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# Study of thin layers of tin oxide for applications in photovoltaic solar cells

Yahya Zakaria

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**Yahya Zakaria**

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**Study of Thin Layers of Tin Oxide for  
Applications in Photovoltaic  
Solar Cells**

**THÈSE dirigée par :**

**M. SLAOUI Abdelilah**  
**M. AHZI Said**

Directeur de Recherche, ICube, CNRS - Université de Strasbourg  
Professeur, ICube, CNRS - Université de Strasbourg

**RAPPORTEURS :**

**Mme. BOUSQUET Angélique**

Maître de Conférences HDR, Institut de Chimie de Clermont-  
Ferrand, Université Clermont Auvergne

**M. PIERSON Jean-François**

Professeur, Institut Jean Lamour, CNRS - Université de Lorraine

**EXAMINATEURS :**

**M. COLIS Silviu**  
**M. JOUIAD Mustapha**  
**M. RUCH David**

Professeur, IPCMS, CNRS - Université de Strasbourg  
Professeur, LPMC, Université de Picardie Jules Verne  
Research Team Leader, Luxembourg Institute of Science and  
Technology (LIST), Luxembourg

**MEMBRES INVITES :**

**M. AISSA Brahim**  
**M. FIX Thomas**

Senior Scientist, Qatar Environment and Energy Research Institute,  
HBKU, Qatar  
Chargé de Recherche HDR, ICube, CNRS - Université de Strasbourg

## Abstract

In the last decades, solar PV has reached the technological maturity as well as high performance levels. However, the current challenge is to make solar PV more cost-effective to be more competitive with conventional sources of energy. One rational way to reduce the cost of solar PV, is through the use of cost-effective materials to reduce the manufacturing cost. Indium tin oxide (ITO) is one of the most used transparent conducting oxides in solar PV, is not a sustainable material as indium, which is the main metal, has a low earth-abundance and a very high cost. On the other hand, tin, which is the metallic element of tin oxide, is much more cost-effective and has a higher earth-abundance compared to indium.

The aim of this thesis is to conduct a detailed study of undoped tin oxide and nitrogen doped tin oxide as a transparent conducting oxide (TCO) using mainly reactive magnetron sputtering with different growth conditions and exploring the material properties with different characterisation techniques. This thesis will also assess the performance of tin oxide as electron transport layer (ETL) for perovskite solar cells.

Chapter 1 will address the state-of-the-art of undoped and doped  $\text{SnO}_x$  films in both chemical states  $\text{SnO}$  and  $\text{SnO}_2$ . Chapter 2 will present the work on undoped  $\text{SnO}_x$  synthesis and post-deposition annealing as well as material characterization related to structural, chemical, optical, and electrical properties. We will show that high conductivity is related to poor oxygen condition during the growth of  $\text{SnO}_x$  films and high optical transmittance is related to high oxidation of  $\text{SnO}_x$  films. Chapter 3 will explore the possibility to dope  $\text{SnO}_x$  with nitrogen using three nitrogen precursor gases namely,  $\text{NH}_3$ ,  $\text{N}_2\text{O}$  and  $\text{N}_2$ . We will show that  $\text{NH}_3$  treatments has etched and damaged the  $\text{SnO}_x$ ,  $\text{N}_2\text{O}$  treatments have enhanced the n-type conductivity and  $\text{N}_2$  treatments have shown the ability to convert the conductivity from n-type to p-type. Chapter 4 will present the implementation of  $\text{SnO}_x$  as electron transport layer in perovskite solar cells (PSCs) where we will show that the post-deposition annealing has improved the electrical properties of  $\text{SnO}_x$  which leads to improve the performance of PSC.

## Résumé

Au cours des dernières décennies, le solaire photovoltaïque a atteint une maturité technologique ainsi que des niveaux de performance élevés. Cependant, le défi actuel consiste à rendre le PV solaire plus rentable pour être plus compétitif par rapport aux sources d'énergie conventionnelles. Une manière rationnelle de réduire le coût du solaire photovoltaïque consiste à utiliser des matériaux rentables pour réduire le coût de fabrication. L'oxyde d'indium et d'étain est l'un des oxydes conducteurs transparents les plus utilisés dans le PV solaire, il n'est pas un matériau durable car l'indium, qui est le métal principal, a une faible abondance et un coût très élevé. D'autre part, l'étain, qui est l'élément métallique de l'oxyde d'étain, a une abondance de terre plus élevée que l'indium et un coût relativement faible.

L'objectif de cette thèse est de mener une étude détaillée de l'oxyde d'étain non dopé et de l'oxyde d'étain dopé à l'azote en tant qu'oxyde conducteur transparent (OCT) en utilisant principalement la pulvérisation magnétron réactive avec différentes conditions de croissance et en explorant les propriétés du matériau par différentes techniques de caractérisation. Cette thèse évaluera également les performances de l'oxyde d'étain comme couche de transport d'électrons (CTE) pour les cellules solaires à pérovskite.

Le chapitre 1 traitera de l'état de l'art des films de  $\text{SnO}_x$  non dopés et dopés dans les deux états chimiques  $\text{SnO}_2$  et  $\text{SnO}$ . Le chapitre 2 présentera les travaux sur la synthèse de  $\text{SnO}_x$  non dopé et le recuit post-dépôt ainsi que la caractérisation des matériaux liée aux propriétés structurales, chimiques, optiques, et électriques. Nous montrons qu'une conductivité élevée est liée à une condition d'oxygène appauvrie pendant la croissance des films de  $\text{SnO}_x$  et qu'une transmission optique élevée est liée à une oxydation élevée des films de  $\text{SnO}_x$ . Le chapitre 3 explorera la possibilité de doper  $\text{SnO}_x$  avec de l'azote en utilisant trois gaz précurseurs de l'azote, à savoir  $\text{NH}_3$ ,  $\text{N}_2\text{O}$  et  $\text{N}_2$ . Nous montrons que les traitements  $\text{NH}_3$  ont gravé et endommagé les films  $\text{SnO}_x$ , les traitements  $\text{N}_2\text{O}$  des films  $\text{SnO}_x$  ont amélioré la conductivité de type n et les traitements  $\text{N}_2$  des films  $\text{SnO}_x$  permettent de convertir la conductivité de type n en type p. Le chapitre 4 présentera la mise en œuvre de  $\text{SnO}_x$  en tant que couche de transport d'électrons dans les cellules solaires à pérovskite où nous montrons que le recuit post-dépôt a amélioré les propriétés électriques de  $\text{SnO}_x$ , ce qui conduit à améliorer les performances des cellules solaires à pérovskite.

## DEDICATION

*To my mother, the greatest mother in the world, to whom I owe everything in my life. To my father, fiancé, sisters, and friends.*

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## Preface

This thesis represents the research work I conducted under the direction of Dr. Abdelilah Slaoui and Pr. Said Ahzi with the support of the research group of ICube - Materials for electronic and photovoltaic devices (MaCEPV) at University of Strasbourg from 2019 to 2023 as well as the research team at Hamad Bin Khalifa University - Core Labs and QEERI Energy Centre at Nanomaterials Laboratory.

Chapter 1 is a detailed literature review which includes around 300 references, will be converted into a review article.

The majority of my work on undoped SnO<sub>x</sub> films in Chapter 2 has been published as open access article “Yahya Zakaria et al. - Study of wide bandgap SnO<sub>x</sub> thin films grown by a reactive magnetron sputtering via a two-step method - Scientific Reports 12, 15294 (2022). <https://doi.org/10.1038/s41598-022-19270-w>” under Creative Commons Attribution 4.0 International License.

Chapter 3 is related to SnO<sub>x</sub> film as electron transport layer was conducted in collaboration with Dr. Brahim Aissa at QEERI-HBKU. This chapter has been published as open access article “Yahya Zakaria et al. - Moderate temperature deposition of RF magnetron sputtered SnO<sub>2</sub>-based electron transporting layer for triple cation perovskite solar cells - Scientific Reports 13, 9100 (2023). <https://doi.org/10.1038/s41598-023-35651-1>” under Creative Commons Attribution 4.0 International License.

Chapter 4 is based on nitrogen doped SnO<sub>x</sub> was conducted with the support of MaCEPV team (mentioned in the acknowledgement). This chapter is based on a finalized manuscript which will be submitted for publication soon.

# General Introduction

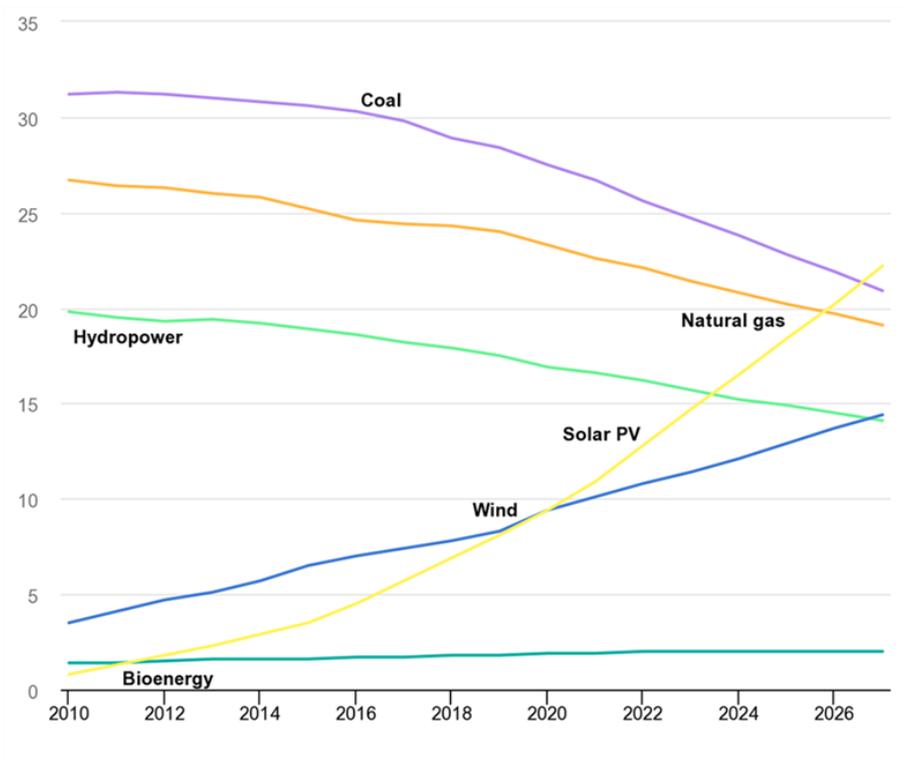
As the electric energy demand is increasing nowadays, the current fossil energy resources are not sufficient for the foreseeable future. Furthermore, despite the appealing advantages the nuclear energy, it is facing various challenges related to the materials handling, transportation, storage and disposal, its high operating cost, and the public view<sup>1</sup>. Adopting and implementing renewable energy, which is originated from sources that can be self-regenerated regularly to meet our energy requirements, is a viable approach to fulfil the human increasing energy need. The renewable energy includes wind energy, hydropower, and solar energy.

Solar energy is an attractive source of energy due to the highly available energy received from the sun. Renewables 2022 report by International Energy Agency reveals that the installed Solar Photovoltaics (PV) power capacity will exceed the power capacity of every other form of energy, particularly the power capacity of coal by 2027 as shown in Figure 1<sup>2</sup>. This increase in cumulative power capacity of solar PV is mainly due to the enhancement of cost-effectiveness that solar technologies have achieved. Like other sources of renewable energy, one of the main global objectives for research and development related to solar PV, is to reach the level of grid parity where the PV-generated power has the levelized cost of energy (LCOE) equal or less than the price of the electricity in the grid from non-renewable sources of energy such as natural gas. Most governments and energy agencies worldwide are encouraging the use of renewable energy by promoting and supporting solar energy and wind energy projects. Some countries are encouraging the use of solar energy by applying incentive plans to households to incorporate solar energy at the individual level.

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<sup>1</sup> <https://www.energy.gov/ne/articles/advantages-and-challenges-nuclear-energy>

<sup>2</sup> IEA (2022), *Renewables 2022*, IEA, Paris <https://www.iea.org/reports/renewables-2022>, License: CC BY 4.0



**Figure 1:** Projected share of cumulative power capacity by technology, 2010-2027. Reprinted (adapted) from [2]. License: CC BY 4.0.

Solar cells or photovoltaic cells, which are the essential components of solar panels, are electronic devices based on semiconducting materials that convert light into electricity directly without any intermediate mean of energy. In terms of performance, the solar cells efficiency is getting closer to the theoretical limit around 30% power conversion energy for single junction<sup>3</sup> using sophisticated processes. There is a need to reduce the cost of solar energy infrastructure, particularly the cost of solar panels. Increasing the efficiency is the rational way to increase the solar PV power capacity and reduce the LCOE which can encourage more investments in solar PV fabrication and implementation. Nevertheless, there is a huge potential in decreasing the fabrication cost by using low-cost abundant materials and lower the cost of deposition processes such as lowering the process temperature to reduce the energy consumption. This can also reduce the LCOE of solar PV which will make it more competitive with conventional non-renewable sources of energy as well as making solar PV more affordable to worldwide markets. Energy storage systems are very complementary to solar PV to provide electricity all the time as solar energy depends on the sun which is only available for several hours per day at most. There are different applications for solar PV as electricity source such as,

<sup>3</sup> Shockley, W. and H.J. Queisser, Detailed Balance Limit of Efficiency of p-n Junction Solar Cells. *Journal of Applied Physics*, 1961. 32(3): p. 510-519.

electricity related the power grid and there are also off grid applications whether in remote areas or in specific isolated sites.

The implementation of cost-effective and abundant elements-based materials to replace expensive and non-abundant elements-based materials in solar PV is a one of the key aspects to reduce the cost of fabrication. Furthermore, Transparent Conducting Oxides (TCO) materials are an essential part of solar cells and currently used in many optoelectronic devices, particularly in PV technologies namely, heterojunction silicon solar cells, thin films solar cells (a-Si, CdTe, CIGS, CZTS), perovskites solar cells, tandem cells. There is a constant research effort to reach this goal by implementing better cost-effective transparent conducting oxides (TCOs). Indium tin oxide (ITO), which is one of the most used TCOs in several applications related to optoelectronics, will be facing some supply challenges in the foreseeable future due to the high cost of indium as well as its limited abundance on earth crust compared to most of post-transition metallic elements. On the other hand, tin is more cost-effective and abundant element compared to indium. Tin, which is the base metallic element of tin oxide ( $\text{SnO}_x$ ) materials, can reach to a high conductivity and electrical mobility by optimizing the growth conditions of  $\text{SnO}_x$  thin films.

This thesis is presenting the work on TCO materials based on tin oxide where growth conditions and materials characterisation will be analysed in depth to explore the potential possibilities to improve the optoelectronic properties of these thin films. The thesis work is presented in four chapters as well as a final conclusion and perspectives.

The first chapter of this thesis will be discussing the literature review of transparent conducting oxides (TCOs) and their respective applications. In particular, the material synthesis and properties, and the applications of  $\text{In}_2\text{O}_3$  and  $\text{ZnO}$  thin films as TCOs will be covered in this chapter. The state-of-the-art of tin oxide ( $\text{SnO}_x$ ), in both chemical states  $\text{SnO}_2$  and  $\text{SnO}$ , will be discussed in terms of material synthesis and properties as well as their applications. The second chapter is a detailed study about deposition, thermal annealing, material characterisation, and optoelectronic performance for undoped  $\text{SnO}_x$ . The third chapter will present the work on  $\text{SnO}_2$  as electron transport layer (ETL) for perovskite solar cells. The material characterisation of  $\text{SnO}_2$  thin film and the performance assessment of perovskite solar cells using  $\text{SnO}_2$  as ETL will be covered in this chapter. The fourth chapter will be focusing on the nitrogen-doped  $\text{SnO}_x$  (N-doped  $\text{SnO}_x$ ) using two-step method of synthesis by magnetron sputtering and plasma assisted thermal annealing in nitrogen precursor atmosphere namely,  $\text{NH}_3$ ,  $\text{N}_2\text{O}$ , and  $\text{N}_2$ . A detailed material characterisation of N-doped  $\text{SnO}_x$  is also presented in this chapter. This thesis will be completed by a general conclusion to present the overall relevant findings and by the perspectives to highlight the potential future work beyond this PhD thesis and the next-level research related to  $\text{SnO}_x$ .

# **Chapter 1: Literature Review on Tin Oxide (SnO<sub>x</sub>) films**

## 1. Literature Review on Tin Oxide (SnO<sub>x</sub>) films

### 1.1. Transparent Conducting Oxides

#### 1.1.1. Introduction

Transparent conducting oxides (TCOs) are oxide materials which have both properties of the electrical conductivity and the high optical transmittance. These two properties can be combined by TCOs as thin films on a highly transparent substrate. While the optical transparency character of materials, which allows the transmittance of the visible and near-infrared of wavelength above 380 nm, is due to the wide bandgap above 3.3 eV. In contrast, the wide bandgap of the materials above 3.3 eV would lead to a high resistivity and make the charge carriers related doping very challenging. Therefore, the transparent conducting oxide materials can be considered very unique and remarkable materials<sup>[1, 2]</sup>. The two important properties of electrical conductivity and optical transparency along with the intrinsic chemical stability of the oxides, make the TCOs essential materials for several optoelectronic applications such as solar cells, light emitting diodes, flat panel displays, touch screens<sup>[1-29]</sup>.

The first widely known report about TCOs was back in 1907 by Karl Bädeker who reported a transparent and conducting film of cadmium oxide <sup>[30]</sup>. The progress in quantum mechanics in 1920s has enabled a better understanding of semiconductors mainly oxides and their doping in terms of electrical properties<sup>[31, 32]</sup>. Since then, the research interest on TCOs has increased till the 1950s when researchers could establish the first Indium Tin Oxide (ITO) which was first patented by J.M. Mochel (Corning Inc)<sup>[33]</sup>. The first academic report about ITO was published in 1954 by Rupperecht<sup>[34]</sup>. Other materials such as SnO<sub>2</sub> and ZnO were also explored as transparent conducting oxide materials. It is relevant to note that one of the first major advance in semiconductors was the ability to dope the amorphous silicon with substitutional pentavalent impurities such as phosphorus or trivalent impurities such as boron to form n-type or p-type conductive films, respectively, as reported by Spear and Le Comber in 1975<sup>[35]</sup>.

Since the 1990s, transparent conducting oxides related properties have had a tremendous progress due to the acquired research knowledge and expertise as well as the major technological advancements which enabled much higher quality in term of material synthesis and characterization, and more powerful computational calculation. Nevertheless, most of the progress was made for n-type conductive TCOs which is due to some of the intrinsic properties of the oxide materials that will be discussed in detail later<sup>[31, 36, 37]</sup>. It is widely established that the main breakthrough in the progress of TCOs in the last three decades, was achieving the first p-type transparent conductive oxide thin films of CuAlO<sub>2</sub> as reported by Kawazoe et al. which have a polycrystalline delafossite microstructure and were prepared by laser ablation<sup>[38]</sup>. This milestone has unlocked a new wider perceptive of TCOs as

well as enabling the possibility of new revolutionary applications for these materials. The other milestone for TCO materials was achieving the transparent amorphous oxide semiconductors which was first reported by Nomura et al. an amorphous Indium Gallium Zinc Oxide (a-IGZO) which was prepared at room temperature on a flexible polyethylene terephthalate substrate by pulsed laser ablation<sup>[3]</sup>. The transparent amorphous oxide semiconductors have more appealing electrical properties compared to hydrogenated amorphous silicon (a-Si:H) and amorphous chalcogenides namely, the Hall mobility which exceeded 10 cm<sup>2</sup>/V·s. This mobility is in the same range of the crystalline TCO materials. Moreover, it is much higher than amorphous thin film materials, namely, a-Si:H and a-chalcogenides<sup>[39, 40]</sup>. Historically, the transparent conductive oxides research was somehow focusing on materials based on In<sub>2</sub>O<sub>3</sub>, SnO<sub>2</sub> and ZnO as well as CdO<sup>[1, 2, 4, 10, 12, 13, 29, 37, 41]</sup>. Therefore, the achievement of new TCO materials using insulating earth-abundant oxides namely, 12CaO·7Al<sub>2</sub>O<sub>3</sub> and converting them to conductive oxides was also a notable breakthrough. This thin film material was achieved through thermal annealing in hydrogen atmosphere followed by ultraviolet irradiation<sup>[42]</sup>. The achievement of new TCO material Niobium doped TiO<sub>2</sub>, which has reached a very low resistivity of 2~3 x10<sup>-4</sup> Ω cm, was also an important breakthrough<sup>[43]</sup>. The advantage of TiO<sub>2</sub> as thin film and as host material for doping is the chemical stability. This achievement has also paved the way for finding alternative materials for ITO<sup>[1, 31]</sup>.

It is worth noting that the global transparent conductive films market was 4.9 billion USD in 2020 and it is expected to reach a market value of 12.9 billion USD in 2030<sup>[44]</sup>. Although Indium Tin Oxide is the widely developed TCO in the industry nowadays, there is still room for improvement in terms of optical transparency, surface electrical conductivity, material's flexibility and low cost<sup>[1-12, 23-29, 31, 45-50]</sup>.

The demand on Indium, which is the main metal in ITO, has been exponentially increasing since the emerging of the ITO as an optimal TCO for many industrial applications. This demand is facing future challenges related to the low abundance of Indium metal in the earth's crust and consequently its potential price increase, or supply chain policies. While there are some research efforts focusing on Indium recycling, this new field is facing some difficulties in terms of lack of accessible and reliable data of mining activities as well as the lack of research efforts of environmental performance related to the Indium recovery. The success of this path would require efficient recycling processes and high-level international agreements between producers/consumers countries, key industrial players and the research entities which is not currently available and not possible for the foreseeable future. Thus, the research efforts and industrial sector are targeting alternatives to ITO due to the challenge of Indium, the main metal of ITO, while keeping the sheet resistances under 5 Ω □<sup>-1</sup> particularly for large surface display and large solid state light emitting devices<sup>[31, 50, 51]</sup>.

The basic requirements for the alternative TCOs are the high electrical conductivity and the high optical transparency. Furthermore, alternative TCOs should have the advantage of low-cost production, non-toxic materials, and the ability to control the interface with the active materials. Alternative TCOs should also have the mechanical flexibility (bendability) which would systematically enable more applications<sup>[31, 50]</sup>.

### 1.1.2. General Material Properties of Transparent Conducting Oxides

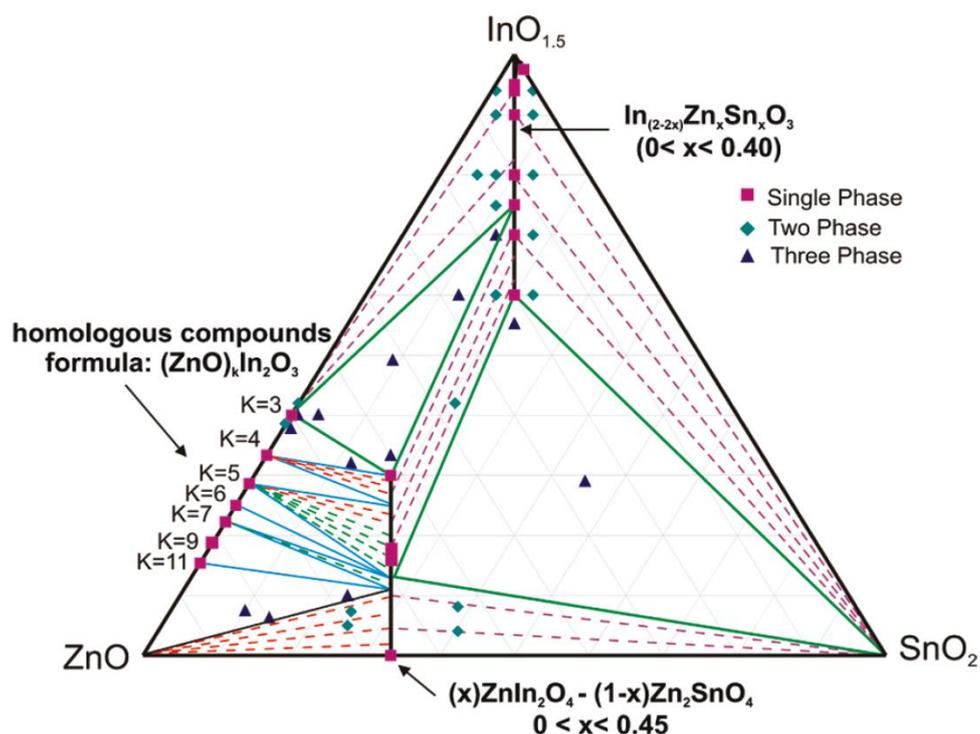
As discussed in the previous section, Transparent Conducting Oxides has gained a lot of progress in the last century, particularly in the last few decades which led to the implementation of TCOs in many applications, mainly in optoelectronics. This progress took place in many aspects, such as deep understanding of the material's properties namely, structural, electronic, optical, and electrical properties as well as improving the synthesis quality by controlling the material's defects, enhancing the crystallinity, fine-tuning the bandgap, and alloying engineering which eventually enabled the substantial improvement in TCOs design and manufacturing in optoelectronic devices.

In general, oxides, which represent an important class of solid materials, are chemical compounds formed by one oxygen atom at least and one other element at least. Furthermore, metal oxides are chemical compounds which are formed by one oxygen atom at least and one metal at least, with many metals existing naturally on earth as metal oxides. The naturally high earth abundance of metal oxides is mainly related to the stable chemical bonds between metals and to the electronegative oxygen atoms. Most of metals can be easily oxidised in presence of oxygen and form metal oxide materials. Metal oxides have largely a tendency to form ionic structure by metal oxygen metal crosslinking. They are often insoluble materials in majority of solvents. However, they can be vulnerable to acids and bases.

Technically, Transparent Conducting Oxides need to satisfy three main criteria. First, the bandgap has to be equal or above 3.1 eV and the optical transmittance has to be superior to 85% in the visible spectrum as thin film material. Secondly, the TCO need to have the ability to form degenerated doping to allow the pure oxide to transition from transparent semiconducting oxide to a metallic conductor. Finally, TCO related metal cations should have typical electronic configuration of  $(n-1)d^{10}ns^2np^a$  where n can be 4 or 5 <sup>[52, 53]</sup>.

Many metal oxides have been implemented in industry, particularly in optoelectronic applications. Certain metal oxides were investigated significantly and are commonly used in various applications. Among these metal oxides, there are  $\text{In}_2\text{O}_3$ ,  $\text{ZnO}$ ,  $\text{SnO}_2$ , and  $\text{CdO}$  where the metal cations are  $\text{In}^{3+}$ ,  $\text{Zn}^{2+}$ ,  $\text{Sn}^{4+}$ , and  $\text{Cd}^{2+}$ , respectively. Nevertheless, the interest in  $\text{CdO}$  has dropped drastically due to the toxicity of Cd which is the main metallic element.  $\text{In}_2\text{O}_3$ ,  $\text{ZnO}$ , and  $\text{SnO}_2$  as TCOs are important

conducting materials for optoelectronic applications as they are all considered wide bandgap as shown in Table 1 and these materials could reach resistivities as low as the order of  $10^{-4} \Omega \cdot \text{cm}$ . Furthermore, ternary and quaternary compounds have been achieved based on these three materials, including ITO<sup>[1, 31, 37, 41, 50, 54]</sup>. Figure 1 shows possible TCO materials based on  $\text{In}_2\text{O}_3$ ,  $\text{SnO}_2$ ,  $\text{ZnO}$ . These materials have a very similar structural, optical, and electronic properties. Their electronic configurations of the metallic elements are  $(n-1)d^{10}ns^2(np^a)$  where electrons from the  $ns$  orbitals interact strongly with the electrons  $2p$  orbital of the oxygen as shown in Table 2. These TCO materials have close-packed structures enabling four coordinates for  $\text{ZnO}$  where the Zn elements have a coordination to four oxygen atoms while forming a corner sharing tetrahedra as well as enabling six coordinates for  $\text{CdO}$ ,  $\text{In}_2\text{O}_3$  and  $\text{SnO}_2$  where the metallic atoms have a coordination to six oxygen atoms while forming a corner sharing octahedra. The valence band (VB) in these oxides is formed by O  $2p$  energy states both bonding and non-bonding, while the conduction band is formed by the antibonding Ms-Op interactions. The energy states of the electrons  $ns$  related to metals and  $2p$  related to oxygen contribute to the formation of the conduction band (CB) to form a three-dimension network for charge carrier transport once additional carriers fill the band<sup>[2, 12, 54-59]</sup>.



**Figure 1:** TCO semiconductors based on  $\text{In}_2\text{O}_3$ ,  $\text{SnO}_2$ , and  $\text{ZnO}$ . Reprinted (adapted) with permission from [53].

Copyright 2010 American Chemical Society.

Property	ZnO	In <sub>2</sub> O <sub>3</sub>	SnO <sub>2</sub>
Mineral name	Zincite	—	Cassiterite
Average metal amount in the Earth's crust, ppm	132	0.1	40
Bandgap $E_g$ (300 K), eV	3.4 (dir.)	2.7 (indir.) 3.75 (dir.)	3.6 (dir.)
Crystal structure	Hexagonal, wurtzite	Cubic, bixbyite	Tetragonal, rutile
Space group (number)	$P6_3mc$ (no. 186)	$Ia\bar{3}$ (no. 206)	$P4_2mmm$ (no. 136)
Lattice parameter (s), nm	$a$ : 0.325 $c$ : 0.5207	$a$ : 1.012	$a$ : 0.474 $c$ : 0.319
Density $\rho$ , g cm <sup>-3</sup>	5.67	7.12	6.99
Thermal conductivity $\lambda$ (300 K), W m <sup>-1</sup> K <sup>-1</sup>	$\parallel c$ : 69 [36] $\perp c$ : 60	14 [37]	$\parallel c$ : 98 $\perp c$ : 55
Thermal expansion $\alpha$ (300 K), 10 <sup>-6</sup> K <sup>-1</sup>	$\parallel c$ : 2.92 $\perp c$ : 4.75	6.7	$\parallel c$ : 3.7 $\perp c$ : 4.0
Melting point, °C [39]	1975	1910	1620 <sup>a)</sup> [40]
Metal melting point, °C	420	157	232
Heat of formation, eV	3.6	9.7	6.0

a) Decomposition into SnO and O<sub>2</sub> at 1500 °C.

**Table 1:** Material properties of ZnO, In<sub>2</sub>O<sub>3</sub> and SnO<sub>2</sub>. Reprinted (adapted) with permission from [12]. Copyright 2018 John Wiley & Sons.

Oxide	Metal atom electron orbitals	Metal ion electron orbitals
CdO	[Kr]4d <sup>10</sup> 5s <sup>2</sup>	[Kr]4d <sup>10</sup>
In <sub>2</sub> O <sub>3</sub>	[Kr]4d <sup>10</sup> 5s <sup>2</sup> 5p <sup>1</sup>	[Kr]4d <sup>10</sup>
SnO <sub>2</sub>	[Kr]4d <sup>10</sup> 5s <sup>2</sup> 5p <sup>2</sup>	[Kr]4d <sup>10</sup>
ZnO	[Ar]3d <sup>10</sup> 4s <sup>2</sup>	[Ar]3d <sup>10</sup>

**Table 2:** metal electron orbitals. Reprinted (adapted) with permission from [12]. Copyright 2018 John Wiley & Sons.

In the next section, Transparent Conducting Oxides as thin films will be discussed mainly, the most used ones namely, In<sub>2</sub>O<sub>3</sub> and ZnO in terms of synthesis, crystal structure, electronic structure and their related properties as well as the possible doping and their applications. The theoretical study of electronic structure of semiconducting thin film materials is one of the most complex calculations in the scientific fields as it requires high performance supercomputing systems, and it involves solving computational chemistry studies which requires quantum chemistry methods. One of the most widely used method is density functional theory (DFT) which is a computational quantum mechanical modelling method. Therefore, computational modelling related studies are of great importance to understand, predict, and design/engineer the electronic structure of semiconducting thin film materials<sup>[60, 61]</sup>. Furthermore, a literature review of SnO<sub>x</sub> (both SnO<sub>2</sub> and SnO) thin films will be

presented thoroughly in a main section to cover its related synthesis techniques, crystalline structure, electronic structure as well as its potential doping to improve its optoelectronic performance as well as to convert the conductivity from p-type to n-type. In the last section of this chapters, metal oxide thin films will be discussed as an electron transport layer (ETL) for perovskite solar cells (PSC).

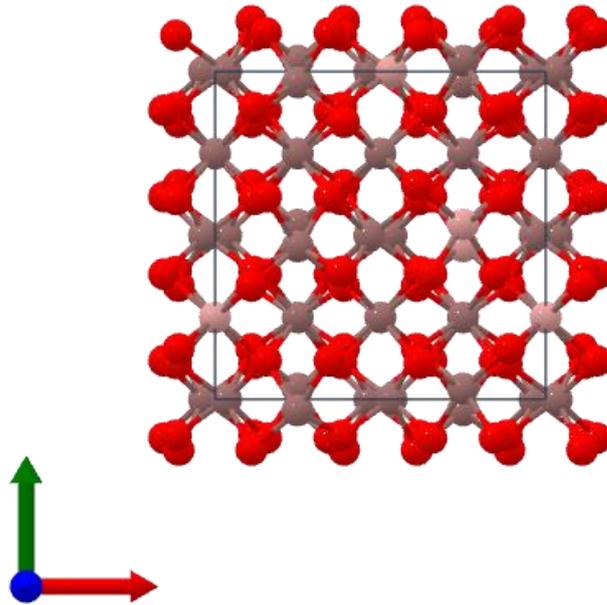
### 1.1.3. Indium Oxide

Indium oxide is a wide bandgap TCO material which was and is still heavily investigated and employed in many applications. It has very appealing optical and electrical properties as an intrinsic material as well as a doped material. It is worth noting again that the most widely used TCO material is ITO is basically an indium oxide material doped with Sn [1-12, 23-29, 31, 45-50]. There many applications for  $\text{In}_2\text{O}_3$  based materials namely, solar cells, light emitting diodes, thin film transistors [54, 62-68]. It can be prepared as single crystal or polycrystalline thin film. In addition to ITO, there are many ternary compound semiconductors based on  $\text{In}_2\text{O}_3$  namely,  $\text{Zn}_2\text{In}_2\text{O}_5$ ,  $\text{Zn}_3\text{In}_2\text{O}_6$ ,  $\text{In}_4\text{Sn}_3\text{O}_{12}$ , In-Sn-Zn-O as well as Zn doped Indium oxide as shown in Figure 1 [24, 53].

#### a) Growth and structural properties

There are many synthesis techniques to grow  $\text{In}_2\text{O}_3$  thin films which are physical deposition techniques namely, Magnetron Sputtering [54, 69], E-beam Evaporation [70], Pulsed Laser Deposition [71, 72], and Molecular Beam Epitaxy [73] as well as chemical deposition techniques namely, Atomic Layer Deposition (ALD) [74-76], Spin Coating [77], Spray Pyrolysis [78], Inkjet Printing [79, 80] and Metal Organic Chemical Vapour Deposition [81, 82].

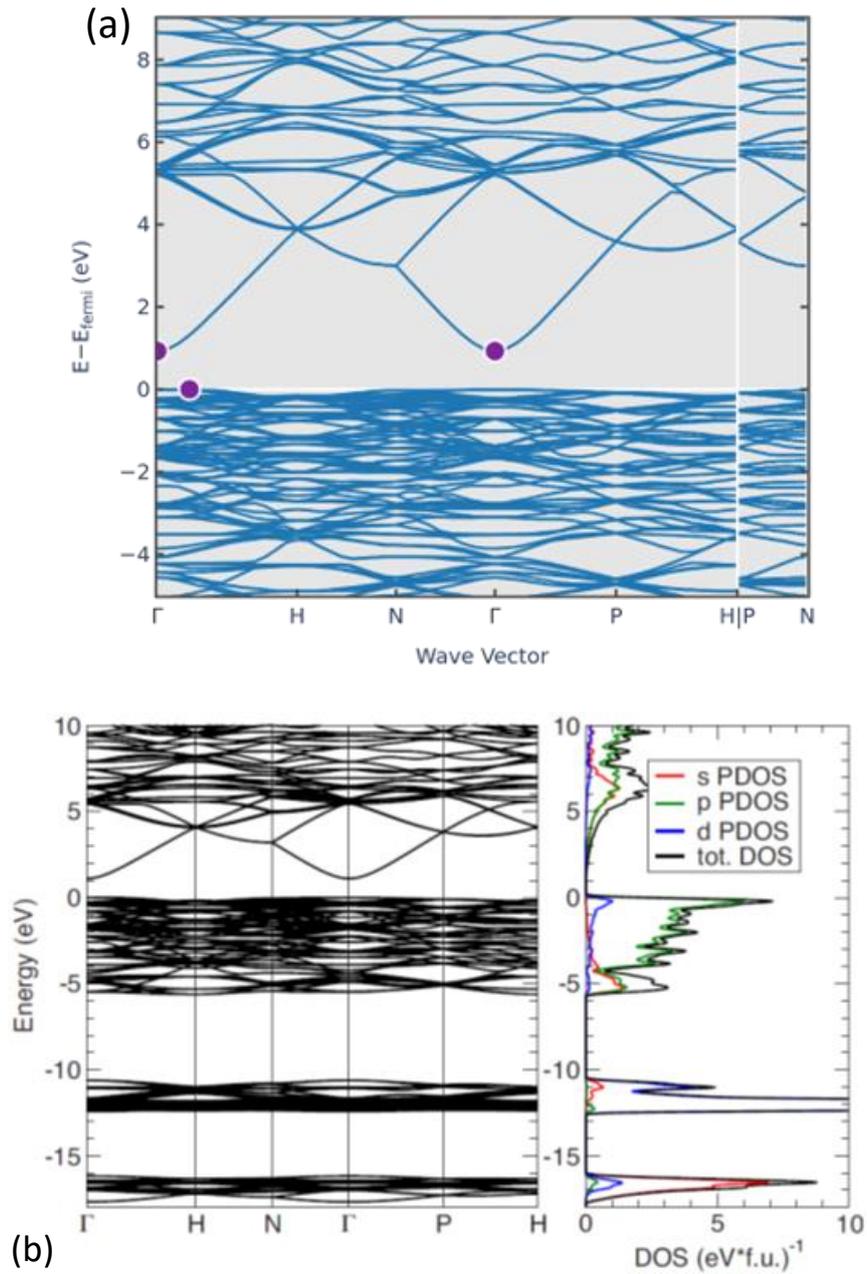
Figure 2 is showing the widely known crystal structure of  $\text{In}_2\text{O}_3$  which is Body-Centred Cubic denoted as bcc- $\text{In}_2\text{O}_3$  where the space group is  $Ia\bar{3}$  (#206) and the lattice parameter  $a = 10.118 \text{ \AA}$ . This crystal structure can be regarded as a bixbyite structure or as a C-type rare-earth sesquioxide structure. As mentioned previously, In atoms have six coordinate to oxygen atoms where they form a corner sharing octahedra [54, 57]. There is another known phase of  $\text{In}_2\text{O}_3$  has a rhombohedral structure denoted rh- $\text{In}_2\text{O}_3$  where the space group is  $R\bar{3}c$  (#167) and the lattice constants are  $a = 5.478 \text{ \AA}$  and  $c = 14.51 \text{ \AA}$  [83]. Rh- $\text{In}_2\text{O}_3$  phase is stabilized under specific conditions particularly, high pressure condition as reported by Zhang et al [84]. While bcc- $\text{In}_2\text{O}_3$  phase has been extensively investigated, there is only a limited literature about the rh- $\text{In}_2\text{O}_3$  phase which is due mostly to the difficult synthesis conditions [83, 85].



