
\cdot \\
\cdot \\
0
\end{array}\right), \quad X_{j}=\left(\begin{array}{c}
0 \\
a_{2, j} \\
\cdot \\
\cdot \\
a_{n, j}
\end{array}\right) \quad \text { for } 1<j<n
$$

for any choice of the coordinates in $W$ and any choice of $X_{2}, \ldots, X_{n}$, under the condition that they provide an orthonormal frame related to the subRiemannian metric.

Notice that thanks to the transversality conditions the vectors fields $X_{2}, \ldots, X_{n}$ are not only orthogonal to $X_{1}$ for the sub-Riemannian metric but also for the canonical inner product of $\mathbb{R}^{n}$ for the chosen coordinates.

Let us assume now that $p=0$ belongs to $\mathcal{Z}_{r}$ with $r \geq 1$ and $r^{2} \leq n$. We want to show that the coordinates and $X_{2}, \ldots, X_{n}$ can be chosen in such a way that they write:

$$
\begin{aligned}
& X_{1}=\left(\begin{array}{c}
1 \\
0 \\
\cdot \\
\cdot \\
\cdot \\
\cdot \\
0
\end{array}\right) X_{2}=\left(\begin{array}{c}
0 \\
1+b_{2}(x) \\
a_{3,2}(x) \\
\cdot \\
\cdot \\
\cdot \\
a_{n, 2}(x)
\end{array}\right) \quad \ldots \quad X_{n-r}=\left(\begin{array}{c}
0 \\
\cdot \\
\cdot \\
1+b_{n-r}(x) \\
a_{n-r+1, n-r} \\
\cdot \\
a_{n, n-r}(x)
\end{array}\right) \quad \text { and } \\
& \left.\left(\begin{array}{c}
X_{n-r+1}
\end{array}\right) \quad . \quad X_{n}\right)=\left(\begin{array}{c}
0 \\
\cdot \\
\cdot \\
0 \\
----- \\
D(x)
\end{array}\right) \quad \text { where } D(x) \in \mathcal{M}_{r}(\mathbb{R})
\end{aligned}
$$

$b_{j}(0)=a_{i, j}(0)=0$, and $D(0)=0$.
Firstly we can assume that $X_{j}(0)=\partial_{x_{j}}$ for $j=2, \ldots, n-r$ and $X_{j}(0)=0$ for $j=n-r+1, \ldots, n$. Indeed $X_{2}, \ldots, X_{n}$ are tangent to $W$ where we can choose freely the cooordinates $\left(x_{2}, \ldots, x_{n}\right)$.

If $r=n-1$, which is generically possible only if $n=2$ and $r=1$, it is finished. Otherwise we can first replace $X_{2}$ by the normalization of $\sum_{j=2}^{n} a_{2 j} X_{j}$. Then we replace $X_{j}$ by $X_{j}-\frac{a_{2 j}}{a_{22}} X_{2}$ for $j>2$. These vector fields belong to $\Delta$, their first two coordinates vanish and they are orthonormal to (the new) $X_{2}$. It remains to orthonormalize these $n-2$ vector fields. This cannot be done directly because some of them vanish in the singular locus. However they are images by the vector bundle morphism $f$ of locally nonvanishing smooth sections of $E$ that can be orthonormalized.

The desired form of the vector fields is obtained by induction.

### 4.3 Nilpotent and solvable approximations of generic distributions

Theorem 7 For a generic distribution holds:
(i) Let $p$ be a tangency point in $\mathcal{Z}_{1}$, that is a point where $T_{p} \mathcal{Z}_{1}=\Delta_{p}$. Then $\widehat{X}_{n}=0$ but $\widetilde{X}_{n} \neq 0$, in normal form.
(ii) At all other points, including all points in $\mathcal{Z}_{r}$ with $r \geq 2$, the nilpotent approximation $\widehat{X}_{1}, \widehat{X}_{2}, \ldots, \widehat{X}_{n}$ is a set of $n$ linearly independant vector fields.

Proof.
Consider a generic distribution, a point $p$ in $\mathcal{Z}_{1}$ and privileged coordinates centered at $p$ such that
$X_{1}=\left(\begin{array}{c}1 \\ 0 \\ \cdot \\ \cdot \\ \cdot \\ 0\end{array}\right) X_{2}=\left(\begin{array}{c}0 \\ 1+b_{2}(x) \\ a_{3,2}(x) \\ \cdot \\ \cdot \\ a_{n, 2}(x)\end{array}\right) \ldots X_{n-1}=\left(\begin{array}{c}0 \\ \cdot \\ \cdot \\ 0 \\ 1+b_{n-1}(x) \\ a_{n, n-1}(x)\end{array}\right) X_{n}=\left(\begin{array}{c}0 \\ \cdot \\ \cdot \\ \cdot \\ 0 \\ a_{n}(x)\end{array}\right)$
where $b_{j}(0)=a_{i, j}(0)=0$. The determinant of $X$ is $a_{n}(x) \Pi_{j=2}^{n-1}\left(1+b_{j}(x)\right)$ and the singular locus is locally $\mathcal{Z}=\mathcal{Z}_{1}=\left\{a_{n}=0\right\}$.

Let $a_{n}(x)=\sum_{i=1}^{n} \alpha_{i} x_{i}+o(\|x\|)$. Since the determinant is a submersion at $p$ there exists $i$ such that $\alpha_{i} \neq 0$.

1. If there exists $i_{0} \leq n-1$ such that $\alpha_{i_{0}} \neq 0$ then $\widehat{X}_{n} \neq 0$. Moreover $d a_{n}(0) \cdot X_{i_{0}}=\alpha_{i_{0}} \neq 0$. Hence $X_{i_{0}}(0) \notin T_{0} \mathcal{Z}_{1}$ and $T_{0} \mathcal{Z}_{1}+\Delta_{0}=\mathbb{R}^{n}$.
2. If $\alpha_{i}=0$ for $i=1, \ldots, n-1$, then $\widehat{X}_{n}=0$, but $\alpha_{n} \neq 0$ and $\widetilde{X}_{n}=\left(\alpha_{n} x_{n}+q\left(x_{1}, \ldots, x_{n-1}\right)\right) \frac{\partial}{\partial x_{n}}$, where $q\left(x_{1}, \ldots, x_{n-1}\right)$ is quadratic. Moreover $T_{0} \mathcal{Z}_{1}=\operatorname{ker} d a_{n}(0)=\left\{x_{n}=0\right\}$, hence $X_{j}(0) \in T_{0} \mathcal{Z}_{1}$ for $j=1, \ldots, n-1$ and $\Delta_{0}=T_{0} \mathcal{Z}_{1}$.

Let now $p \in \mathcal{Z}_{r}$ with $r \geq 2$. We can choose privileged coordinates centered at $p$ such that:

$$
X_{1}=\left(\begin{array}{c}
1 \\
0 \\
\cdot \\
\cdot \\
\cdot \\
\cdot \\
0
\end{array}\right) X_{2}=\left(\begin{array}{c}
0 \\
1+b_{2}(x) \\
a_{3,2}(x) \\
\cdot \\
\cdot \\
\cdot \\
a_{n, 2}(x)
\end{array}\right) \quad \ldots \quad X_{n-r}=\left(\begin{array}{c}
0 \\
\cdot \\
\cdot \\
1+b_{n-r}(x) \\
a_{n-r+1, n-r} \\
\cdot \\
a_{n, n-r}(x)
\end{array}\right) \quad \text { and }
$$

$$
\left(\begin{array}{lll}
X_{n-r+1} & \cdot & X_{n}
\end{array}\right)=\left(\begin{array}{c}
0 \\
\cdot \\
0 \\
0 \\
----- \\
D(x)
\end{array}\right)
$$

where $b_{j}(0)=a_{i, j}(0)=0$ and $D(x) \in \mathcal{M}_{r}(\mathbb{R})$.
Let $x=(y, z)$ where $y=\left(x_{1}, \ldots, x_{n-r}\right)$ and $z=\left(x_{n-r+1}, \ldots, x_{n}\right)$. The assumption $p=0 \in \mathcal{Z}_{r}$ implies $D(0)=0$ and we can write:

$$
D(x)=l(y)+k(z)+\mathcal{Q}(y, z)
$$

where $l$ and $k$ are linear and $\mathcal{Q}$ contains all the terms of degree greater than one. Here $l(y)=\left(l_{i j}(y)\right)_{n-r+1 \leq i, j \leq n}$ belongs to $\mathcal{M}_{r}(\mathbb{R})$ and each entry $l_{i j}$ is linear w.r.t. $y=\left(x_{1}, \ldots, x_{n-r}\right)$.

As in the proof of Theorem 6 we denote by $d(p)$ the dimension of the quotient $\left(T_{p} \mathcal{Z}_{r}+\Delta_{p}\right) / T_{p} \mathcal{Z}_{r}$. It is clear that here $d(0)=\operatorname{dim}\left(d D(0) . \Delta_{0}\right)$, at $p=0$. But $\Delta_{0}=\operatorname{Span}\left\{X_{1}(0), \ldots, X_{n-r}(0)\right\}$ and for $j=1, \ldots, n-r$ :

$$
d D(0) \cdot X_{j}(0)=\frac{\partial D}{\partial x_{j}}(0)=\frac{\partial l}{\partial x_{j}}(0) .
$$

Consequently $d(0)=\operatorname{rank}(l)$ (as a mapping from $\mathbb{R}^{n-r}$ to $\mathcal{M}_{r}(\mathbb{R})$ ).
For $j=n-r+1, \ldots, n$ the nilpotent approximation $\widehat{X}_{j}$ of $X_{j}$ is

$$
\widehat{X}_{j}=\sum_{i=n-r+1}^{n} l_{i j}(y) \frac{\partial}{\partial x_{i}} .
$$

This is due to the fact that the weights of the coordinates are 1 for $i=$ $1, \ldots, n-r$ and 2 for $i=n-r+1, \ldots, n$.

Let us assume that the vector fields $\widehat{X}_{n-r+1}, \ldots, \widehat{X}_{n}$ are not linearly independant. Then the mapping $l$ takes its values in a subspace of $\mathcal{M}_{r}(\mathbb{R})$ of dimension $r^{2}-r$ and $d(p)$ is smaller than or equal to $r^{2}-r$. But $r^{2}-r \leq$ $d(p) \leq \min \left(n-r, r^{2}\right)$ according to the proof of Theorem 6 and the equality can hold for $r=1$ only. Indeed let $s=\min \left(n-r, r^{2}\right)-d(p)$ as in the proof of Theorem 6. Then:

1. If $n-r \geq r^{2}$ then $s=r^{2}-\left(r^{2}-r\right)=r$. But $r \geq 2$ implies $s \geq 2$ and $s^{2}>r$ which is impossible according to Theorem 6 .
2. If $n-r<r^{2}$ then $s=n-r-\left(r^{2}-r\right)=n-r^{2}$. But $n, r$ and $s$ must satisfy $n \geq r^{2}+s\left(r^{2}+r+s-n\right)$ according to the proof of Theorem 6 . For $s=n-r^{2}$ this is $s \geq s r$ which is impossible for $r \geq 2$.

## 5 Examples

The following four examples use the notations of Section 3.
The first one is a standard example of a solvable approximation on the group Heisenberg. Example 2 shows that the elements of the Lie algebra $\mathcal{L}$ of nonholonomic order smaller than -1 are not necessarily in $\mathfrak{g}$. We exhibit a vector field of the distribution that vanishes at $p$ but belongs to the ideal $\mathfrak{g}$ in Example 3 (which implies that the rank of the nilpotent approximation is not full).

To finish the Lie algebra $\mathcal{L}$ of Example 4 is not solvable, it contains a semi-simple subalgebra. Recall that this is not generic.

In these four examples, the vector fields are equal to their nilpotent or solvable approximations at 0 . It is of course possible to add terms of higher nonholonomic order without modifying the conclusions.