**Figure 2:** Crystal structure of Body-Centred Cubic  $\text{In}_2\text{O}_3$ . Reprinted (adapted) from The Materials Project [57]. License: CC BY 4.0.

#### b) Electronic structure

In this section, the electronic properties of  $\text{In}_2\text{O}_3$  will be focused on the most stable and reported phase bcc- $\text{In}_2\text{O}_3$ . Indium metal atom electron orbitals are  $[\text{Kr}] 4d^{10}5s^25p^1$  where the valence electrons are  $5s^2$  and  $5p^1$  as shown in Table 2, while oxygen atom electron orbitals are  $[\text{He}] 2s^22p^4$ . As shown in the band structure (Figure 3(a)), the lowest conduction band is at  $\Gamma$  point. On the other hand, the valence band maximum is on  $\Gamma$ -H line while close to the  $\Gamma$  point. Therefore,  $\text{In}_2\text{O}_3$  is an indirect bandgap semiconductor due to the fact that the minimum CB and the maximum VB are not on the same crystal momentum which is defined by the associate wave vector  $k^{[57]}$ . The band structure is quite complex which is due to the different atomic orbitals' hybridization. As per the partial Density of State in figure 3(b), O 2p has the main contribution to the upper part of the valence band. However, there is also an influence on the upper part of the VB by hybridization with In 4d states. In 5s states contribute primarily to the minimum of the CB, In 5p states also contributes while increasing the k vector.



**Figure 3:** (a) Band structure, density of states (DOS), (b) partial DOS for bcc-In<sub>2</sub>O<sub>3</sub>. Calculations were performed within DFT-LDA. (a) Reprinted (adapted) from The Materials Project [57]. License: CC BY 4.0. (b) Reprinted (adapted) with permission from [86]. Copyright 2008 American Physical Society.

### c) Doped Indium Oxide

It is widely reported that unintentional and/or self-doped indium oxide thin films as well as bulk crystals exhibit a moderate conductivity with n-type charge carrier concentration generally between  $10^{17}$  and  $10^{19}$  cm<sup>-3</sup>. The origin of this significantly high charge carrier concentration, for a wide bandgap semiconducting material, is related to many point defects and unintentional doping by impurities as reported by Bierwagen. The main point defects are related to oxygen vacancies ( $V_o$ ) which are established experimentally as a shallow donor. Hydrogen atoms, which tend to form interstitial

defects, are also another source of shallow donors for indium oxide. This has been demonstrated both theoretically as well as experimentally [87, 88].

Although the previously mentioned unintentional and self-doping for indium oxide enable forming n-type charge carriers, the best electrical performance for indium-based materials can be achieved by a controlled doping process where the charge carrier concentration can increase drastically through various dopant, including Sn, Zr, Ti, Mo, W and F, to be used as substitutes. The doping using tin element has been the most successful and it is used widely in various industrial applications as it provides the highest electrical conductivity as well as providing a very good optical transparency. The stability of the heavily Sn-doped  $\text{In}_2\text{O}_3$ , widely known as ITO, is due to the high solubility of Sn in indium oxide which is up to about 6% atomic. These doping elements provides shallow donors within the  $\text{In}_2\text{O}_3$  bandgap which are very close to the conductive band minimum. ITO can reach to a high charge carrier concentration of  $6 \times 10^{21} \text{ cm}^{-3}$  and resistivities under the  $10^{-4} \Omega\text{-cm}$ . The high doping of Sn in  $\text{In}_2\text{O}_3$  can enable the Fermi level to be in the conductive band. Even though, p-type doping is theoretically possible through elements from group - II, particularly Ca, Mg and Zn. Nevertheless, the p-type conductivity is still not experimentally established as the negative formation energy which compensate the  $V_o$  donors takes place and tend to drive the Fermi level above the valence band maximum [88-90].

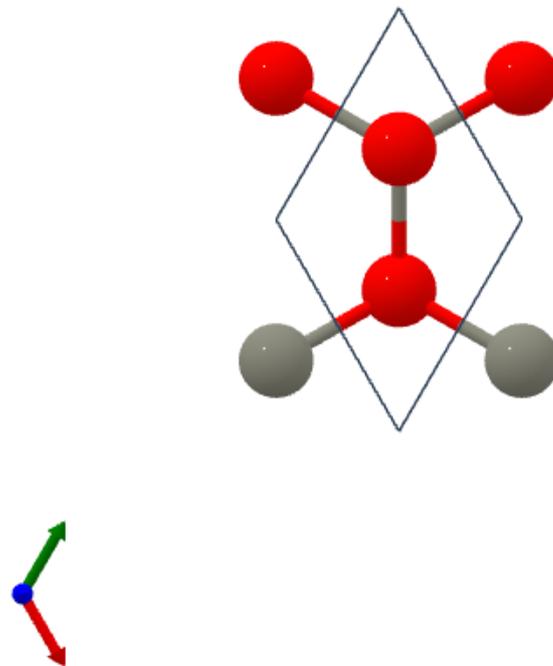
#### 1.1.4. Zinc Oxide

Zinc oxide is also one of the wide and direct bandgap TCO materials which has been heavily investigated and employed in many applications. It has some appealing material properties. It is known to have similar properties compared  $\text{GaN}$  [91]. It is worth to note that  $\text{ZnO}$  TCO material is not suitable to generate p-type charge carriers due to the presence of O 2p orbitals on the top of the valence band. There are many applications for  $\text{ZnO}$  based materials namely, solar cells including Perovskite Solar Cells (PSC),  $\text{CuInGaSe}$  (CIGS) Solar cells, Organic Solar Cells and Dye Sensitized Solar Cells (DSSC) [14, 15, 92-99], light emitting diodes [100-103], gas sensors [104], thin film transistors [105].  $\text{ZnO}$  material exhibits a great material stability. There are many ternary compound semiconductors based on  $\text{ZnO}$  namely,  $\text{Zn}_2\text{In}_2\text{O}_5$ ,  $\text{Zn}_3\text{In}_2\text{O}_6$ ,  $\text{In-Sn-Zn-O}$  [24] as well as Zn doped Indium oxide as shown in Figure 1.

##### a) Growth and structural properties

$\text{ZnO}$  thin films can be grown using different synthesis techniques namely, magnetron sputtering [106], Chemical Vapour Deposition [107, 108] which can provide an average crystalline quality. Advanced techniques such as Molecular Beam Epitaxy [109, 110], Pulsed Laser Deposition [111] and Metal Organic Chemical Vapour Deposition [112] can provide much better crystalline quality due to higher quality of the used materials/precursors as well as the advanced control of the growth conditions.

Zinc oxide has three known crystal structures namely, hexagonal wurtzite, cubic zincblende and Rocksalt. The hexagonal wurtzite phase, which is shown in Figure 4, is the most stable thermodynamically. Zinc blende thin film is normally formed while grown on a cubic structure substrate. However, the Rocksalt phase formation requires a very high-pressure condition. The wurtzite phase has four atoms per unit cell<sup>[113]</sup>, its space group is  $P6_3mc$  (#186) and the lattice parameters reported experimentally using XRD are  $a = 3.2475 \text{ \AA}$  and  $c = 5.2075 \text{ \AA}$ <sup>[114]</sup>. The theoretical lattice parameters using ab initio calculations are  $a = 3.286 \text{ \AA}$  and  $c = 5.241 \text{ \AA}$ <sup>[115]</sup>. The zincblende phase has cubic crystal structure and its space group is  $F\bar{4}3m$  (#216). The experimental lattice parameter is  $4.463 \text{ \AA}$  determined by XRD and TEM<sup>[116]</sup> and the theoretical lattice parameters were found to be slightly higher than the experimental values. The Rocksalt phase has a cubic structure, its group space is  $Fm\bar{3}m$  (#225) and the lattice parameter reported experimentally is  $a = 4.271 \text{ \AA}$ . The theoretical calculated lattice parameter is  $a = 4.30 \text{ \AA}$ <sup>[113, 117, 118]</sup>.

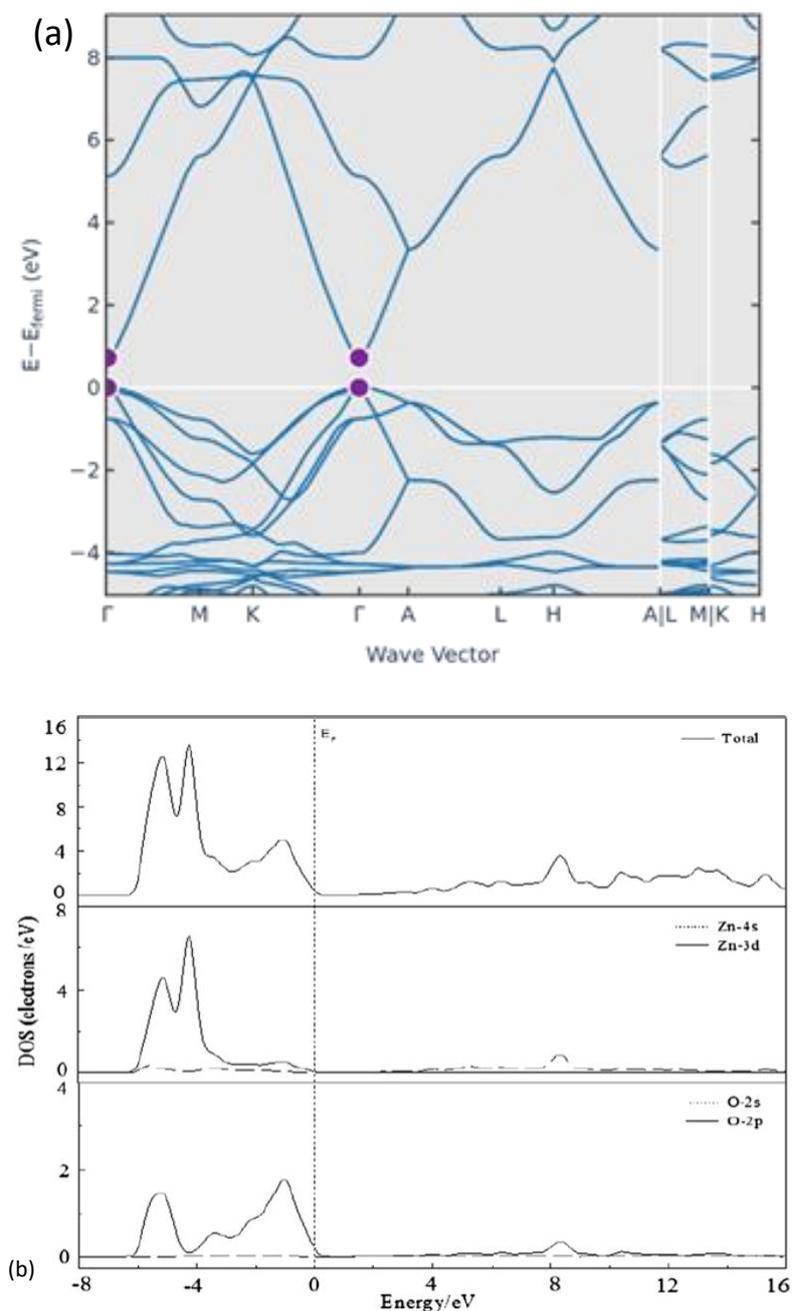


**Figure 4:** Crystal structure of Wurtzite ZnO. Reprinted (adapted) from The Materials Project [58]. License: CC BY 4.0.

#### b) Electronic structure

In this section, the electronic properties of ZnO will be focused on the most stable and reported phase which is hexagonal wurtzite. Zinc metal atom electron orbitals are  $[\text{Ar}] 3d^{10}4s^2$  and valence electron  $4s^2$  as shown in Table 2, while oxygen atom electron orbitals are  $[\text{He}] 2s^22p^4$ . As shown in the band structure (Figure 5(a)), similarly to  $\text{In}_2\text{O}_3$ , the lowest conduction band is at  $\Gamma$  point. However, the valence band maximum is on the symmetric  $\Gamma$  point. Therefore, ZnO is a direct bandgap semiconductor where both the minimum CB and the maximum VB are on the same crystal momentum which is

defined by the associate wave vector  $k^{[58]}$ . As per the partial Density of State in Figure 5(b), oxygen related 2p and 2s states as well as zinc related 3d states have the main contribution to the valence band. However, the upper part of the VB is mainly formed by O 2p states. The minimum of the conductive band is mainly formed by both O 2p and Zn 4s states.



**Figure 5:** (a) Band structure, (b) density of states (DOS), and partial DOS for ZnO. Calculations were performed within DFT-LDA. (a) Reprinted (adapted) from The Materials Project [58]. License: CC BY 4.0. (b) Reprinted (adapted) with permission from [119]. Copyright 2012 Elsevier.

### c) Doped Zinc Oxide

Zinc oxide is also widely reported to be an intrinsic n-type conductive and its charge carrier concentration, as a result of the unintentional and self doping, can reach up to  $10^{19} \text{ cm}^{-3}$ . These n-type

charge carriers are present due to the formation of shallow donors. It is widely known that oxygen vacancies ( $V_o$ ) mainly and zinc interstitial ( $Zn_i$ ) are the source of these shallow donors<sup>[88, 120]</sup>. Nevertheless, there is still a debate around this as it is reported that the donor level formed by the oxygen vacancies is around 1 eV lower than the conductive band minimum (CBM) which is a relatively deep level. Others reported that Hydrogen impurity is the main reason for the n-type conductivity. Janotti and co-authors have reported that hydrogen is likely to form shallow donors due to its interstitial defects and/or by oxygen sites substitution<sup>[121, 122]</sup>.

Even though the shallow donors formed by the presence of oxygen vacancies as well as hydrogen impurities, the n-type doping through elements, mainly group-III atoms, is highly required to assure the higher electrical performance of zinc oxide. Elements like Al, Ga and In are reported to be used as dopants to substitute Zn sites and form n-type charge carriers. Their respective donor energy levels are positioned at 51.55, 54.6, 63.2 meV from their recombination of excitons bounds which are slightly higher than the donor energy level formed by hydrogen of 46.1 meV<sup>[123]</sup>. However, the impurity band can be formed due to the heavy doping of group-III atoms which will tend to get closer to the valence band and reduce the donors' energy level. Therefore, a charge carrier concentration of  $10^{21} \text{ cm}^{-3}$  can be achieved, and the conductivity can be enhanced substantially. Zinc oxide doped with aluminium, which is widely known as AZO, was investigated extensively as a low-cost material in order to replace ITO. It has reached the acceptable performance to be used in many applications such as light emitting diodes and solar cells. Its electrical performance could reach as low as  $3 \times 10^{-4} \Omega\text{-cm}$  as resistivity while the mobility is  $20 \text{ cm}^2/\text{V}\cdot\text{s}$ . This best resistivity of AZO is slightly higher than the one achieved by ITO which is due to the lower mobility in AZO compared to the one in ITO<sup>[124]</sup>.

## **1.2. Tin Oxide thin film as Transparent Conducting Oxide**

Tin oxide based TCO materials have been widely investigated in different research fields, particularly in the optoelectronic ones in the last several decades. Generally, tin oxide ( $\text{SnO}_x$ ) thin films, as one the metal oxide materials, are very appealing materials due to their wide bandgap from 3.3 to 4.3 eV and possibility to achieve a very low and competitive resistivity in the order of the low  $10^{-4} \Omega\text{-cm}$ . Furthermore, tin oxide is earth-abundant and non-toxic material which makes it very suitable for several applications. Tin as raw material has a much lower cost of 27.5 USD/kg compared to indium (210.6 USD/kg) which makes it very appealing for industrial application. In general tin-based materials are important materials to be applied for any global large scale terawatt solar energy supply.  $\text{SnO}_x$  has been employed in various applied fields as devices such as gas sensors, photocatalyst for water treatment, large area electronics such as flat panel displays, thin film transistors, solar cells including perovskite solar cells, light emitting diodes, transparent electrodes, anode materials, coatings, catalysis<sup>[12, 53, 88, 125-141]</sup>.

Sensors related-industry demand is increasing for affordable, thermal and mechanical stable, abundant sensing materials. SnO<sub>2</sub> is one of the most used materials for sensing thanks to its appealing properties<sup>[142]</sup>. Undoped SnO<sub>x</sub> as a gas sensor was reported by Martín et al.<sup>[143]</sup> to detect NO<sub>2</sub> with a limit of detection varying from 0.01 to 0.25 ppm and to detect CO with a limit of detection from 1 to 20 ppm. While Lei Li and co-workers have reported SnO<sub>2</sub> with reduced Graphene Oxide (rGO) based sensor to detect NO<sub>2</sub> with a detection limit down to 1 ppb<sup>[144]</sup>. There are other reports which reveal the employment of SnO<sub>2</sub> based materials to detect CH<sub>4</sub> and H<sub>2</sub><sup>[142, 145, 146]</sup>. Photoelectrochemical water splitting is another application for SnO<sub>2</sub> based materials. Li et al. have reported F-doped SnO<sub>2</sub> nanowire as a light harvesting electrode<sup>[147]</sup>. Hernández and co-workers have also reported a transparent, conductive and porous F-doped SnO<sub>2</sub> coating for quartz sheets as an electrode for photoelectrochemical water splitting<sup>[148]</sup>. The increasingly growing optoelectronic industry leads to a high demand on semiconductors with high performance features including high carrier mobility, high optical transmittance, and its potential mechanical flexibility. Tin oxide can fulfil these criteria, as well as it has the advantage of abundant material, mechanical and chemical stability. In addition, it has a moderate temperature processing range compared to Si and other metal oxides. SnO<sub>2</sub> based thin film transistors (TFT), including its composite based TFT, have reached the maturity to be incorporated in the display device industry<sup>[149]</sup>.

Tin oxide-based materials were explored in various application related to solar cells particularly for Perovskite Solar Cells (PSCs). This latter will be discussed in detail later in this chapter. Research efforts in thin film solar cells have explored tin oxide-based materials for the implementation of tin oxide thin film solar cells. Thomere et al. have reported fluorine-doped tin oxide as a transparent back contact for thin film solar cells of CuGaSe<sub>2</sub> absorber via two-step process and achieved the Power Conversion Efficiency (PCE) of 5.4%. It is worth noting that the relatively higher bandgap of the CuGaSe<sub>2</sub> (pure Ga) of 1.68 eV provides an appealing absorber for a tandem solar cell<sup>[150, 151]</sup>. Jiashuai Li et al. have reported Mg-doped SnO<sub>2</sub> as a buffer layer for Sb<sub>2</sub>S<sub>3</sub>-based solar cells and the best cell has reached 6.3% of PCE<sup>[152]</sup>. Jian-Min Wu and co-workers have reported a thin film GeSe-based solar cell using SnO<sub>2</sub> as a buffer layer with PCE of 0.51%. It was found that SnO<sub>2</sub> is suitable for the band alignment of this cell<sup>[153]</sup>. Ismael et al. have explored a heterojunction solar cell of Pd-doped SnO<sub>2</sub>/porous Si/c-Si and have achieved a PCE higher than 14%<sup>[154]</sup>. It is worth noting that Indium Tin Oxide (ITO) material is widely explored in different solar cells, namely organic solar cells, polymer solar cells, dye synthesized solar cells, silicon based solar cells including polycrystalline, monocrystalline, amorphous, mono-like, and heterojunction solar cells<sup>[16, 18, 155-188]</sup>.

Sn oxide-based materials has been widely investigated in Light Emitting Diodes (LEDs) to improve the device performance. Zinc doped tin oxide can improve light extraction to its high optical transmittance.

The external quantum efficiency is also a key aspect to improve the performance of the LEDs. Tin oxide based-material was also explored as a potential anode material for large area and flexible Organic Light Emitting Diodes (OLEDs). Esro and co-workers have established that Sb-doped SnO<sub>2</sub> at Wt.2% has provided the lowest sheet resistance and highest carrier mobility. Lee et al. have explored tin oxide nanoparticles as an interface layer sandwiched by both ITO cathode and ETL for OLED device. It was found that the incorporation of this interface layer improves the band alignment of layers and improve the electron injection in the OLED device [189-191].

There are many other applications for tin oxide-based materials where its properties were suitable to achieve the required functionalities, coating is among these other applications. Park et al. have reported a fluorine-doped tin oxide as a coating layer of stainless steel 316 plates for development of both corrosion and low contact resistant for polymer electrolyte membrane fuel cells<sup>[192]</sup>. Catalysis is another field where tin oxide-based materials are applied. Cheng et al. have reported the synthesis of a P-modified Sn/SnO<sub>x</sub> as core/shell structure catalyst for electrochemical CO<sub>2</sub> reduction reaction<sup>[193]</sup>.

Tin oxide can crystallize in two stoichiometric phases, namely SnO and SnO<sub>2</sub> which are known to be wide band gap oxide semiconductors with tetragonal litharge and rutile type structures, respectively. Tin, as an oxide, has two chemical states of Sn<sup>2+</sup> for SnO phase and Sn<sup>4+</sup> for the SnO<sub>2</sub> one. The first is known to be intrinsically p-type semiconductor while the second is n-type. SnO<sub>2</sub> material, which is known as tin (IV) dioxide or stannic oxide, is one of the widely known as wide bandgap TCO materials, while SnO is known as tin (II) oxide or stannous oxide<sup>[88, 127]</sup>.

### 1.2.1. Tin(IV) oxide (SnO<sub>2</sub>) thin film

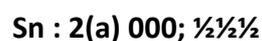
#### a) Thin film growth

There are various techniques to grow SnO<sub>2</sub> thin film and these techniques can be categorized as physical deposition or chemical deposition. The widely used physical deposition techniques are Pulsed Laser Deposition<sup>[194, 195]</sup>, E-beam Evaporator<sup>[22, 196]</sup>, RF sputtering<sup>[197-202]</sup>, Molecular Beam Epitaxy<sup>[203-205]</sup>. The chemical deposition techniques are Metal Organic Chemical Vapour Deposition<sup>[206, 207]</sup>, Atmospheric Pressure Chemical Vapour Deposition<sup>[208]</sup>, Atomic Layer Deposition<sup>[209]</sup>, Sol-Gel process<sup>[210]</sup>, Spin Coating<sup>[211]</sup>, and Spray Pyrolysis<sup>[212, 213]</sup>. Many of the described depositions techniques above are Gaseous state – deposition techniques. Within this class, there are mainly two distinct thin film deposition process categories which are chemical vapour deposition (CVD) and physical vapour deposition (PVD). The basic principle of CVD relies on the precursor vapour which is a used for chemical reaction(s) taking place at the substrate. While PVD's basic principle relies on the condensation at the substrate of the vapour made from atoms and/or molecules. Both CVD and PVD processes form thin films onto the substrate. The other class of deposition techniques is Solution state – deposition

techniques. This class includes spin coating, sol-gel, spray pyrolysis. Among all these techniques, sputtering technique is the most important technique due to its appealing features including the non-thermal ion impact induced vaporization due to the collision between the atoms at the surface of the target. The high deposition rate, the reduced space between the target and the substrate, and the high purity and vacuum deposition are interesting advantages. Furthermore, its main advantage compared to other PVD techniques is the superior adhesion and possibility to achieve uniform and higher thickness thin films. Sputtering has benefited from the technological advancements in the last decades which made it highly efficient and reliable in terms of real-time control, pumping, plasma, and automation systems. Sputtering is already used in many industrial applications, such as coatings, solar cells, and microelectronics. For metal oxides, reactive sputtering or sputtering in reactive mode is an attractive technique due to its intrinsic features of sputtering in addition to the oxygen reactive gas which allows the control of the level of the oxidation and consequentially tuning the optical and electrical properties of the thin film. Substrate temperature and O<sub>2</sub> partial pressure/flow rate parameters are of great importance for reactive sputtering to form metal oxides. These two parameters can have an impact on both crystalline properties and oxygen vacancies within the thin film<sup>[214-218]</sup>.

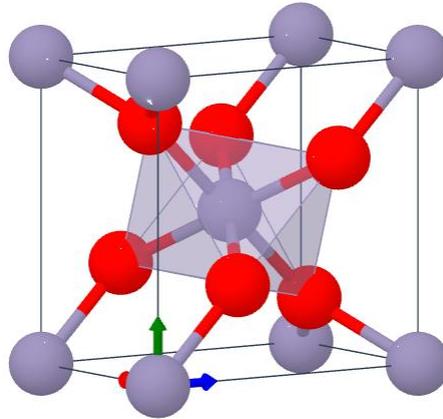
b) Structural properties of SnO<sub>2</sub>

Tin(IV) oxide or stannic oxide (SnO<sub>2</sub>) exists naturally on earth and in mineralogy it is known as Cassiterite which is the main ore, or in other words, the main source of tin metal. It has the rutile crystal structure where the unit cell is tetragonal P4<sub>2</sub>/mnm. The lattice constants are around a = b = 4.74 Å and c = 3.19 Å. There are other oxides, metal dioxides, which have also this type of crystal structure such as MnO<sub>2</sub>, TiO<sub>2</sub>, GeO<sub>2</sub>, VO<sub>2</sub>, CrO<sub>2</sub>, and TeO<sub>2</sub>. This crystal structure is the main structure for SnO<sub>2</sub> as it is very stable at ambient temperature and under atmospheric pressure. The crystal lattice where Sn has a six-fold coordination to the three-fold oxygen atom coordination is illustrated in Figure 6. The lattice positions can be represented by:



Where u is the variable parameter around 0.307 for SnO<sub>2</sub>. Each octahedra of oxygen is attached to two adjacent ones via edge sharing in the c-axis, while it is attached to the other octahedra via corner sharing. Furthermore, there are other polymorphs of SnO<sub>2</sub> which were demonstrated theoretically and/or experimentally by transiting from rutile SnO<sub>2</sub> at a very high pressure in the order of tens of GPa, namely CaCl<sub>2</sub>-type SnO<sub>2</sub> phase with an orthorhombic structure - Pnmm (#58) and lattice constants around a = 4.68 Å, b = 4.63 Å and c = 3.16 Å as well as α-PbO<sub>2</sub>-type SnO<sub>2</sub> phase with also an

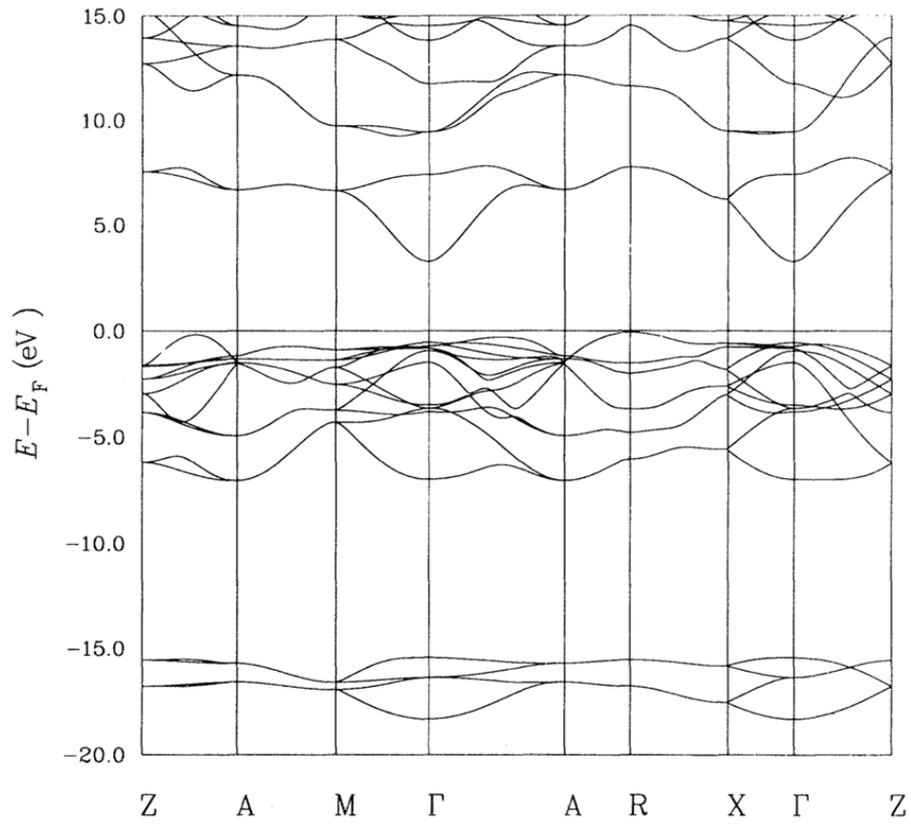
orthorhombic structure -  $Pbcn$  (#60) and lattice constants around  $a = 4.707 \text{ \AA}$ ,  $b = 4.710 \text{ \AA}$  and  $c = 5.246 \text{ \AA}$ . Furthermore, Pyrite-type  $\text{SnO}_2$  phase has also been reported with cubic structure -  $Pa\bar{3}$  (#205) and lattice constant of  $a = 4.888 \text{ \AA}$  and  $a = 5.066 \text{ \AA}$ . However, there is a significant gap between the theoretical and experimental data for pyrite-type  $\text{SnO}_2$  [36, 219].



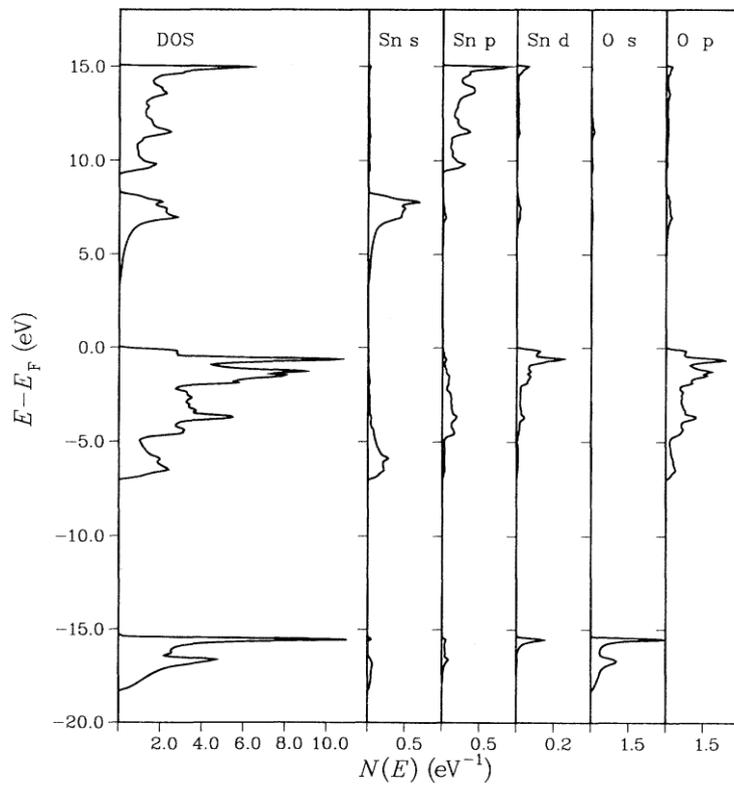
**Figure 6:** Crystal structure of rutile structure  $\text{SnO}_2$ . Tin atoms in purple and oxygen atoms in red. Reprinted (adapted) from The Materials Project [56]. License: CC BY 4.0.

#### c) Electronic structure of $\text{SnO}_2$

Tin(IV) oxide of rutile structure is established to be a direct wide bandgap semiconductor with a bandgap of 3.7 eV at the point  $\Gamma$  in the Brillouin-zone. Like  $\text{In}_2\text{O}_3$  and  $\text{ZnO}$  which were discussed in previous sections,  $\text{SnO}_2$  has a complex band structure as it has many valence electrons in each unit cell. In the Brillouin-zone, the band structure of  $\text{SnO}_2$  show a high degree of symmetry. Figure 7 reveals that the conduction band in both  $\Gamma$ -X and  $\Gamma$ -M directions shows a free electron-like feature. Moreover, the minimum gap between the conduction and valence bands was found at  $\Gamma$  point and it is around 3.7 eV. Figure 8 reveals that  $\text{O}(2s)$  and  $\text{Sn}(4d)$  related states induce four narrow and deep bands within the valence bands and due to their dispersion, they are not considered as core states. There are twelve bands, above these narrow and deep bands, which are considered as the valence band. The lower of part of these twelve bands is dominated by hybridized  $\text{O}(2p)$  with  $\text{Sn}(5s)$  states and the upper part of these bands is dominated by O related 2p and Sn related 5p-4d orbitals. The lower part of the conduction band is dominated by antibonding  $\text{Sn}(5s)$ - $\text{O}(2p)$  orbital while the upper part is dominated by  $\text{Sn}(5p)$  states [54, 61, 220, 221].



**Figure 7:** Band structure for rutile SnO<sub>2</sub>. Reprinted (adapted) with permission from [220]. Copyright 1995 American Physical Society.



**Figure 8:** Density of states (DOS) of rutile SnO<sub>2</sub> along with the partial density of states (pDOS) related Sn and O. Reprinted (adapted) with permission from [220]. Copyright 1995 American Physical Society.

#### d) Literature of undoped SnO<sub>2</sub>

In the literature, recent undoped SnO<sub>x</sub> thin films reports are very limited compared to the reports related to the doped ones. This is mainly due to the attractive properties brought by the doping of SnO<sub>x</sub>. In this section, the literature review is focusing only on the undoped tin oxide particularly, the growth conditions and the optoelectronic performance. This performance is assessed by Haacke figure of merit (FoM) equation<sup>[301]</sup> ( $\phi = T^{10}/R_s$ ), where  $\phi$  is figure of merit in  $\Omega^{-1}$ , T is the average optical transmittance in %, and  $R_s$  is the sheet resistance in  $\Omega \cdot \square^{-1}$ .

Sivakumar et al. have reported the deposition of tin oxide using sol-gel spin coating technique on glass substrate where they could synthesize a polycrystalline tetragonal rutile SnO<sub>2</sub> with a preferred orientation of (110) and the crystallite size around 56 nm. The achieved optoelectronic properties were a high optical average transmittance greater than 85% and the electrical resistivity was around  $3.3 \times 10^{-3} \Omega \cdot \text{cm}$ . The optical bandgap established by Tauc plot was 3.94 eV<sup>[222]</sup>. On the other hand, using PVD technique, Tao et al. have reported SnO<sub>2</sub> films on glass substrate prepared by reactive magnetron sputtering using metallic Sn target while varying the substrate temperature and O<sub>2</sub> flow rate. They found that SnO<sub>2</sub> thin films have an amorphous microstructure if deposited at room temperature regardless the O<sub>2</sub> flow rate. However, the crystallinity improved at higher temperature (150 and 300 °C) even at low O<sub>2</sub> flow rate. It was also reported that the preferred crystalline orientation is also affected by O<sub>2</sub> flow rate. The best conductive sample with resistivity of  $3.65 \times 10^{-3} \Omega \cdot \text{cm}$  was found to be deposited at room temperature and its average transmittance is 81.5% in the visible range. The authors established that electrical conductivity is controlled by oxygen vacancies<sup>[223]</sup>. Soumya reported the synthesis of amorphous tin oxide thin film prepared via sol-gel spin coating technique followed by a heating process to remove the solvents at 100 °C. XRD diffraction has not shown any crystalline features which reveals the amorphous microstructure of the tin oxide thin film. The average optical transmittance in the visible region was 85% and the established optical bandgap by Tauc plot method was 3.64 eV. The calculated resistivity was around  $0.12 \Omega \cdot \text{cm}$  and its FoM at 550 nm was  $2.2 \times 10^{-5} \Omega^{-1}$ <sup>[224]</sup>. Yadava et al. have reported a tin oxide thin film deposition using chemical vapour deposition set-up, namely, reactive thermal evaporation of SnCl<sub>2</sub>, for conducting electrode application. Structural analyses using XRD have shown that increasing the temperature has improved the crystallinity of SnO<sub>2</sub> and lead to different preferential orientation at each temperature. Optical analyses using UV-Visible spectroscopy (UV-Vis) have shown that the optical transmittance increases as the deposition temperature of the substrate increases from 400 to 550 °C, then this optical transmittance sharply decreases at 600 °C. The lowest electrical resistivity was reported for samples deposited at 500 and 550 °C around  $0.05 \Omega \cdot \text{cm}$ <sup>[225]</sup>. Yadav et al. have reported a tin oxide thin film synthesis using reactive thermal evaporation of SnCl<sub>2</sub> in air atmosphere on glass substrate as a cost-effective deposition technique for good quality SnO<sub>2</sub>. Structural analyses using XRD have shown that absence of diffraction

peaks related Sn and SnO phases and that increasing the temperature has improved the crystallinity of SnO<sub>2</sub> and all samples show the presence of polycrystalline cassiterite tetragonal rutile structure. The average crystallite size related to peak SnO<sub>2</sub> (110) ranged from 20 to 32 nm. Optical analyses using UV-Vis has shown that the maximum optical transmittance is around 90% and 85% before and after heating, respectively and the transmittance increase as the temperature of the substrate, while deposition, increases from 250 to 450 °C. The sheet resistance decreases as the deposition temperature of the substrate increases from 250 to 450 °C [226]. Khallaf et al. have reported a deposition of SnO<sub>2</sub> using multi-dip chemical bath deposition followed by post-deposition thermal annealing at 400 °C for 1 hour. Structural analysis by XRD has shown an amorphous structure for the as-deposited tin oxide and a polycrystalline orthorhombic microstructure for the annealed SnO<sub>2</sub>. X-ray Photoelectron Spectroscopy (XPS) results have revealed the presence of Sn(IV) and mainly oxide related O which represents SnO<sub>2</sub> phase. The optical bandgap for the annealed SnO<sub>2</sub> was reported around 4.42 eV. The electrical resistivity for the annealed SnO<sub>2</sub> was very high in the order of 10<sup>3</sup> Ω·cm which is due to the low charge carrier concentration in the order of 10<sup>12</sup> cm<sup>-3</sup> [227]. An older article (reported in 1981) from Muranaka et al. have described the deposition of SnO<sub>2</sub> by reactive Sn thermal evaporation technique along with O<sub>2</sub> flow. Structural analyses have shown that increasing the temperature and the O<sub>2</sub> pressure have formed and/or improved the crystallinity of the SnO<sub>2</sub>. Low O<sub>2</sub> pressure has led to the formation of metallic Sn and SnO phases while low substrate temperature during the growth has led to an amorphous microstructure of tin oxide. Both films deposited at 3 x 10<sup>-3</sup> and 5 x 10<sup>-3</sup> Torr at 420 °C has reached an optical transmittance around 90% and an electrical resistivity below 5 x 10<sup>-2</sup> Ω·cm [228]. Belayachi and co-workers reported the synthesis of polycrystalline cassiterite tetragonal tin oxide thin films on silicon substrate (oriented (100)) prepared by reactive radio-frequency magnetron sputtering at a constant O<sub>2</sub> flow rate and at different substrate temperatures of 100, 200, 300 and 400 °C. The best conductive sample was deposited at 100 °C with an average transmittance around 84%. It was found that the substrate temperature had a slight effect on the optical transmittance and small increase in crystallite size as substrate temperature increases. However, as the substrate temperature increases, the resistivity increases by orders of magnitude from 4.45 x 10<sup>-3</sup> Ω·cm at 100 °C to 3.26 Ω·cm at 400 °C. The main reason for the conductivity decrease, as substrate deposition increases, is substantial decrease of mobility, while charge carrier concentration underwent a limited decrease [200].

Overall, there are a few important points to take into consideration. The first is that increasing the temperature of the substrate during the growth improves the crystalline microstructure of tin oxide thin films and the reports have revealed the polycrystalline tetragonal rutile structure of the crystalline tin oxide. Second, most of the reports reveal a relatively high average optical transmittance superior to 80% and an optical bandgap above 3.5 eV. Finally, the best electrical resistivity for these reports is around 10<sup>-3</sup> Ω·cm.

#### e) Literature of doped SnO<sub>2</sub> thin film

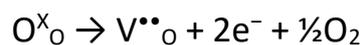
There is a large amount of literature related to SnO<sub>2</sub> doping as TCO thin films which is oriented towards different applications. This literature review will focus on most of the relevant and reports of SnO<sub>2</sub> doping in the optoelectronic field. This section will be presented in two parts, the first will be tackling the n-type conductive doped SnO<sub>2</sub> thin films literature while the second will focus on the literature reporting the conversion of the n-type conductivity of SnO<sub>2</sub> to the p-type one.

##### i. N-type doping of SnO<sub>2</sub>

As SnO<sub>2</sub> is widely known as an intrinsic n-type semiconducting material, most of the research efforts have been focusing on enhancing the optoelectronic properties of tin oxide by introducing doping elements and sometimes even compounds. It is worth noting that most of the literature about SnO<sub>2</sub> doping resulting an n-type conductivity were reporting material synthesis and deposition via chemical processes. This is mainly due to the ability to add and control doping element which the SnO<sub>2</sub> hosts.

Spray Pyrolysis is one of the most utilized techniques to deposit doped SnO<sub>2</sub> due to its process simplicity, precursor solution formulation versatility, cost-effectiveness, fast process as well as it is thermodynamical predictability<sup>[214]</sup>. Moholkar et al. have reported the deposition of fluorine doped tin oxide (FTO) thin films by spray pyrolysis on glass substrate at 475 °C using precursor solution SnCl<sub>4</sub>·5H<sub>2</sub>O (from 0.32 to 1.29 M) and ammonium fluoride. All FTO thin films have a polycrystalline microstructure matching with SnO<sub>2</sub> tetragonal structure with a preferential orientation of (200) plane. Sheet resistivity varies from 61.6 to 3.71 Ω·□<sup>-1</sup> where the lowest value is reported for SnCl<sub>4</sub> concentration of 0.65 M. The optical transmittance at 550 nm was 87%. The best Figure of Merit (FoM), that indicates the optoelectronic performance, was 7.26 x 10<sup>-2</sup> Ω<sup>-1</sup> for FTO with SnCl<sub>4</sub> concentration of 0.81 M<sup>[229]</sup>. Singh and Basu have reported synthesis of antimony doped tin oxide (ATO) and fluorine doped tin oxide on glass slides at temperature varying from 400 to 540 °C using a modified spray pyrolysis system by spraying upwards to assure that all droplets are evaporated while reaching the surface of the substrate. The polycrystalline tetragonal Cassiterite microstructure have been determined by XRD related CuK<sub>α</sub> diffracted pattern of the FTO films with no specific crystalline preferential orientation. Optical measurements have shown that both ATO and FTO thin films at 100 nm thickness have 95% optical transmittance in the visible region. The determined optical bandgap is 3.9 eV for ATO and 4.1 eV for FTO. Moreover, increasing the thickness has led to a reduction in optical transmittance from 95% at 100 nm to 31% at 510 nm thickness for ATO. However, the thickness effect is relatively smaller for FTO as the optical transmittance is reduced from 95% at 100 nm to 70% at 470 nm. Electrical measurements have revealed that increasing the substrate temperature from 400 to 530 °C for FTO and from 430 to 540 °C for ATO has decreased remarkably the sheet resistance of both FTO and ATO thin films. This study also has established that increasing the thin film thickness has

decreased the sheet resistance while varying the thickness from 100 to approximately 600 nm for both ATO and FTO. However, the effect of the thickness is more predominant for ATO compared to FTO. This study has found that the best sheet resistance was recorded at 1.5 wt.% of NH<sub>4</sub>F for FTO and at 3 mol % of SbCl<sub>3</sub> for ATO. The highest Figure of Merit for FTO is  $24.2 \times 10^{-3} \Omega^{-1}$  achieved at 190 nm thickness with 1.2 wt.% of NH<sub>4</sub>F at the growth where the  $R_{sh} = 20 \Omega \cdot \square^{-1}$  and the optical transmittance at 600 nm is 93%. The highest FoM for ATO is  $7 \times 10^{-3} \Omega^{-1}$  achieved at 200 nm thickness with SbCl<sub>3</sub> at 3 mol%, where the  $R_{sh} = 40 \Omega \cdot \square^{-1}$  and the optical transmittance at 600 nm is 88%. As per these optoelectronic performance results, FTO have shown a better optoelectronic performance compared to ATO<sup>[213]</sup>. Babar et al. have reported antimony doped tin(IV) oxide thin film on glass substrate at 475 °C using spray pyrolysis deposition technique, Sb doping was established by adding SbCl<sub>3</sub> at 2 at.%. This concentration was determined as the optimum for high quality Sb-doped SnO<sub>2</sub> (Sb:SnO<sub>2</sub>) thin film. The structural properties, assessed by XRD, have determined the polycrystalline tetragonal structure of the Sb:SnO<sub>2</sub> thin films without presence of any secondary phases. The optical transmittance at 550 nm was 70% and the electrical resistivity was  $8.9 \times 10^{-4} \Omega \cdot \text{cm}$ . The optoelectrical performance of these thin films were assessed via FoM and the best thin film have achieved  $2.1 \times 10^{-3} \Omega^{-1}$  [230]. Md. Faruk Hossain et al. have assessed the impact of Sb doping on tin oxide and they reported the synthesis Sb:SnO<sub>2</sub> on glass substate at 400 °C using nebulized spray pyrolysis via SnCl<sub>2</sub> where the doping concentration varies from 0 to 8 mol.% via SbCl<sub>5</sub>. The crystalline structure of undoped SnO<sub>2</sub> was found to be close to a standard tetragonal rutile SnO<sub>2</sub> structure where the lattice constants “a” and “b” are equal to 4.738 Å and “c” is equal to 3.188 Å. These values were found to be decreasing as the antimony doping increases. The average optical transmittance for the undoped SnO<sub>2</sub> in the visible region was 85% and it slightly decreases as the Sb doping concentration is increasing from 0 to 8%. The optical bandgap is found to be 3.7 eV for undoped SnO<sub>2</sub> while the Sb:SnO<sub>2</sub> increases from 3.8 to 3.9 eV as the doping increases from 2 to 8%, respectively. Electrical measurements have revealed that the resistivity decreases for SnO<sub>2</sub> from  $16.1 \times 10^{-4}$  to  $6.4 \times 10^{-4} \Omega \cdot \text{cm}$  as the Sb doping concentration increases from 0 to 8%. This work has highlighted the importance of oxygen vacancies and interstitial Sn in the formation of charge carriers as described in the below Kroger-Vink notation.



Where ( $\text{V}^{\bullet\bullet}_{\text{O}}$ ) is a doubly ionized vacancy site with an oxygen anion leaving the SnO<sub>2</sub> crystal structure and two conduction electrons generated from an occupied oxygen site ( $\text{O}^{\times}_{\text{O}}$ ). The best optoelectrical performance assessed by FoM is found to be  $1.39 \times 10^{-3} (\Omega/\text{cm}^2)^{-1}$  accomplished in the visible region with 2% antimony doping<sup>[231, 232]</sup>. Ramarajan et al. have reported Sb:SnO<sub>2</sub> thin films by chemical spray pyrolysis and deposited on glass substate at temperatures varying from 360 to 490 °C while keeping the antimony doping at 5 wt.%. The 0.5 M precursor solution is containing SnCl<sub>2</sub>.2H<sub>2</sub>O and SbCl<sub>5</sub>

dissolved and diluted in mixture of deionized water and ethylene glycol. ATO samples have exhibited a tetragonal polycrystalline structure without any secondary phases. Optical transmittance for ATO thin films varies from 62% to 89% and optical bandgap slightly varies from 3.73 to 3.81 eV. Electrical measurements have shown that the lowest reported resistivity for ATO thin films was  $7.9 \times 10^{-4} \Omega\cdot\text{cm}$  for ATO thin film deposited at 450 °C. The best FoM for ATO thin films was  $1.25 \times 10^{-2} \Omega^{-1}$  which was also reported at 450 °C deposition temperature<sup>[233]</sup>. Ramarajan et al. also reported the deposition of Tantalum doped SnO<sub>2</sub> (Ta:SnO<sub>2</sub>) on glass substrate using different doping weight percentage from 0 (undoped SnO<sub>2</sub>) to 5 wt.% Ta:SnO<sub>2</sub>. The precursor solution is prepared by SnCl<sub>2</sub>·2H<sub>2</sub>O and TaCl<sub>5</sub> dissolved and diluted in mixture of deionized water and ethylene glycol. The substrate temperature was kept at 400 °C during the growth. The microstructure of all samples was found to be a polycrystalline tetragonal crystalline structure without any secondary phases. The optical transmittance at 550 nm varies from 60 to 86% and the highest value was found to be 86% for 4 wt.% Ta:SnO<sub>2</sub>. The optical bandgap to slightly varies from 3.95 to 4.01 eV. Electrical measurements have revealed that electrical resistivity ranges from  $10^{-2}$  to  $4.36 \times 10^{-4} \Omega\cdot\text{cm}$ , this lowest resistivity is reported for 4 wt.% Ta:SnO<sub>2</sub>. Thus, this sample has the best achieved FoM of  $4.73 \times 10^{-2} \Omega^{-1}$  <sup>[234]</sup>. Mrabet et al. have reported the deposition of lanthanum doped tin(IV) oxide (La:SnO<sub>2</sub>) thin films on glass substrate via spray pyrolysis at 450 °C where the La doping varies from 0 to 3%. La:SnO<sub>2</sub> have exhibited a polycrystalline tetragonal structure and the preferred orientation for all samples is (110) plane except for La doped SnO<sub>2</sub> at 2% where the preferred orientation is along the (200) plane. For all La:SnO<sub>2</sub> thin films, including the undoped one, the mean optical transmittance is found to be varying between 0.73 and 0.81. The optical bandgap of La:SnO<sub>2</sub> thin films slightly varies from 3.86 to 3.88 eV. Electrical measurements have established that the sheet resistance varies drastically from 358 and 334  $\Omega\cdot\text{cm}^{-1}$  for undoped SnO<sub>2</sub> and 1% La:SnO<sub>2</sub>, respectively, to 2.1 and 3  $\Omega\cdot\text{cm}^{-1}$  for 2% and 3% La:SnO<sub>2</sub>, respectively. Therefore, the best optoelectronic performance assessed by FoM is found to be  $10.6 \times 10^{-2}$  and  $10.7 \times 10^{-2} \Omega^{-1}$  for 2% and 3% La:SnO<sub>2</sub>, respectively<sup>[235]</sup>.

On the other hand, chemical vapour deposition (CVD) was also reported as a synthesis technique for SnO<sub>2</sub>. Vishwakarma et al have reported a TCO of arsenic doped tin oxide deposited on soda glass substrate using CVD through the mixture vapour of SnCl<sub>2</sub> using O<sub>2</sub> as a gas carrier and AsCl<sub>3</sub> using N<sub>2</sub> as a gas carrier. The best electrical conductivity was reported at arsenic doping of 0.15 wt.% where the electrical resistivity has reached  $1.5 \times 10^{-4} \Omega\cdot\text{cm}$  and this resistivity increase as arsenic concentration decreases towards lower concentration and also as arsenic concentration increases towards higher concentrations. The arsenic doping concentration of 0.15 wt.% is also reported for the best mobility and charge carrier concentration. The best reported figure of merit was  $16.5 \times 10^{-3} \Omega^{-1}$  observed at 700 nm for SnO<sub>2</sub> deposited at 400 °C <sup>[236]</sup>. Li et al. have reported ultrasonic-vibration-assisted laser annealed fluorine doped tin oxide (FTO) commercially prepared by CVD where the laser beam central

wavelength at 532 nm with 1 kHz repetition rate. XRD results have revealed the polycrystalline tetragonal structure of tin oxide films without presence of any secondary phases and the crystallite size varies between 22 and 30 nm. The thickness of the films ranges from 695 and 725 nm. The average optical transmittance from 400 to 800 nm of FTO ranges from 81 to 85.5% while the sheet resistance varies from 8.8 to 9.4  $\Omega \cdot \square^{-1}$ . The optoelectronic performance was assessed by figure of merit (FoM) and the best film has reached FoM of  $2.32 \times 10^{-2} \Omega^{-1}$  for using 48 kHz vibrational frequency and 300 W vibration power parameters for the annealing<sup>[237]</sup>. Along with reports related to CVD synthesis, Atmospheric Pressure Chemical Vapour Deposition (APCVD) was also reported as a much more appealing synthesis technique to form SnO<sub>2</sub> based thin films due to its deposition upscaling ability to large area as well as the superior growth rate. For instance, Talaty et al. reported the synthesis and material's characterisation of tin(IV) oxide based thin films on glass substrate via APCVD process. One of the main studies of this report was the growth rate at different temperature. It is established that increasing the substrate temperature increases the growth rate in this technique and the best growth rate was as high as 2300 Å/min at 650 °C which is highly desired growth rates for industrial TCO applications. Structural properties revealed by XRD have found that all samples have a polycrystalline tetragonal structure for both undoped and F-doped SnO<sub>2</sub> (FTO) without any preferred orientation. Furthermore, Scanning Electron Microscopy (SEM) has revealed microstructural features where the SnO<sub>2</sub> related grain size varies for from 50 nm to 500 nm. Optical properties revealed by UV-Visible spectroscopy have shown that the optical transmittance for all samples varies from 76 to 86% from 400 to 800 nm. The fluorine doping for SnO<sub>2</sub> has clearly improved the optical transmittance. Furthermore, the fluorine doping for SnO<sub>2</sub> has also improved the electrical properties as indicated by the decrease in the sheet resistance  $R_{sh}$  from 600 to 10  $\Omega \cdot \square^{-1}$ . These two improvements in optical and electrical properties have a huge impact in optoelectronic performance of SnO<sub>2</sub> thin film after the fluorine doping as they improved FoM from  $0.4 \times 10^{-3}$  to  $16.7 \times 10^{-3} \Omega^{-1}$ . They found that their FTO sample has reached an optoelectronic performance close to the commercial FTO (American Float Glass, Pilkington NSG TEC Glass TEC 15)<sup>[208]</sup>. Tran et al. have reported a FTO thin films on glass substrate using a sol-gel dip coating method using a precursor solution made of SnCl<sub>4</sub>·5H<sub>2</sub>O and SnF<sub>2</sub> (from 0 to 10 mol%) followed by post-annealing process at 500 °C for 2 hours. ATO thin films have all a polycrystalline tetragonal rutile structure. Optical transmittance at 550 nm varies from 91.7 to 89% while increasing SnF<sub>2</sub> concentration from 0 to 10 mol% and the optical bandgap varies between 3.66 to 3.91 eV. Electrical resistivity varies from  $7 \times 10^{-4}$  to  $29 \times 10^{-4} \Omega \cdot \text{cm}$  where the lowest resistivity corresponds to ATO made with SnF<sub>2</sub> concentration of 6 mol%. This ATO sample has also the best FoM of  $2.43 \times 10^{-2} \Omega^{-1}$  [238].

Finally, the best Sb:SnO<sub>2</sub> thin film, as a transparent conducting oxide, was also reported by Fauzia et al. which was deposited by ultrasonic assisted spray pyrolysis on glass substrates. The precursor solution

was prepared by 1M SnCl<sub>2</sub>.H<sub>2</sub>O with Sb doping using SbCl<sub>3</sub> at 2 wt.% and diluted in ethanol. The structural properties revealed by XRD and the crystal structure was found to be tetragonal rutile structure. The average optical transmittance of all thin films was determined to be higher than 80% from 450 to 900 nm and all ATO samples exhibit an optical bandgap of 3.83 eV. In this research work and to the best of our knowledge, the co-authors have reported the lowest recorded electrical resistivity for ATO of  $0.58 \times 10^{-4} \Omega\cdot\text{cm}$  ( $5.8 \times 10^{-5} \Omega\cdot\text{cm}$ ). Thus, the achieved optoelectronic performance by FoM of  $47.22 \times 10^{-2} (\Omega/\text{cm}^2)^{-1}$  is found to be the largest compared to the literature<sup>[239]</sup>.

Reports about tin(IV) oxide doping using physical process are relatively limited compared to chemical processes. While it is very straightforward to add doping in solution precursor with every desired concentration, it is fairly more challenging to add doping elements in physical processes as it requires a long optimization process for each desired doping concentration. There is a report from Dian and co-workers about deposition of tellurium doped tin(IV) oxide (Te:SnO<sub>2</sub>) thin films on glass substrate at 500 °C using pulsed laser deposition (PLD) followed by post-deposition annealing at 500 °C for 30 minutes at atmospheric pressure. All Te:SnO<sub>2</sub> thin films have a polycrystalline tetragonal structure matching with the one of SnO<sub>2</sub>. However, the crystallinity is more predominant as the O<sub>2</sub> pressure increases from 1 to 150 mTorr during the growth. The post-deposition annealing has improved the crystallinity of Te:SnO<sub>2</sub> thin films. The highest optical transmittance of Te:SnO<sub>2</sub> thin films is 87% for PO<sub>2</sub> of 100 mTorr and the optical bandgap related to this thin film is 3.5 eV. The lowest resistivity of  $4 \times 10^{-2} \Omega\cdot\text{cm}$  is recorded for the as-deposited thin film as PO<sub>2</sub> of 100 mTorr. The highest FoM is around  $1.2 \times 10^{-5} \Omega^{-1}$  [240].

## ii. P-type doping of SnO<sub>2</sub>

In the last two decades and upon the major achievement in metal oxide p-type conductivity, there are several reports related to theoretical studies and experimental results, showing the possibility of forming a doped SnO<sub>2</sub> with a p-type conductivity. Among other achievements in p-type metal oxide, this has opened a new horizon for optoelectronic devices where an all-oxide device is possible, if these p-type metal oxides can achieve the required performance for industrial applications. Nevertheless, there is still huge research and development to be done on p-type SnO<sub>2</sub> thin films, as well as most of p-type metal oxides, to reach the standards for such large-scale industrial applications.

To reach p-type there are two ways of doping SnO<sub>2</sub>, the first approach is to incorporate metals with lower valence cations such as Al, Ga, and In, to substitute Sn in its SnO<sub>2</sub> lattice site to generate p-type charge carriers. Theoretically, this approach might be a subject to some challenges mainly due to the deep localised defect's state which can lead to pinning the Fermi level above the valence band maximum (VBM) and prevent the formation of p-type doping due to self-compensation. The second approach is to incorporate a higher anion doping element mainly nitrogen due to many advantages

including similar electronegativity, ionic radius compared to oxygen. Nitrogen doping can also reduce the oxygen vacancies which are the main donor for electron charge carriers as nitrogen atoms can occupy the oxygen lattice site within the SnO<sub>2</sub> crystal structure. In addition, N2p-O2p hybridisation orbitals, induced by nitrogen doping, is expected to increase the valence band maximum (VBM) position. Therefore, nitrogen doping for SnO<sub>2</sub> is a potential doping to convert the n-type conductivity of SnO<sub>2</sub> to p-type<sup>[241]</sup>.

There are numerous reports about theoretical studies related metallic doping of SnO<sub>2</sub>. K. Lai et al. have reported the effect of Al doping on the electronic and optical properties of Al-doped SnO<sub>2</sub> along with the presence and the absence of oxygen vacancies via first-principle calculations. The co-authors have revealed that Al doping enable p-type conductivity while substituting Sn in SnO<sub>2</sub> and the presence of O vacancies can shift the Fermi level closer to conduction band. This can convert the conductivity back to n-type even with presence of Al doping<sup>[242]</sup>. Furthermore, there are fewer reports about experimental studies of p-type metallic doping for SnO<sub>2</sub>. This is mainly due to the difficulty and challenges to achieve a p-type conductivity based SnO<sub>2</sub> via metallic doping. J. Ni et al have reported a transparent and p-type conducting Sb doped SnO<sub>2</sub> (Sb:SnO<sub>2</sub>) thin films on quartz glass prepared by RF magnetron sputtering at 200°C using a pre-made target of 20 mol% Sb:SnO<sub>2</sub> followed by thermal annealing at different temperatures ranging 550 to 800°C. Structural properties determined by XRD have revealed that all samples have a rutile polycrystalline structure with fine grains without any secondary phases related to tin oxide or to antimony oxide. Optical measurements determined by UV-Visible spectroscopy have shown that the thermal annealing has improved the optical transmittance on the visible spectrum and that the absorption edge did not change for all samples with or without annealing. The calculated optical bandgap was around 3.6 eV. They reported that 600°C is a critical temperature for conductivity conversion from n-type to p-type, and all samples annealed above 600°C are p-type conductive. However, there is a reverse effect at higher temperatures, particularly at 750°C and above, where the hole related charge carrier concentration decreases to a very low level and eventually the Sb:SnO<sub>2</sub> conductivity switches back to n-type. The highest hole concentration was achieved using the thermal annealing at 700°C for 4 hours. The electrical resistivity has decreased significant after the annealing at 550°C from  $9.9 \times 10^{-2}$  to  $2.9 \times 10^{-3}$  Ω·cm, both are n-type. The electrical resistivity increases drastically at 600°C to 0.28 Ω·cm without a clear conductivity type (n/p). However, it slightly decreases to 0.25 Ω·cm as the annealing temperature increases to 650°C with a clear p-type conductivity. Then, the electrical resistivity decreases 0.17 Ω·cm at 700°C which is the highest p-type conductivity. After that, the electrical resistivity increases 0.15 Ω·cm at 750°C while keeping the p-type character. Finally, the electrical resistivity slightly decreases, and the conductivity is returned to n-type. This conversion from p-type to n-type is due to the decrease of p-type charge carriers at very high temperature. The co-authors have established a pn junction with both p-type and n-type Sb:SnO<sub>2</sub> thin

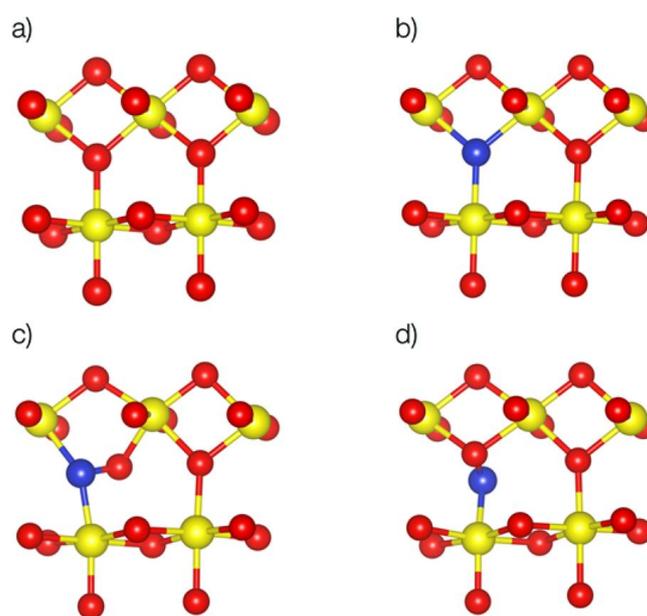
films in a diode structure using the highest p-type charge carrier concentration. The I-V results show that an ohmic contact is achieved as revealed by the linear behaviour<sup>[243]</sup>. T. Le et al. have reported the synthesis and characterization of Zn-doped SnO<sub>2</sub> (Zn:SnO<sub>2</sub>) deposited and annealed at different temperatures on quartz glass substrates via direct current (DC) magnetron sputtering. All the Zn:SnO<sub>2</sub> thin films have revealed a tetragonal rutile polycrystalline structure except the one deposited at room temperature and the crystallinity increases as a function of deposition temperature. Zn doping has favoured more the orientation towards (101) compared to undoped SnO<sub>2</sub> where the preferred orientation was on (110) plane. The average optical transmittance for Zn:SnO<sub>2</sub> thin films whether as-deposited or annealed is 84% which is very appealing for optoelectronic applications as TCO. The electrical resistivity has increased from 0.2 to 14.7 Ω·cm as the deposition temperature of Zn:SnO<sub>2</sub> increases from room temperature to 300°C. However, the resistivity decreases at higher temperatures from 14.7 to 4 Ω·cm as the deposition temperature increases from 300°C to 500°C and more importantly the conductivity has been converted from n-type to p-type at temperatures of 400°C and 500°C. It is worth noting that the thermal annealing has kept the p-type conductivity while significantly reducing the resistivity as well as that longer post-deposition annealing time increases the resistivity of the Zn:SnO<sub>2</sub> thin films. In order to confirm the p-type conductivity of Zn:SnO<sub>2</sub> thin films, a pn junction was made by Zn:SnO<sub>2</sub> on an n-type Si wafer using In as contact electrodes. I-V characteristics have confirmed the p-type character of Zn:SnO<sub>2</sub> thin film by showing an Ωic contact behaviour<sup>[244]</sup>. C.Y. Tsay et al. have also reported a p-type SnO<sub>2</sub> achieved by Ga doping ranging from [Ga]/[Sn]+[Ga] 0 to 20%. The Ga-doped SnO<sub>2</sub> (Ga:SnO<sub>2</sub>) thin films have been prepared using sol-gel spin coating followed first by drying and then by thermal annealing at 520°C for 1 hour under oxygen atmosphere. The crystalline structure of Ga:SnO<sub>2</sub> is found to be polycrystalline tetragonal structure closely matching with SnO<sub>2</sub> rutile structure and the crystallite size decreases as the Ga doping increases. The average optical transmittance from 400 to 800 nm and is relatively high for all samples without a clear trend and it is slightly varying between 87.4% and 89.5%. The optical bandgap is also slightly varying between 3.83 eV and 3.92 eV without a clear trend. The electrical conductivity was found to be n-type for undoped and 5% Ga:SnO<sub>2</sub> and it is p-type for SnO<sub>2</sub> with higher Ga doping. The electrical resistivity is 0.8 Ω·cm and it slightly decreases to around 0.9 Ω·cm for Ga doping from 0 to 5% which is expected due to the charge compensation by the generated acceptors. This is clearly observed by the reduction of n-type charge carrier concentration. For 10% Ga:SnO<sub>2</sub> thin film, the electrical resistivity decreases slightly to 0.8 Ω·cm, and the conductivity becomes p-type with a relatively low charge carrier concentration. As the Ga doping for SnO<sub>2</sub> increases to 15%, the electrical resistivity increases as well as the p-type charge carrier concentration increases which is due to the decrease in mobility. Finally, the electrical resistivity increases significantly for 20% Ga:SnO<sub>2</sub> and the p-type charge carrier concentration decreases. It is worth noting that this decrease in p-type charge carrier concentration at high doping levels is related

to the reverse effect which has been reported previously with different doping elements and it is reported previously in this section. The co-authors have confirmed the p-type conductivity of Ga:SnO<sub>2</sub> by achieving a pn junction made by p-type 15% Ga:SnO<sub>2</sub> and n-type Al:ZnO on glass substrate using In as contact electrode. P-type conductivity of Ga:SnO<sub>2</sub> has been confirmed by the Ωic contact behaviour in I-V characteristics<sup>[245]</sup>.

On the other hand, there are several reports about nitrogen doping for SnO<sub>2</sub>. A theoretical study by E. Albanese and co-authors have discussed the possible N doping in SnO<sub>2</sub>. They reported possible N doping within SnO<sub>2</sub> whether the substitutional doping or interstitial doping, Figure 9 shows these possibilities using realistic models. This study also highlighted the effect of nitrogen doping on the electronic structure<sup>[246]</sup>. K. Watanabe et al have reported nitrogen doped SnO<sub>2</sub> via two steps where the first step is the deposition of SnO<sub>2</sub> through reactive radiofrequency (RF) magnetron sputtering under O<sub>2</sub>/Ar (0.1) on Si wafer with SiO<sub>2</sub> surface, and the second step is the thermal annealing using rapid thermal process at different temperature ranging from 150 to 600°C under N<sub>2</sub> for 30 minutes. It was found that SnO<sub>x</sub> annealed at 600°C have been crystallised to rutile structure from the amorphous as-deposited SnO<sub>x</sub> and the conductivity type conversion from n-type to p-type was only enabled after post-deposition annealing at 600°C in high purity N<sub>2</sub> atmosphere. These results indicate that p-type SnO<sub>2</sub> thin film was achieved through nitrogen doping by filling oxygen lattice sites of the existing oxygen vacancies with nitrogen atoms<sup>[247]</sup>. Pan et al. have reported the synthesis of nitrogen doped SnO<sub>2</sub> using one step process and via reactive RF magnetron sputtering at 400°C with a mixture of O<sub>2</sub>/N<sub>2</sub> atmosphere. All samples have a typical pure cassiterite phase without any secondary phases namely metallic tin, tin nitride, or tin(II) oxide. Hall effect measurements have shown that SnO<sub>2</sub> without N doping have a mobility between 0.21 and 6.9 cm<sup>2</sup>/V·s and charge carrier concentration between 10<sup>18</sup> and 10<sup>19</sup> cm<sup>-3</sup>. By introducing a low N<sub>2</sub> partial pressure (PP), the nitrogen doping was not enough to neutralize n-type charge carriers, electrons, formed by the intrinsic defects of SnO<sub>2</sub>. Therefore, it has kept the n-type conductivity character. By introducing a higher N<sub>2</sub> PP, the nitrogen doping has increased and formed more, acceptors charge carrier, holes which led to conductivity type conversion from n-type to p-type. Introducing a high N<sub>2</sub> PP, the nitrogen doping was relatively high in SnO<sub>2</sub> which resulted in further decrease in resistivity, and it kept the n-type conductivity. This high N<sub>2</sub> PP had a reverse effect. This might be due to the formation of intrinsic compensating donors which became comparable to the generated acceptors at high level doping. In this letter, the authors have demonstrated a process toward achieving atomic nitrogen doping in SnO<sub>2</sub> under nitrogen different partial pressure toward achieving p-type SnO<sub>2</sub> films conductivity<sup>[248]</sup>. Y. Kim et al. have described his work on N-doped SnO<sub>2</sub> synthesis by reactive RF sputtering at 100°C and characterization on borosilicate glass substrate to probe effect of N-doping on material properties of N:SnO<sub>2</sub>, particularly, on the electrical conductivity. The nitrogen doping was controlled by varying the N<sub>2</sub> flow rate during the

growth from 0% for undoped SnO<sub>2</sub> to 40% while O<sub>2</sub> was fixed at 15% and the balance was Ar flow rate. XRD results show that there is no peaks related to Sn, Sn nitrides or any secondary phases which indicates nitrogen has been incorporated successfully within SnO<sub>2</sub> and (101) peak slightly shifted to a lower angle. This indicates that the d-spacing/interplanar spacing has slightly increases from 2.643 Å to 2.658 Å due to nitrogen doping. N-doping has improved the crystallinity of the host SnO<sub>2</sub>. The co-authors reported that nitrogen doping has reduced the conductivity of SnO<sub>2</sub> which is due to the reduction in O vacancies and possible self-compensation effect by the formation of low acceptor carrier concentration. Therefore, p-type conversion was not possible<sup>[249]</sup>. X. Ding et al. have reported the electrical and optical properties of undoped SnO<sub>2</sub> and N-doped SnO<sub>2</sub> thin film deposited via reactive magnetron sputtering at 400°C. SnO<sub>2</sub> was deposited under a mixture of O<sub>2</sub>/Ar of gas flow while N-doped SnO<sub>2</sub> was deposited under O<sub>2</sub>/N<sub>2</sub>. Tuning the O<sub>2</sub> PP as well as N<sub>2</sub> PP, has enabled to investigate the effect of N doping on the structural, optical and electrical properties N-doped SnO<sub>2</sub>. Microstructure was revealed using SEM imaging and it clearly showed the presence of cracks between the SnO<sub>x</sub> grains. These cracks became larger for N-doped SnO<sub>2</sub> compared to undoped SnO<sub>2</sub>. XRD results have found that all samples have a polycrystalline rutile structure and all N-doped SnO<sub>2</sub> samples have also a polycrystalline rutile structure. However, there is a relatively preferred growth direction towards (101) plane. Moreover, nitrogen doping for SnO<sub>2</sub> thin film did not affect its crystalline structure. The average optical transmittance in the visible region has been reduced after the nitrogen doping which is due to the absorption edge redshift for each O<sub>2</sub> PP. The optical bandgap is low for N-doped SnO<sub>2</sub> ranging from 3.60 to 3.98 eV compared to undoped SnO<sub>2</sub> ranging from 4.02 to 4.08 eV. It is worth noting that the thickness of N-doped SnO<sub>2</sub> thin films is higher than the undoped SnO<sub>2</sub> one for each O<sub>2</sub> PP. N doping has improved the electrical resistivity compared to undoped SnO<sub>2</sub> and it is established that increasing O<sub>2</sub> PP has reduced the resistivity from 8.5 x 10<sup>-1</sup> to 3.1 x 10<sup>-2</sup> Ω·cm for undoped SnO<sub>2</sub> and from 2.8 x 10<sup>-2</sup> to 5.9 x 10<sup>-3</sup> Ω·cm for N-doped SnO<sub>2</sub>. These optoelectronic performance results have reflected FoM where the highest value for undoped SnO<sub>2</sub> was 2.5 x 10<sup>-4</sup> Ω<sup>-1</sup> and the highest value for N-doped SnO<sub>2</sub> was 7 x 10<sup>-4</sup> Ω<sup>-1</sup>. This best FoM for N-doped SnO<sub>2</sub> is achieved due to the low resistivity of 5.9 x 10<sup>-3</sup> Ω·cm as well as the average optical transmittance of 80% in the visible range<sup>[250]</sup>. Depositing metal oxide thin films on flexible substrates has attracted a lot of research attention in the last decades due to the potential new applications as well as the flexibility in terms of fabrication process upscaling, namely roll-to-roll. Some of these flexible substrates have a good transmittance in the visible range which makes them an essential component for ideal goal of all oxide optoelectronic devices. In this context, F. Fang et al. have demonstrated N-doped SnO<sub>2</sub> thin film using reactive magnetron sputtering. SnO<sub>2</sub>:N thin film was deposited at room temperature on flexible substrate (PET) using lower process temperature. This study has explored the effect of O<sub>2</sub> partial pressure where O<sub>2</sub>:N<sub>2</sub> varies from 0% to 4% on structural, electrical and optical properties of N:SnO<sub>2</sub> to achieve low cost, eco-friendly, higher

performance flexible electronic materials. SEM imaging shows that nitrogen doping has increased the surface roughness. XRD has shown that undoped SnO<sub>2</sub> and N:SnO<sub>2</sub> thin film are amorphous which is likely due to the deposition at room temperature that is not enough to enable crystallization. However, N:SnO<sub>2</sub> is found to be crystalline when the substrate temperature is at 400°C. X-ray Photoelectron Spectroscopy has revealed that the tin oxide has not a perfect stoichiometry SnO<sub>x</sub> where x = 2. However, the stoichiometry coefficient is 1.61, 1.68, 1.70 for the following O<sub>2</sub> partial pressure 0, 2, 4%. This indicate that there is a large concentration of O vacancies which makes the thin film conductive. In this report, the co-authors stated that in rich growth conditions related to both Sn and O, Interstitial nitrogen has a larger formation energy which means that nitrogen substitution is more favourable compared to interstitial. All samples of undoped SnO<sub>2</sub> and N:SnO<sub>2</sub> have revealed an average optical transmittance above 75%. However, N:SnO<sub>2</sub> have lower optical transmittance compared to undoped SnO<sub>2</sub> which is due to the optical bandgap reduction from 3.46 to 3.19 eV as the doping of nitrogen increases from 7.5 to 19.9 at.%. The electrical resistivity decreases from 2 x 10<sup>-3</sup> to 9 x 10<sup>-4</sup> Ω·cm as the doping of nitrogen increases from 7.5 to 11 at.%, and then it increases to 0.01 Ω·cm as the doping of N increases to 19.9 at.%. The best FoM for N:SnO<sub>2</sub> is found to be slightly above 10<sup>-3</sup> Ω<sup>-1</sup> [251, 252].



**Figure 9:** Structural model for: (a) undoped SnO<sub>2</sub>, (b) substitutional (occupying O lattice site) N-doped SnO<sub>2</sub>, and (c, d) Interstitial N-doped SnO<sub>2</sub> (N in blue, Sn in yellow, O in red). Reprinted (adapted) with permission from [246]. Copyright 2015 American Chemical Society.