## Example 1

Consider in $\mathbb{R}^{3}$ the almost-Riemannian structure defined by the vector fields:

$$
X_{1}=\partial_{x}, \quad X_{2}=\partial_{y}+x \partial_{z}, \quad X_{3}=x \partial_{y}+\frac{1}{2} x^{2} \partial_{z}
$$

At $p=(0,0,0)$ the coordinates $(x, y, z)$ are privileged with weights $(1,1,2)$, the vector fields $X_{1}$ and $X_{2}$ are homogeneous of order -1 and $X_{3}$ is homogeneous of order 0 , so that $\widehat{X}_{1}=X_{1}, \widehat{X}_{2}=X_{2}, \widehat{X}_{3}=0$ and $\widetilde{X}_{3}=X_{3}$. The nilpotent approximation at $p$ is not an almost-Riemannian structure, it is the constant rank 2 sub-Riemannian structure defined by $X_{1}$ and $X_{2}$.

The Lie algebra generated by $\widehat{X}_{1}, \widehat{X}_{2}, \widetilde{X}_{3}$ is $\mathcal{L}=\operatorname{Span}\left\{\widehat{X}_{1}, \widehat{X}_{2}, \widetilde{X}_{3}, \partial_{z}\right\}$, the ideal $\mathfrak{g}=\operatorname{Span}\left\{\widehat{X}_{1}, \widehat{X}_{2}, \partial_{z}\right\}$ is here the Heisenberg Lie algebra, and $\widetilde{X}_{3}$ is a linear vector field on $\mathfrak{g}$. Finally $\widehat{X}_{1}, \widehat{X}_{2}, \widetilde{X}_{3}$ is a linear ARS on the Heisenberg group.

## Example 2

The almost-Riemannian structure is here defined in $\mathbb{R}^{3}$ by:

$$
X_{1}=\partial_{x}, \quad X_{2}=x \partial_{y}, \quad X_{3}=y^{2} \partial_{z}
$$

The Lie algebra $\mathcal{L}$ contains $X_{1}, X_{2}, X_{3}$ and

$$
\begin{aligned}
& X_{4}=\left[X_{1}, X_{2}\right]=\partial_{y}, X_{5}=\frac{1}{2}\left[X_{4}, X_{3}\right]=y \partial_{z} \\
& X_{6}=\left[X_{2}, X_{5}\right]=x \partial_{z}, X_{7}=\left[X_{1}, X_{6}\right]=\left[X_{4}, X_{5}\right]=\partial_{z} \\
& X_{8}=\frac{1}{2}\left[X_{2}, X_{3}\right]=x y \partial_{z}, X_{9}=\left[X_{2}, X_{8}\right]=x^{2} \partial_{z}
\end{aligned}
$$

Therefore:

$$
\begin{gathered}
\Delta^{1}=\Delta=\left\{X_{1}, X_{2}, X_{3}\right\}, \quad \Delta^{2}=\Delta^{1}+\left\{X_{4}, X_{8}\right\} \\
\Delta^{3}=\Delta^{2}+\left\{X_{5}, X_{9}\right\}, \quad \Delta^{4}=\Delta^{3}+\left\{X_{6}\right\}, \quad \Delta^{5}=\Delta^{4}+\left\{X_{7}\right\} .
\end{gathered}
$$

The canonical coordinates are privileged with weights $(1,2,5)$ and the vector fields $X_{1}, X_{2}$, and $X_{3}$ are homogeneous of order -1 hence equal to their nilpotent approximations.

The algebra $\mathfrak{g}$ is here the ideal of $\mathcal{L}$ generated by $X_{1}$ that is

$$
\mathfrak{g}=\operatorname{Span}\left\{X_{1}, X_{4}, X_{5}, X_{6}, X_{7}\right\}
$$

The vector fields $X_{2}$ and $X_{3}$ are linear, as well as $X_{8}=\frac{1}{2}\left[X_{2}, X_{3}\right]$ and $X_{9}=\left[X_{2}, X_{8}\right]$.

The orders of $X_{8}$ and $X_{9}$ are respectively -2 and -3 which shows that the vector fields of order smaller than -1 are not necessarily in $\mathfrak{g}$.

Notice that the singular locus is here $\mathcal{Z}=\{x y=0\}$.

## Example 3

Consider in $\mathbb{R}^{4}$ the almost-Riemannian structure defined by the vector fields:

$$
X_{1}=\partial_{x}, \quad X_{2}=\partial_{y}+x \partial_{z}, \quad X_{3}=y \partial_{w}, \quad X_{4}=\frac{1}{2} x^{2} \partial_{z}+\frac{1}{2} y^{2} \partial_{w}
$$

Since $\left[X_{1}, X_{2}\right]=\partial_{z}$ and $\left[X_{2}, X_{3}\right]=\partial_{w}$, the coordinates $(x, y, z, w)$ are privileged with weights $(1,1,2,2)$ at $p=(0,0,0,0)$. The vector fields $X_{1}, X_{2}$ are homogeneous of order -1 and independent at 0 , and the vector field $X_{3}$ is homogeneous of order -1 but vanishes at 0 . The last field $X_{4}$ is homogeneous of order 0 . Consequently the first three are equal to their nilpotent approximation and $X_{4}=\widetilde{X}_{4}$. According to the notations of Section 3 we have $k=2$ and $m=3$.

The Lie algebra $\mathcal{L}$ is spanned by $X_{1}, X_{2}, X_{3}, X_{4}$ and

$$
X_{5}=\left[X_{1}, X_{2}\right]=\partial_{z}, X_{6}=\left[X_{2}, X_{3}\right]=\partial_{w}, X_{7}=\left[X_{1}, X_{4}\right]=x \partial_{z}
$$

Despite the fact that $\widehat{X}_{3}(0)=X_{3}(0)=0$ we cannot assert as in the nilpotent case that $X_{3}$ does not belong to $\mathfrak{g}$ (see Section 3.3 after Theorem 3). Indeed the ideal generated in $\mathcal{L}$ by $X_{1}$ and $X_{2}$ is here

$$
\mathfrak{g}=\operatorname{Span}\left\{X_{1}, X_{2}, X_{3}, X_{5}, X_{6}, X_{7}\right\}
$$

because $X_{3}=\left[X_{2}, X_{4}\right]$.
As explained in Section 3.3 this may happen when $k<m<n$.
Notice that the determinant of $X_{1}, X_{2}, X_{3}, X_{4}$ is $-\frac{1}{2} x^{2} y$. Therefore the singular locus is $\mathcal{Z}=\{x y=0\}$ which shows that the structure is not generic.

In the general case the Lie algebra $\mathcal{L}$ need not be solvable. Indeed it is a subalgebra of the semi-direct product of $\mathfrak{g}$ by its algebra of derivations. But the algebra of derivations of a nilpotent Lie algebra is not solvable in general. For instance the derivations of the Heisenberg algebra is the set of endomorphisms the matrix of which writes in the canonical basis:

$$
D=\left(\begin{array}{ccc}
a & b & 0 \\
c & d & 0 \\
e & f & a+d
\end{array}\right)
$$

The subalgebra of such derivations that moreover satisfy $e=f=a+d=0$ is equal to $\mathfrak{s l}_{2}$ hence semisimple.

Example 4 illustrates that possibility.

## Example 4

Consider in $\mathbb{R}^{5}$, with coordinates $(x, y, z, w, t)$ :

$$
\begin{aligned}
& X_{1}=\partial_{x}, \quad X_{2}=\partial_{y}+x \partial_{w}+z \partial_{t}, \quad X_{3}=\partial_{z}+x \partial_{t} \\
& X_{4}=x \partial_{y}+\frac{1}{2} x^{2} \partial_{w}, \quad X_{5}=y \partial_{x}+\frac{1}{2} y^{2} \partial_{w}
\end{aligned}
$$

The Lie algebra $\mathcal{L}$ contains also

$$
\begin{aligned}
& X_{6}=\left[X_{1}, X_{2}\right]=\partial_{w}, \quad X_{7}=\left[X_{1}, X_{3}\right]=\left[X_{3}, X_{2}\right]=\partial_{t}, \\
& X_{8}=\left[X_{3}, X_{5}\right]=-y \partial_{t}, \quad X_{9}=\left[X_{8}, X_{4}\right]=x \partial_{t}, \\
& X_{10}=\left[X_{4}, X_{5}\right]=x \partial_{x}-y \partial_{y} .
\end{aligned}
$$

The coordinates $(x, y, z, w, t)$ are privileged with weights $(1,1,1,2,2)$ at the origin. At this point the vector fields $X_{1}, X_{2}, X_{3}$ (resp. $X_{4}, X_{5}$ ) are homogeneous of order -1 (resp. 0), hence equal to their nilpotent approximations (resp. $\widetilde{X}_{4}=X_{4}$ and $\widetilde{X}_{5}=X_{5}$ ).

Since $X_{10}=\left[X_{4}, X_{5}\right],\left[X_{10}, X_{4}\right]=2 X_{4}$ and $\left[X_{10}, X_{5}\right]=-2 X_{5}$ the vector fields $X_{4}, X_{5}$ and $X_{10}$, that do not belong to $\mathfrak{g}$, generate a semi-simple Lie algebra isomorphic to $\mathfrak{s l}_{2}$. Consequently the algebra $\mathcal{L}$ is not solvable.

The singular locus is here $\mathcal{Z}=\{x y z=0\}$, and again the structure is not generic.

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## APPENDIX D

## Paper 3

Distance Induced by the Solvable approximation of $n$-dimensional Almost-Riemannian Structures, article in progress.

# Distance Induced by the Solvable Approximation of $n$-dimensional Almost-Riemannian Structures 

Ronald Manríquez*


#### Abstract

The nilpotent and solvable approximations of a $n$-dimensional almostRiemannian structure (ARS) induce two different distances denoted by $\widehat{d}$ and $\widetilde{d}$ respectively. In this paper, under generic conditions, we prove that $\widetilde{d}$ is closer to to the original distance associated with the ARS than $\widehat{d}$.


Key words: Almost-Riemannian geometry, Nilpotent approximation, Solvable approximation.
AMS subject classifications: 53C15, 53C17, 22E25, 53B99.

## 1 Introduction

This paper aims to prove that generically, the distance induced by the solvable approximation is closer than the one induced by the nilpotent approximation to the original distance of a $n$-dimensional almost-Riemannian structure.

An almost-Riemannian structure (ARS) on a $n$-dimensional differential manifold $M$, locally can be defined by a set of $n$ smooth vector fields on $M$, $\left\{X_{1}, X_{2}, \ldots, X_{n}\right\}$, satisfying the Lie algebra rank condition (Larc in short). This set of $n$ vector fields is considered an orthonormal frame. The set of points where the linear span of the vector fields is not full-rank is called the singular locus or the singular set and denoted by $\mathcal{Z}$. If $\mathcal{Z}$ is empty, then the almost-Riemannian structure is a Riemannian one (more details in [1]).

In the generic 2 and 3 -dimensional cases, $\mathcal{Z}$ is a codimension one embedded submanifold. Furthermore, if $\Delta_{p}=\operatorname{span}\left\{X_{1}(p), X_{2}(p), \ldots, X_{n}(p)\right\}$, the points where $\Delta_{p}=T_{p} \mathcal{Z}$ are isolated (see Figure 1). Such points are called tangency points in [2] and type-2 points in [5]. In [6], the authors show that for generic $n$-dimensional ARSs, the singular locus $\mathcal{Z}$ is a union of submanifolds $\mathcal{Z}_{r}$ of codimension $r^{2}$ where the rank of $\Delta$ is $n-r$ and the tangency points in $\mathcal{Z}_{1}$ are isolated.

[^2]

Figure 1: Representation of a tangency point.

On the other hand, nilpotent approximations are used to locally study the behavior of almost-Riemannian structures due to their significant similarity to the original dynamics. However, there are cases where the nilpotent approximation of an ARS turns out to be a constant rank sub-Riemannian structure. That is, some vector fields may vanish. The above is exactly what happens in the generic 3-dimensional case, dealt in [5], where in privileged coordinates, the local representation of a 3-dimensional ARS at tangency points has the following nilpotent approximation (see section 2.2)

$$
\hat{X}_{1}=\left(\begin{array}{l}
1 \\
0 \\
0
\end{array}\right), \quad \hat{X}_{2}=\left(\begin{array}{l}
0 \\
1 \\
x
\end{array}\right), \quad \hat{X}_{3}=0
$$

which is the Heisenberg sub-Riemannian structure, and hence is not an ARS.
To recover the almost-Riemannian structure lost in the nilpotent approximation, the solvable approximation is introduced in [10]. In that paper, it was considered the case where only one of the vector fields of the nilpotent approximation vanishes and the other ones are independents. A generalization of this approximation is given in [6], where a complete description of the nilpotent and solvable approximations is addressed, and the generic almost-Riemannian structures are described.