In general, one of the main challenges for p-type doping of SnO<sub>2</sub> is the self-compensation where the doping related p-type charge carriers formed within the SnO<sub>2</sub>, are compensated by the intrinsic n-type charge carriers formed by the host semiconducting material itself, SnO<sub>2</sub>. Another challenge of N-doping of SnO<sub>2</sub> is the compromise between a sufficient N-doping to reach a good acceptor charge carrier

concentration and a limited N-doping to avoid altering the structure of SnO<sub>2</sub> rutile and forming secondary phases mainly tin nitrides which can play a significant role in changing the electronic structure and overall electrical properties. There are some reports suggesting novel material designs and doping to tackle this issue. H.P. Dang et al. have suggested a secondary nitrogen doping for the Ga-doped SnO<sub>2</sub> material to prevent the self-compensation effect. The co-authors have reported a Ga:SnO<sub>2</sub> (GTO) thin film by reactive magnetron sputtering in N<sub>2</sub> atmosphere using Ga<sub>2</sub>O<sub>3</sub>-doped SnO<sub>2</sub> sputtering target to allow them to customize separately N-doping and Ga-doping. They prepared N-doped GTO thin films at different temperatures varying from room temperature to 500°C. All samples N-doped GTO have achieved and maintained p-type conductivity except the one deposited at room temperature, while the best p-type conductive sample was the one deposited at 400°C. This low resistivity is owed to the high p-type charge carrier concentration of  $6.99 \times 10^{19} \text{ cm}^{-3}$  [253].

### 1.2.2. Tin(II) oxide (SnO) thin film

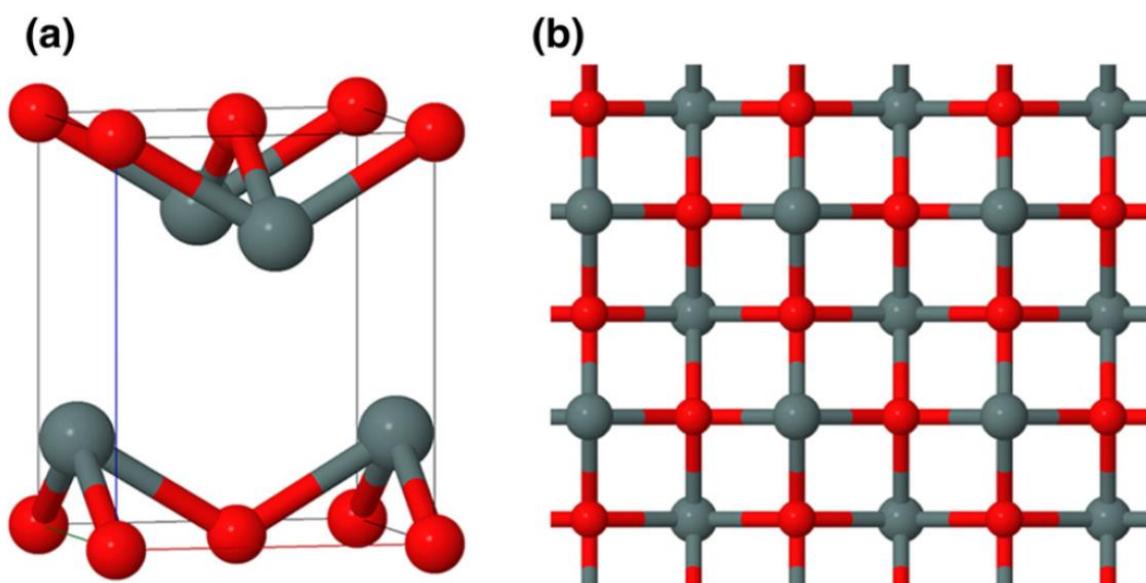
Tin(II) oxide (SnO) which is also known as stannous oxide is another form of tin oxide where the tin oxidation state is +2 and stoichiometry SnO<sub>x</sub> is x = 1. Tin monoxide thin film performance has still a long way to achieve the highest required performance, particularly in electrical properties and chemical stability. Nevertheless, tin monoxide has already many applications such as sensors<sup>[254]</sup>, transparent conducting oxide<sup>[255]</sup>, Li ion battery<sup>[256]</sup>. However, most of the reports about SnO thin films are related to Thin Film Transistor (TFT). This is due to the pressing requirement to achieve p-type oxide thin film which can be synthesised at relatively low-to-moderate temperature. The process to fabricate p-type thin film has to be scalable, and using abundant and cost-effective elements. Therefore, tin monoxide, which can fulfil all these requirements, is an attractive material to be employed in large scale industrial process. There are still some issues with chemical stability due to the oxidation towards SnO<sub>2</sub>. In some literature, the use of a capping layer or a passivation on the top of SnO is reported as a remedy to this issue<sup>[134, 135, 257-263]</sup>.

#### a) Thin film growth

Among the experimental literature, there are several synthesis techniques related to SnO whether through physical or chemical processes. Regarding the physical processes, the reactive magnetron sputtering has more reports compared to other synthesis techniques<sup>[254, 255, 258, 260, 263-270]</sup>. However, SnO synthesis is also reported via Pulsed Laser Deposition<sup>[135, 259, 262, 271, 272]</sup>, Thermal Evaporation<sup>[257, 261, 273, 274]</sup>, electron beam evaporation<sup>[134, 275-278]</sup>. In terms of chemical processes, there are several reports about SnO synthesis techniques, namely Plasma Enhanced Chemical Vapour Deposition (PECVD)<sup>[279]</sup>, Atomic Layer Deposition (ALD)<sup>[280]</sup>, hydrothermal process<sup>[281]</sup>, spray pyrolysis<sup>[282]</sup>, spin coating<sup>[283]</sup>.

### b) Structural properties

Tin(II) oxide crystal structure at room temperature and ambient pressure is a tetragonal litharge PbO type layered structure ( $\alpha$ -SnO) which is part of the tetragonal  $P4/nmm$  space group. This SnO structure is enabling four-coordinate geometry where the Sn elements have a coordination to four oxygen atoms in a pyramidal shape while forming a two-dimensional SnO sheet oriented towards (001) direction. In this four-coordinate geometry, each  $\text{Sn}^{2+}$  cation is bonded to four  $\text{O}^{2-}$  anions and vice-versa as shown in Figure 10. All bonds of O-Sn have a length around 2.23 Å.  $\text{Sn}^{2+}$  cations, while bonding to four  $\text{O}^{2-}$  anions, are forming a both corner-sharing and edge-sharing  $\text{SnO}_4$  tetrahedra. There are several theoretical or experimental reports about the SnO crystal structure, where the lattice constants vary theoretically between  $a = b = 3.79$  Å and  $c = 3.65$  Å and  $a = b = 3.90$  Å and  $c = 3.98$  Å, and are experimentally around  $a = b = 3.80$  Å and  $c = 3.84$  Å. J.P. Allen et al. have revealed using hybrid DFT calculation that while oxygen related defects do not alter the crystal structure of SnO, instead the Sn related interstitial defects ( $\text{Sn}_i$ ) induce more effect on the crystal structure by other atoms away from the  $\text{Sn}_i$  defect<sup>[284-286]</sup>.

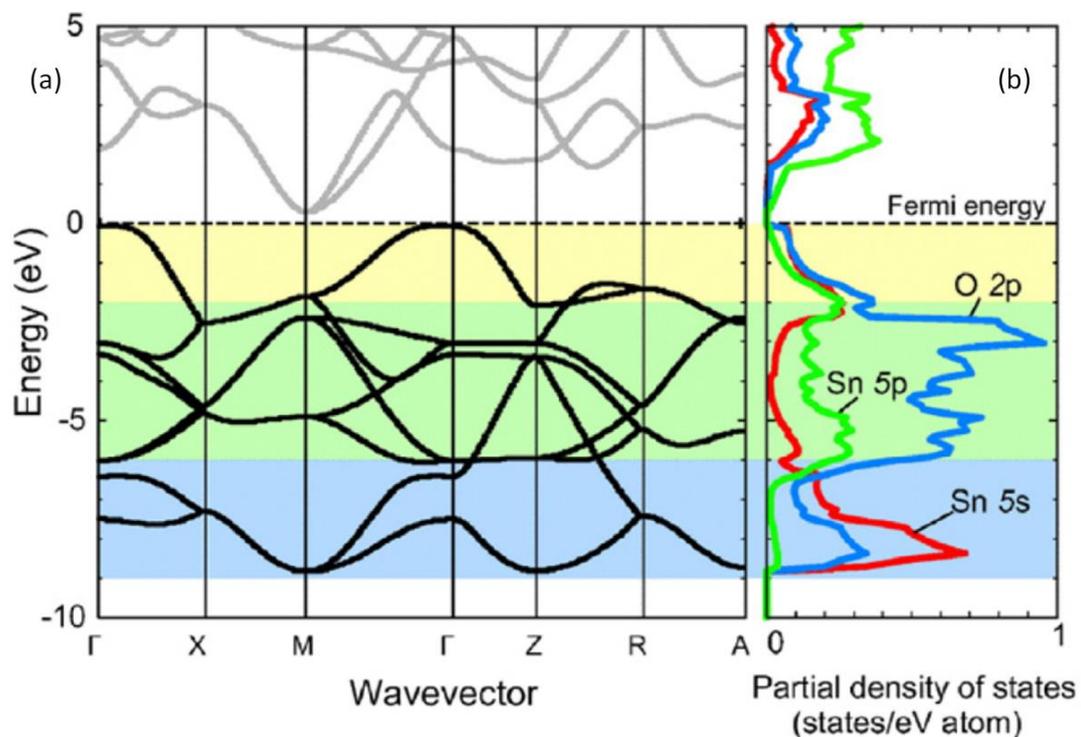


**Figure 10:** (a) SnO crystal structure and (b) SnO crystal structure top-view (gray spheres represent Sn and red spheres represent O ). Reprinted (adapted) with permission from [285]. Copyright 2016 Elsevier.

### c) Electronic structure

There are several reports about the electronic structure of SnO to describe assess the electronic bandgap as well as the contributions of electrons from Sn and O to the valence band and to the conduction band. The electronic structure shows that SnO has an indirect bandgap as the valence band maximum (VBM) and conduction band minimum (CBM) are location in a different point in the Brillouin-

zone. The VBM is located at  $\Gamma$  point while the CBM is located at M point as shown in Figure 11. The calculated indirect bandgap of SnO 0.41 eV as reported by The Materials Project. SnO, at complete oxidation, exhibits ground state  $5s^2$ . The main contribution of the top of the valence band is from hybridised O2p and Sn5s. The core of the valence band is dominated by O2p and Sn5p has a lower contribution. The bottom of the valence band is dominated by Sn5s and there is a relatively smaller contribution from O2p. On the other hand, the conduction band shows a strong contribution of Sn5p orbitals with smaller O2p and Sn5s contribution as shown in the partial Density of State (pDOS) in Figure 11<sup>[284-287]</sup>.



**Figure 11:** SnO (a) electronic band structure, (b) partial density of states. Reprinted (adapted) with permission from [285]. Copyright 2016 Elsevier.

#### d) Undoped SnO thin film

There is a limited number of experimental reports in terms of SnO synthesis compared to the modelling and theoretical literature as there is still a lot to explore about this material. The theoretical studies are important to investigate the electronic structure in order to understand the optical and electronic properties, particularly, the highly desired p-type conductivity which is intrinsic to the undoped SnO<sup>[286-299]</sup>. Among the experimental reports, U. Myeonghun et al. have reported the SnO synthesis by Plasma Enhanced Chemical Vapour Deposition (PECVD) at low temperature, the best SnO thin film was the one deposited at 150°C with a high field effect mobility of 4.86 cm<sup>2</sup>/V·s with a small acceptor related charge carrier concentration. The high electrical performance, including the mobility and lower charge carrier concentration, is suggested to be due to the fine grain microstructure<sup>[279]</sup>. M. Rockele et al.

have reported the synthesis of a p-type tin oxide via thermal evaporation at room temperature on Al<sub>2</sub>O<sub>3</sub> using SnO powder followed by rapid thermal annealing at 350°C under N<sub>2</sub> for 30 minutes. XRD results revealed that the as-deposited SnO is amorphous while the annealed SnO has secondary phases of Sn and SnO<sub>2</sub>. The co-authors could still achieve a p-channel tin monoxide TFT via thermal evaporation<sup>[261]</sup>. H.J. Kim et al. have reported the deposition of 10 nm SnO thin film using reactive RF magnetron sputtering at room temperature. They have confirmed the p-type conductivity of SnO using TFT device<sup>[260]</sup>. J.A. Caraveo-Frescas et al. have reported a high mobility of transparent p-type of pure SnO films using reactive Direct Current (DC) magnetron sputtering at room temperature with metallic Sn target and mixture of O<sub>2</sub>:Ar. XRD analysis shows the presence of SnO phase with a minor Sn secondary phase for low O<sub>2</sub> partial pressure (PP). As the O<sub>2</sub> PP increases a pure tetragonal SnO phase was achieved. For higher O<sub>2</sub> PP the SnO film can be whether amorphous or SnO<sub>2</sub> phase. Hall effect probe analysis shows that most of samples exhibit p-type conductivity while n-type conductivity is only found at very low O<sub>2</sub> PP with a low deposition pressure. The highest deposition pressure and highest O<sub>2</sub> PP has enabled a high resistive tin oxide film. The summary of the sputtering conditions is shown in table 3. It is worth noting that highest mobility and the highest conductivity are achieved at 9% O<sub>2</sub> PP and 1.8 mTorr. The sample which exhibits the highest mobility/conductivity is not a pure SnO phase as it has a minor presence of metallic β-Sn<sup>[258]</sup>.

O <sub>PP</sub> \ Dep. Pressure		1.5 mTorr	1.7 mTorr	1.8 mTorr	1.9 mTorr	2.0 mTorr
7%	Mobility (cm <sup>2</sup> /Vs)	n	9.60	14.48	12.22	7.01
	Carrier Density (cm <sup>-3</sup> )	n	2.42E+17	2.57E+17	3.12E+17	3.33E+17
	Conductivity (S/cm)	n	0.3738	0.5986	0.6133	0.3757
9%	Mobility (cm <sup>2</sup> /Vs)	10.43	13.57	18.71	8.78	4.45
	Carrier Density (cm <sup>-3</sup> )	2.29E+17	2.38E+17	2.18E+17	2.37E+17	3.33E+17
	Conductivity (S/cm)	0.3838	0.5198	0.6560	0.3350	0.2386
11%	Mobility (cm <sup>2</sup> /Vs)	9.83	10.16	15.21	4.45	2.25
	Carrier Density (cm <sup>-3</sup> )	7.55E+16	1.44E+17	7.62E+16	5.11E+16	2.31E+17
	Conductivity (S/cm)	0.1195	0.2363	0.1865	0.0366	0.0836
13%	Mobility (cm <sup>2</sup> /Vs)	9.33	7.68	5.08	3.82	1.95
	Carrier Density (cm <sup>-3</sup> )	7.62E+16	1.91E+17	4.83E+16	9.77E+16	6.88E+16
	Conductivity (S/cm)	0.1144	0.2359	0.0395	0.0601	0.0216
15%	Mobility (cm <sup>2</sup> /Vs)	4.82	4.50	4.00	X	X
	Carrier Density (cm <sup>-3</sup> )	7.83E+16	6.23E+16	5.90E+16	X	X
	Conductivity (S/cm)	0.0607	0.0451	0.0380	X	X

Pure phase SnO
n - n-type conduction
X - high resistance

**Table 3:** Hall mobility, charge carrier concentration and electrical conductivity for room temperature reactive DC magnetron sputtered SnO thin films (the reported results are average of 3 for each sample). Reprinted (adapted) with permission from [258]. Copyright 2013 American Chemical Society.

E. Eqbal and E.I. Anila have reported the synthesis and characterization of tin monoxide as transparent conducting oxide using spray pyrolysis by preparing two precursor solutions of 0.1 M and 0.25 M out of  $\text{SnCl}_2 \cdot \text{H}_2\text{O}$  in water in addition to some drops concentrated HCl. The spray pyrolysis was conducted at 350°C substrate temperature. XRD analysis revealed the presence of orthorhombic crystal structure for the prepared SnO thin films with a minor presence of  $\text{SnO}_2$  phase. The optical transmittance of SnO was found to be 82% with a wide bandgap of 3.58 eV for precursor solution of 0.1 M and 30% for 0.25 M. Hall effect probe analysis has established that both thin films exhibit p-type conductivity, and the resistivity is 0.059  $\Omega \cdot \text{cm}$  for 0.25 M precursor solution compared to 0.12  $\Omega \cdot \text{cm}$  0.1 M. Both samples have the charge carrier concentration around  $7 \times 10^{18} \text{ cm}^{-3}$ . However, the mobility is 17.1  $\text{cm}^2/\text{V} \cdot \text{s}$  for 0.25 M precursor solution compared to 7.8  $\text{cm}^2/\text{V} \cdot \text{s}$  0.1 M. This report shows that p-type conductivity can be achieved in SnO even with presence of secondary phase namely, Sn and  $\text{SnO}_2$ <sup>[282]</sup>.

#### e) SnO thin film doping

There are only few experimental reports about SnO doped thin films which is likely due to the complexity to achieve high purity undoped SnO. Among these reports, N-doped SnO was reported by Y. Kim et al. via reactive RF magnetron sputtering with metallic Sn target and mixture of  $\text{Ar}/\text{O}_2/\text{N}_2$  at 300°C on borosilicate substrates. Structural properties by XRD have revealed the presence of a polycrystalline SnO phase with preferred growth on (101) direction and the absence of any secondary phase related to metallic Sn, other tin oxides or tin nitrides. However, it is noticed that the nitrogen doping has decreased the crystallinity of SnO thin films. Scanning Electron Microscopy (SEM) imaging has revealed the microstructural properties and it showed that nitrogen doping has made the thin film grain more regular and the film more much dense. Hall effect probe analysis has revealed the electrical properties and that generating Sn-N bonds has increased the resistivity. The nitrogen doping of SnO has generated donors and contributed significantly increasing the resistivity due to self-compensation. The mobility is improved by nitrogen doping of SnO due to the increase of the density and the regular microstructure of the thin film. Vacuum annealing has increased the resistivity of both nitrogen doped and undoped SnO. UV-Visible spectroscopy has determined the optical properties and it revealed that the optical bandgap decreased due to the defects formation by nitrogen doping which formed new states within the bandgap. The bandgap decreased from 2.43 eV for undoped SnO to 2.17 eV N doped SnO. Vacuum annealing has slightly improved the optical transmittance<sup>[267]</sup>. G. Dai et al. have reported ball and sheet-like Cr-doped SnO synthesised via hydrothermal process. XRD results have shown that all samples are all closely matching to the polycrystalline tetragonal of SnO without presence of secondary phases. The optical bandgap is slightly changing between 3.1 to 3.3 eV. It was established that Cr atoms have been incorporated as substitutional to Sn ( $\text{Cr}_{\text{Sn}}$ ) and as interstitial ( $\text{Cr}_i$ )<sup>[281]</sup>. H. Hosono et al. have reported the synthesis of tin monoxide thin film doped with antimony (Sb:SnO) using Pulsed Laser Deposition (PLD) with doping level from 0 to 10 at.% (Sb/Sb+Sn) on Yttria-stabilised

ZrO<sub>2</sub> single crystal (001) substrates at temperatures ranging from 300°C to 600°C. It is worth noting that the substrate selection took into consideration the lattice parameter matching with SnO. The best film was achieved at substrate temperature of 550°C and it was determined that increasing Sb doping decreases the crystallinity of SnO. They revealed that 5 at.% Sb:SnO was not enough to switch the p-type conductivity to n-type due to the large concentration of acceptors<sup>[272]</sup>.

### **1.3. Summary**

Metal oxides are interesting materials due to their unique ability to combine both high electrical conductivity and high optical transmittance with a wide bandgap. SnO<sub>x</sub> has shown to be an excellent material that allows tuning its structural, optical and electrical properties. Literature review has shown the possibility to form tin oxide thin films with both conductivity n-type and p-type. N-type conductive SnO<sub>2</sub> thin film is well established material and reached the maturity as TCO in various applications. Nevertheless, p-type SnO<sub>x</sub> based material whether intrinsic tin monoxide SnO or doped SnO<sub>2</sub> to form a stable p-type thin film is still an interesting topic of research and development with a potential prospect for improvement. Achieving a high quality and stability p-type tin oxide based material will be an important step toward all oxide devices such as thin film transistors<sup>[300]</sup>. In the next chapter, the tuning of tin oxide based thin films will be explored to achieved both conductivity types via magnetron sputtering which is vacuum-based technique. This technique can assure a high quality and reproducibility SnO<sub>x</sub> thin films. Furthermore, the post-deposition thermal annealing will be investigated to explore the possibility of tuning the structural properties which can lead to customized optical and electrical properties.

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# **Chapter 2: Growth and Characterization of undoped tin oxide (SnO<sub>x</sub>) films**

## 2. Growth and Characterization of undoped tin oxide (SnO<sub>x</sub>) films

### 2.1. Introduction

Transparent conductive oxides (TCOs) are imperative materials in various technologies requiring optical transparency and electrical conductivity. In addition to these two characteristics, TCOs combine a third property, namely a high chemical stability. More specifically, owing to its optoelectronic properties and its production scalability, tin-doped indium oxide (ITO) is a preferred TCO for industrial applications and it has been extensively employed as transparent electrodes for various solar cell technologies, light emitting diodes and flat panel displays<sup>[1-11]</sup>. However, researchers are working on the development of alternatives materials to ITO due to the Indium (In) supply challenges in the future. Tin, instead, is a low-cost material of much higher earth abundance than Indium. Two stoichiometric tin oxide compounds, namely SnO and SnO<sub>2</sub>, are known to be wide band gap oxide semiconductors with tetragonal litharge and rutile type structures, respectively. Tin in SnO<sub>x</sub> has two chemical states of Sn<sup>2+</sup> for SnO phase and Sn<sup>4+</sup> for the SnO<sub>2</sub> one. The first is intrinsically p-type semiconductor while the second is n-type<sup>[12, 13]</sup>.

In particular, SnO<sub>x</sub> has been largely investigated in applications of gas sensors, solar cells, transparent electrodes, and thin film transistors<sup>[11, 14-20]</sup>. Moreover, in the past decades, SnO has been the key material for anode materials<sup>[21]</sup>, coatings<sup>[22]</sup>, catalysis<sup>[23]</sup>, and precursors for the production of SnO<sub>2</sub><sup>[24, 25]</sup>, because of its properties of gas-sensitivity and metastability to transform into SnO<sub>2</sub> at O<sub>2</sub>-rich ambient. Recently, SnO has been drawn back into attention mainly because of the difficulty in obtaining high-quality p-type such as doped ZnO<sup>[26, 27]</sup>, NiO<sup>[28, 29]</sup>, Cu<sub>2</sub>O<sup>[30]</sup>. Previous studies show that the maximum hole mobility of SnO films is about 2.6 cm<sup>2</sup>/V·s, fairly high among p-type conductive oxides, and it can be further improved via proper doping<sup>[19]</sup>. Those properties render SnO a promising candidate to be a next p-type oxide semiconductor for advanced optoelectronic devices. Several techniques have been used to grow SnO<sub>x</sub> films on various substrates, including reactive RF magnetron sputtering (MS)<sup>[31]</sup>, e-beam evaporation<sup>[18]</sup>, laser ablation<sup>[25]</sup>, and atomic layer deposition<sup>[32]</sup>. However, the prepared SnO<sub>x</sub> films are often mixed with some impurity phases, including metallic Sn, SnO<sub>2</sub>, and intermediate oxides containing both 2<sup>+</sup> and 4<sup>+</sup> valences<sup>[31, 33, 34]</sup>. The reason is that SnO can decompose according to the disproportionation at ion reaction even in the absence of oxygen at suitable temperature<sup>[16, 17]</sup>. Consequently, deposition conditions and growth temperature are of great importance in the fabrication of single phase SnO<sub>x</sub> films.

In this chapter, the synthesis of tin oxide thin films to achieve both phases of Sn(IV)O<sub>2</sub> and Sn(II)O is presented in detail. The SnO<sub>x</sub> synthesis is carried out by reactive magnetron sputtering via two routes. The first is to deposit SnO<sub>x</sub> by reactive sputtering of oxidized target of SnO while the second route is

the deposition by reactive sputtering of metallic target. Through these two routes, different growth conditions and post-deposition annealing effects are investigated with material characterization techniques. In addition to the sputtering target material, the growth conditions, to be varied, are substrate temperature, O<sub>2</sub>/Ar flow rate ratio, and post-deposition thermal annealing. All samples are deposited on amorphous quartz which is highly resistive and has a constant high optical transmittance in infrared, visible and a significant part of the ultraviolet to avoid any electrical or optical interference with the grown tin oxide films.

In the first part, different polycrystalline SnO<sub>x</sub> thin films were fabricated on quartz substrates by a two-step method, i.e., RF MS from high-purity SnO source target first and subsequent air and vacuum annealing treatments. The optimized conditions were clearly defined and discussed thoroughly. In the second part, different polycrystalline SnO<sub>x</sub> thin film were synthesised using one step-method of magnetron sputtering where the effect of O<sub>2</sub>/Ar flow rate ratio and substrate temperature at poor-oxygen condition were investigated. In the third part, the effect of thermal annealing on grown SnO<sub>x</sub> thin film by magnetron sputtering was explored to achieve and understand the mechanism to form a p-type conductivity. In the last part, the ageing and heating effects were assessed to investigate certain aspects of instability related to grown SnO<sub>x</sub> thin films.

In this chapter, material characterization will be conducted on all samples to explore their morphology, crystal phase, chemical composition, optical, and electrical properties of the obtained SnO<sub>x</sub> thin films were characterized by Grazing Incidence X-Ray Diffraction (GIXRD), Scanning Electron Microscopy (SEM), Transmitting Electron Microscopy (TEM), UV-Visible spectroscopy (UV-Vis), and Hall effect measurement, X-ray Photoelectron Spectroscopy (XPS), Transmission Electron Microscopy (TEM) as detailed in the experimental section. SnO<sub>x</sub> sputtered using an oxidized target (SnO)

## **2.2. Materials and Methods**

### **2.2.1. Material Synthesis**

As part of this thesis work, various deposition techniques were explored to synthesize, tune and optimize SnO<sub>x</sub> thin films, namely, hydrothermal process, spin coating, ion layer gas reaction (ILGAR) set-up, e-beam evaporator, and magnetron sputtering. While this exercise was worthwhile to be exposed to the advantages and challenges of each deposition techniques, magnetron sputtering stood out as the most effective technique as it is an affordable and scalable physical vacuum-based technique which enables the deposition of consistent, reproducible, high purity and quality films and it has several parameters of adjustment which allows the tuning of the material properties. In the following, the synthesis of SnO<sub>x</sub> films was achieved using two magnetron sputtering instruments namely, Orion 3

device from AJA International Co. and Torr International system. The parameters of magnetron sputtering for each study will be presented in detail in each section.

### 2.2.2. Material Characterization

Several characterization techniques were employed to investigate the micro-structure, crystalline structure, optical transmittance and bandgap, resistivity, charge carrier concentration and mobility, and elemental depth profiles.

GIXRD has been conducted using Rigaku – Smartlab along with ICDD – PDF crystallography data base. The X-ray source is Cu K-alpha at 1.54 Å, the  $2\theta$  scans of the detector are from 15 to 65°, the step and the speed are 0.02° and 2°/min while the X-ray incident beam was kept at 0.55°. UV-Visible Spectroscopy was performed using Perkin Elmer - Lambda 1050, using 4 nm steps. Electrical properties were conducted for all samples using benchtop four-point probe system and Hall effect - Lakeshore 8400. Depth profiles were obtained using ToF-SIMS<sup>5</sup> - IONTOF. Scanning/Transmission Electron Microscopy (SEM/TEM) images were obtained using FEI - Quanta 650/Talos, respectively. TEM lamella has been prepared by Focused Ion Beam (FIB)/SEM. The process consists of coating the SnO<sub>x</sub> film by a protective Pt layer. Then, the whole stack Pt/SnO<sub>2</sub>/Substrate undergoes the thinning process to reach a convenient vertical thickness for the TEM imaging. TOF-SIMS analysis was performed by positive Bi<sup>+</sup> ion primary beam at 30 keV and ~1.3 pA current over a 100 × 100 μm<sup>2</sup> analysed area using random rastering mode while the sputtering was achieved using Cs<sup>+</sup> ion beam at 30 keV over 400 × 400 μm<sup>2</sup>. The depth profile has been conducted in positive polarity which targets the positive ions emerging from the surface at each etched level. The chemical state analysis of all samples was conducted by X-ray Photoelectron Spectroscopy (XPS) (ESCALAB250Xi Thermo-Fisher-Scientific - UK) using Avantage software and database as well as XPS NIST database. Pass Energy for high resolution scans is 20 eV and for survey scans is 100 eV. XPS equipment was calibrated using triple high purity standards (Gold, Silver, and Copper). All samples were referenced using C1s at 284.8 eV.

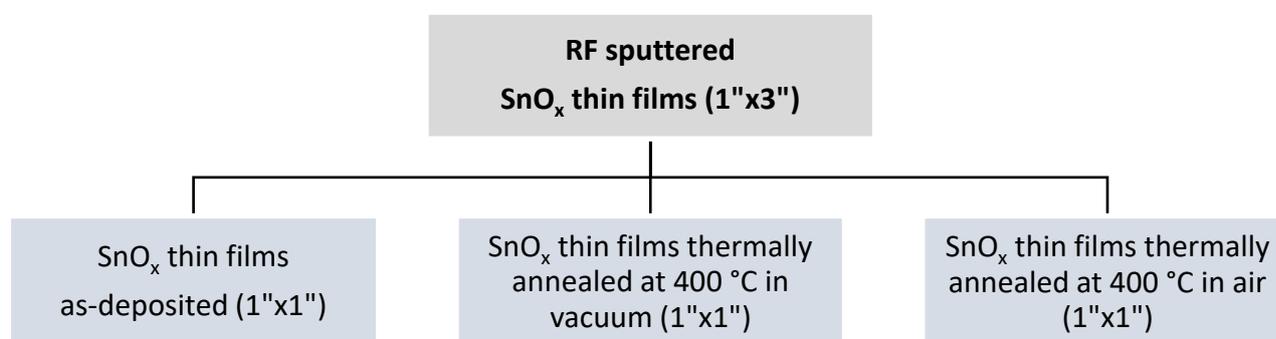
## 2.3. SnO<sub>x</sub> sputtered using an oxidized target (SnO)

### 2.3.1. SnO<sub>x</sub> thin film synthesis

The SnO<sub>x</sub> thin films were prepared in two subsequent steps. The first step consists of the thin film deposition by Radio Frequency (RF) magnetron sputtering, manufactured by Torr (described in the appendix in section 1.3.), with a high purity SnO (99.99%) 2" - diameter sputtering target (Manufactured by Codex International) on 1" × 3" cleaned quartz substrates at two different temperatures, namely 100 and 250 °C, under different O<sub>2</sub> to Ar flow rates ratios O<sub>2</sub>/Ar = 0, 0.005, 0.015, 0.025, 0.045 and 0.075. The Ar flow rate was fixed at 200 sccm. Each sample was cut into three 1" × 1" smaller samples.

The second step consists of two different post deposition annealing at 400 °C for 1 h, one in air and the other one in a controlled moderate vacuum (with a constant Ar flow of 5 sccm) at  $6.66 \times 10^{-2}$  Pa as shown in Diagram 1 and table 1.

The depositions conditions using RF magnetron sputtering were: 50 W power, 200 sccm of Ar flow and 30 min deposition time for all samples. The first batch was deposited at 100 °C, and the second batch was deposited at 250 °C. Prior to the deposition, the base pressure of  $6.66 \times 10^{-3}$  Pa was first achieved, and the deposition pressure, which depends on the variable oxygen flow rate, varied from 0.59 to 0.68 Pa for the chamber's pumping.



**Diagram 1:** Samples preparation, growth and annealing conditions.

Sample	Substrate Temp (°C)	O <sub>2</sub> /Ar (Ar:O <sub>2</sub> sccm)	Depos. Pressure (Pa)	Thickness (nm)
S0	100	0 (200:0)	0.59	100
S1		0.005 (200:1)	0.60	242
S2		0.015 (200:3)	0.61	255
S3		0.025 (200:5)	0.63	277
S4		0.045 (200:9)	0.67	199
S5		0.075 (200:15)	0.68	153
E0	250	0 (200:0)	0.59	143
E1		0.005 (200:1)	0.60	170
E2		0.015 (200:3)	0.61	149
E3		0.025 (200:5)	0.63	181
E4		0.045 (200:9)	0.65	108
E5		0.075 (200:15)	0.68	163

**Table 1:** Summary of the RF sputtering conditions and thickness of samples.

Following a first visual observations just after the SnO<sub>x</sub> deposition, all the samples looked homogenous and most of the films had the yellowish colour which is a typical characteristic of SnO<sub>x</sub> material.

### 2.3.2. Structural properties by GIXRD

For the as-deposited  $\text{SnO}_x$  at 100 °C, the samples do not show any XRD sharp peak, thereby indicating the low crystalline quality and/or the amorphous structure of this series (Figure 1a). The sample deposited without  $\text{O}_2$  flow has a broad peak at around 30° revealing the presence of a very low crystalline structure of SnO. The XRD pattern for  $\text{SnO}_x$  samples deposited in presence of  $\text{O}_2$  show a low crystalline  $\text{SnO}_2$  phase due to the shift of the broad peak from 30 to 34°. It is worth to note that the XRD pattern did not change while increasing the  $\text{O}_2/\text{Ar}$  flow rate ratio from 0.005 to 0.075, which may indicate that increasing the  $\text{O}_2$  flow does not improve the crystalline microstructure of  $\text{SnO}_x$  at this deposition temperature of 100 °C.

For  $\text{SnO}_x$  samples deposited at 100 °C and annealed at 400 °C under air, all samples show a remarkable improved crystallinity (Figure 1e). For air annealed  $\text{SnO}_x$  deposited without  $\text{O}_2$  flow, there are two highly crystalline phases of SnO and  $\text{SnO}_2$  as revealed by the presence of sharp peaks located at 29.9° indicating SnO and at 33.9° indicating  $\text{SnO}_2$ . The crystallite size related to  $\text{SnO}_2$  (110) peak drastically decreases from 116 Å to 51 Å when  $\text{O}_2/\text{Ar}$  goes from 0 to 0.005. Then, it increases slightly up to 62 Å when  $\text{O}_2/\text{Ar}$  varies from 0.005 to 0.025 and it decreases to 57 Å when  $\text{O}_2/\text{Ar}$  goes from 0.025 to 0.075. The sharp peak located at 26.6° and corresponding to the (110) plane of  $\text{SnO}_2$  deposited without  $\text{O}_2$ , represents the highest crystallite size of all the deposited  $\text{SnO}_x$  samples.

For as-deposited  $\text{SnO}_x$  at 250 °C (Figure 1b), the samples deposited at  $\text{O}_2/\text{Ar}$  ratios equal or above 0.005 have an XRD pattern which shows at least one peak, and the most intense peak located at 34° and corresponding to the (101) plane is observed for deposition at  $\text{O}_2/\text{Ar} = 0.015$ . This sample deposited at  $\text{O}_2/\text{Ar} = 0.015$  also exhibits the highest crystallite size at 56 Å for the (110) plane. The crystallite size decreases to 45 nm while increasing the  $\text{O}_2/\text{Ar}$  from 0.015 to 0.045. The crystallite size was not calculated for the other samples due to poor peak definition as shown in Table 2. It is also expected that the growth conditions in the absence of  $\text{O}_2$  flow would favour the presence of SnO due to the material of the sputtering target. Furthermore, the growth under  $\text{O}_2$  flow has formed a  $\text{SnO}_2$  phase and hence, increasing the growth temperature has clearly improved the crystalline quality of  $\text{SnO}_2$  as concluded from XRD analysis.

For  $\text{SnO}_x$  samples deposited at 100 °C and annealed at 400 °C under moderate vacuum, similarly to as-deposited samples at 100 °C without  $\text{O}_2$  flow, the XRD pattern shows a broad peak at around 30° revealing the presence of low crystalline SnO for the sample deposited without  $\text{O}_2$  flow (Figure 1c). The XRD pattern also shows a broad peak at around 34° related to the presence of low crystalline  $\text{SnO}_2$  phase for the samples deposited in the presence of  $\text{O}_2$ . For  $\text{O}_2/\text{Ar}$  above 0.005, there is a small peak

located at  $34^\circ$  and associated to the (101) plane. No noticeable change was then observed while varying the  $O_2/Ar$ .

For  $SnO_x$  samples deposited at  $250^\circ C$  and annealed at  $400^\circ C$  under moderate vacuum, similarly to previous samples deposited without  $O_2$  flow, there is a clear indication of a low crystalline  $SnO$  phase (Figure 1d). While varying the  $O_2/Ar$  from 0.005 to 0.075, there is a clear crystalline microstructure as indicated by XRD pattern. The crystallite size related to  $SnO_2$  (110) peak decreases from  $61 \text{ \AA}$  to  $56 \text{ \AA}$  when  $O_2/Ar$  goes from 0.015 to 0.025 and then it remains constant for higher  $O_2/Ar$  ratios as revealed by Table 2. It was noticed that the vacuum annealing has slightly improved the crystallinity of the  $SnO_x$  samples deposited in the presence of  $O_2$ .

For  $SnO_x$  samples deposited at  $250^\circ C$  and annealed at  $400^\circ C$  under air, all samples show an improved crystallinity compared to as-deposited  $SnO_x$  (Figure 1f). For air annealed  $SnO_x$  deposited without  $O_2$  flow, there are also two high crystalline phases of  $SnO$  and  $SnO_2$  as revealed by the presence of sharp peaks located at  $29.9^\circ$  indicating  $SnO$  and at  $33.9^\circ$  indicating  $SnO_2$ . The crystallite size of  $SnO_2$  drastically decreases from  $109 \text{ \AA}$  to  $54 \text{ \AA}$  when  $O_2/Ar$  goes from 0 to 0.005. Then, it slightly stabilizes around the value of  $53 \text{ \AA}$  when  $O_2/Ar$  goes from 0.005 to 0.075. Air annealing has clearly improved the crystallinity of  $SnO_x$  samples which were deposited at  $100^\circ C$  compared to vacuum annealing process. It is worth to note that  $SnO$  peaks were only observed in both air annealed samples deposited without  $O_2$  flow.

To sum-up, the as-deposited samples have shown that the relatively high temperature of  $250^\circ C$  (i.e. as compared to  $100^\circ C$ ) has improved the crystalline microstructure for all the samples deposited with the presence of  $O_2$ . This is due to the improved crystallization of  $SnO_x$  and the reduction of the microstructure disorder<sup>[31]</sup>. Vacuum annealing has then slightly improved the crystallinity for samples deposited at  $100^\circ C$  due to the higher annealing temperature of  $400^\circ C$ . However, there is no clear improvement of the crystalline microstructure of the samples deposited at  $250^\circ C$ . The air annealing has remarkably improved the  $SnO_x$  microstructure for all samples deposited at  $100^\circ C$  due to the higher annealing temperature of  $400^\circ C$ , compared to the deposition temperature, and to the abundant presence of  $O_2$  which enabled the crystallization of  $SnO_x$ . Air annealing for  $SnO_x$  sample deposited at  $250^\circ C$  without  $O_2$  has substantially improved the crystallinity of  $SnO_x$  due to the presence of  $O_2$  at higher temperature of  $400^\circ C$ . Moreover, further crystallization was less extensive under air annealing for other samples deposited in presence of  $O_2$ . This is likely due to the small temperature gradient between the deposition and the annealing processes, as well as the presence of  $O_2$  in both processes.  $SnO$  crystallinity was neither achieved when depositing  $SnO$  without  $O_2$  at both temperatures of  $100$  and  $250^\circ C$ , nor after annealing these samples under vacuum. However,  $SnO$  phase appeared in both samples deposited at  $100$  and  $250^\circ C$  without  $O_2$  and annealed under air. This indicates that  $SnO$  crystallinity may be improved using a controlled annealing process under inert gas at atmospheric

pressure and in absence of O<sub>2</sub> which can lead to crystallization of SnO without any further oxidation to SnO<sub>2</sub> phase<sup>[31, 32, 35-38]</sup>. Experimental work related to SnO<sub>x</sub> annealing under inert gas will be discussed later in this chapter. The different observations extracted from XRD are summarized in Table 3.

O <sub>2</sub> /Ar ratio	Crystallite Size (Å)					
	Deposition Temperature (°C)					
	100	250	100	250	100	250
	As deposited		Vacuum Annealing at 400 °C		Air Annealing at 400 °C	
0	*	*	*	*	116	109
0.005	*	*	*	*	51	54
0.015	*	56	*	61	56	52
0.025	*	54	*	56	62	53
0.045	*	45	*	56	61	55
0.075	*	*	*	56	57	55

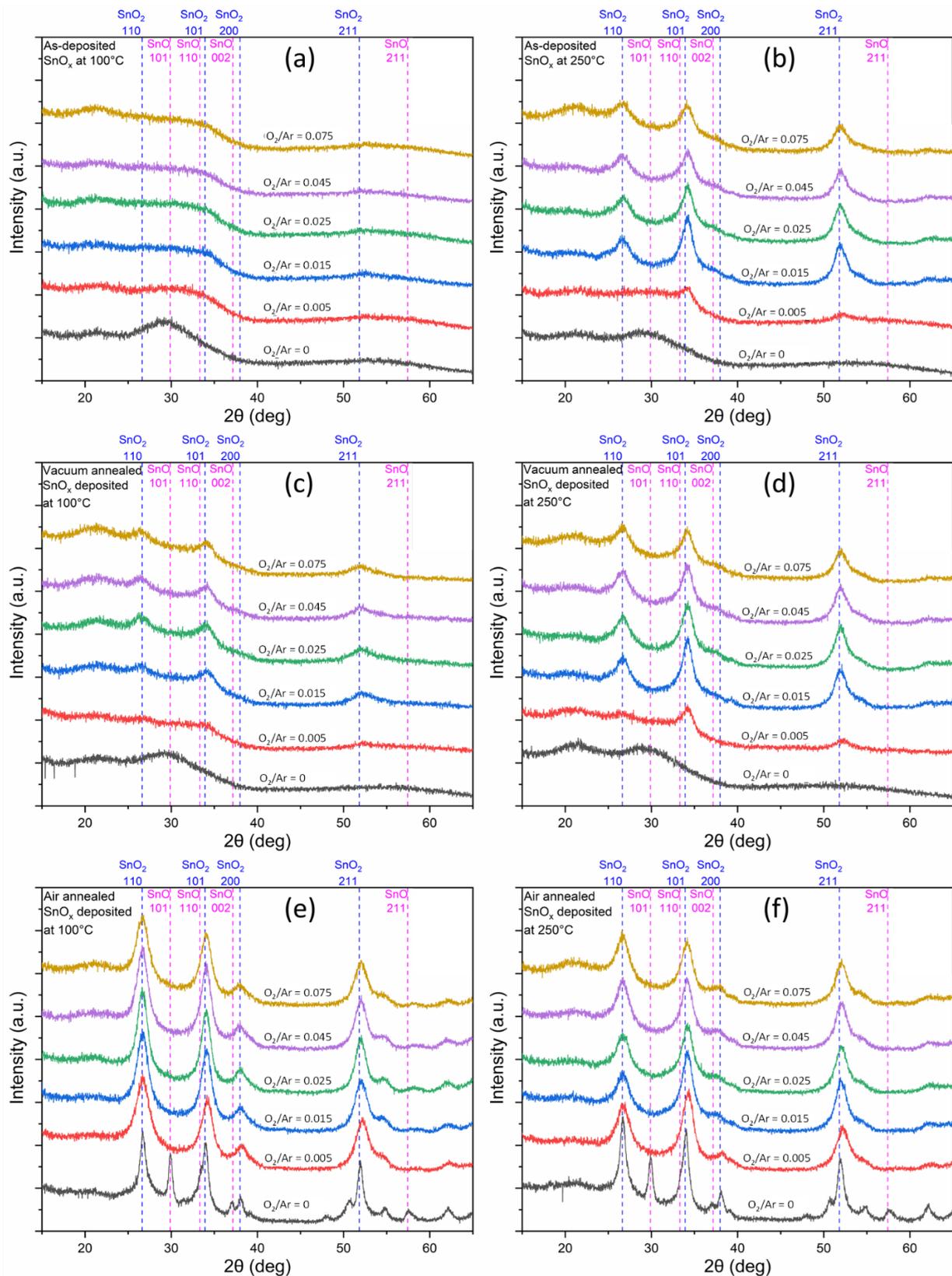
**Table 2:** Crystallite size for SnO<sub>x</sub> sample using Scherrer equation for the peak SnO<sub>2</sub> (110).

O <sub>2</sub> /Ar ratio	Deposition Temperature (°C)					
	100	250	100	250	100	250
	As deposited		Vacuum Annealing at 400 °C		Air Annealing at 400 °C	
0	(SnO <sub>x</sub> ) Low crystallinity	* Higher crystallinity	Higher crystallinity			
0.005	Low crystallinity	Poorly crystallized	Poorly crystallized	Lower crystallinity	Lower crystallinity	Average crystallinity
0.015		Average crystallinity		+ High crystallinity	↓	
0.025		↓		↓	High crystallinity	
0.045		↓		↓	↓	
0.075		Lower crystallinity		Lower crystallinity	Lower crystallinity	

**Table 3:** Summary of SnO<sub>x</sub> crystallinity samples using XRD and Scherrer equation for peak (101).

+: best conductivity of SnO<sub>x</sub> deposited at 250 °C and 0.015 O<sub>2</sub>/Ar ratio followed by vacuum annealing.

\*:best crystallinity of SnO<sub>x</sub> deposited at 100 °C and without O<sub>2</sub> flow followed by air annealing.



**Figure 1:** GIXRD patterns of SnO<sub>x</sub> samples (a) as deposited (as-dep) at 100 °C, (b) as-dep at 250 °C, (c) as-dep at 100°C and annealed in vacuum at 400 °C, (d) as-dep at 250°C and annealed in vacuum at 400 °C, (e) as-dep at 100 °C and annealed in air at 400 °C, and (f) as-dep at 250 °C and annealed in air at 400 °C.

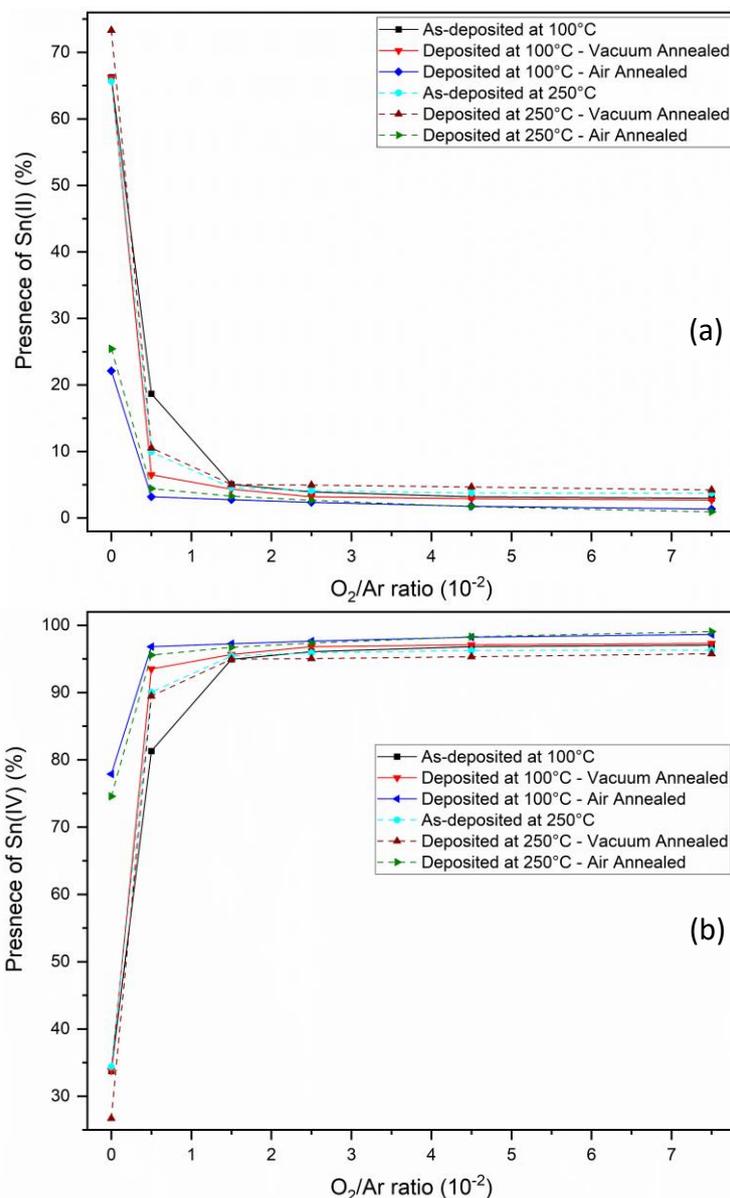
### 2.3.3. Surface chemical state analysis by XPS

Chemical state analysis was performed for all 36 SnO<sub>x</sub> thin film samples using X-ray photoelectron spectroscopy which is a surface sensitive technique analysing the surface of around 5 nm thickness. As expected, the surface of thin films gets oxidized at the surface once exposed to the air. However, this technique can provide an important indication of the present chemical states of the overall film. As expected, all XPS spectra for Sn3d<sub>5/2</sub> show only the presence of the two chemical states Sn(II) and Sn(IV). The peak position of Sn3d<sub>5/2</sub> related to Sn(II) is around 486.5 (+/- 0.15) eV while the one for Sn(IV) is around 486.9 (+/- 0.15) as shown in Figures 2, 3. It is worth noting that all samples were subject to low energy Ar ion cluster cleaning where the adventitious carbon was reduced drastically which improved the Sn and O from lattice and defect oxide. XPS equipment was calibrated using triple high purity metals Au/Ag/Cu and the XPS spectra were referenced using C 1s main peak at 284.8 eV.

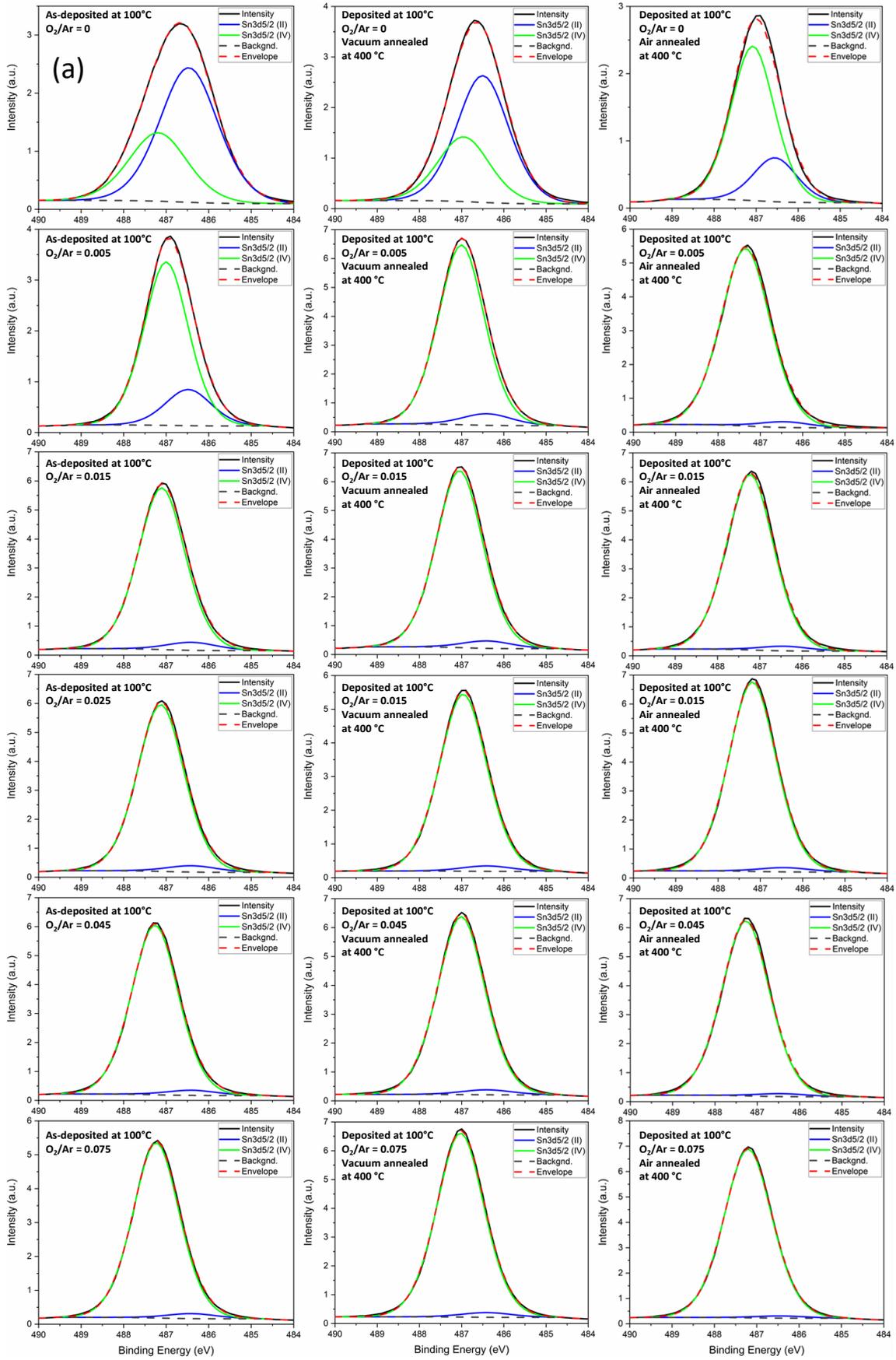
For SnO<sub>x</sub> deposited at 100 °C, the chemical state analysis revealed that the presence of Sn(II) in SnO<sub>x</sub> deposited without O<sub>2</sub> flow is around 66.3% while Sn(IV) is around 33.7%. The presence of Sn(II) decreases drastically once we introduce O<sub>2</sub> flow and it reaches 18.7% for O<sub>2</sub>/Ar ratio of 0.005 while Sn(IV) increases to 81.3%. Sn(II) presence decreases to 5.1% for O<sub>2</sub>/Ar of 0.015 while Sn(IV) increases to 94.9%. Sn(II) continues to decrease while increasing the O<sub>2</sub>/Ar flow rate ratio till it reaches 3% while O<sub>2</sub>/Ar ratio reaches 0.075. Sn(IV) increases and reaches 97% for O<sub>2</sub>/Ar ratio of 0.075. This decrease in Sn(II) (increase in Sn(IV)) was expected due to the increase of the oxidation state while introducing and increasing the O<sub>2</sub> flow rate. Vacuum annealing has kept the Sn(II) presence very similar to the as-deposited samples. However, the Sn(II) percentages slightly decreases except for O<sub>2</sub>/Ar ratio of 0.005 where vacuum annealing led to a significant decrease 6.5% compared to 18.7% for as-deposited. The vacuum annealing at 400°C has proven to slightly improve the crystallinity for SnO<sub>x</sub> deposited in presence of O<sub>2</sub> and favours more SnO<sub>2</sub> phase as revealed by GIXRD which explains the higher percentage of Sn(IV). Air annealing has led to a substantial increase of Sn(IV) for SnO<sub>x</sub> deposited without O<sub>2</sub> flow and for O<sub>2</sub>/Ar ratio of 0.005 compared to as-deposited SnO<sub>x</sub> films while all other SnO<sub>x</sub> were subject to a slight increase. Air annealing has led to a high crystalline SnO<sub>2</sub> phase as shown by GIXRD. These results are matching the structural findings discussed in the previous section.

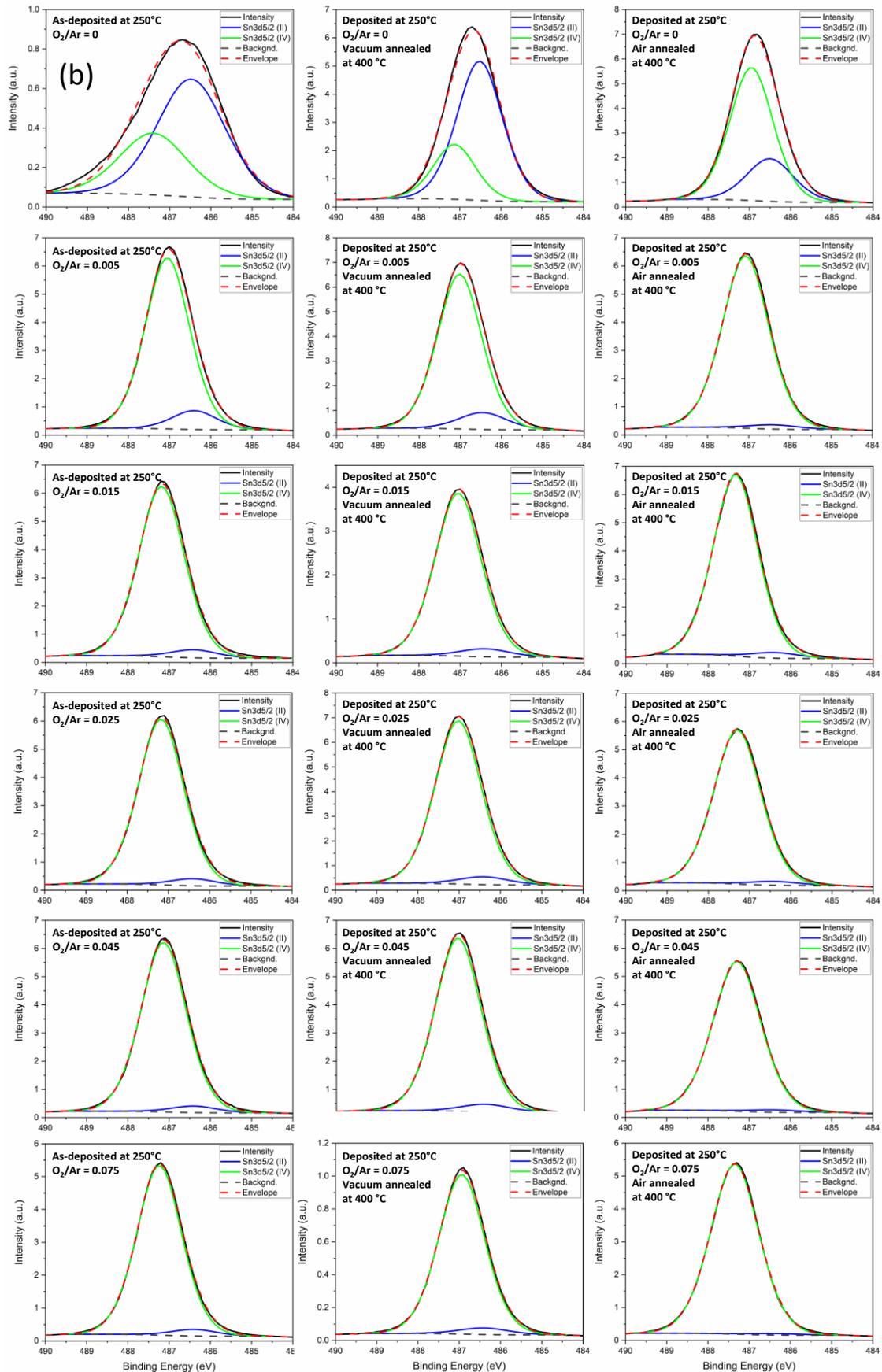
For SnO<sub>x</sub> deposited at 250 °C, the chemical state analysis revealed that the presence of Sn(II) in SnO<sub>x</sub> deposited without O<sub>2</sub> flow is around 65.6% while Sn(IV) is around 34.3%. The presence of Sn(II) decreases drastically once we introduce O<sub>2</sub> flow and it reaches 10% for O<sub>2</sub>/Ar ratio of 0.005 while Sn(IV) increases to 90%. Sn(II) presence decreases to 5.5% for O<sub>2</sub>/Ar ratio of 0.015 while Sn(IV) increases to 94.5%. Sn(II) continue to decrease while increasing the O<sub>2</sub>/Ar flow rate ratio till it reaches 3.7% while O<sub>2</sub>/Ar ratio reaches 0.075. Sn(IV) increases and reaches 96.3% for O<sub>2</sub>/Ar ratio of 0.075. Similar to the previous SnO<sub>x</sub> series, this decrease in Sn(II) (increases in Sn(IV)) was expected due to the increase of

the oxidation state while introducing and increasing the O<sub>2</sub> flow rate. Vacuum annealing has kept the Sn(II) presence trend very similar to the as-deposited samples. However, the Sn(II) percentages slightly increases which might be due to vacuum annealing and the small temperature gradient. The vacuum annealing at 400 °C has proven to slightly improve the presence of Sn(II) phase.



**Figure 2:** XPS chemical state (a) Sn(II) and (b) Sn(IV) quantification using Sn3d<sub>5/2</sub> spectra fitting for SnO<sub>x</sub> as a function of O<sub>2</sub>/Ar ratio and deposition temperature in different annealing conditions.





**Figure 3:** XPS spectra for Sn $3d_{5/2}$  for (a) SnO $_x$  deposited at 100°C as a function of O $_2$ :Ar in different annealing conditions, (b) SnO $_x$  deposited at 250 °C as a function of O $_2$ :Ar in different annealing conditions.

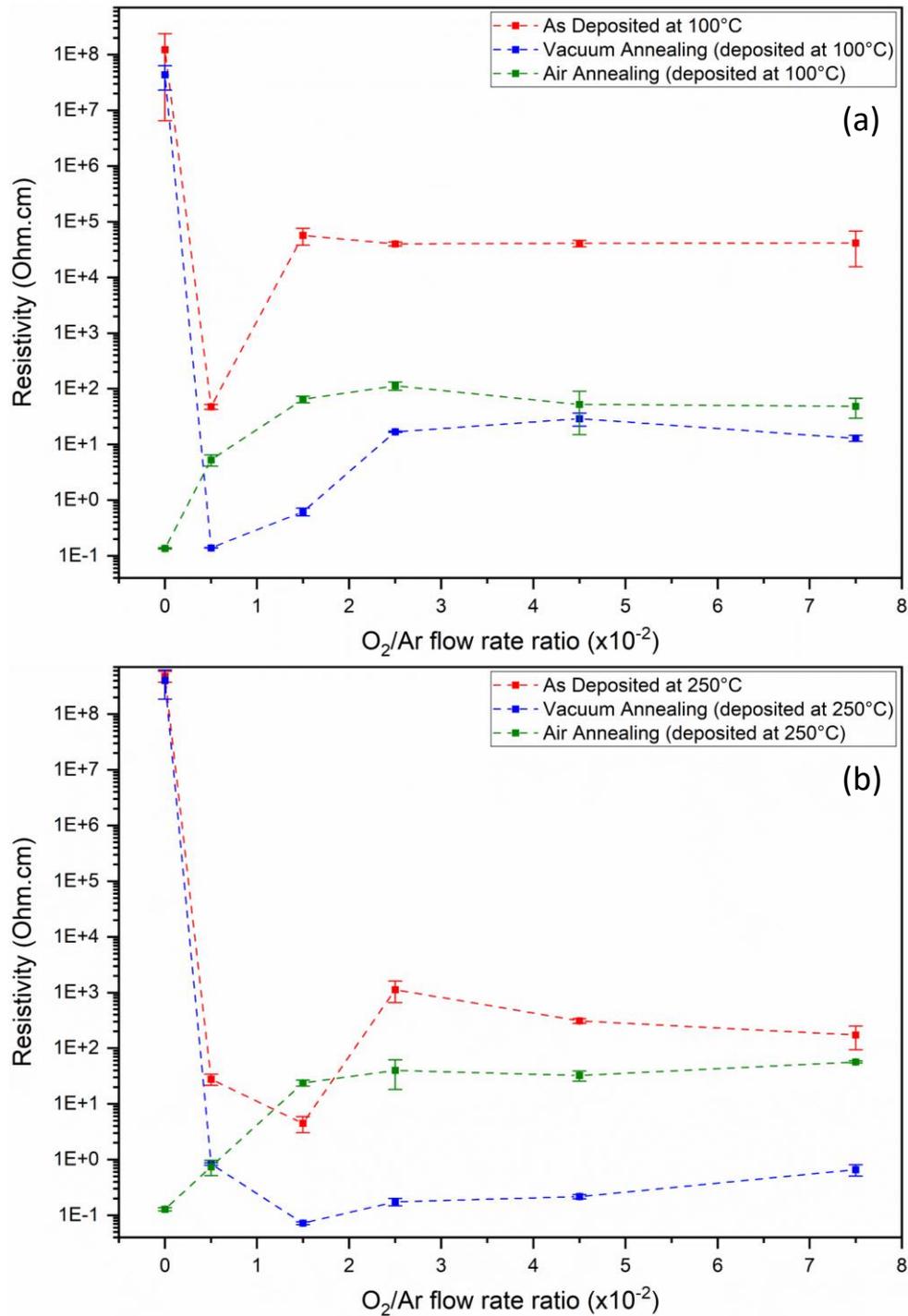
#### 2.3.4. Electrical transport properties

The electrical properties were investigated initially using the four-point probe IV tool then through the Hall effect measurements. Four-point probe measurements were performed on all samples using three-point statistics method. The obtained results revealed a very high resistivity around  $10^8 \Omega \cdot \text{cm}$  for  $\text{SnO}_x$  deposited without the presence of  $\text{O}_2$  for both as-deposited series at  $100^\circ\text{C}$  and  $250^\circ\text{C}$ . As revealed by XRD results, this may be attributed to the low crystallinity of the  $\text{SnO}_x$  films. The resistivity decreases to its lowest value of  $47 \Omega \cdot \text{cm}$  at  $\text{O}_2/\text{Ar} = 0.005$  for  $\text{SnO}_x$  samples deposited at  $100^\circ\text{C}$  and to its lowest value of  $4.5 \Omega \cdot \text{cm}$  at  $\text{O}_2/\text{Ar} = 0.015$  for  $\text{SnO}_x$  samples deposited at  $250^\circ\text{C}$ . The lower value of resistivity for samples deposited at  $250^\circ\text{C}$  is matching with the highest crystallite size which indicated as discussed earlier the improvement of the crystalline microstructure.

For  $\text{SnO}_x$  deposited at  $100^\circ\text{C}$ , the resistivity increases significantly from  $\text{O}_2/\text{Ar} = 0.005$  to  $0.015$  and it relatively stabilizes at higher  $\text{O}_2/\text{Ar}$  starting from  $0.015$  (Figure 2). For  $\text{SnO}_x$  as deposited at  $250^\circ\text{C}$ , the same behaviour is observed, where the resistivity increases significantly from  $\text{O}_2/\text{Ar} = 0.015$  to  $0.025$  and it relatively decreases towards higher  $\text{O}_2/\text{Ar}$  starting from  $0.025$ . In the presence of  $\text{O}_2$ , all deposited  $\text{SnO}_x$  samples at  $250^\circ\text{C}$  show a better electrical conductivity compared to films grown at  $100^\circ\text{C}$ .

Vacuum annealed samples were found to follow similar pattern as the as-deposited  $\text{SnO}_x$  at both deposited temperatures where the best crystalline as-deposited  $\text{SnO}_x$  samples are still the most conductive ones. Moreover, the lowest recorded resistivity values are  $0.14 \Omega \cdot \text{cm}$  for the best  $\text{SnO}_x$  deposited at  $100^\circ\text{C}$  and  $0.07 \Omega \cdot \text{cm}$  for the best  $\text{SnO}_x$  deposited at  $250^\circ\text{C}$  of all the fabricated batches after vacuum annealing. Furthermore, as expected after annealing process, the best recorded resistivity values are significantly lower compared to the as-deposited ones. This improvement is likely attributed to the microstructure improvement leading to the conservation of charge carrier after the vacuum annealing of the samples.

For both air annealed series, the lowest resistivity of around  $0.13 \Omega \cdot \text{cm}$  was observed for samples deposited without  $\text{O}_2$  flow. The resistivity increases then significantly until  $\text{O}_2/\text{Ar} = 0.015$  and slightly stabilizes for higher  $\text{O}_2/\text{Ar}$  starting from  $0.015$ . The trend of resistivity with respect to  $\text{O}_2/\text{Ar}$  (of both air-annealed series) are very closely matching regardless the deposition temperature of  $\text{SnO}_x$ . However, the resistivity for  $\text{SnO}_x$  samples deposited at  $250^\circ\text{C}$  is slightly lower for all  $\text{O}_2/\text{Ar}$  ratios except the two extreme values of  $0$  and  $0.075$  as shown in Figure 4.



**Figure 4:** Resistivity as a function of O<sub>2</sub>/Ar flow rate ratio in growth conditions for as-deposited, vacuum-annealed, and air-annealed SnO<sub>x</sub> (a) deposited at 100°C and (b) deposited at 250°C.

The samples deposited at 250 °C followed by a moderate vacuum annealing at 400 °C showed the lowest resistivity of 0.07 Ω·cm at O<sub>2</sub>/Ar = 0.015. Furthermore, the other SnO<sub>x</sub> samples in the same series deposited with O<sub>2</sub>/Ar ratio above 0.015 show substantially a lower resistivity compared to samples from other series with the same conditions of O<sub>2</sub>/Ar. This can be attributed to the relatively higher crystallite size due to the annealing process, as well as the expected charge carrier concentrations due to the vacuum annealing [31, 37, 38].

Further analysis using Hall effect measurement were conducted solely on the best conductive samples identified by four-point probe method, both for vacuum and air annealed samples series. The electron mobility for the best conductive samples for each series were  $1.07 \text{ cm}^2/\text{V}\cdot\text{s}$  for vacuum annealed series deposited at  $100^\circ\text{C}$ ,  $7.77 \text{ cm}^2/\text{V}\cdot\text{s}$  for vacuum annealed series deposited at  $250^\circ\text{C}$ ,  $2.11 \text{ cm}^2/\text{V}\cdot\text{s}$  for air annealed series deposited at  $100^\circ\text{C}$ , and  $2.58 \text{ cm}^2/\text{V}\cdot\text{s}$  for air annealed series deposited at  $250^\circ\text{C}$ . Their respective charge carrier concentrations were  $1.47 \times 10^{19} \text{ cm}^{-3}$  for vacuum annealed series deposited at  $100^\circ\text{C}$ ,  $5.84 \times 10^{18} \text{ cm}^{-3}$  for vacuum annealed series deposited  $250^\circ\text{C}$ ,  $1.39 \times 10^{19} \text{ cm}^{-3}$  for air annealed series deposited at  $100^\circ\text{C}$ , and  $1.26 \times 10^{19} \text{ cm}^{-3}$  for air annealed series deposited at  $250^\circ\text{C}$  as shown in Table 4. All these conductive samples were n-type semiconductors which is expected for vacuum annealed samples due to the presence of  $\text{SnO}_2$  phase.

Series' best conductive sample	O <sub>2</sub> /Ar ratio	Mobility (cm <sup>2</sup> /V·s)	electron carrier (cm <sup>-3</sup> )
As deposited at 100 °C / Vacuum annealed	0.005	1.07	$1.47 \times 10^{19}$
As deposited at 250 °C / Vacuum annealed	0.015	7.77	$5.84 \times 10^{18}$
As deposited at 100 °C / Air annealed	0	2.11	$1.39 \times 10^{19}$
As deposited at 250 °C / Air annealed	0	2.58	$1.26 \times 10^{19}$

**Table 4:** Hall effect measurements for  $\text{SnO}_x$  samples of each annealed series.

However, the n-type conductivity for air annealed samples reveals that the majority charge carriers are related to the  $\text{SnO}_2$  phase. The best conductive  $\text{SnO}_x$ , which is deposited at  $250^\circ\text{C}$  /  $\text{O}_2/\text{Ar} = 0.015$  and annealed under vacuum at  $400^\circ\text{C}$  is a result of a much higher mobility and an average charge carrier concentration compared to other samples.

Normally, scattering mechanisms are the main explanation for electron mobility. Furthermore, as per the grain-boundary scattering mechanism, the mobility increases while increasing the carrier concentration or the crystallite size. However, for higher values than  $10^{20} \text{ cm}^{-3}$  of charge carrier concentration, the mobility decreases due to the domination of ionized scattering mechanism [36, 39, 40]. Therefore, the grain-boundary scattering mechanism is likely the main mechanism responsible for the higher electron mobility of the best conductive  $\text{SnO}_x$  sample. As a matter of fact, Kim et al. [41] established that the grain boundary scattering was the dominant scattering mechanism for  $\text{SnO}_{2-x}$  thin films prepared by magnetron sputtering [41]. In another work performed on polycrystalline Gallium doped Zinc Oxide (GZO), Hall mobility measurements indicated that the mobility of electron carriers across many grains and grain boundaries within conduction path was limited by both scattering effects intra-grains and at grain boundaries [42]. The dominance of the scattering effects varies with electron concentration. In the case of higher electron concentration above about  $10^{20} \text{ cm}^{-3}$ , the dominant

scattering effect for  $\mu_{\text{Hall}}$  has been considered to be ingrain scattering (ionized-impurity scattering)<sup>[43, 44]</sup>. On the other hand, it is well established that  $\mu_{\text{opt}}$  shows the electron mobility in ingrain. Thus,  $\mu_{\text{opt}}$  is limited by the effect of ingrain scattering. In advanced investigations, comparing electron and optical mobilities ( $\mu_{\text{Hall}}$  and  $\mu_{\text{opt}}$ ) has been employed to highlight the contribution of grain boundary scattering on electron mobility, and this approach has frequently been used as a means to study the effect of ingrain and grain boundary scattering on electron transport properties in many TCOs<sup>[45-50]</sup>.

### 2.3.5. Optical properties

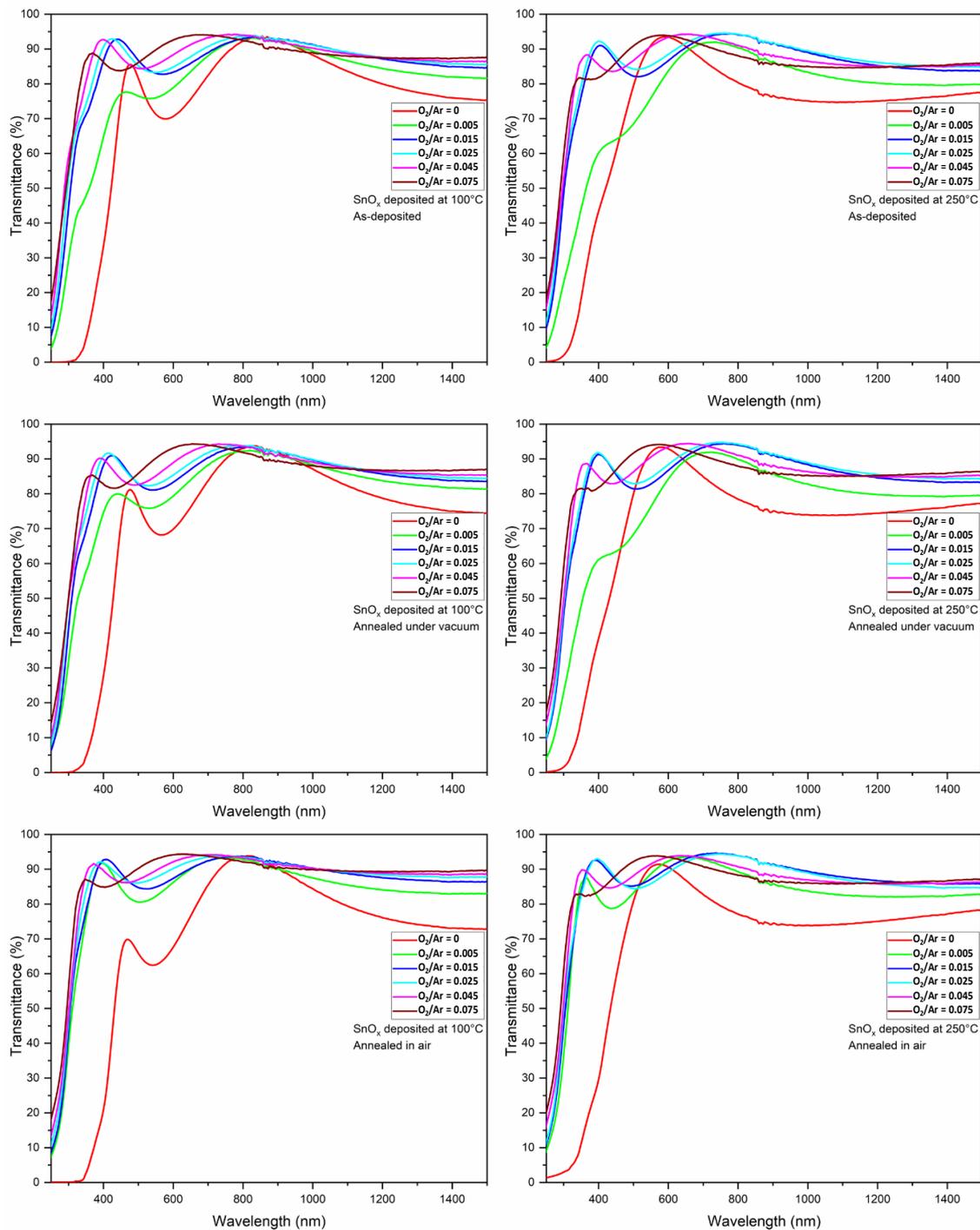
The optical properties have been studied using UV-Visible spectroscopy. The optical transmittance measurements (Figure 4) were conducted on all the grown samples and the average optical transmittance between 400 and 700 nm is summarized in Table 5.