This paper deals with the distance induced by the solvable approximation at tangency points, of an $n$-dimensional ARS considering generic assumptions, which are: the rank of $\Delta_{p}+[\Delta, \Delta]_{p}$ is $n$ for all $p \in M$, the singular set is a union of submanifolds $\mathcal{Z}_{r}$ of codimension $r^{2}$ where the rank of $\Delta$ is $n-r$, the points $p \in \mathcal{Z}_{1}$ where $T_{p} \mathcal{Z}_{1}=\Delta_{p}$ are isolated in $\mathcal{Z}_{1}$, and the solvable approximation can be defined only at a tangency point in $\mathcal{Z}_{1}$. Under the above conditions, we get a normal form of a generic $n$-dimensional ARS at a tangency point belonging to the singular set, and thus we obtain the nilpotent and solvable approximations with this local representation.

The original system, the nilpotent and solvable approximations give rise to three different distances: $d, \hat{d}$ and $\widetilde{d}$ respectively. Section 4 deals with the
almost-Riemannian distance defined by the solvable approximation. The main result is Theorem 5 which states that generically, the distance $\widetilde{d}$ is closer to $d$ than $\widehat{d}$ for pairs of points translated in an appropriate direction (Section 4.3). To prove this result, it is important to determine two facts. First, to state the order of approximation of $d$ by $\widetilde{d}$ (Theorem 2), and second, to find translation directions such that the distance $\widetilde{d}$ of a pair of translated points is decreasing (see Section 4.2).

To state the order of approximation of $d$ by $\tilde{d}$, we analyze the divergence of curves admissible for $d$ and $\widetilde{d}$, defined by the same control functions and starting at the same point (see Proposition 2). By using this fact, we obtain that the distance induced by the solvable approximation improves the order of approximation of $d$ given by $\hat{d}$ (see Theorem 2).

To find translation directions, we consider a vector field $Y \in \mathfrak{g}^{1}$ and then $Y \in \mathfrak{g}^{2}$, where $\mathfrak{g}=\mathfrak{g}^{1} \oplus \mathfrak{g}^{2}$ is the ideal generated in $\mathcal{L}$ by $\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}$, where $\mathfrak{g}^{s}$ is the set of homogeneous vector fields of order $-s$, and $\mathcal{L}$ is the Lie algebra generated by the solvable approximation.

The paper is organized as follows. Section 2 contains generalities about ARSs, nonholonomic order, privileged coordinates, the nilpotent approximation, and solvable one.

Section 3 is devoted to provide the normal form of a generic $n$-dimensional ARS at tangency points. Moreover, subsection 3.2 provides the solvable approximation, and the Hamiltonian associated with the flow defined by the solvable approximation in the normal form.

In Section 4, we analyze the divergence of curves admissible for $d$ and $\widetilde{d}$, we state the result related to the comparison of distances, and we address the translation, more accurately, we state conditions under which the distance defined by the solvable approximation is decreasing. Finally, subsection 4.3 is devoted to prove that the distance induced by the solvable approximation is strictly closer than $\hat{d}$ to the original distance $d$.

## 2 Preliminaries

This section is dedicated to present some definitions and results used in this paper related to the almost-Riemannian structures, nonholonomic order, privileged coordinates, the nilpotent approximation, and solvable one.

### 2.1 Almost-Riemannian structures

For all that concern general sub-Riemannian geometry, including almost- Riemannian one, the reader is referred to [1].

Let $M$ be a $n$-dimensional, connected, $C^{\infty}$ manifold. The $C^{\infty}$-module of $C^{\infty}$ vector fields on $M$ is denoted by $\Gamma(M)$. Let $\Delta$ be a sub-module of $\Gamma(M)$.

The flag of submodules

$$
\begin{equation*}
\Delta=\Delta^{1} \subseteq \Delta^{2} \subseteq \cdots \subseteq \Delta^{k-1} \subseteq \Delta^{k} \subseteq \ldots \tag{1}
\end{equation*}
$$

is defined by induction: $\Delta^{2}=\Delta+[\Delta, \Delta]$ is the sub-module of $\Gamma(M)$ generated by $\Delta$ and the Lie algebra of its elements, and $\Delta^{k+1}=\Delta^{k}+\left[\Delta, \Delta^{k}\right]$. The Lie algebra generated by $\Delta$ is $\mathcal{L}(\Delta)=\bigcup_{k \geqslant 1} \Delta^{k}$. The sub-module $\Delta$ satisfies the rank condition if the evaluaction of $\mathcal{L}(\Delta)$ at each point $p$ is equal to $T_{p} M$.

Definition 1. An almost-Riemannian structure (resp. distribution) on a smooth n-dimensional manifold $M$ is a triple $(E, f,\langle\cdot, \cdot\rangle)$ (resp. a pair $(E, f)$ ) where $E$ is a rank $n$ vector bundle over $M, f: E \rightarrow T M$ is a morphism of vector bundles, and $(E,\langle\cdot, \cdot\rangle)$ is an Euclidean bundle, that is $\langle\cdot, \cdot\rangle_{p}$ is an inner product on the fiber $E_{p}$ of $E$, smoothly varying w.r.t. $p$, assumed to satisfy the following properties:

1. The set of points $p \in M$ such that the restriction of $f$ to $E_{p}$ is onto is a proper open and dense subset of $M$;
2. the module $\Delta$ of vector fields on $M$ defined as the image by $f$ of the module of smooth sections of $E$ satisfies the rank condition.

The set of points of $M$ where the rank of $f\left(E_{p}\right)=\Delta_{p}$ is less than $n$ is called the singular locus (or singular set) of the almost-Riemannian structure and denoted by $\mathcal{Z}$.

The inner product on $E$ induces a bilinear symmetric and positive definite mapping, also denoted by $\langle\cdot, \cdot\rangle$, from $\Delta \times \Delta$ to $C^{\infty}(M)$. Indeed, and element $X \in \Delta$ (resp. $Y \in \Delta$ ) is the image by $f$ of a unique section $\sigma$ (resp. $\eta$ ) of $E$ and we can set $\langle X, Y\rangle_{p}=\langle\sigma, \eta\rangle_{p}$. Consequently andalmost-Riemannian structure (ARS in short) can be alternately defined as follows.

Definition 2. An almost-Riemannian structure on a smooth n-dimensional manifold $M$ is a pair $(\Delta,\langle\cdot, \cdot\rangle)$ where $\Delta$ is a sub-module of $\Gamma(M)$ that can be locally defined by $n$ vector fields and satisfies the rank condition, and $\langle\cdot, \cdot\rangle$ is a bilinear symmetric and positive definite mapping from $\Delta \times \Delta$ to $C^{\infty}(M)$, such that the set $\mathcal{Z}$ of points $p$ where the dimension of $\Delta_{p}$ is less than $n$ is nonempty but with empty interior.

Around any point $p \in M$ the sub-module $\Delta$ can be locally defined by an orthonormal frame $\left(X_{1}, X_{2}, \ldots, X_{n}\right)$. It is enough to select a set of $n$ sections $\left(e_{1}, e_{2}, \ldots, e_{n}\right)$ of $E$ orthonormal in a neighborhood of $p$ and define $X_{i}=f_{*} e_{i}$, where $f_{*} e_{i}=f \circ e_{i}$.

Remark 1. The structure is Riemannian out of $\mathcal{Z}$.

The almost-Riemannian norm on $\Delta_{p}$ is defined by

$$
\|v\|=\min \left\{\|u\|^{2}: u \in E_{p}, f(u)=v\right\} .
$$

An absolutely continuous curve $\gamma:[0, T] \longrightarrow M$ is admissible for $E$ if there exists a measurable essentially bounded function $t \mapsto u(t)$ from $[0, T]$ into $E_{\gamma(t)}$ called control function such that $\dot{\gamma}(t)=f(u(t))$ for almost every $t \in$ $[0, T]$. Locally this means that $\dot{\gamma}(t)=u_{1}(t) X_{1}(\gamma(t))+u_{2}(t) X_{2}(\gamma(t))+\ldots+$ $u_{n}(t) X_{n}(\gamma(t))$ for almost every $t \in[0, T]$, where $X_{1}, X_{2}, \ldots, X_{n} \in \Delta$.

Given an admissible curve $\gamma:[0, T] \longrightarrow M$, the length of $\gamma$ is defined by

$$
l(\gamma)=\int_{0}^{T}\|\dot{\gamma}(t)\| d t
$$

The almost-Riemannian distance (or Carnot-Caratheodory distance) on $M$ associated with the $n$-ARS is defined as

$$
d\left(p_{0}, p_{1}\right)=\inf \left\{l(\gamma): \gamma(0)=p_{0}, \gamma(T)=p_{1}, \gamma \text { admissible }\right\} .
$$

It induces the usual topology on $M$.

### 2.2 Privileged coordinates and approximations

Everything related to privileged coordinates and nilpotent approximation comes from [8] and [3]. The definition of solvable approximation was introduced in [10].

Let $p$ be a point of $M$ and let $\Delta_{p}^{k}, k \geqslant 1$ be the evaluation of the submodule $\Delta^{k}$ at $p$. Thanks to the rank condition these submodules verify $\Delta_{p}^{1} \subset \Delta_{p}^{2} \subset$ $\cdots \subset \Delta_{p}^{r-1} \nsubseteq \Delta_{p}^{r}=T_{p} M$, for some $r$ refered to as the degree of nonholonomy at $p$. Let $n_{j}(p)=\operatorname{dim} \Delta_{p}^{j}$. The nonholonomic weights $w_{1}, w_{2}, \ldots, w_{n}$ at $p$ are defined by $w_{i}=j \Leftrightarrow n_{j-1}(p)<i \leqslant n_{j}(p)$. Let $\left(x_{1}, x_{2}, \ldots, x_{n}\right)$ be a system of coordinates centered at $p$. These coordinates are privileged if for each $i$ there exist $w_{i}$ vector fields in $\Delta$ such that the Lie derivative $X_{j_{1}} X_{j_{2}} \ldots X_{j_{w_{i}}} x_{i}$ does not vanish at $p=0$ but that any such Lie derivative of length smaller than $w_{i}$ vanishes at 0 .

Systems of privileged coordinates always exist (under the rank condition) and in such a system the weighted degree (homogeneous nonholonomic order) of the monomial $x_{1}^{\alpha_{1}} x_{2}^{\alpha_{2}} \cdots x_{n}^{\alpha_{n}}$ is $w_{1} \alpha_{1}+\cdots+w_{n} \alpha_{n}$, the the weighted degree of the vector field $\frac{\partial}{\partial x_{j}}$ is $-w_{j}$, and the weighted degree of the vector field $x_{1}^{\alpha_{1}} x_{2}^{\alpha_{2}} \cdots x_{n}^{\alpha_{n}} \frac{\partial}{\partial x_{j}}$ is $w_{1} \alpha_{1}+\cdots+w_{n} \alpha_{n}-w_{j}$.

More generally, the nonholonomic order at $p$ of a function $f$ (resp. a vector field $X$ ) is the minimum of the homogeneous nonholonomic orders of the monomials of its Taylor series.

Definition 3. The function defined by $x \mapsto\|x\|_{p}=\sum_{i=1}^{n}\left|x_{i}\right|^{\frac{1}{w_{i}}}$ is the so-called pseudo-norm at $p$.

Remark 2. Let $x=\left(x_{1}, \ldots, x_{n}\right)$ be a system of privileged coordinates defined on an open neighborhood $U$ of the point $p$. When composed with the coordinate functions, the pseudo-norm at $p$ is (non smooth) homogeneous of order 1, that is, $\|x(q)\|_{p}=O(d(p, q))$, where $x(q)$ are the coordinates of $q \in U$.