O <sub>2</sub> /Ar ratio	Optical transmittance (%)					
	Deposition Temperature (°C)					
	100	250	100	250	100	250
	As-deposited		Vacuum Annealing at 400 °C		Air Annealing at 400 °C	
0	72	80	70	79	65	77
0.005	78	76	80	76	86	87
0.015	87	87	86	87	88	89
0.025	87	88	87	88	89	89
0.045	88	90	87	89	90	90
0.075	89	91	89	91	91	91

**Table 5:** Average optical transmittance between 400 and 700 nm for SnO<sub>x</sub> samples with substrate.

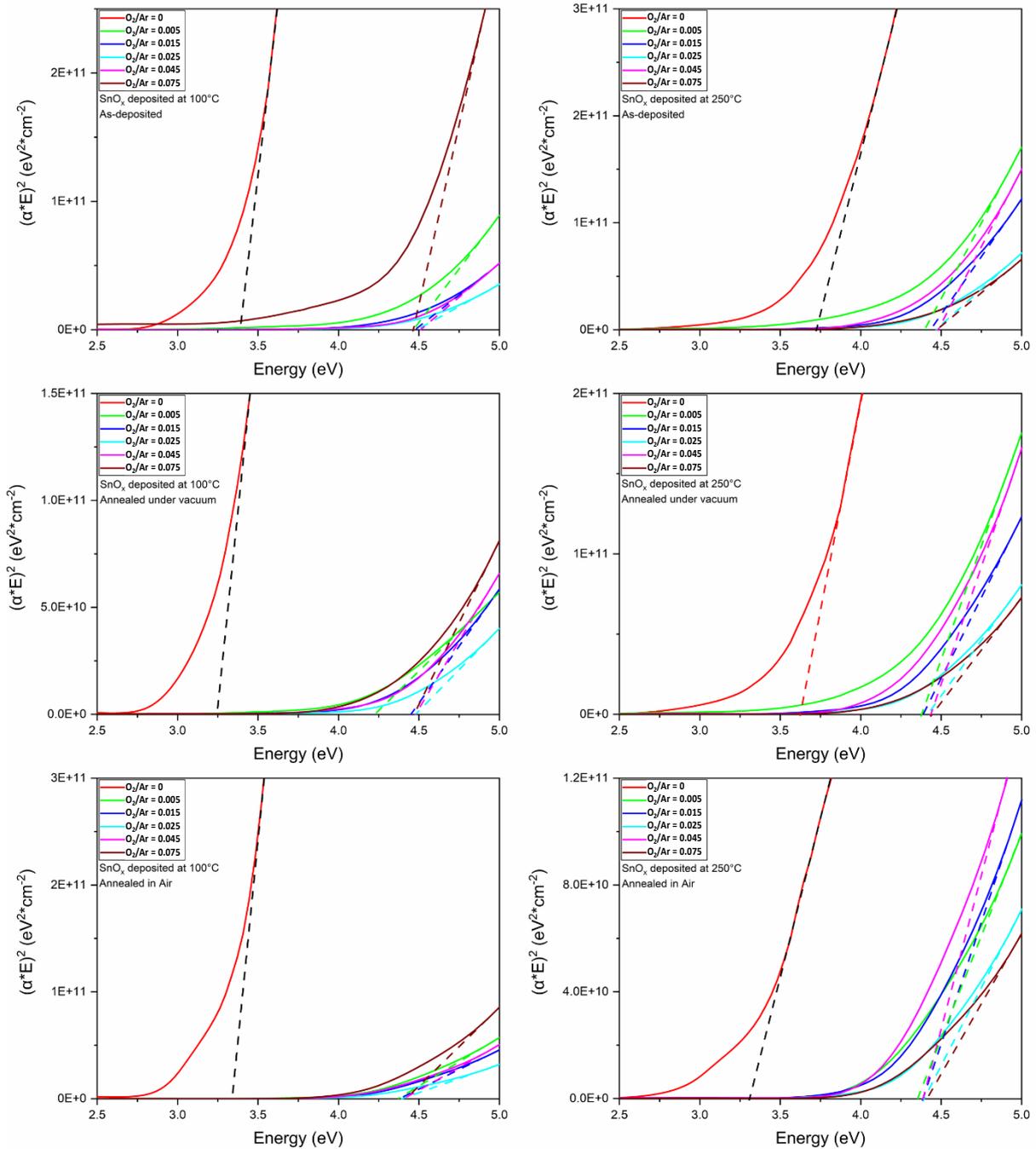
The average optical transmittance of the reference Quartz substrate was initially measured at 93%. For SnO<sub>x</sub> samples deposited at 100 °C, it is observed that the average optical transmittance from 400 to 700 nm (labelled transmittance) increases from 72 up to 89% when O<sub>2</sub>/Ar increases from 0 to 0.075. Moreover, the transmittance increases from 70 to 89% after vacuum annealing and increases from 65 % to 91 % when O<sub>2</sub>/Ar varies from 0 to 0.075 after air annealing. For SnO<sub>x</sub> samples deposited at 250 °C, the transmittance in as-deposited and vacuum annealed samples decreases from around 80% to 76% when O<sub>2</sub>/Ar varies from 0 to 0.005. Furthermore, the transmittance increases up to 91% for O<sub>2</sub>/Ar increasing from 0.005 to 0.075. The transmittance increases from 77 % to 91 % for O<sub>2</sub>/Ar varying from 0 to 0.075 after air annealing. These results are summarized in table 5.

It is established that the oxygen flow has a strong effect on the optical properties of SnO<sub>x</sub><sup>[36]</sup>. This can be directly observed by the blue shift of the absorption edge as well as the relative increase of transmittance when O<sub>2</sub>/Ar is increasing as shown in Figure 5.



**Figure 5:** UV-Visible transmittance as a function of the wavelength for all the measured SnO<sub>x</sub> samples.

Figure 6 shows Tauc plots for all samples considering the SnO<sub>x</sub> thin films as a direct bandgap semiconductor. It can be concluded from Tauc plots that all the SnO<sub>x</sub> films have a wide optical band gap varying from 3.3 to 4.5 eV.



**Figure 6:** Tauc plots for all the measured SnO<sub>x</sub> samples showing the respective optical bandgaps.

Table 6 summarised the values of the optical band gap for all the samples, which are found to increase and then stabilise when O<sub>2</sub>/Ar is increasing. The low optical bandgap for the SnO<sub>x</sub> samples deposited without the presence of O<sub>2</sub> compared to the other samples, is a typical characteristic for SnO and it was already reported in the literature varying from 2.6 to 3.4 eV [36]. These results corroborate well the XRD analysis and are confirming the major presence and the low crystallinity behaviour of the SnO phase for both the as-deposited and vacuum annealed samples. The measured bandgap above 4 eV for the rest of the samples is related to SnO<sub>2</sub> phase which is very close to the values reported in the literature. The band gap variation is mainly related to the difference of the stoichiometry of the SnO<sub>x</sub>

films. However, the effect of the sample disorder can also decrease the bandgap of the SnO<sub>x</sub> thin films<sup>[35, 36, 51]</sup>.

The band gap of SnO<sub>2</sub> samples prepared in presence of O<sub>2</sub> is relatively very high compared to the literature where all samples reached 4.4~4.5 eV except one sample of SnO<sub>2</sub> deposited at 100 °C / 0.005 O<sub>2</sub>/Ar and annealed under vacuum. These band gap values are related to the high oxidation of SnO<sub>2</sub> due to the oxidized sputtering target and the presence of oxygen during the growth.

O <sub>2</sub> /Ar ratio	Optical Bandgap (eV)					
	SnO <sub>x</sub> Deposition Temperature (°C)					
	100	250	100	250	100	250
	As-deposited		Vacuum Annealing at 400 °C		Air Annealing at 400 °C	
0	3.4	3.7	3.3	3.6	3.3	3.3
0.005	4.5	4.4	4.2	4.4	4.4	4.4
0.015			4.5			
0.025						
0.045						
0.075						

**Table 6:** Optical bandgap for SnO<sub>x</sub> samples.

### 2.3.6. Optoelectronic performance

The optoelectronic performance for all SnO<sub>x</sub> samples was evaluated using Haacke figure of merit (FoM) equation (1)<sup>[52]</sup>. The results are shown in Table 7. SnO<sub>x</sub> samples have a relatively low figure of merit due to the moderate resistivity of the thin films. The highest figure of merit  $\phi$  of  $5.14 \times 10^{-2}$  ( $10^{-3} \Omega^{-1}$ ) is related to our best conductive SnO<sub>x</sub> sample. Furthermore, the SnO<sub>x</sub> samples deposited at 250 °C where the O<sub>2</sub>/Ar ratio between 0.015 and 0.045 as well as SnO<sub>x</sub> sample deposited at 100 °C and at 0.005 O<sub>2</sub>/Ar ratio have shown figure of merit above  $10^{-2}$  ( $\times 10^{-3} \Omega^{-1}$ ). These results reveal that the vacuum annealing has improved the optoelectronic performance of the SnO<sub>x</sub> thin films.

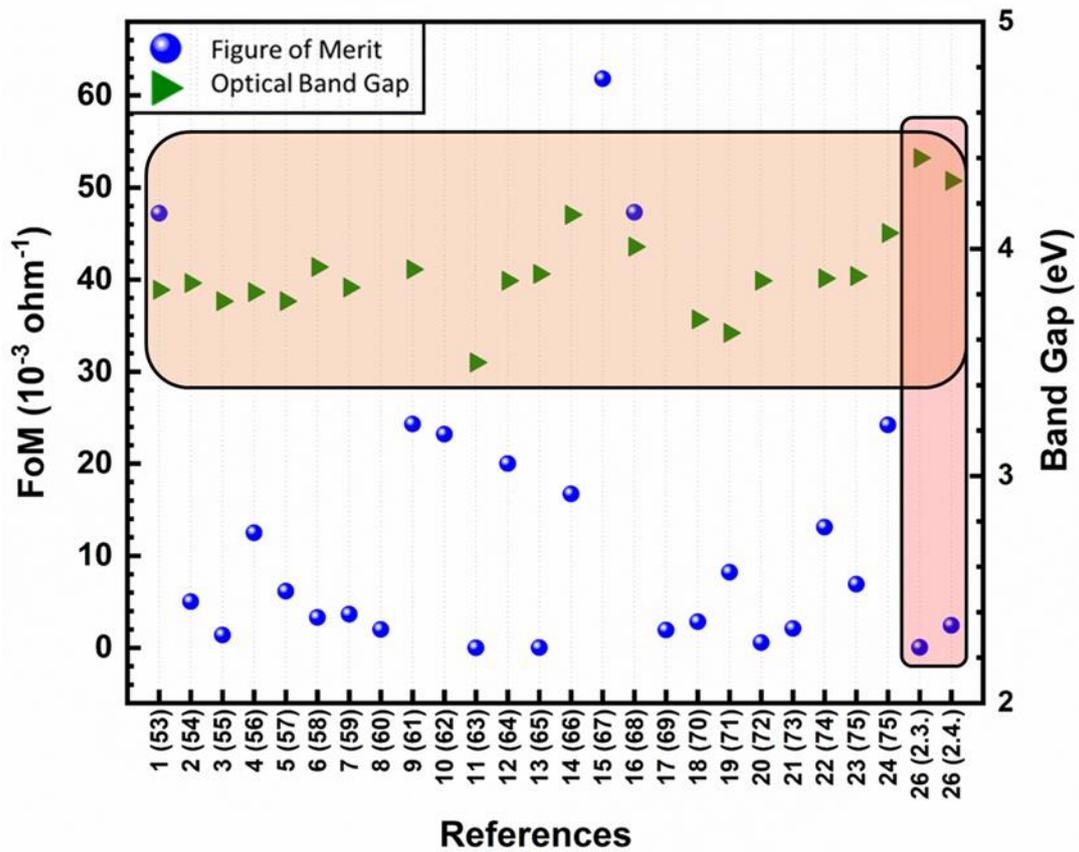
$$\phi = T^{10} / R_s \quad (1)$$

where  $\phi$  is figure of merit in  $\Omega^{-1}$ , T is the average optical transmittance from 400 to 700 nm in %, and R<sub>s</sub> is the sheet resistance in  $\Omega \cdot \square^{-1}$ .

O <sub>2</sub> /Ar ratio	Figure of merit ( $10^{-3} \Omega^{-1}$ )					
	Deposition Temperature ( $^{\circ}\text{C}$ )					
	100	250	100	250	100	250
	As-deposited		Vacuum Annealing at 400 $^{\circ}\text{C}$		Air Annealing at 400 $^{\circ}\text{C}$	
0	$3.06 \times 10^{-12}$	$3.21 \times 10^{-12}$	$6.52 \times 10^{-12}$	$3.38 \times 10^{-12}$	$9.91 \times 10^{-04}$	$8.17 \times 10^{-03}$
0.005	$4.27 \times 10^{-05}$	$3.93 \times 10^{-05}$	<b><math>1.87 \times 10^{-02}</math></b>	$1.32 \times 10^{-03}$	$1.02 \times 10^{-03}$	$5.70 \times 10^{-03}$
0.015	$1.12 \times 10^{-07}$	$8.31 \times 10^{-04}$	$9.13 \times 10^{-03}$	<b><math>5.14 \times 10^{-02}</math></b>	$1.10 \times 10^{-04}$	$1.96 \times 10^{-04}$
0.025	$1.72 \times 10^{-07}$	$4.45 \times 10^{-06}$	$4.10 \times 10^{-04}$	<b><math>2.90 \times 10^{-02}</math></b>	$7.66 \times 10^{-05}$	$1.42 \times 10^{-04}$
0.045	$1.35 \times 10^{-07}$	$1.22 \times 10^{-05}$	$1.72 \times 10^{-04}$	<b><math>1.56 \times 10^{-02}</math></b>	$1.33 \times 10^{-04}$	$1.17 \times 10^{-04}$
0.075	$1.15 \times 10^{-07}$	$3.70 \times 10^{-05}$	$3.70 \times 10^{-04}$	$9.66 \times 10^{-03}$	$1.23 \times 10^{-04}$	$1.12 \times 10^{-04}$

**Table 7:** Figure of merit for all SnO<sub>x</sub> samples.

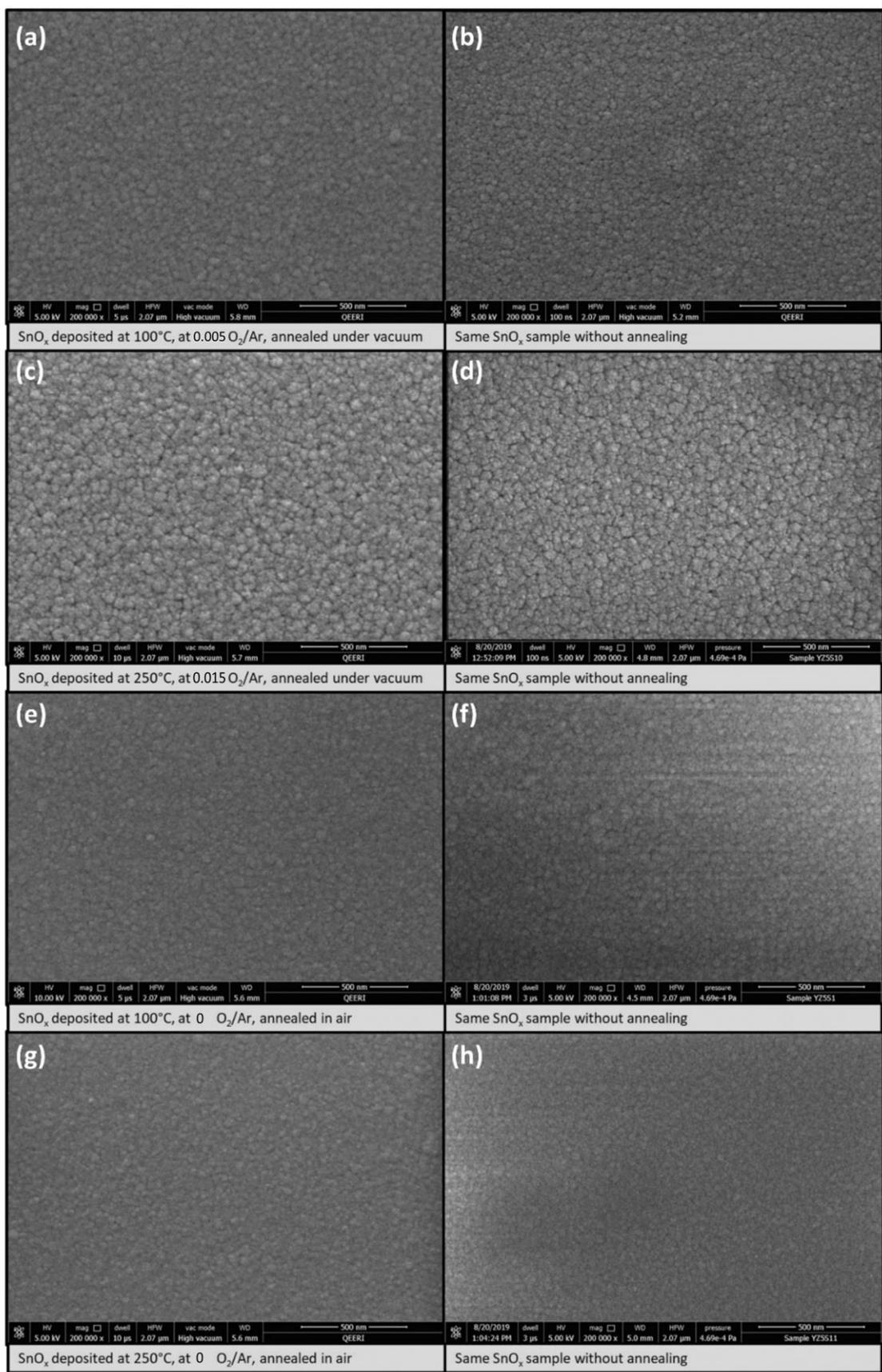
Table 1 in the appendix summarizes selected values from relevant literature of FoM for different doped SnO<sub>2</sub> thin films along with the doping type, synthesis method, band gap value, electrical resistivity and sheet resistivity, and optical transmittance<sup>[53-75]</sup>. Only two references of undoped films have been found in addition to our present work. Figure 7 highlights these FoM and band gap values as a function of the various references along with our present work in this section 2.3 and the next one 2.4. Highest FoM has been recorded for SnO<sub>2</sub> grown with spray pyrolysis and doped with fluorine, while the lowest value characterized SnO<sub>2</sub> was deposited by Pulsed Laser Deposition and doped with Tellurium. While our measured FoM belongs rather to the category of low values in Table 1 in the appendix, it is rather expected since our SnO<sub>2</sub> films are undoped and did not reach a low resistivity around  $10^{-4} \Omega\cdot\text{cm}$ . However, the optical band gap was among the highest reported in the literature which is, as discussed previously, due to the high oxidation state of SnO<sub>x</sub>.



**Figure 7:** Summary of Figure of Merit values and bang gap as a function of the various references <sup>[53-75]</sup> as well as the results of this thesis work in sections 2.3. and 2.4.

### 2.3.7. Morphology, TEM imaging & diffraction, and mapping

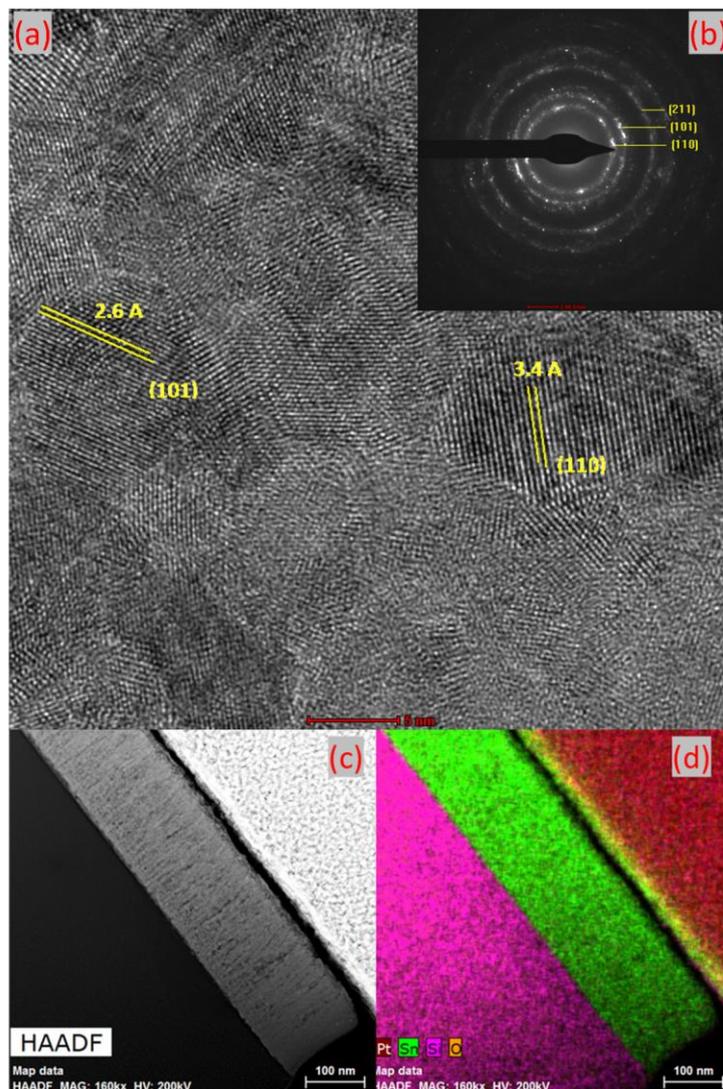
SEM observations were conducted on the four samples from the annealed series showing the highest conductivity as well as their related SnO<sub>x</sub> samples without annealing. All four samples show a crack-free SnO<sub>x</sub> films. The annealing process did not change drastically the morphology of the SnO<sub>x</sub> thin films. It can be clearly observed that the best conductive sample deposited at 250 °C with O<sub>2</sub>/Ar = 0.015 and annealed under vacuum (Figure 8c) has the largest grain size compared to the other samples as shown in Figure 8. The large grain size has improved the conductivity of SnO<sub>x</sub> as previously reported <sup>[36]</sup>. Following the discussion regarding the Hall effect results, the high mobility associated to these specific growth conditions is also attributed to the large grain size <sup>[33, 36]</sup>. The sample deposited at 100 °C and with O<sub>2</sub>/Ar = 0.005 and annealed under vacuum (Figure 8a) has shown the smallest grain size compared to the other three SnO<sub>x</sub> samples. This result is corroborating well the Hall effect measurement as it has shown the lowest mobility among the four selected samples.



**Figure 8:** Representative SEM images for the most electrically conductive annealed samples and their related SnO<sub>x</sub> samples without annealing: (a) SnO<sub>x</sub> deposited at 100°C, O<sub>2</sub>/Ar = 0.005, annealed under vacuum; (b) SnO<sub>x</sub> deposited at 100°C, O<sub>2</sub>/Ar = 0.005, without annealing; (c) SnO<sub>x</sub> deposited at 250°C, O<sub>2</sub>/Ar = 0.015, annealed under vacuum; (d) SnO<sub>x</sub> deposited at 250°C, O<sub>2</sub>/Ar = 0.015, without annealing, (e) SnO<sub>x</sub> deposited at 100°C,

O<sub>2</sub>/Ar = 0, annealed in air; (f) SnO<sub>x</sub> deposited at 100°C, O<sub>2</sub>/Ar = 0, without annealing; (g) SnO<sub>x</sub> deposited at 250°C, O<sub>2</sub>/Ar = 0, annealed in air; (h) SnO<sub>x</sub> deposited at 250°C, O<sub>2</sub>/Ar = 0, without annealing.

In order to confirm the multi-crystalline structure of the best conductive SnO<sub>2</sub> sample, we performed TEM imaging and mapping (Figure 9). The interplanar spacing could be measured directly from the image (Figure 9a) namely (110) and (101) planes, which is matching with the results revealed by XRD. Figure 9b shows the TEM diffraction pattern indexation revealing the intense patterns are related to (110) and (101) planes. High-angle annular dark-field imaging (HAADF) shown in Figure 9c has revealed a dense SnO<sub>2</sub> film with elongated column-shape crystalline structure towards the growth direction. EDS mapping (Figure 9d) has revealed the presence of a uniform layer of SnO<sub>2</sub> which is forming a sharp and clear interface with the quartz substrate grain size [33, 36].



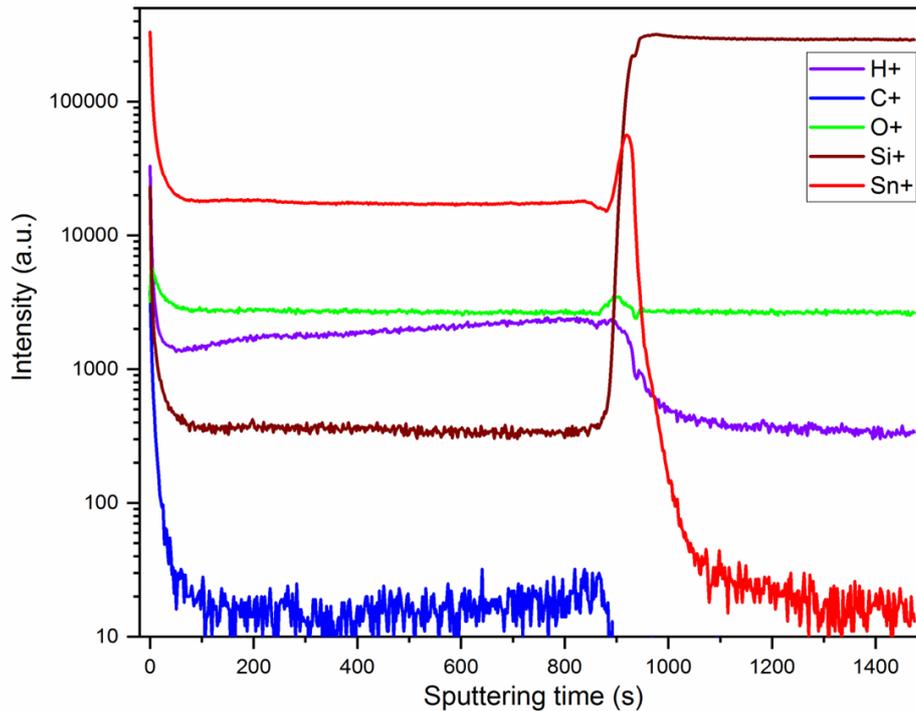
**Figure 9:** Cross-section TEM images for the best conductive SnO<sub>x</sub> thin film deposited at 250°C and O<sub>2</sub>/Ar = 0.015, annealed under vacuum (a) high resolution TEM image, (b) TEM diffraction pattern, (c) High-angle annular dark-field imaging (HAADF), (d) EDS mapping.

### 2.3.8. Thickness variation and elemental depth profiling

As reported in literature, the variation in thickness due to oxygen flow rate and temperature has been assessed using theoretical and experimental methods, and this variation depends on many factors. One of these factors is oxygen adsorption as oxygen flow rate increases from poor oxygen condition to the optimal condition which improves the deposition rate. On the other hand, one of the factors related to the decrease in deposition rate for higher O<sub>2</sub> flow rate (oxygen rich condition) is the further oxidation of the sputtering target surface<sup>[76]</sup>.

After having investigated the structural properties with XRD and TEM, we further confirm the homogeneity of the samples by performing TOF-SIMS on the best conductive samples to show the presence of high quality SnO<sub>x</sub> by the two constant intensities of Sn and O in the ToF-SIMS steady state conditions between the surface and the interface as shown in Figure 10. This analysis reveals the constant stoichiometry throughout the depth. The ion yield is much higher for Sn at the surface and the interface due to the matrix effect where the chemical environment changes as the secondary ion yields are strongly dependent on the chemical environment, which explains the high intensity of Sn at the surface and the interface. SIMS in general is inherently not a quantitative measurement technique. The secondary ion yields are strongly dependent on the chemical environment (matrix effect) and therefore, there is no direct correlation of elemental/compound intensity vs concentration. This technique also revealed the presence of H which slightly increases from the surface to the interface<sup>[77]</sup>.

TOF-SIMS also confirmed the absence of organic or inorganic contamination throughout the depth and it shows also perfect interfaces between the SnO<sub>x</sub> thin film and the quartz substrate.



**Figure 10:** TOF-SIMS profiles for the best conductive sample ( $\text{SnO}_x$  thin film deposited at  $250\text{ }^\circ\text{C}$ , at  $\text{O}_2/\text{Ar} = 0.015$ , and annealed under vacuum at  $400\text{ }^\circ\text{C}$ ).

### 2.3.9. Summary

For as-deposited  $\text{SnO}_x$  samples sputtered using a  $\text{SnO}$  target and based on their characterizations, it is suggested that such  $\text{SnO}_x$  is predominantly amorphous and/or showing low crystalline  $\text{SnO}$  in absence of additional  $\text{O}_2$  in the chamber, owing to the nature of the sputtering target ( $\text{SnO}$ ). This has led to lower electrical conductivity. At low  $\text{O}_2/\text{Ar}$  ratios, the predominating phase becomes  $\text{SnO}_2$  and the crystallinity tended to improve at higher temperature of  $250\text{ }^\circ\text{C}$  as the  $\text{O}_2/\text{Ar}$  ratio increased from 0.005 to 0.015. The poor oxygen condition is suggested to form defects within  $\text{SnO}_x$  thin film, particularly oxygen vacancies, and these defects are expected to increase while improving the crystalline structure of  $\text{SnO}_2$  by increasing the  $\text{O}_2/\text{Ar}$  ratio in the oxygen poor condition range (from 0.005 to 0.015  $\text{O}_2/\text{Ar}$ ). This has led to higher electrical conductivity. At higher  $\text{O}_2/\text{Ar}$  ratios, the crystallinity decreased at higher temperature of  $250\text{ }^\circ\text{C}$ . The rich  $\text{O}_2$  condition is expected to reduce the oxygen vacancy defects and eventually decreases the charge carrier concentration. This has led to lower electrical conductivity<sup>[78]</sup>. Both annealing processes are expected to improve the crystallinity of the films due to a thermal treatment at  $400\text{ }^\circ\text{C}$  for 1 hour. However, vacuum annealing is expected to conserve the charge carrier concentration by preventing annihilation of the oxygen vacancies due to the lack of  $\text{O}_2$ . On the other hand, air annealing is found to reduce the charge carrier concentration by filling the oxygen vacancies with oxygen supplied from air.

The air annealing has shown a better crystallinity compared to vacuum annealing as all the samples have reported high crystallinity and two SnO<sub>x</sub> samples deposited without O<sub>2</sub> recorded the highest crystallite size. Both SnO<sub>x</sub> deposited without O<sub>2</sub> and annealed in air have revealed that the presence of both phases SnO and SnO<sub>2</sub>, and it is clear that O<sub>2</sub> from the air atmosphere has oxidized significantly SnO to SnO<sub>2</sub>. Moreover, as per crystallite size reported in Table 2, it is found that the crystallite size is relatively higher after vacuum annealing for the SnO<sub>x</sub> samples (deposited at 250°C). While for the annealing under vacuum, the charge carrier concentration which is mainly due to oxygen vacancies is found to be conserved. The combination of improved crystallite size and the conservation of charge carrier are the key factors for improving the electrical conductivity as compared to SnO<sub>x</sub> samples from other series with same O<sub>2</sub>/Ar ratio.

The first part of the chapter highlighted the structure/performance correlations of SnO<sub>x</sub> thin films grown by RF MS. High quality SnO<sub>x</sub> samples were prepared using magnetron sputtering deposition method followed by thermal annealing processes. Crystalline microstructure, electrical and optical properties were characterised in-depth. Both SnO<sub>2</sub> and mixed SnO/SnO<sub>2</sub> thin films were synthesized using RF sputtering. The **most electrically conductive sample** was obtained by using O<sub>2</sub>/Ar = 0.015 during the growth at 250 °C and followed by a moderate vacuum post annealing at 400 °C /  $6.66 \times 10^{-2}$  Pa. This SnO<sub>x</sub> film has shown a compact and dense morphology without presence of pinholes or cracks, and its grain sizes were relatively larger compared to other samples, which clearly improved the electron mobility. Its average optical transmittance between 400 and 700 nm was measured to be above 80%. The **best optical transmittance** of 91% is achieved only using the highest O<sub>2</sub>/Ar ratio of 0.075 for deposited SnO<sub>x</sub> at 250°C without annealing and with vacuum annealing as well as both air annealed samples. Vacuum annealing provided a higher electrical conductivity compared to the as-deposited and air-annealed processes. This is attributed to the improvement of crystalline microstructure as well as the presence of oxygen lattice vacancies which has led to a high charge carrier concentration. These growth conditions summarise a good compromise between a high grain size, higher crystalline structure, and high charge carrier concentration. The most conductive SnO<sub>x</sub> sample obtained by using O<sub>2</sub>/Ar = 0.015 during the growth at 250 °C and followed by a moderate vacuum post annealing have also enabled the best optoelectronic performance as assessed by FoM.

## 2.4. SnO<sub>x</sub> sputtered using a Sn target

In this study, the SnO<sub>x</sub> thin films were prepared using reactive magnetron sputtering, manufactured by AJA International Co. (described in the appendix in section 1.3.), with a high purity Sn (99.999%) 2".

### 2.4.1. Effect of O<sub>2</sub> flow rate during the growth on SnO<sub>x</sub>

The aim of this study is to investigate the effect of O<sub>2</sub> flow rate while using magnetron sputtering on deposited SnO<sub>x</sub> using Sn metallic sputtering target while setting the other parameters, namely power at 50 W, Ar flow rate at 12 sccm, deposition temperature at 100 °C and deposition pressure at 0.45 Pa. The deposition time was kept constant for low O<sub>2</sub> flow rate at 12 minutes. However, it was slightly increased to 15 minutes at higher O<sub>2</sub> flow rates to adjust the expected decrease in deposition rate. The targeted thickness is slightly above 100 nm. The ultimate objective is to explore the SnO<sub>x</sub> thin film properties from low O<sub>2</sub> flow rates to higher one to study the structural and optoelectronic properties while varying the O<sub>2</sub> at low flow rate. The samples synthesis parameters are presented in table 8.

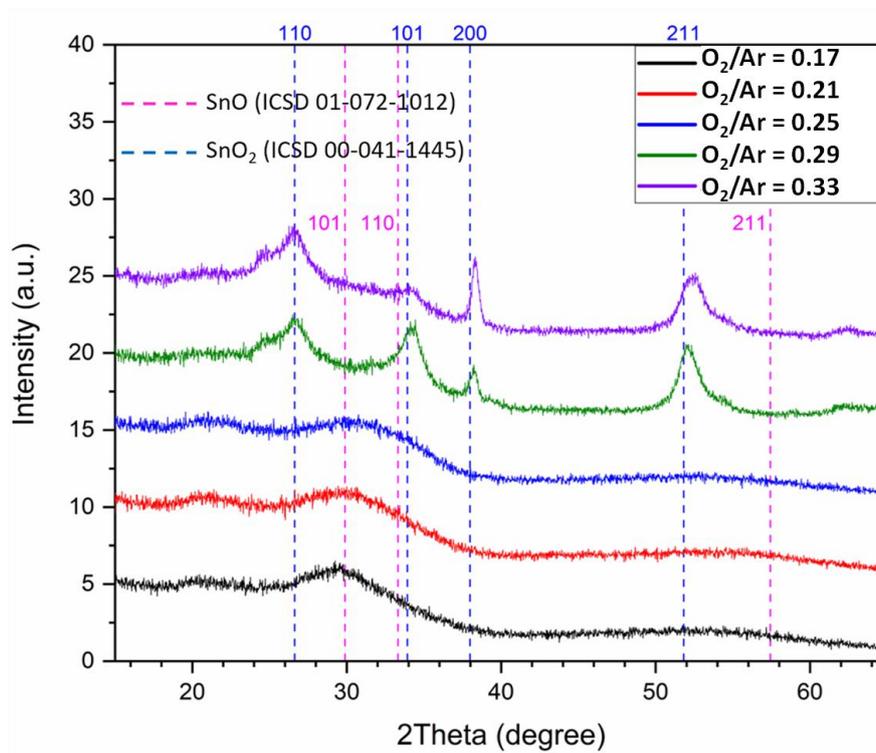
Sample name	Power (W)	Substrate Temperature (°C)	Deposition time (min)	O <sub>2</sub> /Ar (O <sub>2</sub> :Ar sccm)	Deposition Pressure (Pa)
6S1D	50	100	12 (100nm target)	0.17 (2:12)	0.45
6S2D				0.21 (2.5:12)	
6S3D				0.25 (3:12)	
6S4D			15 (100nm target)	0.29 (3.5:12)	
6S5D				0.33 (4:12)	

**Table 8:** SnO<sub>x</sub> thin film deposition parameters using radio-frequency magnetron sputtering.

#### a) Structural properties and chemical state analysis

Grazing incidence X-ray diffraction was performed to probe the crystalline microstructure of the SnO<sub>x</sub> thin films in different O<sub>2</sub> partial pressures during the growth. Both SnO<sub>x</sub> samples deposited at 0.17 and 0.21 O<sub>2</sub>/Ar flow rate ratios have shown a broad peak around 30° which indicate the presence of low crystalline or amorphous SnO phase. Similar results were reported in a previous section of this chapter. For SnO<sub>x</sub> deposited at 0.25 O<sub>2</sub>/Ar flow rate ratio the peak become broader which indicate the presence of either or both low crystalline or amorphous phases SnO and SnO<sub>2</sub>. Samples deposited at 0.29 and 0.33 O<sub>2</sub>/Ar ratios have revealed a clear crystalline microstructure by the presence of XRD peaks (110), (101), (200) and (211) related to SnO<sub>2</sub> phase only. It is worth noting that the sharpest peak was (200) observed for SnO<sub>x</sub> deposited at 0.29 O<sub>2</sub>/Ar flow rate ratio as shown in Figure 11. The poor oxygen condition during the growth has enabled the formation of SnO which the lower oxidation state of SnO<sub>x</sub>. Increasing the O<sub>2</sub> partial pressure in the growth atmosphere has increased the oxidation state of SnO<sub>x</sub> which led to a dominant phase of SnO<sub>2</sub>. The crystallization of the SnO phase was not achieved in this

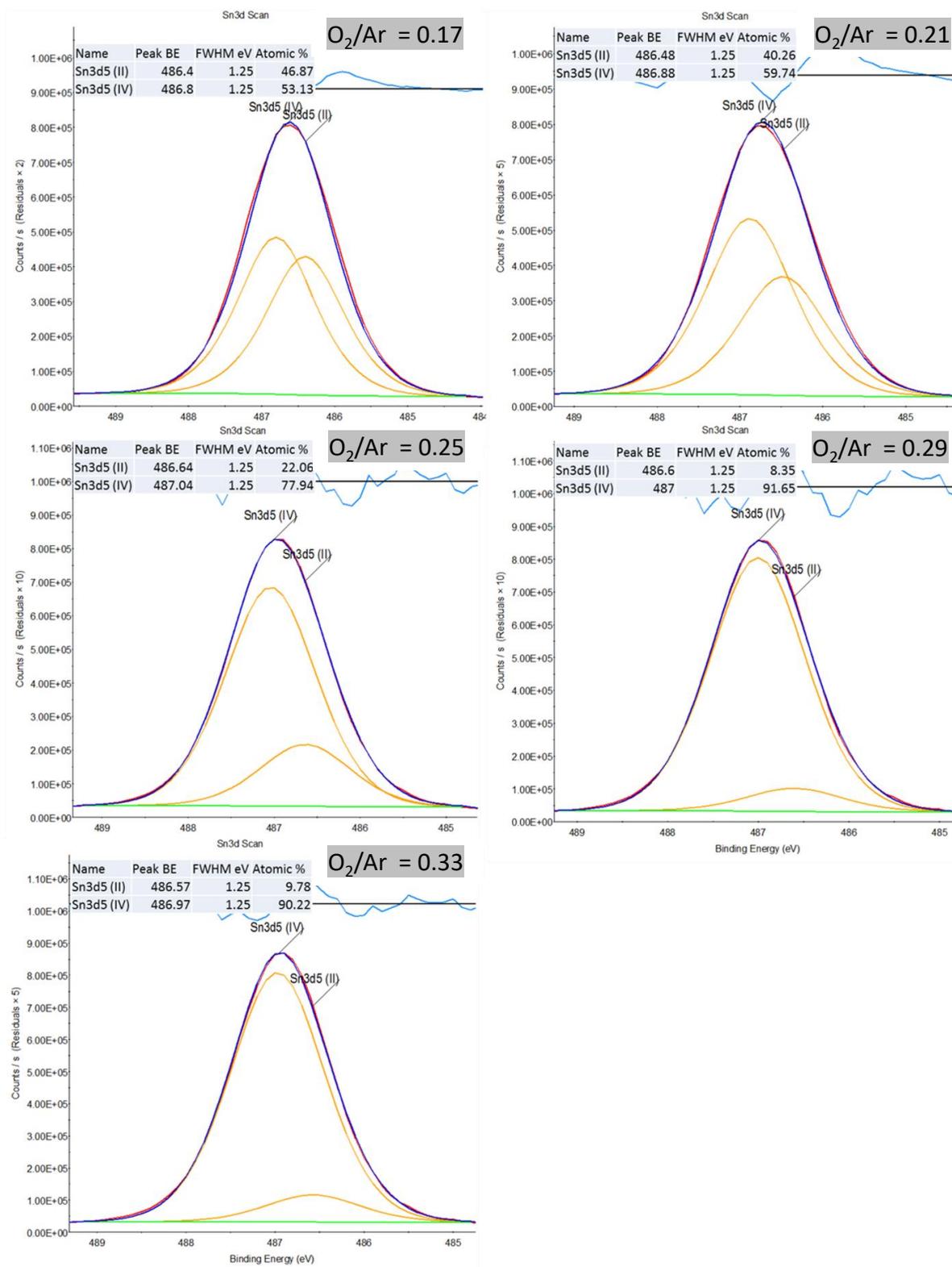
study as the  $\text{SnO}_x$  tend to change the phase from SnO to  $\text{SnO}_2$  while increasing the  $\text{O}_2$  partial pressure. On the other hand, increasing the  $\text{O}_2$  partial pressure to richer condition enabled the crystallization of  $\text{SnO}_2$  for samples grown at 0.29 and 0.33  $\text{O}_2/\text{Ar}$  flow rate ratios.



**Figure 11:** GIXRD pattern for  $\text{SnO}_x$  sample using different  $\text{O}_2$  flow rates during the growth.

Chemical state analysis was performed using X-ray photoelectron spectroscopy which is a surface sensitive technique analysing the surface of around 5 nm thickness. As expected, the surface of thin films gets oxidized at the surface once exposed to the air. However, the technique can provide an important indication of the present chemical states of the overall film. All XPS spectra for  $\text{Sn}3d_{5/2}$  show only the presence of the two chemical states Sn(II) and Sn(IV) without any presence of Sn(0) which indicates that the Sn sputtering deposition is completely oxidized and there is no indication of metallic Sn presence within the surface of  $\text{SnO}_x$ . The peak position of  $\text{Sn}3d_{5/2}$  related to Sn(II) is around 486.5 (+/- 0.1) eV while the one for Sn(IV) is around 486.9 (+/- 0.1) as shown in Figure 12. The chemical state analysis revealed that the presence of Sn(II) in  $\text{SnO}_x$  at 0.17  $\text{O}_2/\text{Ar}$  flow rate ratio is around 47% while Sn(IV) is 53%. Sn(II) percentage decreases to 40% while the  $\text{O}_2/\text{Ar}$  flow rate ratio increases to 0.29 and Sn(IV) increases to 60%. Sn(II) percentage decreases drastically to 22% while the  $\text{O}_2/\text{Ar}$  flow rate ratio increases to 0.25 and Sn(IV) increases to 78%. Sn(II) percentage continue decreasing to 8% while the  $\text{O}_2/\text{Ar}$  flow rate ratio reaches 0.29 and Sn(IV) increases to 92%. Sn(II) percentage slightly increasing to 10% at  $\text{O}_2/\text{Ar}$  flow rate ratio of 0.33 while Sn(IV) decreases to 90%. XPS results shows that the chemical state of  $\text{SnO}_x$  thin film made from metallic Sn sputtering target follows the same pattern of  $\text{SnO}_x$  made from oxidised SnO target. It can be summarised by the presence of mainly SnO as revealed by Sn(II) at

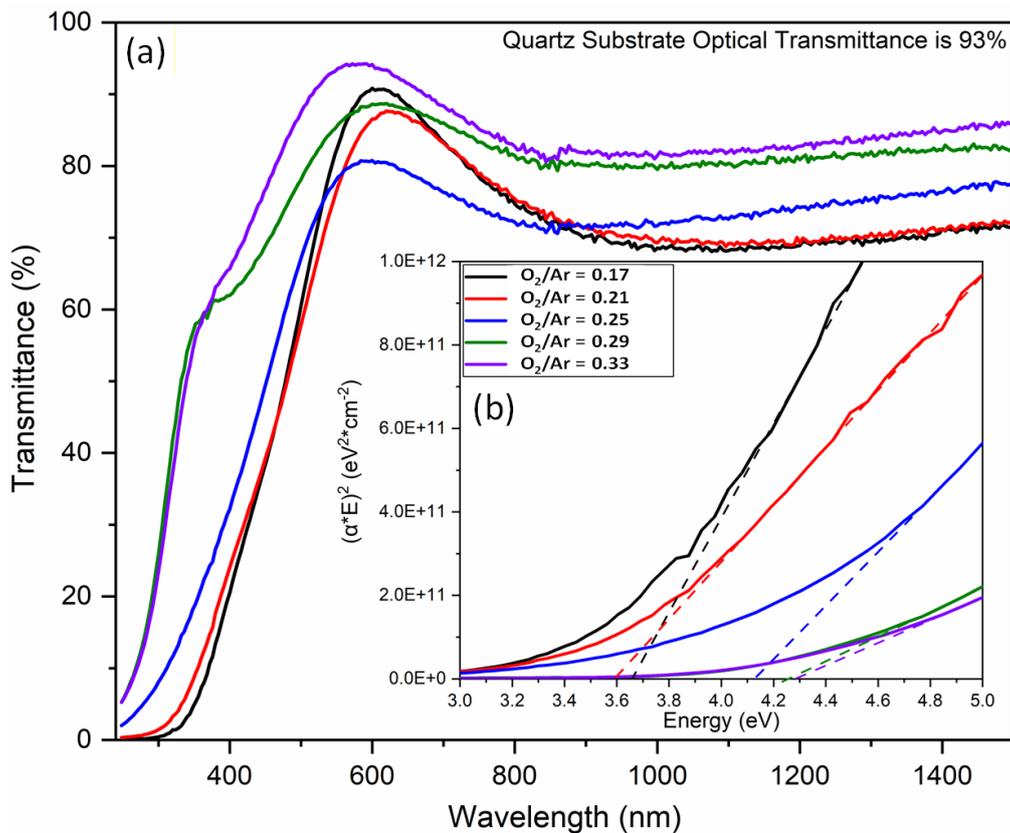
low  $O_2/Ar$  flow rates along with the presence of Sn(IV) at the surface. It is worth noting that the surface of thin films would always oxidise to a higher chemical state due to the presence of oxygen in air. The main chemical state of  $SnO_x$  becomes  $SnO_2$  as revealed by Sn(IV) in  $O_2$  rich conditions for  $O_2/Ar$  ratio at 0.25 or higher. These results are matching with the XRD results discussed previously.



**Figure 12:** XPS peak fitting for  $SnO_x$  samples using different  $O_2$  flow rates during the growth.

b) Optical properties

The optical transmittance in UV-Visible-Infrared range for SnO<sub>x</sub> thin film samples is shown in Figure 13(a). Both average of optical transmittance from 400 to 700 nm and from 500 to 800 nm as well as the optical transmittance at 550 nm and the optical bandgap are defined in Table 9. The average optical transmittance from 500 to 800 nm is 82% for SnO<sub>x</sub> thin film sample with O<sub>2</sub>/Ar flow rate ratio at 0.17. This transmittance decreases to 80% at O<sub>2</sub>/Ar of 0.21 and continues to decrease to 76.5% at O<sub>2</sub>/Ar of 0.25. The optical transmittance increases to 85.3% when the O<sub>2</sub>/Ar flow rate ratio increases to 0.29 and it continues to increase to 89.7% for O<sub>2</sub>/Ar ratio at 0.33. The optical transmittance spectra show that cut-on edge is towards higher wavelength with oxygen poor growth conditions for SnO<sub>x</sub> thin film grown samples and it is towards lower wavelength for the relatively oxygen rich conditions grown SnO<sub>x</sub> thin film samples. The optical transmittance is relatively high which is above than 80% for four samples out of five without excluding the absorbance/reflection contribution of the quartz substrate. The highest optical transmittance is achieved by the SnO<sub>x</sub> thin sample which has the highest O<sub>2</sub> flow rate during the growth which is related to highest oxidized SnO<sub>x</sub>.



**Figure 13:** (a) UV-Vis transmittance spectra and (b) Tauc plot for SnO<sub>x</sub> samples using different O<sub>2</sub>/Ar flow rate ratio during the growth.

Tauc plot for SnO<sub>x</sub> thin film samples is shown in Figure 13(b) and their related optical bandgap is defined in Table 10. Tauc plot was established for both direct and indirect bandgap. Only the direct

bandgap plot had a clear linear region which could be extrapolated to the abscissa. Therefore, the SnO<sub>x</sub> thin film samples are direct bandgap. The SnO<sub>x</sub> sample of oxygen poor condition during the growth, namely for O<sub>2</sub>/Ar ratio at 0.17 and 0.21, have a lower bandgap which is 3.7 and 3.6 eV respectively. The oxygen rich condition during the growth, namely for O<sub>2</sub>/Ar ratio at 0.25, 0.29 and 0.33, have a higher bandgap which is 4.1, 4.3 and 4.3 eV, respectively. The optical bandgap results, and as per the literature, are interpreted as the SnO<sub>x</sub> thin films grown at O<sub>2</sub>/Ar ratio of 0.17 and 0.21 are mainly SnO and the SnO<sub>x</sub> thin films grown at O<sub>2</sub>/Ar ratio of 0.25, 0.29 and 0.33 are mainly SnO<sub>2</sub>. It is worth mentioning that 4.3 eV one of the highest results reported in the literature as shown previously in the summary (Figure 7).

O <sub>2</sub> /Ar (O <sub>2</sub> :Ar sccm)	Optical Transmittance (%)			Optical Bandgap (eV)
	Average 400–700 nm	Average 500–800 nm	At 550 nm	
0.17 (2:12)	68.8	82.0	82.3	3.7
0.21 (2.5:12)	66.9	80.0	76.6	3.6
0.25 (3:12)	68.2	76.5	78.1	4.1
0.29 (3.5:12)	80.5	85.3	85.5	4.3
0.33 (4:12)	86.7	89.7	93.4	4.3

**Table 9:** Average optical transmittance from 400 to 700 nm and from 500 to 800 nm, optical transmittance for 550 nm, and optical bandgap for SnO<sub>x</sub> thin film samples.

### c) Electrical properties

Hall effect measurements on SnO<sub>x</sub> thin films shown in Table 10 have revealed that poor oxygen condition led to a very high electrical resistivity and further measurements were not possible. The relatively oxygen richer condition, namely for SnO<sub>x</sub> thin films grown at O<sub>2</sub>/Ar ratio of 0.29 and 0.33, has provided a very low electrical resistivity below  $5 \times 10^{-3} \Omega \cdot \text{cm}$ . The SnO<sub>x</sub> sample grown at O<sub>2</sub>/Ar ratio of 0.25 has an average electrical resistivity of  $2.84 \times 10^{-1} \Omega \cdot \text{cm}$ .

The high electrical resistivity of SnO<sub>x</sub> thin films, grown at O<sub>2</sub>/Ar ratio of 0.17 and 0.21, is mainly related to the amorphous character of the dominant phase SnO as revealed by XRD, XPS and UV-Vis. The low electrical resistivity of SnO<sub>x</sub> thin films, grown at O<sub>2</sub>/Ar ratio of 0.29 and 0.33, is related to both high mobility which is mainly a result of the crystalline microstructure of SnO<sub>2</sub> dominant phase and intrinsic charge carrier of SnO<sub>2</sub> due to the oxygen vacancies<sup>[36]</sup>. The moderate electrical resistivity of SnO<sub>x</sub> thin film, grown at O<sub>2</sub>/Ar ratio of 0.25, is related to the amorphous character of the dominant SnO<sub>2</sub> phase which has led to a relatively smaller electron mobility of  $1.51 \text{ cm}^2/\text{V}\cdot\text{s}$  compared to both  $18.1$  and  $24.4 \text{ cm}^2/\text{V}\cdot\text{s}$  for rich oxygen conditions.

O <sub>2</sub> /Ar (O <sub>2</sub> :Ar sccm)	Thickness (nm)	Rate (nm/min)	Resistivity (Ω.cm)	Sheet resistance (Ω/□)	Mobility (cm <sup>2</sup> /V.s)	Charge concentration (cm <sup>-3</sup> ) – (carrier type)
0.17 (2:12)	132	11	N/A*	N/A	N/A	N/A
0.21 (2.5:12)	125	10.42	N/A*	N/A	N/A	N/A
0.25 (3:12)	113	9.42	2.84 x 10 <sup>-1</sup>	2.51 x 10 <sup>4</sup>	1.51	1.64 x 10 <sup>19</sup> - (n-type)
0.29 (3.5:12)	136	9.07	4.93 x 10 <sup>-3</sup>	3.63 x 10 <sup>2</sup>	18.1	9.52 x 10 <sup>19</sup> - (n-type)
0.33 (4:12)	145	9.67	3.00 x 10 <sup>-3</sup>	2.07 x 10 <sup>2</sup>	24.4	1.24 x 10 <sup>20</sup> - (n-type)

**Table 10:** Electrical properties of SnO<sub>x</sub> thin film samples using different O<sub>2</sub> flow rate ratios during the growth.

\* Very high resistivity.

#### d) Optoelectronic performance

Table 11 shows the Figure of Merit (FoM) for SnO<sub>x</sub> thin film samples. The poor oxygen conditions of grown SnO<sub>x</sub> thin film at O<sub>2</sub>/Ar ratio of 0.17 and 0.21 led to a high resistivity which is due to the amorphous character of the dominant phase SnO. Thus, their FoM could not be computed, and it is expected to be very low. However, richer oxygen condition grown SnO<sub>x</sub> thin film at O<sub>2</sub>/Ar ratio of 0.29 and 0.33 have shown relatively high optoelectrical performance by their high FoM. The SnO<sub>x</sub> thin film grown at O<sub>2</sub>/Ar ratio of 0.25, has a relatively low optoelectrical performance due to the moderate electrical resistivity. It is worth noting that as per the literature and to the best of our knowledge, SnO<sub>x</sub> thin film grown O<sub>2</sub>/Ar ratio of 0.33 of 2.44 x 10<sup>-3</sup> Ω<sup>-1</sup> has one of the highest reported optoelectrical performance for undoped SnO<sub>x</sub> as indicated in the summary of Figure 7.

O <sub>2</sub> /Ar (O <sub>2</sub> :Ar sccm)	Figure of merit (10 <sup>-3</sup> Ω <sup>-1</sup> )		
	Average Transmittance 400 – 700 nm	Average Transmittance 500 – 800 nm	Transmittance at 550 nm
0.17 (2:12)	N/A	N/A	N/A
0.21 (2.5:12)	N/A	N/A	N/A
0.25 (3:12)	8.67 x 10 <sup>-4</sup>	2.73 x 10 <sup>-3</sup>	3.36 x 10 <sup>-3</sup>
0.29 (3.5:12)	0.315	0.562	0.575
0.33 (4:12)	1.16	1.63	2.44

**Table 11:** Figure of Merit for SnO<sub>x</sub> thin film samples using different O<sub>2</sub> flow rates during the growth.

#### 2.4.2. Effect of deposition temperature on SnO<sub>x</sub> using Sn metallic target

After assessing the effect O<sub>2</sub> flow rate during the growth on SnO<sub>x</sub>, this study is mainly focusing on the effect of deposition temperature during the growth on SnO<sub>x</sub> thin films in poor oxygen condition. All parameters were kept fixed except the deposition temperature. The fixed parameters are power at 50 W, Ar flow rate at 12 sccm, deposition pressure at 0.45 Pa, and O<sub>2</sub>/Ar flow rate ratio at 0.21. This last parameter was chosen because it is just below the poor oxygen condition point for the SnO<sub>x</sub> thin film grown at O<sub>2</sub>/Ar of 0.25 that provided a predominant SnO<sub>2</sub> phase. The deposition temperature used in

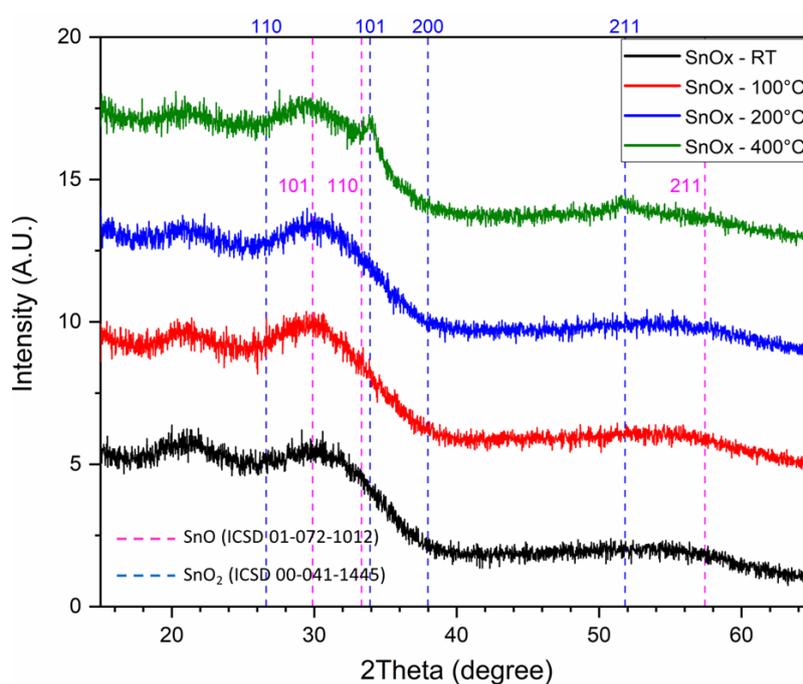
this study, are room temperature (RT), 100, 200 and 400°C. The deposition time was kept constant at 12 minutes for all samples. The targeted thickness is slightly above 100 nm. The objective of this study is also to assess the possibility to improve the crystallinity of the dominant SnO phase and eventually improve the electrical conductivity. The samples synthesis parameters are presented in Table 12.

Sample name	Mode/target	Power (W)	Substrate Temperature	Deposition time (min)	O <sub>2</sub> /Ar (O <sub>2</sub> :Ar sccm)	Depos. Pressure (Pa)
7S1D	RF/Sn (2")	50	RT	12 (100nm target)	0.21 (2.5:12)	0.45
7S2D			100°C			
7S3D			200°C			
7S4D			400°C			

**Table 12:** SnO<sub>x</sub> thin film deposition parameters using radio-frequency magnetron sputtering.

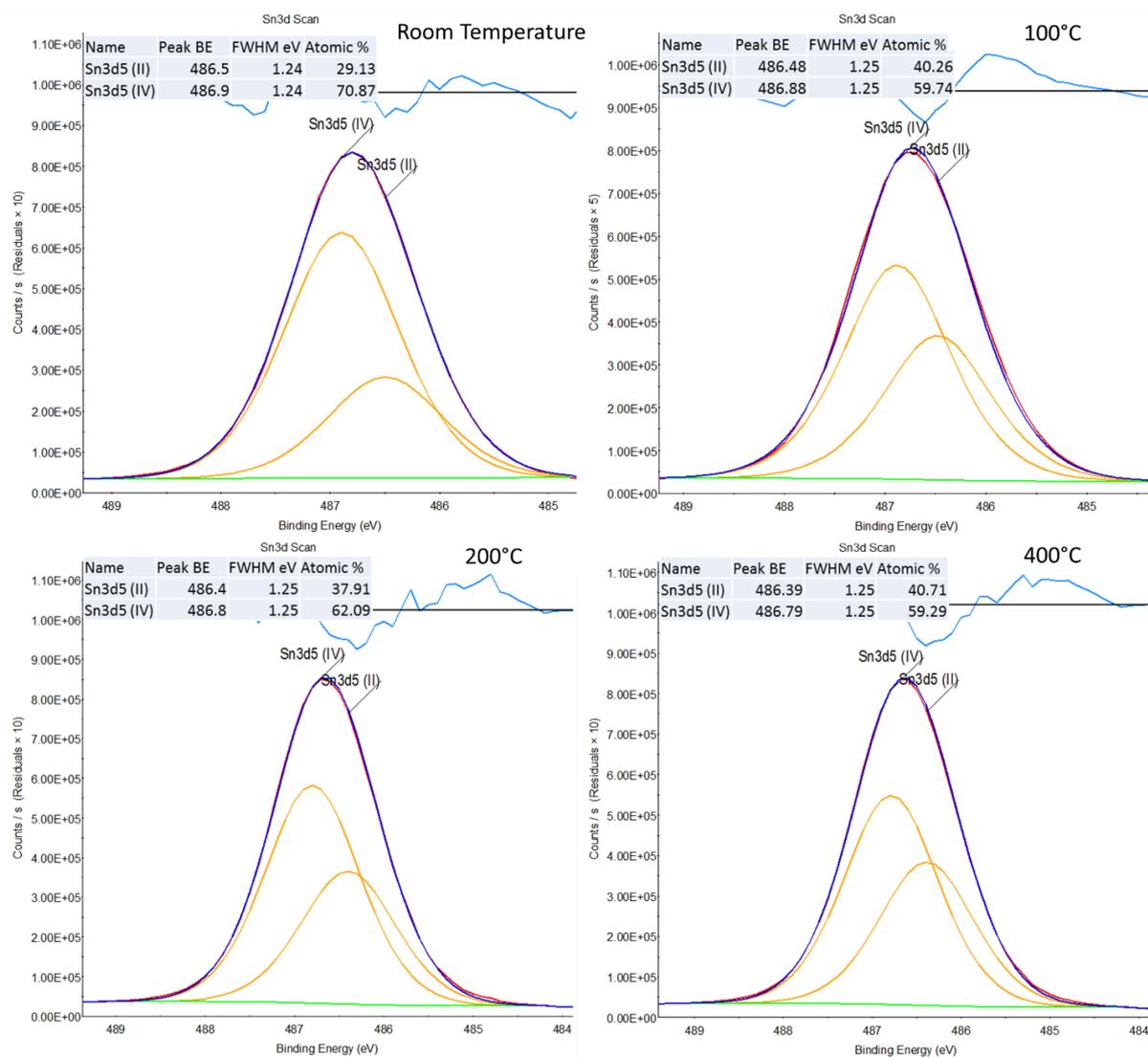
a) Structural properties and chemical state analysis

Grazing incidence X-ray diffraction, shown in Figure 14, was performed to probe the crystalline structure of the SnO<sub>x</sub> thin films in different deposition temperatures during the growth. All SnO<sub>x</sub> samples have shown a broad peak around 30° which indicate the presence of low crystalline or amorphous SnO phase. In addition, GIXRD pattern of SnO<sub>x</sub> deposited at 400°C has shown a small peak at 34° related to SnO<sub>2</sub> (101) and a smaller peak at 52° related to SnO<sub>2</sub> (211). Both these peaks indicate a small presence of SnO<sub>2</sub> crystalline phase. Varying the deposition temperature did not change the amorphous microstructure of the SnO at RT, 100°C and 200°C. However, the deposition temperature at 400°C enabled the formation of a small crystalline phase of SnO<sub>2</sub> along with predominant amorphous SnO.



**Figure 14:** GIXRD pattern for SnO<sub>x</sub> sample using different deposition temperatures during the growth.

Surface sensitive chemical state analysis was performed using X-ray photoelectron spectroscopy. Similar to the previous study, it is expected to get higher oxidation state at the surface of thin films due to the air exposure. All XPS spectra for Sn3d<sub>5/2</sub> show only the presence of the two chemical states Sn(II) and Sn(IV) without any presence of Sn(0) which indicates that the Sn sputtering deposition is completely oxidized and there is no indication of metallic Sn presence within the surface of SnO<sub>x</sub>. The peak position of Sn3d<sub>5/2</sub> related to Sn(II) is around 486.5 (+/- 0.1) eV while the one for Sn(IV) is around 486.9 (+/- 0.1) as shown in Figure 15. The chemical state analysis revealed that the presence of Sn(II) in SnO<sub>x</sub> deposited at room temperature is around 29% while Sn(IV) is 71%. Sn(II) percentage increases to 40% and Sn(IV) increases to 60% for SnO<sub>x</sub> deposited at 100°C. Sn(II) percentage slightly decreases to 38% and Sn(IV) increases to 62% for SnO<sub>x</sub> deposited at 200°C. For SnO<sub>x</sub> deposited at 400°C, Sn(II) percentage slightly increasing to 41% and Sn(IV) decreases to 59%. XPS results the chemical states distribution of Sn(II) and Sn(IV) were very closely matching in samples deposited at 100, 200 and 400°C. While SnO<sub>x</sub> deposited at room temperature was more oxidized to the higher oxidation state as revealed by the lower Sn(II) and the higher Sn(IV) chemical states. As stated before, the presence of higher oxidation states at the surface is likely related to the post deposition oxidation of SnO<sub>x</sub> once exposed to the air. These results along the XRD ones have revealed that there is a major presence of Sn(II) chemical state as SnO except at the surface, and the temperature deposition did not change drastically the chemical state except for the room temperature deposition which favours more the Sn(IV) chemical state. It is also worth noting that the higher temperature 400°C enables the crystallization of the phase SnO<sub>2</sub>.

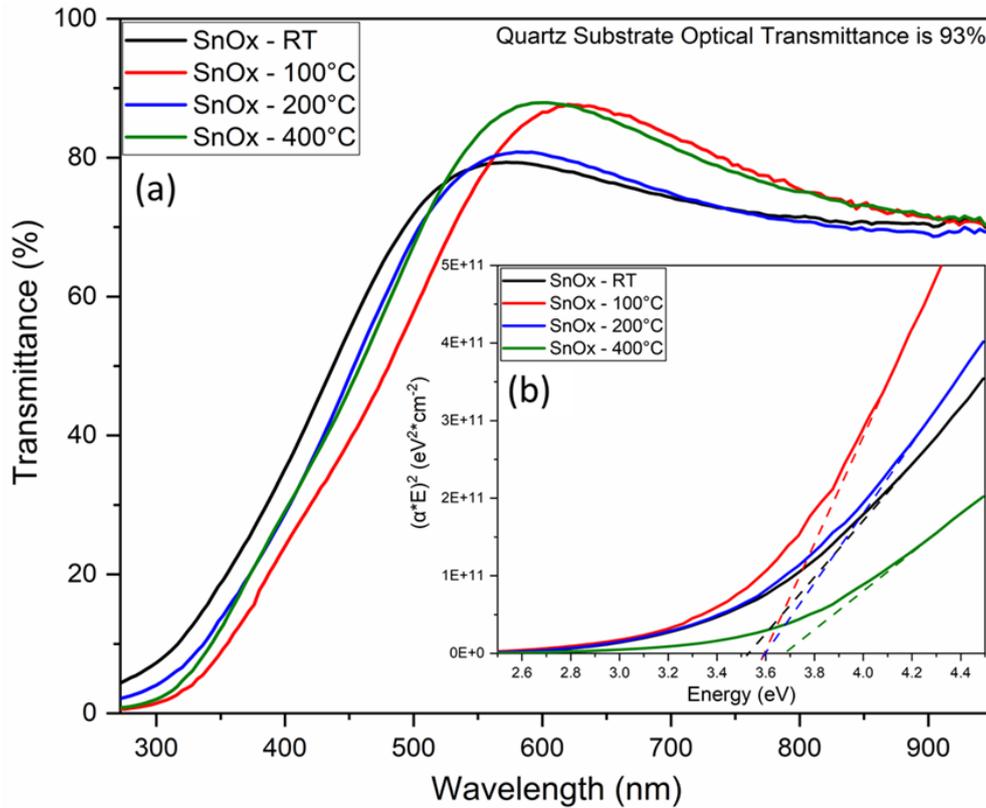


**Figure 15:** XPS peak fitting for SnO<sub>x</sub> samples using different deposition temperatures during the growth.

### b) Optical properties

The optical properties were probed by UV-Visible spectrometer. The optical transmittance in UV-Visible-Infrared range for SnO<sub>x</sub> thin film samples is shown in Figure 16(a). The average of optical transmittance from 400 to 700nm, from 500 to 800nm, the optical bandgap are defined in Table 13. The average optical transmittance from 500 to 800nm is 75.5% for SnO<sub>x</sub> thin film sample deposited at room temperature. This transmittance increases to 80% for SnO<sub>x</sub> deposited at 100°C and it decreases to 75.9% for SnO<sub>x</sub> deposited at 200°C. The optical transmittance increases to 83.2% when the deposition temperature of SnO<sub>x</sub> increases to 400°C. The optical transmittance spectra show that cut-on edges for all samples is very close to each other. The optical transmittance is relatively high which is above than 75% for all samples without excluding the absorbance/reflection contribution of the quartz substate. The highest optical transmittance is

obtained with the SnO<sub>x</sub> thin film sample grown at 400°C. This is due to the better microstructural ordering and disorder reduction<sup>[35, 36, 51, 79]</sup>.



**Figure 16:** (a) UV-Vis transmittance spectra and (b) Tauc plot for SnO<sub>x</sub> samples using different deposition temperatures during the growth.

Tauc plot for SnO<sub>x</sub> thin film samples is shown in Figure 16(b) and their related optical bandgap is defined in Table 13. Tauc plot was established for both direct and indirect bandgap. Similar to the previous study, only the direct bandgap plot had a clear linear region for all samples which could be extrapolated to the abscissa. Therefore, the SnO<sub>x</sub> thin film samples are direct bandgap. All SnO<sub>x</sub> samples have very similar optical bandgap of about 3.6 (+/- 0.1) eV. SnO<sub>x</sub> deposited at room temperature have shown a bandgap of 3.5 eV. Both SnO<sub>x</sub> samples deposited at 100 and 200°C have exhibited an optical bandgap of 3.6 eV while SnO<sub>x</sub> deposited at 400°C show an optical bandgap of 3.7 eV. The optical bandgap difference between the four samples is very small. However, this study shows that increasing the temperature can increase even slightly the optical bandgap which is more due to the microstructural ordering and defect reduction. It is worth mentioning that the presence of defects can form shallow energy levels which reduces the optical bandgap. As per the literature, the optical bandgaps of 3.6 (+/- 0.1) eV reveal that predominant phase in the SnO<sub>x</sub> thin film samples is SnO which is expected to the poor oxygen condition during the growth of SnO<sub>x</sub>. This predominant SnO phase was not significantly affected by the deposition temperature variation from RT to 400°C<sup>[35, 36, 51, 79]</sup>.

Deposition Temperature	Optical Transmittance (%)			Optical Bandgap (eV)
	Average 400-700nm	Average 500-800 nm	At 550nm	
RT	69.4	75.5	78.7	3.5
100°C	66.9	80.0	76.6	3.6
200°C	67.8	75.9	79.1	3.6
400°C	71.1	81.4	83.2	3.7

**Table 13:** Average optical transmittance from 400 to 700nm and from 500 to 800nm, optical transmittance for 550nm, and optical bandgap for SnO<sub>x</sub> thin film samples using different deposition temperatures during the growth.

#### c) Electrical properties

Hall effect measurements were not possible due to the high resistivity of the SnO<sub>x</sub> thin film samples as shown in table 14. The deposition temperature did not improve the electrical conductivity at poor oxygen condition during the growth. The high resistivity is likely related to the amorphous character of the SnO. Improving the SnO crystallinity can be a key factor to improve the electrical and optoelectrical performance of the SnO samples. Further studies related to this topic will be presented and discussed in this thesis.

Deposition Temperature	Thickness (nm)	Rate (nm/min)	Resistivity (Ω*cm)	Sheet resistance (Ω/sq)	Mobility (cm <sup>2</sup> /Vs)	Charge concentration (cm <sup>-3</sup> )
RT	100	8.33	N/A*	N/A	N/A	N/A
100°C	125	10.42	N/A*	N/A	N/A	N/A
200°C	116	9.67	N/A*	N/A	N/A	N/A
400°C	205	17.08	N/A*	N/A	N/A	N/A

**Table 14:** Electrical properties of SnO<sub>x</sub> thin film samples using different deposition temperatures during the growth. \* Very high resistivity.

#### d) Optoelectronic performance

The Figure of Merit which reveals the optoelectrical performance of the SnO<sub>x</sub> thin film samples were not calculated due to the very high resistivity of the SnO<sub>x</sub> thin films. Figure of Merit is expected to be extremely low.

#### 2.4.3. Summary

SnO<sub>x</sub> thin film samples, with different O<sub>2</sub>/Ar flow rate ratios during the growth, have shown very different structural, optical and electrical properties. At poor oxygen condition during the growth, SnO<sub>x</sub> were mainly amorphous SnO phase, the optical transmittance was equal or above 80%. However, the electrical resistivity was very high. At higher oxygen condition during the growth, the predominant

phase was SnO<sub>2</sub>, the optical transmittance varied from 76.8 to 89.7 % and the resistivity varied from  $2.84 \times 10^{-1}$  to  $3 \times 10^{-3} \Omega \cdot \text{cm}$ . All samples deposited at higher O<sub>2</sub>/Ar flow rate ratios were n-type conductive. The optoelectrical performance has improved while increasing the O<sub>2</sub>/Ar flow rate ratios.

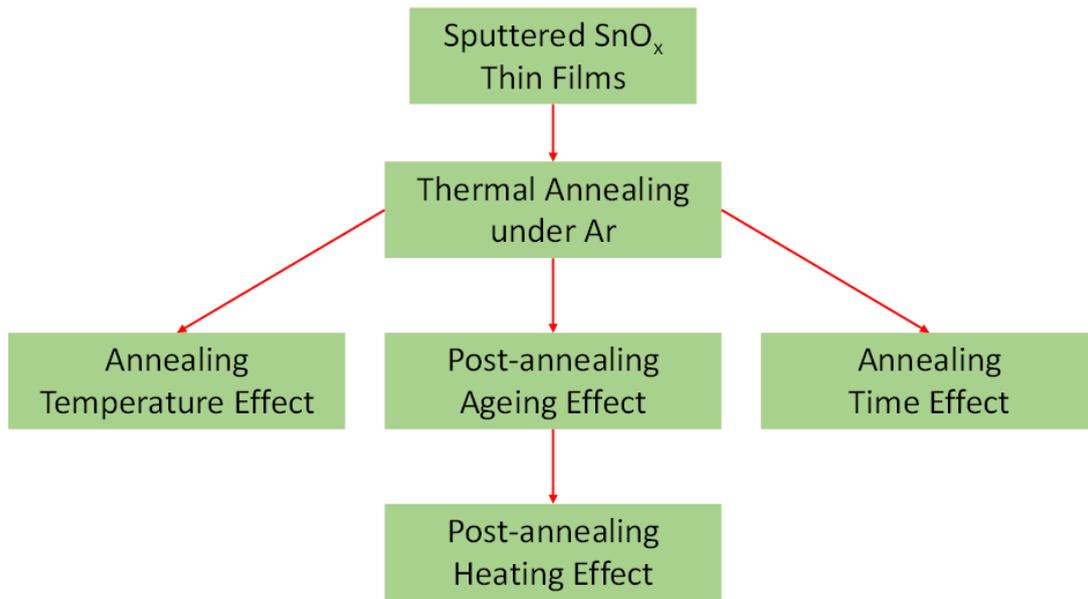
At relatively poor oxygen condition, SnO<sub>x</sub> thin film samples deposited at different temperature have shown the presence of a predominant amorphous SnO phase. Furthermore, SnO<sub>x</sub> deposited at 400°C has shown a small presence of SnO<sub>2</sub>. The optical properties have shown a relatively high optical transmittance and very similar bandgaps for all the SnO<sub>x</sub> which related SnO. However, the electrical resistivity was very high for all samples at different deposition temperatures<sup>[36]</sup>.