It is important to notice that the nonholonomic degree $\operatorname{ord}_{p} X$ of a vector field $X$ at $p$ cannon be less than $-r$, and that $\operatorname{ord}_{p}[X, Y]=\operatorname{ord}_{p} X+\operatorname{ord}_{p} Y$ if $X$ and $Y$ are homogeneous and $[X, Y] \neq 0$ (more details see [3] and [8]).

The nonholonomic order of a vector field $X$ belonging to $\Delta$ is at least equal to -1 . Consider a set of vector fields $\left\{X_{1}, X_{2}, \ldots, X_{n}\right\}$ that generates $\Delta$ around $p$. In privileged coordinates each $X_{i}$ can be decomposed into

$$
X_{i}=X_{i}^{(-1)}+X_{i}^{(0)}+X_{i}^{(1)}+X_{i}^{(2)}+\cdots,
$$

where $X_{i}^{(s)}$ is the component of $X_{i}$ of homogeneous order $s$.
The nilpotent approximation of $X_{i} \in \Delta$ is $\widehat{X}_{i}=X_{i}^{(-1)}$.
The Lie algebra generated by $\hat{X}_{1}, \widehat{X}_{2}, \ldots, \hat{X}_{n}$ is nilpotent and finite dimensional. The rank condition is satisfied by $\widehat{X}_{1}, \widehat{X}_{2}, \ldots, \widehat{X}_{n}$ as soon it is satisfied by $X_{1}, X_{2}, \ldots, X_{n}$.

We consider the nilpotent approximation $\left\{\hat{X}_{1}, \ldots, \hat{X}_{n}\right\}$ of $\left\{X_{1}, \ldots, X_{n}\right\}$ at $p=0$ such that $\hat{X}_{n}=0$. Denoting $\tilde{X}_{n}=X_{n}^{(0)} \neq 0$, we have the following definition:
Definition 4 (Solvable approximation). The family $\left\{\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}, \widetilde{X}_{n}\right\}$ is the solvable approximation of $\left\{X_{1}, \ldots, X_{n}\right\}$.

More details see [10].

## 3 Generic $n$-dimensional almost-Riemannian structures

This section provides a normal form of a $n$-dimensional ARS around the point $p \in \mathcal{Z}$ such that $T_{p} \mathcal{Z}=\Delta_{p}$, under generic conditions. Moreover, we compute its solvable approximation. This normal form is essential in the following sections.

When we work with $n$-ARS, we can find different and complex structures. For this reason, in the following, we deal only with generic ARSs. We will say that a property of almost-Riemannian distributions (resp. structures) on a manifold $M$ is generic if for any rank $n$ vector bundle $E$ (resp. Euclidean vector bundle $(E,\langle\cdot, \cdot\rangle)$ over $M$ the set of smooth morphisms of vector bundles from $E$ to $T M$ for which this property is satisfied is open and dense in the $\mathcal{C}^{2}$ Whitney topology.

We take from [6] the following results.

Theorem 1 ([6]). The following properties are generic:

1. Let $R$ be the largest integer such that $R^{2} \leqslant n$. For $1 \leqslant r \leqslant R$ let $\mathcal{Z}_{r}$ be the set of points where the rank of $\left\{X_{1}, \ldots, X_{n}\right\}$ is $n-r$. Each $\mathcal{Z}_{r}$ is a codimension $r^{2}$ submanifold and the singular locus $\mathcal{Z}$ is the union of these disjoint submanifolds.
2. The points $p \in \mathcal{Z}_{1}$ where $T_{p} \mathcal{Z}_{1}=\Delta_{p}$ are isolated in $\mathcal{Z}_{1}$.
3. The mapping $x \mapsto \operatorname{det}(X(x))$ is a submersion at all points $x \in \mathcal{Z}_{1}$, where $X$ is the $n \times n$ matrix whose $j$ th column is $X_{j}$.
4. Let $p$ be a tangency point in $\mathcal{Z}_{1}$, that is a point where $T_{p} \mathcal{Z}_{1}=\Delta_{p}$. Then $\widehat{X}_{n}=0$ but $\widetilde{X}_{n} \neq 0$, in normal form.
5. At all other points, including the points in $\mathcal{Z}_{r}, r \geqslant 2$, the nilpotent approximation $\widehat{X}_{1}, \ldots, \widehat{X}_{n}$ is a set of $n$ linearly independant vector fields.
6. The rank of $\Delta_{p}+[\Delta, \Delta]_{p}$ is full at all points.

In conclusion, generically: (1) the points $p \in \mathcal{Z}_{1}$ where $T_{p} \mathcal{Z}_{1}=\Delta_{p}$ are isolated in $\mathcal{Z}_{1} ;(2)$ the solvable approximation can be defined only at a tangency point in $\mathcal{Z}_{1} ;(3)$ the degree of nonholonomy is 2 .

### 3.1 Normal forms: $n$-dimensional case

We consider an ARS defined by a Euclidean vector bundle $(E,\langle\cdot, \cdot\rangle)$ and a vector bundle morphism $f$. We are interested in local normal forms of orthonormal vector fields defining the structure in a neighborhood of a point $p$ that we can assume $p=0 \in \mathcal{Z}$ in local coordinates.

By Theorem 1, generically there are points where the rank is $n-r$, as long as $r^{2}<n$. In this paper, we consider only points belonging to $\mathcal{Z}_{1}$ because the only generic points where the solvable approximation is useful are tangency points in $\mathcal{Z}_{1}$ (see Theorem 1 item 4 and 5). Hence the distribution at $p=0$ has always dimension $n-1$, then assume that $X_{j}(0)=\frac{\partial}{\partial x_{j}}$ for $j=2, \ldots, n-1$. It is shown in [6] that the coordinates and $X_{1}, \ldots, X_{n}$ can be chosen in such a way that:

$$
X_{1}=\left(\begin{array}{c}
1 \\
0 \\
\vdots \\
\vdots \\
\vdots \\
0
\end{array}\right), X_{2}=\left(\begin{array}{c}
0 \\
1+\beta_{2}(x) \\
\alpha_{3,2}(x) \\
\vdots \\
\vdots \\
\alpha_{n, 2}(x)
\end{array}\right), \ldots, X_{n-1}=\left(\begin{array}{c}
0 \\
\vdots \\
\vdots \\
0 \\
1+\beta_{n-1}(x) \\
\alpha_{n, j}(x)
\end{array}\right), X_{n}=\left(\begin{array}{c}
0 \\
\vdots \\
\vdots \\
\vdots \\
0 \\
\alpha_{n}(x)
\end{array}\right)
$$

where $\beta_{j}(0)=\alpha_{i, j}(0)=0 i, j=2, \ldots, n$.
By the normal form of the vector fields, the singular locus is locally $\mathcal{Z}=$ $\mathcal{Z}_{1}=\left\{\alpha_{n}(x)=0\right\}$.

Let $\alpha_{n}(x)=\widehat{\alpha}_{n}(x)+\widetilde{\alpha}_{n}(x)+\bar{\alpha}_{n}(x)$, that is $\alpha_{n}$ decomposed into its components of nonholonomic order 1, 2 and greater than or equal to 3 respectively. Let $p$ be a tangency point in $\mathcal{Z}_{1}$, then by Theorem $1, \widehat{X}_{n}=0$ and $\widetilde{X}_{n} \neq 0$, hence $\alpha_{n}(x)=\widetilde{\alpha}_{n}(x)+\bar{\alpha}_{n}(x)$. Moreover, since $\operatorname{ord}_{p}\left(\widetilde{\alpha}_{n}\right)=2$ we get $\alpha_{n}(x)=a x_{n}+Q\left(x_{1}, x_{2}, \ldots, x_{n-1}\right)+\bar{\alpha}_{n}(x)$, where $Q\left(x_{1}, x_{2}, \ldots, x_{n-1}\right)$ is quadratic. Notice that, since the determinant is a submersion at $p$ (Theorem 1 item 3) $a \neq 0$.

### 3.2 Solvable approximation

We can find the nilpotent and solvable approximation in the coordinate system constructed in the normal form. Since $p=0$ is a tangency point and the weights of the coordinates are $w_{1}=\cdots=w_{n-1}=1$, and $w_{n}=2$, the nilpotent approximation is defined by

$$
\begin{aligned}
& \hat{X}_{1}=X_{1}=\frac{\partial}{\partial x_{1}} \\
& \hat{X}_{j}=\frac{\partial}{\partial x_{j}}+\hat{\alpha}_{n, j}(x) \frac{\partial}{\partial x_{n}} \quad \text { for } j=2, \ldots, n-1, \text { and } \\
& \hat{X}_{n}=0
\end{aligned}
$$

where $\widehat{\alpha}_{n, j}$ is the component of $\alpha_{n, j}$ of nonholonomic order 1 , hence it is linear in $x_{1}, x_{2}, \ldots, x_{n-1}$.

Therefore, the solvable approximation in the tangent case is defined by $\left\{\hat{X}_{1}, \ldots, \hat{X}_{n-1}, \widetilde{X}_{n}\right\}$ where

$$
\tilde{X}_{n}=\left(a x_{n}+Q\left(x_{1}, x_{2}, \ldots, x_{n-1}\right)\right) \frac{\partial}{\partial x_{n}}=\widetilde{a}_{n}(x) \frac{\partial}{\partial x_{n}},
$$

### 3.2.1 Hamiltonian equations to the solvable approximation

The following is helpful to prove one of the main results of this paper (see Section 4.3).

With the normal forms of the vector fields, the Hamiltonian for the normal flow is given by

$$
\begin{aligned}
H(\lambda(t)) & \left.=\frac{1}{2}\left(\sum_{i=1}^{n-1}\left\langle\lambda(t), \widehat{X}_{i}(x)\right)\right\rangle^{2}+\left\langle\lambda(t), \widetilde{X}_{n}(x)\right\rangle^{2}\right), \\
& =\frac{1}{2}\left(\lambda_{1}^{2}+\sum_{i=2}^{n-1}\left(\lambda_{i}+\lambda_{n} \widehat{\alpha}_{n, j}(x)\right)^{2}+\lambda_{n}^{2} \widetilde{a}_{n}(x)^{2}\right),
\end{aligned}
$$

where $\lambda(t)=\left(\lambda_{1}, \ldots, \lambda_{n}\right) \in T_{\gamma(t)}^{*} \mathbb{R}^{n}$. Hence
$\dot{x_{1}}(t)=\lambda_{1}$
$\dot{x_{j}}(t)=\lambda_{j}+\lambda_{n} \widehat{\alpha}_{n, j}(x) \quad($ for $j=2, \ldots, n-1)$
$\dot{x_{n}}(t)=\sum_{i=2}^{n-1} \widehat{\alpha}_{n, j}(x)\left(\lambda_{i}+\lambda_{n} \widehat{\alpha}_{n, j}(x)\right)+\lambda_{n} \widetilde{a}_{n}(x)^{2}$
$\dot{\lambda}_{j}(t)=-\sum_{i=2}^{n-1} \lambda_{n} \frac{\partial}{\partial x_{j}} \widehat{\alpha}_{n, j}(x)\left(\lambda_{i}+\lambda_{n} \widehat{\alpha}_{n, j}(x)\right)-\frac{\partial}{\partial x_{j}} Q\left(x_{1}, \ldots, x_{n-1}\right) \lambda_{n}^{2} \widetilde{a}_{n}(x)$, for $j=1, \ldots, n-1$
$\dot{\lambda}_{n}(t)=-a \lambda_{n}^{2} \widetilde{a}_{n}(x)$
are the associated Hamiltonian equations to the solvable approximation in the normal form.
Remark 3. Notice that $H(\lambda(0))=\frac{1}{2} \sum_{i=1}^{n-1} \lambda_{i}(0)^{2}$, and then it does not depend of $\lambda_{n}(0)$. In consequence, $\lambda_{n}(0)$ can be chosen arbitrarily.

## 4 Distance induced by the solvable approximation

The original system, the nilpotent and solvable approximations give rise to three different distances: $d, \widehat{d}$ and $\widetilde{d}$ respectively. This section deals with the almost-Riemannian distance defined by the solvable approximation. The main result is Theorem 5 which states that generically, the distance $\widetilde{d}$ is closer to $d$ than $\hat{d}$ for pairs of points translated in an appropriate direction (Section 4.3). This translation condition is significant because the distance $\tilde{d}$ is not closer to $d$ than the distance induced by the nilpotent approximation for any pair of points. Then to prove the main result, it is essential to determine two facts. First, to state the order of approximation of $d$ by $\widetilde{d}$ (Theorem 2), and second, to find translation directions such that the distance $\tilde{d}$ of a pair of translated points is decreasing (see Section 4.2).

A first proposition provides a relation between the solvable distance and the nilpotent one. This result is important because is useful to prove the main result of subsection 4.3.
Proposition 1. For all $x, y \in \mathbb{R}^{n}, \widetilde{d}(x, y) \leqslant \widehat{d}(x, y)$.
Proof. See [10].

### 4.1 Comparison of distances

In this subsection, we state the order of approximation of the original distance by $\widetilde{d}$. An important conclusion is that $\widetilde{d}$ improves the order of approximation
of $d$ given by the nilpotent approximation.
Let $q$ and $q^{\prime}$ belong to the ball centered at $p$ and radius $\epsilon$, denoted by $B(p, \epsilon)$. We start by analyzing the divergence of curves respectively admissible for $d$ and $\tilde{d}$, defined by the same control functions and starting at the same point $q$. Let us consider the vector fields $X_{1}, \ldots, X_{n}$ in normal form as in the above section. Hence we can express each vector field $X_{j}$ for $j=2, \ldots, n-1$, as

$$
X_{j}=\left(\begin{array}{c}
0 \\
\vdots \\
1+\beta_{j}(x) \\
\alpha_{j+1, j}(x) \\
\vdots \\
\alpha_{n, j}(x)
\end{array}\right) .
$$

Notice that $\alpha_{n, j}(x)$ can be split into components of order 1 and the remainder i.e $\alpha_{n, j}(x)=\widehat{\alpha}_{n, j}(x)+\alpha_{n, j}^{+}(x)$. We denote by $\rho_{n, j}^{+}$the order of $\alpha_{n, j}^{+}$for $j=$ $2, \ldots, n-1$.

Proposition 2. Let $\gamma$ be the geodesic for $d$ such that $\gamma(0)=q, \gamma(T)=q^{\prime}$ and associated with the control function $u(\cdot)$ satisfying $\|u\|=1$. Let $\widetilde{\gamma}$ be the admissible curve for $\widetilde{d}$ defined by the same control functions as $\gamma$ and $\widetilde{\gamma}(0)=q$. If $\rho_{n, j}^{+} \geqslant 3$ for $j=2, \ldots, n-1$, then

$$
\begin{equation*}
\|\gamma(t)-\widetilde{\gamma}(t)\|_{p} \leqslant C s t \cdot \tau^{\frac{3}{2}} \cdot t^{\frac{1}{2}}, \tag{2}
\end{equation*}
$$

where $\tau=\max \left(\|q\|_{p}, t\right)$.
Proof. We have the following:
$\dot{\gamma}(t)-\dot{\tilde{\gamma}}(t)=\sum_{j=1}^{n-1} u_{j}(t)\left(X_{j}(\gamma(t))-\widehat{X}_{j}(\widetilde{\gamma}(t))\right)+u_{n}(t)\left(X_{n}(\gamma(t))-\widetilde{X}_{n}(\widetilde{\gamma}(t))\right)$.
In details,

1. $\dot{\gamma}_{1}(t)-\dot{\tilde{\gamma}}_{1}(t)=0$,
2. $\dot{\gamma}_{i}(t)-\dot{\widetilde{\gamma}}_{i}(t)=\sum_{j=2}^{i-1} u_{j}(t) \alpha_{i, j}(\gamma(t))+u_{i}(t) \beta_{i}(\gamma(t))$, for $i=2, \ldots, n-1$,
3. 

$$
\begin{gathered}
\dot{\gamma}_{n}(t)-\dot{\widetilde{\gamma}}_{n}(t)=\sum_{j=2}^{n-1} u_{j}(t) \alpha_{n, j}^{+}(\gamma(t))+\sum_{j=2}^{n-1} u_{j}(t)\left(\widehat{\alpha}_{n, j}(\gamma(t))-\widehat{\alpha}_{n, j}(\widetilde{\gamma}(t))\right)+ \\
u_{n}\left(a\left(\gamma_{n}(t)-\widetilde{\gamma}_{n}(t)\right)+Q\left(\gamma_{1}(t), \ldots, \gamma_{n-1}(t)\right)-\right. \\
\left.Q\left(\widetilde{\gamma}_{1}(t), \ldots, \widetilde{\gamma}_{n-1}(t)\right)+\bar{\alpha}_{n}(\gamma(t))\right) .
\end{gathered}
$$

We have successively,

1. $\gamma_{1}(t)=\widetilde{\gamma}_{1}(t)$.
2. $\dot{\gamma}_{i}(t)-\dot{\widetilde{\gamma}}_{i}(t)=\sum_{j=2}^{i} u_{j}(t) \alpha_{i, j}(\gamma(t))+u_{i}(t) \beta_{i}(\gamma(t))$, hence $\left|\dot{\gamma}_{i}(t)-\dot{\widetilde{\gamma}}_{i}(t)\right| \leqslant$ $\sum_{j=2}^{i}\left|\alpha_{i, j}(\gamma(t))\right|$, for $i=2, \ldots, n-1$.
Since $\operatorname{ord}_{p}\left(\alpha_{i, j}\right) \geqslant 1$, for $j=2, \ldots, n-1$, then $\left|\alpha_{i, j}(\gamma(t))\right| \leqslant \operatorname{Cst} \cdot\|\gamma(t)\| \leqslant$ Cst $\cdot \tau$ because $\|\gamma(t)\|_{p} \leqslant$ Cst $\cdot \tau$, where $\tau=\max \left(\|q\|_{p}, t\right)$ (the proof of the above inequality is given in [8] in the proof of Lemma 2.2). Similarly for $\beta_{i}$, hence

$$
\begin{equation*}
\left|\gamma_{i}(t)-\widetilde{\gamma}_{i}(t)\right| \leqslant \operatorname{Cst} \cdot \tau \cdot t \quad \text { for } i=2, \ldots, n-1 \tag{3}
\end{equation*}
$$

3. 

$$
\begin{aligned}
& \dot{\gamma}_{n}(t)-\dot{\tilde{\gamma}}_{n}(t)=\sum_{j=2}^{n-1} u_{j}(t) \alpha_{n, j}^{+}(\gamma(t))+\sum_{j=2}^{n-1} u_{j}(t)\left(\widehat{\alpha}_{n, j}(\gamma(t))-\widehat{\alpha}_{n, j}(\widetilde{\gamma}(t))\right)+ \\
& +u_{n}\left(a\left(\gamma_{n}(t)-\widetilde{\gamma}_{n}(t)\right)+\sum_{k=2}^{n-1}\left(\gamma_{k}(t)-\widetilde{\gamma}_{k}(t)\right) P_{n, k}(\gamma, \widetilde{\gamma})+\right. \\
& \left.\quad \bar{\alpha}_{n}(\gamma(t))\right)
\end{aligned}
$$

where $P_{n, k}(\gamma, \widetilde{\gamma})$ is a homogeneous polynomial of weighted degree 1 , hence

$$
\begin{aligned}
\left|\gamma_{n}(t)-\widetilde{\gamma}_{n}(t)\right| \leqslant & \int_{0}^{t}\left|\alpha_{n, j}^{+}(s)\right| d s+\int_{0}^{t} \sum_{k=2}^{n-1}\left|\widehat{\alpha}_{n, j}(\gamma(s))-\widehat{\alpha}_{n, j}(\widetilde{\gamma}(s))\right| d s+ \\
& \int_{0}^{t}|a|\left|\gamma_{n}(s)-\widetilde{\gamma}_{n}(s)\right| d s+\int_{0}^{t}\left|\bar{\alpha}_{n}(\gamma(s))\right| d s+ \\
& \int_{0}^{t} \sum_{k=2}^{n-1}\left|\gamma_{k}(s)-\widetilde{\gamma}_{k}(s)\right|\left|P_{n, k}(\gamma, \widetilde{\gamma})\right| d s
\end{aligned}
$$

Since $\operatorname{ord}_{p}\left(\bar{\alpha}_{n}\right) \geqslant 3$ then $\left|\bar{\alpha}_{n}(\gamma(s))\right| \leqslant$ Cst $\cdot \tau^{3}$. Moreover, since $\rho_{n, j}^{+} \geqslant 3$ is the order of $\alpha_{n, j}^{+}$, for $j=2, \ldots, n-1$, then $\left|\alpha_{n, j}^{+}(\gamma(s))\right| \leqslant$ Cst $\cdot \tau^{3}$. Furthermore, by inequality (3) $\left|\gamma_{k}(s)-\widetilde{\gamma}_{k}(s)\right| \leqslant$ Cst $\cdot \tau^{2} \cdot t$, then

$$
\left|\gamma_{k}(t)-\widetilde{\gamma}_{k}(t)\right|\left|P_{n, k}(\gamma, \widetilde{\gamma})\right| \leqslant \operatorname{Cst} \cdot \tau^{2} \cdot t \cdot \tau=\mathrm{Cst} \cdot \tau^{3} \cdot t
$$

Also, we know that $\hat{\alpha}_{n, j}$ is linear on $x_{1}, \ldots, x_{n-1}$, hence by inequality (3) $\left|\widehat{\alpha}_{n, j}(\gamma(s))-\widehat{\alpha}_{n, j}(\widetilde{\gamma}(s))\right| \leqslant \operatorname{Cst} \cdot \tau \cdot t$. Then

$$
\left|\gamma_{n}(t)-\widetilde{\gamma}_{n}(t)\right| \leqslant \int_{0}^{t}|a|\left|\gamma_{n}(s)-\widetilde{\gamma}_{n}(s)\right| d s+\mathrm{Cst} \cdot \tau^{3} \cdot t .
$$

By Gronwall lemma we get

$$
\begin{aligned}
& \left|\gamma_{n}(t)-\widetilde{\gamma}_{n}(t)\right| \leqslant \operatorname{Cst} \cdot \tau^{3} \cdot t \cdot e^{|a| t} \\
& \left|\gamma_{n}(t)-\widetilde{\gamma}_{n}(t)\right| \leqslant \text { Cst } \cdot \tau^{3} \cdot t .
\end{aligned}
$$

Finally, we have

$$
\begin{equation*}
\|\gamma(t)-\widetilde{\gamma}(t)\|_{p}=\sum_{i=1}^{n}\left|\gamma_{i}(t)-\widetilde{\gamma}_{i}(t)\right|^{\frac{1}{w_{i}}} \leqslant \mathrm{Cst} \cdot \tau^{\frac{3}{2}} \cdot t^{\frac{1}{2}} . \tag{4}
\end{equation*}
$$

In order to state the result related to the comparison of distances, we need upper bounds for the distances $d$ and $\widetilde{d}$. So, from Theorems 7.31 and 7.26 of [3] we get

$$
\begin{aligned}
& d\left(q, q^{\prime}\right) \leqslant \mathrm{Cst} \sum_{k, j \mid w_{k} \leqslant w_{j}}\|q\|_{p}^{1-\frac{w_{k}}{w_{j}}}\left|q_{k}-q_{k}^{\prime}\right|^{\frac{1}{w_{j}}} \text {, and } \\
& \widehat{d}\left(q, q^{\prime}\right) \leqslant \mathrm{Cst} \sum_{k, j \mid w_{k} \leqslant w_{j}}\|q\|_{p}^{1-\frac{w_{k}}{w_{j}}}\left|q_{k}-q_{k}^{\prime}\right|^{\frac{1}{w_{j}}} .
\end{aligned}
$$

Since $\widetilde{d}\left(q, q^{\prime}\right) \leqslant \widehat{d}\left(q, q^{\prime}\right)$, we get

$$
\widetilde{d}\left(q, q^{\prime}\right) \leqslant \mathrm{Cst} \sum_{k, j \mid w_{k} \leqslant w_{j}}\|q\|_{p}^{1-\frac{w_{k}}{w_{j}}}\left|q_{k}-q_{k}^{\prime}\right|^{\frac{1}{w_{j}}} .
$$

Since the weights of the coordinates are $w_{1}=\cdots=w_{n-1}=1$, and $w_{n}=2$, we obtain

$$
\begin{align*}
& d\left(q, q^{\prime}\right) \leqslant \operatorname{Cst}\left(\left\|q-q^{\prime}\right\|_{p}+\|q\|_{p}^{\frac{1}{2}}\left(\sum_{k=1}^{n-1}\left|q_{k}-q_{k}^{\prime}\right|^{\frac{1}{2}}\right)\right),  \tag{5}\\
& \tilde{d}\left(q, q^{\prime}\right) \leqslant \operatorname{Cst}\left(\left\|q-q^{\prime}\right\|_{p}+\|q\|_{p}^{\frac{1}{2}}\left(\sum_{k=1}^{n-1}\left|q_{k}-q_{k}^{\prime}\right|^{\frac{1}{2}}\right)\right) . \tag{6}
\end{align*}
$$

The following notation and proposition are required for the comparison of distances result.

We denote by $\rho_{i, j}$ the order of $\alpha_{i, j}$ for $j=2, \ldots, n-1$ and $i=j, \ldots, n-1$ with the convention that $\alpha_{i, j}=\beta_{j}$ if $i=j$.

Proposition 3. Let $\gamma$ be the geodesic for $d$ such that $\gamma(0)=q, \gamma(T)=q^{\prime}$ and associated with the control function $u(\cdot)$ satisfying $\|u\|=1$. Let $\widetilde{\gamma}$ be the admissible curve for $\widetilde{d}$ defined by the same control functions as $\gamma$ and $\widetilde{\gamma}(0)=q$. If $\rho_{i, j} \geqslant 2$ and $\rho_{n, j}^{+} \geqslant 3$ for $j=2, \ldots, n-1$ and $i=j, \ldots, n-1$, then

1. $\tilde{d}(\gamma(t), \widetilde{\gamma}(t)) \leqslant C s t \cdot t^{\frac{1}{2}} \cdot \tau^{\frac{3}{2}}$.
2. $d(\gamma(t), \widetilde{\gamma}(t)) \leqslant C s t \cdot t^{\frac{1}{2}} \cdot \tau^{\frac{3}{2}}$,
where $\tau=\max \left(\|q\|_{p}, t\right)$.
Proof. 1. If $\rho_{i, j} \geqslant 2$ is the order of $\alpha_{i, j}$ for $j=2, \ldots, n-1$ and $i=j, \ldots, n-1$ then

$$
\begin{equation*}
\left|\gamma_{k}(t)-\widetilde{\gamma}_{k}(t)\right| \leqslant \operatorname{Cst} \cdot \tau^{2} \cdot t \quad \text { for } k=2, \ldots, n-1 \tag{7}
\end{equation*}
$$

Finally, by inequalities (7) and (6), and Proposition 2

$$
\widetilde{d}\left(q, q^{\prime}\right) \leqslant \operatorname{Cst}\left(t^{\frac{1}{2}} \cdot \tau^{\frac{3}{2}}+\tau^{\frac{1}{2}}\left(\tau^{2} \cdot t\right)^{\frac{1}{2}}\right) \leqslant \operatorname{Cst} \cdot t^{\frac{1}{2}} \cdot \tau^{\frac{3}{2}}
$$

2. The proof of this item is similar to the previous one, only exchanging the roles of $d$ and $\widetilde{d}$, and considering inequality (5).

Theorem 2 (Comparison of distances). If $\rho_{i, j} \geqslant 2$ and $\rho_{n, j}^{+} \geqslant 3$ for $j=$ $2, \ldots, n-1$ and $i=j, \ldots, n-1$, then there exist constants $C$ and $\epsilon>0$, such that, for all $q, q^{\prime} \in B(p, \epsilon)$, we have

$$
\begin{equation*}
-C \tau^{\frac{3}{2}} d\left(q, q^{\prime}\right)^{\frac{1}{2}} \leqslant d\left(q, q^{\prime}\right)-\widetilde{d}\left(q, q^{\prime}\right) \leqslant C \cdot \widetilde{\tau}^{\frac{3}{2}} \widetilde{d}\left(q, q^{\prime}\right)^{\frac{1}{2}}, \tag{8}
\end{equation*}
$$

where $\tau=\max \left(\|q\|_{p}, d\left(q, q^{\prime}\right)\right), \widetilde{\tau}=\max \left(\|q\|_{p}, \widetilde{d}\left(q, q^{\prime}\right)\right)$.
Proof. Let $q$ belonging to $B(p, \epsilon)$. Let us consider the geodesic $\gamma:[0, T] \rightarrow M$ for the distance $d$ such that $\gamma(0)=q, \gamma(T)=q^{\prime}$ and associated with the control function $u(\cdot)$ satisfying $\|u(t)\|=1$ and $\widetilde{\gamma}$ the admissible curve for $\widetilde{d}$ defined by the same control functions that $\gamma$ with $\widetilde{\gamma}(0)=q$. By Proposition 3 item 1

$$
\begin{equation*}
\widetilde{d}(\gamma(T), \widetilde{\gamma}(T)) \leqslant \mathrm{Cst} \cdot T^{\frac{1}{2}} \cdot \tau^{\frac{3}{2}} \tag{9}
\end{equation*}
$$

On the other hand, note that

$$
d\left(q, q^{\prime}\right)=l(\gamma)=l(\widetilde{\gamma}) \geqslant \tilde{d}(q, \widetilde{\gamma}(T))
$$

Moreover, by triangle inequality, we have

$$
\widetilde{d}(q, \widetilde{\gamma}(T)) \geqslant \widetilde{d}\left(q, q^{\prime}\right)-\widetilde{d}\left(q^{\prime}, \tilde{\gamma}(T)\right),
$$

Then, from (9), transitivity and since $\gamma(T)=q^{\prime}$, we get

$$
\begin{align*}
d\left(q, q^{\prime}\right) & \geqslant \widetilde{d}\left(q, q^{\prime}\right)-\mathrm{Cst} \cdot T^{\frac{1}{2}} \cdot \tau^{\frac{3}{2}} \\
d\left(q, q^{\prime}\right)-\widetilde{d}\left(q, q^{\prime}\right) & \geqslant-\mathrm{Cst} \cdot T^{\frac{1}{2}} \cdot \tau^{\frac{3}{2}} \tag{10}
\end{align*}
$$

Now, we change the roles of $d$ and $\widetilde{d}$ and by Proposition 3 item 2, we obtain

$$
\begin{equation*}
d\left(q, q^{\prime}\right)-\widetilde{d}\left(q, q^{\prime}\right) \leqslant \mathrm{Cst} \cdot \widetilde{T}^{\frac{1}{2}} \cdot \widetilde{\tau}^{\frac{3}{2}}, \tag{11}
\end{equation*}
$$

where $\widetilde{T}=\widetilde{d}\left(q, q^{\prime}\right)$.
Therefore from (10) and (11)

$$
-C \tau^{\frac{3}{2}} d\left(q, q^{\prime}\right)^{\frac{1}{2}} \leqslant d\left(q, q^{\prime}\right)-\widetilde{d}\left(q, q^{\prime}\right) \leqslant C \cdot \widetilde{\tau}^{\frac{3}{2}} \widetilde{d}\left(q, q^{\prime}\right)^{\frac{1}{2}} .
$$

The proof is complete.
Notice that, if $d\left(q, q^{\prime}\right) \geqslant d(p, q)$ we get $\left|d\left(q, q^{\prime}\right)-\widetilde{d}\left(q, q^{\prime}\right)\right| \leqslant \operatorname{Cd}\left(q, q^{\prime}\right)^{2}$. The similar inequality for $\hat{d}$ is $\left|d\left(q, q^{\prime}\right)-\widehat{d}\left(q, q^{\prime}\right)\right| \leqslant \mathrm{C} d\left(q, q^{\prime}\right)^{\frac{3}{2}}$. This show that the order of bound of $\left|d\left(q, q^{\prime}\right)-\widetilde{d}\left(q, q^{\prime}\right)\right|$ is stricty better than the one of $\left|d\left(q, q^{\prime}\right)-\hat{d}\left(q, q^{\prime}\right)\right|$.

### 4.2 Translation

In this subsection, we address the second fact needed to prove the main result of this paper, that is, to find directions where the distance of a pair of translated points is decreasing. These directions are the appropriate ones where $\widetilde{d}$ is closer to $d$ than $\hat{d}$.

It is well known that the distance defined by the nilpotent approximation is left-invariant (cf. [8]) while $\widetilde{d}$ is not. Let $p_{2}$ be a point in a neighborhood of $p=0$ and $g \in \mathbb{R}^{n}$. We are interested in conditions under which $\widetilde{d}\left(g, g \cdot p_{2}\right) \leqslant$ $\widetilde{d}\left(0, p_{2}\right)$ (this means decreasing), where the product is the Lie group one. For this, some elements are required.
Let $\mathcal{L}=\operatorname{Lie}\left(\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}, \tilde{X}_{n}\right), \mathfrak{g}$ the ideal generated in $\mathcal{L}$ by $\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}$ and $G$ the simply connected Lie group whose Lie algebra is $\mathfrak{g}$ that is, the set of left-invariant vector fields on $G$. We know that $\mathfrak{g}$ is a nilpotent Lie algebra. This Lie algebra $\mathfrak{g}$ can be split into homogeneous components

$$
\mathfrak{g}=\mathfrak{g}^{1} \oplus \mathfrak{g}^{2}
$$

where $\mathfrak{g}^{s}$ is the set of homogeneous vector fields of order $-s$.
The translation will be dealt with by considering a vector field $Y \in \mathfrak{g}^{1}$ and then $Y \in \mathfrak{g}^{2}$. We start by considering $Y \in \mathfrak{g}^{1}$. For the above, Definition 5 and Theorem 3 are necessary and they come from [11] and [7].

Definition 5. Let $\gamma:[0, T] \longrightarrow \mathbb{R}^{n}$ be a smooth curve, and $\epsilon>0$. A variation of $\gamma$ is a smooth map $F:[0, T] \times(-\epsilon, \epsilon) \longrightarrow \mathbb{R}^{n}$ such that

$$
F(t, 0)=\gamma(t)
$$

for all $t \in[0, T]$.
For each $s \in(-\epsilon, \epsilon)$, the curve $\gamma_{s}:[0, T] \rightarrow \mathbb{R}^{n}$ given by $\gamma_{s}(t)=F(s, t)$, is called a curve of the variation $F$.

A variation $F$ of $\gamma$ determines a differentiable vector field $V(t)$ along $\gamma$ by $V(t)=\frac{\partial F}{\partial s}(t, 0)$.

We denote by $l\left(\gamma_{s}\right)$ the length of the curve $\gamma_{s}$.
Theorem 3 (First variation of length, [11]). Let $\gamma:[0, T] \longrightarrow \mathbb{R}^{n}$ be any unit speed admissible curve and $F(t, s)$ a smooth variation of $\gamma$. Then

$$
\frac{d}{d s} l\left(\gamma_{s}\right)(0)=-\int_{0}^{T}\left\langle V(t), \nabla_{\dot{\gamma}} \dot{\gamma}\right\rangle d t+\langle V(T), \dot{\gamma}(T)\rangle-\langle V(0), \dot{\gamma}(0)\rangle,
$$

where $\nabla$ is the Levi-Civita connection.
From [6], we have that $\mathbb{R}^{n}$ is diffeomorphic to the quotient $G / H$ where $H$ stands for the connected subgroup of $G$ whose Lie algebra is the set of elements of $\mathfrak{g}$ that vanish at 0 . Moreover, the homogeneous space $G / H$ is the manifold of the right cosets of $H$. We denote by $\Pi$ the canonical projection of $G$ onto $G / H$.

Let $\gamma:[0, T] \longrightarrow \mathbb{R}^{n}$ be a geodesic of $\tilde{d}$ such that $\gamma(t) \notin \mathcal{Z}$ for $\left.\left.t \in\right] 0, T\right]$. Let $Y \in \mathfrak{g}^{1}$, and $F:[0, T] \times(-\epsilon, \epsilon) \longrightarrow \mathbb{R}^{n}$ a variation of $\gamma$ such that

$$
F(t, s)=\gamma_{s}(t)=\gamma(t) \Pi(\exp (s Y))=L_{\gamma(t)}(\Pi(\exp (s Y))),
$$

where $L_{\gamma(t)}$ is the left translation by $\gamma(t)$.
For each $Y \in \mathfrak{g}^{1}$, the projection of $Y$ onto $G / H$ is denoted by $\Pi_{*} Y$, the latter is an invariant vector field on $G / H$ (cf. [9]) Denoting by $V(t)$ the variation field of $F$, we get

$$
V(t)=\frac{\partial F}{\partial s}(t, 0)=\left.\frac{\partial}{\partial s}\right|_{s=0} \gamma_{s}(t)=T L_{\gamma(t)} \cdot \Pi_{*} Y=\left(\Pi_{*} Y\right)(\gamma(t)) .
$$

Then by Theorem 3,
$\frac{d}{d s} l\left(\gamma_{s}\right)(0)=\langle V(T), \dot{\gamma}(T)\rangle-\langle V(0), \dot{\gamma}(0)\rangle=\left\langle\left(\Pi_{*} Y\right)(\gamma(T)), \dot{\gamma}(T)\right\rangle-\left\langle\left(\Pi_{*} Y\right)(0), \dot{\gamma}(0)\right\rangle$.
Notice that the integral vanishes because $\gamma$ is a geodesic for $\tilde{d}$ and hence $\nabla_{\dot{\gamma}} \dot{\gamma}=0$.

Since we want to study the translation of the curve $\gamma$ then $Y(0)$ must not be zero. Indeed, if $Y(0)=0$ then $\gamma(t) \Pi(\exp (0))=\gamma(t)$. So, we must look for a vector field $Y \in \mathfrak{g}^{1}$ such that $Y(0) \neq 0$.

To obtain conditions on $Y$ such that the distance $\widetilde{d}$ be decreasing in the direction of $Y$, we will analyze $\langle V(t), \dot{\gamma}(t)\rangle$ considering $Y \in \operatorname{span}\left\{\widehat{X}_{1} \ldots, \widehat{X}_{n-1}\right\}$.

We assume that $Y \in \operatorname{span}\left\{\hat{X}_{1} \ldots, \hat{X}_{n-1}\right\}$. Hence there exist constants $\alpha_{1}, \alpha_{2}, \ldots, \alpha_{n-1}$ such that $Y(\gamma(t))=\sum_{i=1}^{n-1} \alpha_{i} \hat{X}_{i}(\gamma(t))$. Since $\gamma$ is an admisible curve for $\widetilde{d}$, this is to say,

$$
\dot{\gamma}(t)=\sum_{i=1}^{n-1} u_{i} \hat{X}_{i}(\gamma(t))+u_{n} \widetilde{X}_{n}(\gamma(t)),
$$

and $\left\{\widehat{X}_{1}, \ldots, \widehat{X}_{n-1}, \tilde{X}_{n}\right\}$ is an orthonormal frame for the metric, we have

$$
\langle Y(\gamma(t)), \dot{\gamma}(t)\rangle=\sum_{i=1}^{n-1} \alpha_{i} u_{i}(t)
$$

Then

$$
\langle Y(\gamma(T)), \dot{\gamma}(T)\rangle-\langle Y(0), \dot{\gamma}(0)\rangle=\sum_{i=1}^{n-1} \alpha_{i}\left(u_{i}(T)-u_{i}(0)\right) .
$$

So, we have obtained the following proposition.
Proposition 4. Let $\gamma:[0, T] \longrightarrow \mathbb{R}^{n}$ be a length minimizer of $\tilde{d}$ with control functions $u_{1}, \ldots, u_{n}$, and $Y \in \mathfrak{g}^{1}$ such that $Y(\gamma(t))=\sum_{i=1}^{n-1} \alpha_{i} \widehat{X}_{i}(\gamma(t))$.

$$
\text { If } \sum_{i=1}^{n-1} \alpha_{i}\left(u_{i}(T)-u_{i}(0)\right)<0 \text { then } \frac{d}{d s} l\left(\gamma_{s}\right)(0)<0 .
$$

To deal with the case where the translation is in direction of a vector field $Y \in \mathfrak{g}^{2}$, i.e., a vector field $Y$ such that its evaluation does not belong to the tangent space at 0 , we must change the above strategy since Lemma ?? depends on $\langle Y(0), \dot{\gamma}(0)\rangle$, so if $Y(0)$ does not belong to the tangent space at 0 , $\langle Y(0), \dot{\gamma}(0)\rangle$ does not make sense. The below (Proposition 5) is necessary and comes from [4].

Let $G$ be a connected, simply connected Lie group of dimension $n$ such that $G$ is a Carnot (or stratified) group of step $r$ (see more details in [4]) and $\mathfrak{g}$ its Lie algebra. After the choice of a basis $X_{1}, \ldots, X_{n}$ for $\mathfrak{g}$, the group $G$ is identified with $\mathbb{R}^{n}$ via the exponential mapping; this means that a point $x=\left(x_{1}, \ldots, x_{n}\right) \in \mathbb{R}^{n}$ is identified with the point $\exp \left(x_{1} X_{1}+\ldots+x_{n} X_{n}\right)$ of the group. Hence we have the following result obtained from [4] and [12].

Proposition 5. The group product has the form

$$
\begin{equation*}
x \cdot y=x+y+Q(x, y) \quad \forall x, y \in \mathbb{R}^{n}, \tag{12}
\end{equation*}
$$

where $Q=\left(Q_{1}, \ldots, Q_{n}\right): \mathbb{R}^{n} \times \mathbb{R}^{n} \rightarrow \mathbb{R}^{n}$ and $Q_{j}$ is a homogeneous polynomial of degree $w_{j}$. Moreover, for all $x, y \in \mathbb{R}^{n} Q_{1}(x, y)=\cdots=Q_{n_{1}}(x, y)=0$, where $n_{1}$ is such that $w_{1}=\cdots=w_{n_{1}}=1$, and

$$
Q_{i}(x, y)=\sum_{h, k} \mathcal{R}_{k, h}^{i}(x, y)\left(x_{k} y_{h}-x_{h} y_{k}\right),
$$

where the functions $\mathcal{R}_{k, h}^{i}$ are polynomials, homogenous of degree $w_{i}-w_{h}-w_{k}$ with respect to group dilations, and the sum is extended to all $h, k$ such that $w_{k}+w_{h} \leqslant w_{i}$.
Remark 4. In the context of the generic case, which is the case that interests us in this paper, $n_{1}=n-1$, and since rank $\left(\Delta_{p}+[\Delta, \Delta]_{p}\right)=n$ at all points then $\mathcal{R}_{k, h}^{n}$ is a constant, and $w_{k}=w_{h}=1$.

Let $Y \in \mathfrak{g}^{2}$. In local coordinates $Y(x)=\sum_{i=1}^{n} f_{i}(x) \frac{\partial}{\partial x_{i}}$. Since $w_{1}=\ldots=$ $w_{n-1}=1$ and $w_{n}=2$ then $f_{i}(x) \equiv 0$ for $i=1, \ldots, n-1$, and $f_{n}(x)$ is a constant different from 0 , hence $Y(x)=\nu \frac{\partial}{\partial x_{n}}$, with $\nu \neq 0$ (see Figure 2).


Figure 2: Singular set and its tangent space at $p$ with the translation vector $Y(x)$.

Let $\mu_{Y}$ be the integral curve of $Y$ passing through the identity of $G$ when $t=0$, then

$$
\dot{\mu}_{Y}(t)=Y\left(\mu_{Y}(t)\right)=\nu \frac{\partial}{\partial x_{n}},
$$

hence

$$
\mu_{Y}(t)=\left(\begin{array}{c}
0 \\
\vdots \\
0 \\
t \nu
\end{array}\right)
$$

Since $\exp (Y)=\mu_{Y}(1)$ (see more details in [4]), then $\exp (Y)=\left(\begin{array}{c}0 \\ \vdots \\ 0 \\ \nu\end{array}\right)$.
Let $\gamma:[0, T] \longrightarrow \mathbb{R}^{n}$ be a geodesic of $\tilde{d}$ such that $\gamma(t)=\left(\gamma_{1}(t), \ldots, \gamma_{n}(t)\right)$, $\gamma(t) \notin \mathcal{Z}$ for $t \in] 0, T]$, and $\gamma(0)=0$, with control functions $u_{1}, \ldots, u_{n}$. We consider $Y \in \mathfrak{g}^{2}$ such that, $Y(x)=\nu \frac{\partial}{\partial x_{n}}$. Let $\gamma_{L Y}(t)=L_{\exp (Y)}(\gamma(t))=$ $\left(\gamma_{1 Y}(t), \ldots, \gamma_{n Y}(t)\right)$ and $\bar{u}_{1}, \ldots, \bar{u}_{n}$ its control functions.

Recall that $\widetilde{a}_{n}(x)=a x_{n}+Q\left(x_{1}, x_{2}, \ldots, x_{n-1}\right)$, where $Q\left(x_{1}, x_{2}, \ldots, x_{n-1}\right)$ is quadratic. We set $\widetilde{a}_{n}(\gamma)=\widetilde{a}_{n}(\gamma(t))$.

The following result provides conditions on $Y$ such that $\gamma_{L Y}$ has a length less than $\gamma$.

Theorem 4. Let $\gamma:[0, T] \longrightarrow \mathbb{R}^{n}$ be a length minimizer of $\tilde{d}$ with control functions $u_{1}(t), \ldots, u_{n}(t)$ with $u_{n}(t) \neq 0$ almost everywhere, and $\widetilde{a}_{n}(\gamma)+a \nu \neq$ 0. If $\left|\widetilde{a}_{n}(\gamma)\right|<\left|\widetilde{a}_{n}(\gamma)+a \nu\right|$ then $\widetilde{d}\left(\gamma_{L Y}(0), \gamma_{L Y}(T)\right)<\widetilde{d}(\gamma(0), \gamma(T))$.

Proof. By Proposition 5, the curve $\gamma_{L Y}(t)$ is simply $\gamma_{L Y}(t)=\left(\gamma_{1}(t), \ldots, \gamma_{n}(t)+\nu\right)$. Then $\dot{\gamma}_{L Y}(t)=\dot{\gamma}(t)$ i.e. $\dot{\gamma}_{i Y}(t)=\dot{\gamma}_{i}(t)$ for $i=1, \ldots, n$. The latter has two implications: (1) Since $\dot{\gamma}_{i Y}(t)=\dot{\gamma}_{i}(t)=u_{i}(t)$ then $\bar{u}_{i}(t)=u_{i}(t)$ for $i=1, \ldots, n-1$. (2) Since $\dot{\gamma}_{n Y}(t)=\dot{\gamma}_{n}(t)$ then

$$
\begin{aligned}
\bar{u}_{n}(t)\left(\widetilde{a}_{n}(\gamma)+a \nu\right) & =u_{n}(t) \widetilde{a}_{n}(\gamma) \\
\bar{u}_{n}(t) & =u_{n}(t) \frac{\widetilde{a}_{n}(\gamma)}{\widetilde{a}_{n}(\gamma)+a \nu}
\end{aligned}
$$

The condition $\left|\widetilde{a}_{n}(\gamma)\right|<\left|\widetilde{a}_{n}(\gamma)+a \nu\right|$ implies that $\left|\bar{u}_{n}(t)\right|<\left|u_{n}(t)\right|$, hence $\bar{u}_{n}(t)^{2}<u_{n}(t)^{2}$. In consequence the length of $\gamma_{L Y}$ is less than the length of $\gamma$. Therefore $\tilde{d}\left(\gamma_{L Y}(0), \gamma_{L Y}(T)\right)<\widetilde{d}(\gamma(0), \gamma(T))$.

### 4.3 Solvable distance is better than the nilpotent one

It is known that the almost-Riemannian distance $d$ of the original system, close to $p=0$, behaves at the first-order as the distance defined by the nilpotent approximation at $p=0$. However, thanks to Theorem 2 we know that the solvable approximation improves the order of approximation of $d$ given by the nilpotent approximation. Despite the above, we can not state that the solvable distance is closer than the nilpotent one to the original distance for all pairs of points.

In this section, we prove that the approximation by $\tilde{d}$ is better than the one by $\widehat{d}$ for a pair of points translated in a direction where the distance $\widetilde{d}$ is decreasing.

Before that, notice that by Proposition 1 we know that $\widetilde{d}\left(q, q^{\prime}\right) \leqslant \widehat{d}\left(q, q^{\prime}\right)$ then we can conclude that for $\tilde{d}$ to be better than $\hat{d}$ it must be satisfied that $\frac{\widetilde{d}\left(q, q^{\prime}\right)+\widehat{d}\left(q, q^{\prime}\right)}{2}>d\left(q, q^{\prime}\right)$.

Let $\gamma(t)=\left(\gamma_{1}(t), \ldots, \gamma_{n}(t)\right)$ be a (normal) geodesic of $\widetilde{d}$ such that $\gamma(0)=0$ with control functions $u_{1}, u_{2}, \ldots, u_{n}$ parametrized by arc length on $[0, T]$, and the length of the curve $\gamma$ is denoted by $l(\gamma)$. We consider $g \in \mathbb{R}^{n}$ such that $g=\exp (Y)$ with $Y \in \mathfrak{g}$ and $Y$ satisfiying Theorem 4. Let $\gamma_{g}(t)=L_{g}(\gamma(t))$ and $\bar{u}_{1}, \bar{u}_{2}, \ldots, \bar{u}_{n}$ its control functions. Note that $\gamma_{g}$ is admissible for $\tilde{d}$ as long as it does not meet $\mathcal{Z}$. Indeed, all absolutely continuous curves are admissible out of the singular locus since the metric is Riemannian.

Let $\epsilon>0$ such that $C_{2} \cdot \widetilde{d}(0, \gamma(T))=C_{2} \cdot T \leqslant \epsilon$, where $C_{2}$ is the constant of inequality 8 of Theorem 2 .
On the other hand, by Proposition 5 we have that $u_{i}(t)=\bar{u}_{i}$ for $i=1, \ldots, n-1$ an since $l\left(\gamma_{g}\right)<l(\gamma)$, we can assume that there exists $C:[0, T] \rightarrow[0,1[$ such that

$$
\begin{equation*}
\left|u_{n}(t)\right| C(t)=\left|\bar{u}_{n}(t)\right| . \tag{13}
\end{equation*}
$$

Moreover, from Pontryagin's maximum principle (more details see [1]) we know that

$$
\begin{equation*}
u_{n}(t)=\left\langle\lambda(t), \tilde{X}_{i}\right\rangle=\lambda_{n}(t) \cdot\left(a x_{n}+Q\left(x_{1}, x_{2}, \ldots, x_{n-1}\right)\right) \tag{14}
\end{equation*}
$$

where $\lambda(t) \in T_{\gamma(t)}^{*} \mathbb{R}^{n}$.
Let $b=\max (C(t), t \in[0, T])$ and $S=\frac{\left(1-b^{2}\right) \lambda_{n}(0)^{2}(n-2)^{2} a^{2}}{40(n-1)^{2}} T^{4}$.
The following Theorem is the main result in this section.
Theorem 5. With the previous notations. If $\frac{2 \epsilon}{1+2 \epsilon}<S$ then

$$
\left|\widehat{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right)-d\left(\gamma_{g}(0), \gamma_{g}(T)\right)\right|>\left|d\left(\gamma_{g}(0), \gamma_{g}(T)\right)-\tilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right)\right|,
$$

and in consequence $\tilde{d}$ is closer than $\hat{d}$ to $d$.
Proof. The different distances between $\gamma_{g}(0)$ and $\gamma_{g}(T)$ are analyzed in several steps.

1. Since the controls associated to $\gamma_{g}$ are $u_{1}, \ldots, u_{n-1}, \bar{u}_{n}$, then by equation (13) the length of $\gamma_{g}$ related to $\tilde{d}$ is

$$
\widetilde{l}\left(\gamma_{g}\right)=\int_{0}^{T}\left(\sum_{i=1}^{n-1} u_{i}^{2}+u_{n}^{2} C(t)^{2}\right)^{\frac{1}{2}} d t
$$

Let

$$
\begin{equation*}
I_{b}=\int_{0}^{T}\left(\sum_{i=1}^{n-1} u_{i}^{2}+u_{n}^{2} b^{2}\right)^{\frac{1}{2}} d t \tag{15}
\end{equation*}
$$

then $\widetilde{l}\left(\gamma_{g}\right) \leqslant I_{b}$. Since $\sum_{i=1}^{n} u_{i}^{2}=1$ then $\sum_{i=1}^{n-1} u_{i}^{2}=1-u_{n}^{2}$. Replacing in equation (15)

$$
\begin{aligned}
I_{b} & =\int_{0}^{T}\left(1-u_{n}^{2}+u_{n}^{2} b^{2}\right)^{\frac{1}{2}} d t \\
& =\int_{0}^{T}\left(1+\left(b^{2}-1\right) u_{n}^{2}\right)^{\frac{1}{2}} d t
\end{aligned}
$$

Since $b \in\left[0,1\left[\right.\right.$ then $b^{2}-1<0$. We set $-\beta=b^{2}-1$. Hence

$$
\begin{equation*}
I_{b}=\int_{0}^{T}\left(1-\beta u_{n}^{2}\right)^{\frac{1}{2}} d t \tag{16}
\end{equation*}
$$

2. Approximation of $u_{n}$.

From the Hamiltonian equations (section 3.2.1), we make the choice $\lambda_{1}(0)=\lambda_{2}(0)=\ldots=\lambda_{n-1}(0)=\lambda_{0}$ and the following approximations hold:

$$
\begin{align*}
& x_{i} \approx \lambda_{0} \cdot t, \quad \text { for } i=1, \ldots, n-1, \\
& x_{n} \approx \frac{t^{2}(n-2)}{2(n-1)} . \tag{17}
\end{align*}
$$

Replacing approximations (17) in equation (14) we obtain

$$
\begin{equation*}
u_{n} \approx \frac{\lambda_{n}(0)(n-2)}{2(n-1)} a t^{2} . \tag{18}
\end{equation*}
$$

3. Considering the above approximation of $u_{n}$, and replacing in equation (16) we get:

$$
\begin{aligned}
I_{b}=\int_{0}^{T}\left(1-\beta u_{n}^{2}\right)^{\frac{1}{2}} d t & \leqslant \int_{0}^{T}\left(1-\beta \cdot \frac{\lambda_{n}(0)^{2}(n-2)^{2}}{4(n-1)^{2}} a^{2} t^{4}\right)^{\frac{1}{2}} d t \\
& \leqslant \int_{0}^{T}\left(1-\beta \cdot \frac{\lambda_{n}(0)^{2}(n-2)^{2}}{8(n-1)^{2}} a^{2} t^{4}\right) d t
\end{aligned}
$$

the latter inequality is thanks to $(1-c)^{\frac{1}{2}} \leqslant 1-0.5 c$ whenever $0 \leqslant c \leqslant 1$. Finally,
$\widetilde{l}\left(\gamma_{g}\right) \leqslant I_{b} \leqslant \int_{0}^{T}\left(1-\beta \cdot \frac{\lambda_{n}(0)^{2}(n-2)^{2}}{8(n-1)^{2}} a^{2} t^{4}\right) d t=T-\frac{\beta \lambda_{n}(0)^{2}(n-2)^{2} a^{2}}{40(n-1)^{2}} T^{5}=T-S T$.

We assume $S<1$. This is possible, even if $T$ is small, because by Remark $3 H(\lambda(0))$ does not depend of $\lambda_{n}(0)$. In consequence, $\lambda_{n}(0)$ can be chosen arbitrarily large. Hence

$$
\begin{equation*}
\widetilde{l}\left(\gamma_{g}\right) \leqslant I_{b} \leqslant(1-S) T \tag{19}
\end{equation*}
$$

4. We assume $\|g\|_{p} \leqslant \widetilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right)$. We apply now inequality 8 of Theorem 2 and since $C_{2} \cdot \widetilde{d}(0, \gamma(T))=C_{2} \cdot T \leqslant \epsilon$ we get

$$
\begin{aligned}
d\left(\gamma_{g}(0), \gamma_{g}(T)\right)-\widetilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right) & \leqslant C_{2} \cdot \tilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right)^{2} \\
d\left(\gamma_{g}(0), \gamma_{g}(T)\right)-\widetilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right) & \leqslant C_{2} \cdot \tilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right) \widetilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right)
\end{aligned}
$$

since $\tilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right) \leqslant \tilde{d}(\gamma(0), \gamma(T))$ then

$$
\begin{align*}
d\left(\gamma_{g}(0), \gamma_{g}(T)\right)-\tilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right) & \leqslant C_{2} \cdot \tilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right) \tilde{d}(\gamma(0), \gamma(T)) \\
d\left(\gamma_{g}(0), \gamma_{g}(T)\right)-\widetilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right) & \leqslant \epsilon \cdot \widetilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right)  \tag{20}\\
d\left(\gamma_{g}(0), \gamma_{g}(T)\right) & \leqslant \widetilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right)(1+\epsilon) \tag{21}
\end{align*}
$$

From inequality (21)

$$
d\left(\gamma_{g}(0), \gamma_{g}(T)\right) \leqslant(1+\epsilon) \cdot T \cdot \frac{I_{b}}{T} \leqslant(1+\epsilon) \cdot \hat{d}(\gamma(0), \gamma(T)) \frac{I_{b}}{T}
$$

then

$$
\begin{align*}
\hat{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right)-d\left(\gamma_{g}(0), \gamma_{g}(T)\right) & \geqslant-(1+\epsilon) \cdot \hat{d}(\gamma(0), \gamma(T)) \frac{I_{b}}{T}+\widehat{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right) \\
& \geqslant \widetilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right) \cdot \frac{\widehat{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right)}{\widetilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right)}\left(1-(1+\epsilon) \frac{I_{b}}{T}\right) \\
& \geqslant \widetilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right) \cdot \frac{T}{I_{b}}\left(1-(1+\epsilon) \frac{I_{b}}{T}\right) \\
\hat{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right)-d\left(\gamma_{g}(0), \gamma_{g}(T)\right) & \geqslant \widetilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right) \cdot\left(\frac{T}{I_{b}}-(1+\epsilon)\right) . \tag{22}
\end{align*}
$$

Therefore from inequalities (20) and (22)

$$
\widehat{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right)-d\left(\gamma_{g}(0), \gamma_{g}(T)\right)>d\left(\gamma_{g}(0), \gamma_{g}(T)\right)-\widetilde{d}\left(\gamma_{g}(0), \gamma_{g}(T)\right)
$$

as soon as

$$
\begin{aligned}
\frac{T}{I_{b}}-(1+\epsilon) & >\epsilon, \\
\frac{T}{I_{b}} & >1+2 \epsilon \\
I_{b} & <\frac{1}{1+2 \epsilon} T .
\end{aligned}
$$

According to inequality (19): $I_{b} \leqslant(1-S) T$, and since $\frac{2 \epsilon}{1+2 \epsilon}<S$, then $1-S<\frac{1}{1+2 \epsilon}$ and in consequence $\widetilde{d}$ is closer to $d$ than $\hat{d}$.

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[^1]:    ${ }^{1}$ i.e. $(X, Y) \longmapsto\langle X, Y\rangle$ is $\mathcal{C}^{\infty}(M)$-bilinear; $\langle X, X\rangle_{q} \geq 0$ for all $q \in M$ and $\langle X, X\rangle_{q}=$ $0 \Longrightarrow X_{q}=0$

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