As per both studies related to O<sub>2</sub> flow effect and growth temperature, it can be concluded that the high electrical resistivity of the predominant SnO phase thin film is due to the amorphous microstructure regardless of the deposition temperature during the growth. The predominant amorphous SnO<sub>2</sub> phase is moderately resistive due to the amorphous microstructure. However, predominant crystalline SnO<sub>2</sub> thin film is highly conductive. Tuning the deposition temperature at poor oxygen condition during the growth of SnO<sub>x</sub> did not improve the crystalline microstructure of SnO<sub>x</sub> thin film samples. Thus, the electrical properties could not be improved. Improving additional parameters during the growth and/or post annealing might be required to achieve a crystalline SnO phase.

## **2.5. Effect of thermal annealing of sputtered SnO<sub>x</sub> under Ar atmosphere**

The objective of this part is to explore the possibility of achieving SnO<sub>x</sub> p-type conductivity through thermal annealing using inert gas, namely Ar, to avoid any reaction with the SnO<sub>x</sub> during the process. For this study, the SnO<sub>x</sub> films were synthesized using the method described in part 2.3. More specifically, the SnO<sub>x</sub> thin films were prepared using magnetron sputtering, manufactured by Torr Int. (described in the appendix in section 1.3.) with a high purity SnO (99.99%) 2"- diameter target (Manufactured by Codex International) on cleaned quartz substrates at a temperature of 250 °C, under different O<sub>2</sub>/Ar flow rates ratios where O<sub>2</sub>/Ar = 0.015, 0.035, 0.055, 0.065 and 0.075. The other deposition conditions were fixed, at 50 W as a power, 200 sccm of Ar flow rate and 30 min deposition duration time for all samples. Prior deposition, the base pressure was around  $6.66 \times 10^{-3}$  Pa while the deposition pressure was approximately 0.66 Pa as summarized in Table 15. The as-deposited samples were subject to different annealing processes as summarized in Diagram 2.

The thermal annealing process was carried out using a tube furnace (made by Carbolite, UK) and high purity Ar gas which is used as process gas. Diagram 2 below shows the summary of the applied thermal annealing processes. The annealing temperature varied from 400 to 600°C and duration time from 20 to 120 min. As the objective of this study is to explore the possibility of p-type conductivity, the main used characterization technique is the Hall effect probe.



**Diagram 2:** Different experiments conducted using Ar thermal annealing.

Sample	Substrate Temp (°C)	Power (W)	O <sub>2</sub> /Ar (O <sub>2</sub> :Ar sccm)	Depos. Pressure (Pa)	Deposition time (min)
S1	250	50	0.015 (3:200)	0.66	30
S2			0.035 (7:200)		
S3			0.055 (11:200)		
S4			0.065 (13:200)		
S5			0.075 (15:200)		

**Table 15:** Summary of the magnetron sputtering conditions of SnO<sub>x</sub> thin film samples.

### 2.5.1. Annealing temperature effect on SnO<sub>x</sub> thin films

The aim of this experiment using different annealing temperatures is to define if there is an effect of the annealing temperature on the electrical properties of SnO<sub>x</sub> thin films. The annealing processes temperatures were 400, 500 and 600 °C and all other parameters were kept fixed, namely Ar flow rate and annealing time as shown in Table 16.

Annealing	Flow rate (sccm)	Temperature (°C)	Time (min)
A400	57 (Ar)	400	50
A500	57 (Ar)	500	50
A600	57 (Ar)	600	50

**Table 16:** Summary of the annealing conditions of SnO<sub>x</sub> thin film samples.

The as-deposited samples have shown a high resistivity for O<sub>2</sub>/Ar ratio of 0.015 and 0.035. However, the SnO<sub>x</sub> samples were very high resistive for higher O<sub>2</sub>/Ar ratios and their conductivity type and electrical parameters could not be determined as shown in Table 17.

Furthermore, the thermal annealing at 400 °C in Ar atmosphere has improved remarkably the electrical conductivity for all samples to a moderate resistivity except for SnO<sub>x</sub> deposited at O<sub>2</sub>/Ar of 0.035 which has slightly improved the resistivity from  $2.4 \times 10^2$  to  $1.6 \times 10^2$  Ω.cm. Similarly, while the mobility of all samples are in the range of  $10^{-1}$  cm<sup>2</sup>/V.s, that of SnO<sub>x</sub> deposited at O<sub>2</sub>/Ar ratio of 0.035 has a high mobility of 2.6 cm<sup>2</sup>/V.s. The charge carrier concentration for all samples is relatively high in the range of  $10^{18}$  cm<sup>-3</sup> whereas SnO<sub>x</sub> deposited at O<sub>2</sub>/Ar ratio of 0.035 exhibited a lower concentration of about  $1.5 \times 10^{16}$  cm<sup>-3</sup>. All thermally annealed SnO<sub>x</sub> samples at 400 °C have n-type conductivity. The improvement in conductivity is mainly related to the improvement of the n-type charge carrier concentration at least for the SnO<sub>x</sub> thin films deposited at O<sub>2</sub>/Ar ratio of 0.015 and 0.035.

The thermal annealing at 500 °C has improved the electrical resistivity for all samples. However, this improvement is less important compared to the thermal annealing at 400 °C. The SnO<sub>x</sub> samples deposited at O<sub>2</sub>/Ar ratio of 0.015 and 0.035, and annealed at 500 °C, have demonstrated a n-type conductivity with electrical resistivity of 22 and 69 Ω.cm, respectively. The SnO<sub>x</sub> thin films deposited at O<sub>2</sub>/Ar ratio of 0.055 and 0.065, and annealed at 500 °C, have shown a p-type conductivity with electrical resistivity of  $2 \times 10^2$  and  $2.2 \times 10^2$  Ω.cm, respectively. Increasing further the O<sub>2</sub>/Ar ratio to 0.075 followed by an annealing at 500 °C led to a mixture of n-type and p-type conductivity with electrical resistivity of 42 Ω.cm. As for the carriers mobility of annealed samples, it decreased for SnO<sub>x</sub> film deposited at O<sub>2</sub>/Ar ratio of 0.015 while it slightly increased for the SnO<sub>x</sub> deposited at O<sub>2</sub>/Ar of 0.035. Increasing the O<sub>2</sub>/Ar ratio to 0.055 and 0.065 led to relatively good mobility of 8.3 and 14 cm<sup>2</sup>/V.s, respectively. The mobility is found to be as low as 0.11 cm<sup>2</sup>/V.s for SnO<sub>x</sub> deposited at O<sub>2</sub>/Ar ratio of 0.075 followed by thermal annealing. The n-type charge carrier concentration has relatively improved for SnO<sub>x</sub> deposited at O<sub>2</sub>/Ar ratio of 0.015 and 0.035 after thermal annealing at 500 °C while the conductivity type of SnO<sub>x</sub> deposited at O<sub>2</sub>/Ar ratio of 0.055 and 0.065 have changed from n-type to p-type after this thermal annealing.

The thermal annealing at 600 °C has slightly improved the electrical resistivity for SnO<sub>x</sub> thin films deposited at O<sub>2</sub>/Ar of 0.015 and 0.035, and significantly improved the electrical resistivity for the rest of the samples. However, this thermal annealing has enabled p-type conductivity for all samples except for the oxygen rich grown SnO<sub>x</sub> deposited at O<sub>2</sub>/Ar of 0.075. The electrical resistivity for all samples is between  $1.1 \times 10^2$  and  $2 \times 10^2$  Ω.cm which is relatively high. The mobility is ranging between 0.13 to 1.4 cm<sup>2</sup>/V.s. The charge carrier concentration ranges from  $3 \times 10^{16}$  to  $4.5 \times 10^{17}$  cm<sup>-3</sup> for the p-type

conductive SnO<sub>x</sub> thin films, while it is  $2.5 \times 10^{16} \text{ cm}^{-3}$  for the n-type conductive SnO<sub>x</sub>. All electrical measurements data for this post-annealing study are reported in Table 17.

Charge carrier concentration ( $\text{cm}^{-3}$ )

O <sub>2</sub> /Ar	As-deposited	Annealed at 400°C 50 minutes	Annealed at 500°C 50 minutes	Annealed at 600°C 50 minutes
0.015	-1.34E+15	-2.60E+18	-7.30E+17	2.30E+17
0.035	N/A**	-1.50E+16	-3.40E+16	3.00E+16
0.055	N/A*	-3.90E+18	4.00E+15	4.50E+17
0.065	N/A*	-3.20E+18	2.00E+15	4.40E+16
0.075	N/A*	-3.00E+18	N/A**	-2.50E+16

Resistivity ( $\Omega\cdot\text{cm}$ )

O <sub>2</sub> /Ar	As-deposited	Annealed at 400°C 50 minutes	Annealed at 500°C 50 minutes	Annealed at 600°C 50 minutes
0.015	2.00E+02	2.70E+00	2.20E+01	1.90E+02
0.035	2.39E+02	1.60E+02	6.90E+01	1.60E+02
0.055	N/A*	3.50E+00	2.00E+02	1.40E+02
0.065	N/A*	3.20E+00	2.20E+02	1.10E+02
0.075	N/A*	3.20E+00	4.20E+01	2.00E+02

Mobility ( $\text{cm}^2/\text{V}\cdot\text{s}$ )

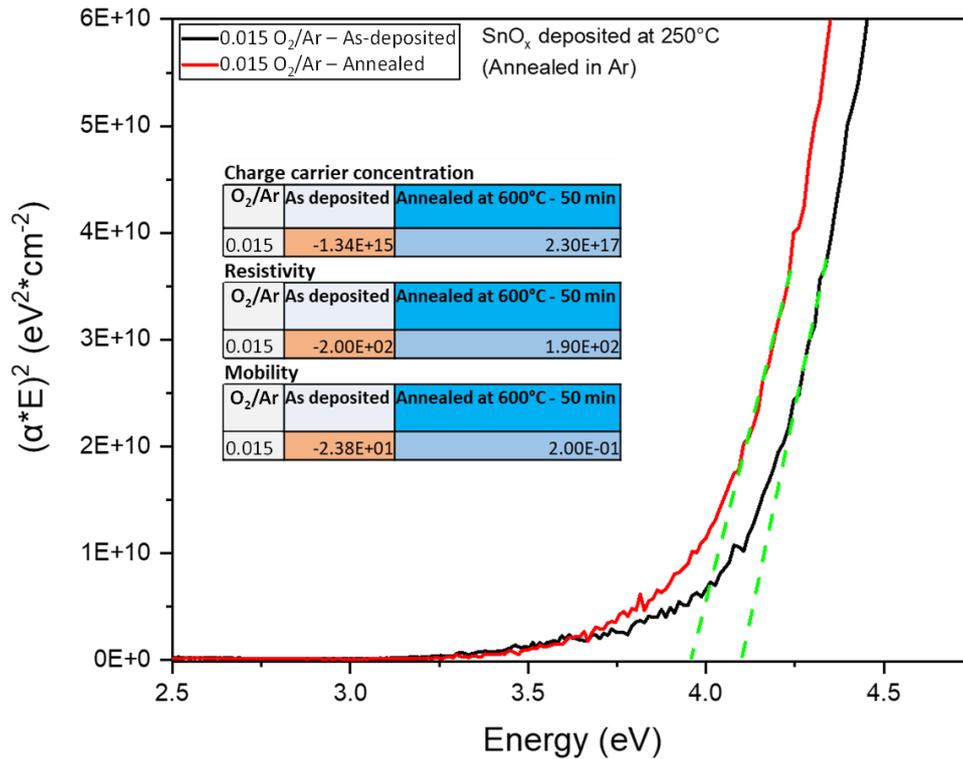
O <sub>2</sub> /Ar	As-deposited	Annealed at 400°C 50 minutes	Annealed at 500°C 50 minutes	Annealed at 600°C 50 minutes
0.015	2.38E+01	8.90E-01	4.00E-01	2.00E-01
0.035	4.48E+00**	2.60E+00	4.40E+00	1.40E+00
0.055	N/A*	4.70E-01	8.30E+00	1.30E-01
0.065	N/A*	6.20E-01	1.40E+01	1.30E+00
0.075	N/A*	6.50E-01	1.10E-01**	1.30E+00

**Table 17:** Hall effect measurement for as deposited and thermally annealed SnO<sub>x</sub>.

\* Values with orange background represent n-type conductive samples and blue represent p-type conductive samples. Values with white background represent the samples where the conductivity type could not be determined. The samples with the mark. N/A\*: high resistivity SnO<sub>x</sub>. \*\*: conductivity type could not be determined.

To sum up, the thermal annealing under inert atmosphere namely Ar, has improved the n-type conductivity at 400 °C and converted the n-type conductivity to p-type for four samples out of five at the annealing temperature of 600 °C. It is clear that increasing the annealing temperature has enabled the conductivity type conversion from n-type to p-type through enabling the formation of hole charge carrier due to the phase change from SnO<sub>2</sub> to SnO. The low charge carrier concentration is likely due to the charge compensation effect<sup>[13, 18, 19, 34, 80-82]</sup>.

SnO<sub>x</sub> thin film deposited at O<sub>2</sub>/Ar ratio of 0.015 has shown a clear n-type conductivity and has changed its conductivity type to p-type upon thermal annealing at 600 °C for 50 min. In order to understand the changes, UV-Visible spectra related Tauc plot calculation was conducted on both as deposited and thermally annealed SnO<sub>x</sub> deposited O<sub>2</sub>/Ar ratio of 0.015 as depicted in Figure 17. Tauc plot shows that the as-deposited SnO<sub>x</sub> is mainly SnO<sub>2</sub> as the optical bandgap is above 4 eV and the optical bandgap annealed SnO<sub>x</sub> is less than 4 eV which indicates the presence of SnO phase that decreases the optical bandgap. After thermal annealing at 600 °C, there is a clear redshift of the SnO<sub>x</sub> optical bandgap. As discussed previously in Chapter 1, thermal annealing at high temperature is expected to reduce the disorder which increases the optical bandgap. However, in this case the optical bandgap has reduced to a lower value below 4 eV which indicates the formation of SnO phase during the thermal annealing<sup>[35, 36, 51, 79]</sup>.

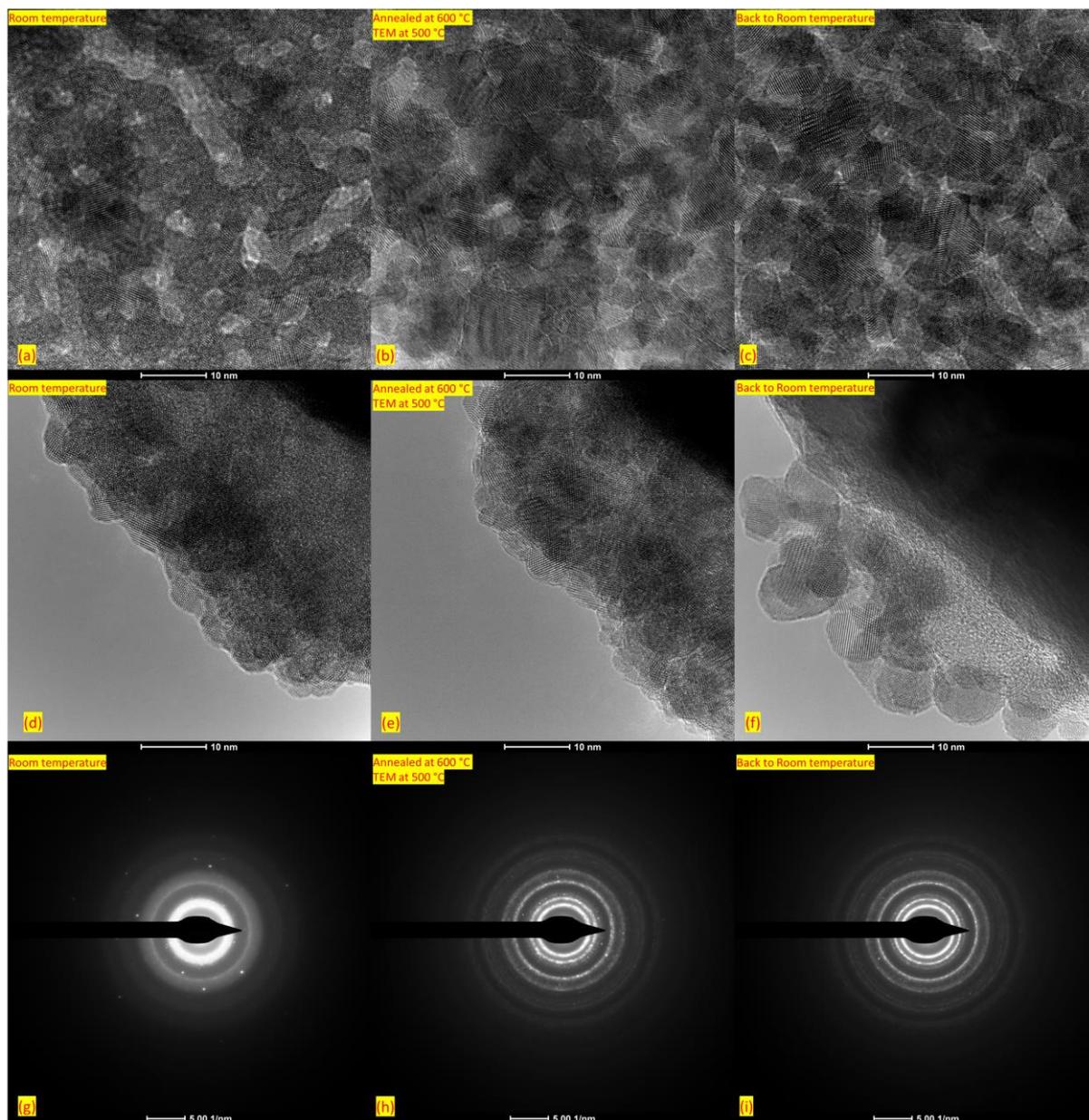


**Figure 17:** Tauc plot of as deposited and thermally annealed at 600°C SnO<sub>x</sub> deposited O<sub>2</sub>/Ar ratio of 0.015.

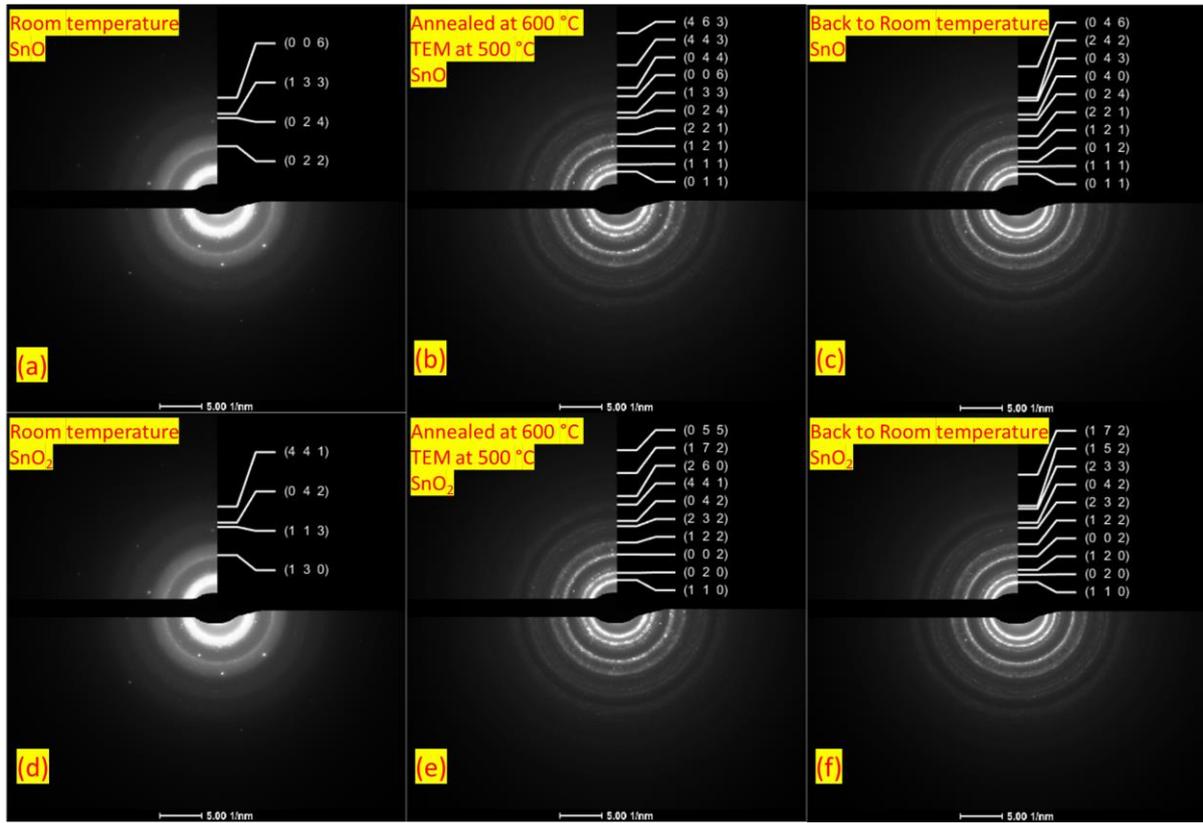
To understand the changes at the crystalline structure, High Resolution - Transmission Electron Microscopy (TEM) in-situ thermal annealing was performed on the SnO<sub>x</sub> prepared at 0.015 O<sub>2</sub>/Ar ratio where the substrate temperature was set at 250°C. The in-situ study consists of performing the TEM imaging and diffraction for as-deposited SnO<sub>x</sub> followed by increasing the temperature from room temperature to 600°C for 10 minutes to complete the annealing process, and then reducing the temperature to 500°C to be able to perform the TEM imaging and diffraction at high temperature (maximum 500°C), and finally slowing down to room temperature to access the changes after the annealing process. The TEM images are shown in Figure 18. It is worth noting that high resolution imaging and diffraction of this TEM can be performed at maximum temperature of 500°C without the cooling system. However, the cooling system is required for higher temperatures, and it induces a certain level of vibration that reduces the resolution of TEM.

As expected, the as-deposited SnO<sub>x</sub> sample shows an amorphous microstructure which is confirmed by the TEM images in Figure 18 (a, d) as well as the diffraction pattern in Figure 18 (g) due to the presence of large diffraction bands along. Nevertheless, the presence of some diffraction spots indicates the low crystallinity of SnO<sub>x</sub> thin film. Upon thermal annealing and after cooling back to 500°C, there is clearly an improvement in crystallization of SnO<sub>x</sub> due to the presence of much larger crystallites in Figure 18 (b) and by clear atomic planes in Figure 18 (e). The diffraction pattern has revealed a polycrystalline structure of SnO<sub>x</sub> as witnessed by the presence of multiple sharp rings as well as some

bright spots mainly within these rings. Upon returning to room temperature, the crystalline structure has been kept after the thermal annealing as shown in Figure 18 (c, f). TEM diffraction pattern has also revealed a polycrystalline structure by the presence of sharp rings. These rings are slightly thinner compared to the ones recorded at 500°C<sup>[83, 84]</sup>. The diffraction pattern analysis was performed using diffractGUI to define the analysis zone axis and ringGUI to determine the crystallographic planes related to the rings in the diffraction pattern as shown in figure 19 and table 18. The diffraction pattern of SnO<sub>x</sub> thin film at 500°C (after thermal annealing at 600°C) and after returning to room temperature is closely matching with SnO phase theoretical positions. This indicates the presence of polycrystalline SnO. These results also show that the amorphous as-deposited SnO<sub>x</sub> has been crystallized to SnO and it has kept this SnO phase after the thermal annealing process. One of the lessons from this In-Situ TEM study is that the formed SnO at high temperature is unlikely to be subject to further oxidation while returning back to room temperature in the presence of O<sub>2</sub> as SnO<sub>x</sub> sample was constantly under high vacuum in TEM system. This might also explain the non-stability of SnO<sub>x</sub> after thermal annealing as well as the ageing affect<sup>[85-88]</sup>.



**Figure 18:** In-Situ Transmission Electron Microscopy imaging and diffraction: (a, d, g) as deposited SnO<sub>x</sub>, (b, e, h) in-situ at 500 °C, (c, f, i) back to room temperature.



**Figure 19:** In-Situ Transmission Electron Microscopy diffraction analysis: for SnO indexation (a) as deposited SnO<sub>x</sub>, (b) annealed SnO<sub>x</sub> at 600 °C and returned to 500 °C, (c) annealed SnO<sub>x</sub> at 600 °C and back to room temperature. for SnO<sub>2</sub> indexation (d) as deposited SnO<sub>x</sub>, (e) annealed SnO<sub>x</sub> at 600 °C and returned to 500 °C, (f) annealed SnO<sub>x</sub> at 600 °C and back to room temperature.

		SnO Phase Indexation				SnO <sub>2</sub> Phase Indexation						
Room temperature	Ring identification								Ring identification			
	Plane	Radius [1/nm]		d-spacing [nm]		Plane	Radius [1/nm]		d-spacing [nm]			
		theor.	measured	theor.	measured		theor.	measured	theor.	measured		
	(0 2 2)	6.691	6.499	0.149	0.154	(1 3 0)	6.675	6.499	0.15	0.154		
	(0 2 4)	9.802	9.892	0.102	0.101	(1 1 3)	9.877	9.892	0.101	0.101		
(1 3 3)	10.376	10.418	0.096	0.096	(0 4 2)	10.521	10.418	0.095	0.096			
(0 0 6)	12.407	12.33	0.081	0.081	(4 4 1)	12.347	12.33	0.081	0.081			
Annealed at 600 °C & TEM Diffraction at 500 °C	Ring identification								Ring identification			
	Plane	Radius [1/nm]		d-spacing [nm]		Plane	Radius [1/nm]		d-spacing [nm]			
		theor.	measured	theor.	measured		theor.	measured	theor.	measured		
	(0 1 1)	3.346	3.345	0.299	0.299	(1 1 0)	2.985	3.345	0.335	0.299		
	(1 1 1)	4.256	4.205	0.235	0.238	(0 2 0)	4.222	4.205	0.237	0.238		
(1 2 1)	6.234	6.404	0.16	0.156	(0 0 2)	6.277	6.404	0.159	0.156			
(2 2 1)	7.721	7.837	0.13	0.128	(1 2 2)	7.853	7.837	0.127	0.128			

	(0 2 4)	9.802	9.797	0.102	0.102	(2 3 2)	9.865	9.797	0.101	0.102
	(1 3 3)	10.376	10.418	0.096	0.096	(0 4 2)	10.521	10.418	0.095	0.096
	(0 0 6)	12.407	12.377	0.081	0.081	(4 4 1)	12.347	12.377	0.081	0.081
	(0 4 4)	13.383	13.381	0.075	0.075	(2 6 0)	13.351	13.381	0.075	0.075
	(4 4 3)	16.12	16.105	0.062	0.062	(1 7 2)	16.192	16.153	0.062	0.062
	(4 6 3)	19.955	19.928	0.05	0.05	(0 5 5)	18.911	18.877	0.053	0.053
Back to Room temperature	Ring identification					Ring identification				
	Plane	Radius [1/nm]		d-spacing [nm]		Plane	Radius [1/nm]		d-spacing [nm]	
		theor.	measured	theor.	measured		theor.	measured	theor.	measured
	(0 1 1)	3.346	3.345	0.299	0.299	(1 1 0)	2.985	3.345	0.335	0.299
	(1 1 1)	4.256	4.253	0.235	0.235	(0 2 0)	4.222	4.253	0.237	0.235
	(0 1 2)	4.901	4.827	0.204	0.207	(1 2 0)	4.72	4.827	0.212	0.207
	(1 2 1)	6.234	6.452	0.16	0.155	(0 0 2)	6.277	6.452	0.159	0.155
	(2 2 1)	7.721	7.885	0.13	0.127	(1 2 2)	7.853	7.885	0.127	0.127
	(0 2 4)	9.802	9.845	0.102	0.102	(2 3 2)	9.865	9.845	0.101	0.102
	(0 4 0)	10.521	10.466	0.095	0.096	(0 4 2)	10.521	10.466	0.095	0.096
	(0 4 3)	12.214	12.138	0.082	0.082	(2 3 3)	12.107	12.138	0.083	0.082
	(2 4 2)	12.468	12.473	0.08	0.08	(1 5 2)	12.46	12.473	0.08	0.08
(0 4 6)	16.267	16.2	0.061	0.062	(1 7 2)	16.192	16.2	0.062	0.062	

**Tables 18:** In-Situ Transmission Electron Microscopy diffraction rings indexation for SnO and SnO<sub>2</sub>.

The p-type conduction character reached at higher annealing temperature, revealed by Hall effect measurements, can be attributed to the formation of SnO phase which is known to form hole related charge carriers and enable the p-type conductivity. In the next experiment, the duration of the thermal annealing at 600 °C will be assessed.

### 2.5.2. Annealing time effect on SnO<sub>x</sub> thin films

The aim of this experiment is to define if there is an effect of the annealing time on the electrical properties of SnO<sub>x</sub> thin films. The thermal annealing durations were 20, 35, 50 (already established) and 120 minutes while all other parameters were kept fixed, namely Ar flow rate and annealing temperature at 600 °C as shown in Table 19.

Annealing	Flow rate (sccm)	Temperature (°C)	Time (min)
AT020	57 (Ar)	600	20
AT035	57 (Ar)	600	35
AT050	57 (Ar)	600	50
AT120	57 (Ar)	600	120

**Table 19:** Summary of the annealing conditions of SnO<sub>x</sub> thin film samples.

The SnO<sub>x</sub> samples deposited at O<sub>2</sub>/Ar flow rate ratio of 0.015 and 0.035, and thermally annealed at 600 °C for 20 minutes have shown a relatively high resistivity similar to the as-deposited SnO<sub>x</sub> ones with a very minor change. In contrast, the SnO<sub>x</sub> deposited at O<sub>2</sub>/Ar ratio of 0.065 and 0.075, and thermal annealed at 600°C for 20 minutes have shown an improvement in electrical conductivity with resistivity values around  $2.2 \times 10^2 \Omega\cdot\text{cm}$  for both samples. The SnO<sub>x</sub> thin film deposited at O<sub>2</sub>/Ar ratio of 0.035, has kept its high resistivity after the thermal annealing at 600 °C for 20 minutes and its conductivity type could not be determined as well as the mobility and the charge carrier concentration. The SnO<sub>x</sub> thin films deposited at O<sub>2</sub>/Ar ratio of 0.015, 0.065 and 0.075, and thermally annealed at 600 °C are found to have p-type conductivity, with hole carriers' mobility ranging from 2.2 to 6.3 cm<sup>2</sup>/V·s and charge carrier concentration ranging from  $5.1 \times 10^{15}$  to  $2.2 \times 10^{16} \text{ cm}^{-3}$ . Surprisingly, the sample deposited at 0.035 and annealed at 600 °C is found to have n-type conductivity, an electron mobility of 10 cm<sup>2</sup>/V·s and a charge carrier concentration of  $2.9 \times 10^{15} \text{ cm}^{-3}$ .

For SnO<sub>x</sub> thin films thermally annealed at 600 °C for 35 minutes and deposited at O<sub>2</sub>/Ar ratio of 0.015, 0.055 and 0.075 were found to be p-type conductive, with charge carrier concentration of  $1.5 \times 10^{16}$ ,  $8.2 \times 10^{17}$  and  $10^{16} \text{ cm}^{-3}$ , respectively. Their resistivity values were found to range between  $1.7 \times 10^2$  and  $1.9 \times 10^2 \Omega\cdot\text{cm}$ . The mobility is 3.3 and 3.4 cm<sup>2</sup>/V·s for the SnO<sub>x</sub> films deposited at O<sub>2</sub>/Ar ratio of 0.015 and 0.075 while it is as low as 0.07 cm<sup>2</sup>/V·s when the O<sub>2</sub>/Ar ratio is 0.055. Surprisingly, the SnO<sub>x</sub> thin films deposited at O<sub>2</sub>/Ar ratio of 0.035 and 0.065 followed by the thermal annealing process, are found to have a n-type conductivity, their electrical resistivity is respectively  $1.9 \times 10^2$  and  $1.3 \times 10^2 \Omega\cdot\text{cm}$ , their charge carrier concentration is respectively  $2.1 \times 10^{16}$  and  $1.4 \times 10^{17} \text{ cm}^{-3}$ , and their mobility is respectively 1.7 and 0.35 cm<sup>2</sup>/V·s.

For the thermal annealing at 600 °C for 120 minutes, all SnO<sub>x</sub> samples have shown a n-type conductivity except the SnO<sub>x</sub> grown at O<sub>2</sub>/Ar ratio of 0.075 whose the conductivity type could not be determined. The measured electrical resistivity for SnO<sub>x</sub> thin films deposited at O<sub>2</sub>/Ar flow ratio of 0.015, 0.035, 0.055, and 0.065 followed by the thermal annealing is 78.7, 98.7, 26.5, and 29.2 Ω·cm, the charge carrier concentration is  $2.3 \times 10^{16}$ ,  $8.4 \times 10^{15}$ ,  $2.4 \times 10^{18}$  and  $2 \times 10^{18} \text{ cm}^{-3}$ , and the mobility is 3.7, 7.6, 0.099 and 0.14 cm<sup>2</sup>/V·s, respectively. Thus, the prolonged thermal annealing at 600 °C of two hours has led to a counter effect and kept the n-type conductivity while shorter duration annealing less than 50 minutes did not have a clear effect. All the results are shown in Table 20.

Charge carrier concentration (cm<sup>-3</sup>)

O <sub>2</sub> /Ar	As deposited	Annealed at 600°C 20 minutes	Annealed at 600°C 35 minutes	Annealed at 600°C 50 minutes	Annealed at 600°C 120 minutes
0.015	-1.34E+15	2.17E+16	1.51E+16	2.30E+17	-2.27E+16
0.035	N/A**	-2.91E+15	-2.13E+16	3.00E+16	-8.41E+15
0.055	N/A*	N/A*	8.24E+17	4.50E+17	-2.41E+18
0.065	N/A*	5.13E+15	-1.37E+17	4.40E+16	-1.95E+18
0.075	N/A*	1.38E+16	1.01E+16	-2.50E+16	N/A**

Resistivity (Ω·cm)

O <sub>2</sub> /Ar	As deposited	Annealed at 600°C 20 minutes	Annealed at 600°C 35 minutes	Annealed at 600°C 50 minutes	Annealed at 600°C 120 minutes
0.015	2.00E+02	2.43E+02	1.68E+02	1.90E+02	7.87E+01
0.035	2.39E+02**	2.42E+02	1.91E+02	1.60E+02	9.87E+01
0.055	N/A*	N/A*	1.66E+02	1.40E+02	2.65E+01
0.065	N/A*	2.18E+02	1.31E+02	1.10E+02	2.92E+01
0.075	N/A*	2.24E+02	1.86E+02	2.00E+02	3.32E+01

Mobility (cm<sup>2</sup>/V·s)

O <sub>2</sub> /Ar	As deposited	Annealed at 600°C 20 minutes	Annealed at 600°C 35 minutes	Annealed at 600°C 50 minutes	Annealed at 600°C 120 minutes
0.015	2.38E+01	3.07E+00	3.27E+00	2.00E-01	3.69E+00
0.035	4.48E+00**	1.03E+01	1.67E+00	1.40E+00	7.56E+00
0.055	N/A*	N/A*	6.90E-02	1.30E-01	9.94E-02
0.065	N/A*	6.27E+00	3.53E-01	1.30E+00	1.44E-01
0.075	N/A*	2.18E+00	3.40E+00	1.30E+00	6.84E-01

**Table 20:** Hall effect measurement for as deposited and thermally annealed SnO<sub>x</sub>.

\* Values with orange background represent n-type conductive samples and blue represent p-type conductive samples. Values with white background represent the samples where the conductivity type could not be determined. The samples with the mark. N/A\*: high resistivity SnO<sub>x</sub>. \*\*: conductivity type could not be determined.

### 2.5.3. Post-annealing ageing and heating effect on SnO<sub>x</sub> thin films

The stability of SnO<sub>x</sub> in terms of electrical properties is very important to assess optoelectrical performance of the SnO<sub>x</sub> samples. Therefore, a set of three samples was kept in air for seven days before performing a second Hall effect measurement as shown in Table 21.

All three samples have shown a better electrical conductivity after seven days of ageing while being exposed to air. The electrical resistivity is found to be 16.7, 22.5 and 26.7 Ω·cm for SnO<sub>x</sub> samples deposited at O<sub>2</sub>/Ar ratio of 0.055, 0.065 and 0.075, thermally annealed at 600 °C for 50 minutes and aged in air for seven days, respectively. All samples have shown an n-type conductivity including two SnO<sub>x</sub> thin films which were p-type before the ageing. The corresponding electron mobility is 0.16, 0.17 and 0.17 cm<sup>2</sup>/V·s and the charge carrier concentration is found to be 2.4 × 10<sup>18</sup>, 1.7 × 10<sup>18</sup> and 1.5 × 10<sup>18</sup> cm<sup>-3</sup>, respectively.

It is worth noting that the charge carrier concentration has increased significantly for all samples which is the main factor for improving the electrical conductivity. The mobility has decreased for the SnO<sub>x</sub> thin films deposited at O<sub>2</sub>/Ar ratio of 0.065 and 0.075, thermally annealed at 600 °C for 50 minutes and aged in air for seven days while it increased slightly for the other sample grown at 0.055 and thermally annealed at 600°C. The increase in the electrical conductivity can be related to the further oxidation which took place in the SnO<sub>x</sub> in presence of O<sub>2</sub> in air. This oxidation is expected to change the crystalline structure towards mainly SnO<sub>2</sub>. The intrinsic formation of oxygen vacancies has led to form electron charge carriers within the SnO<sub>x</sub> which compensate the presence of hole charge carriers and eventually improve the electrical conductivity by the increasing of n-type carriers.

The air ambient also contains moisture which can affect the chemistry of the SnO<sub>x</sub> thin films. Therefore, we carried out an experiment to evaluate if there is a reversible effect of exposure to moisture within the SnO<sub>x</sub> thin films. Thus, the aged three samples were subject to a heating process at 150 °C for 50 minutes. Similar to previous studies, the resistivity, mobility and charge carrier concentration for SnO<sub>x</sub> thin films deposited at O<sub>2</sub>/Ar ratio of 0.055, 0.065 and 0.075, annealed at 600 °C for 50 minutes, aged for seven days in air and finally heated at 150 °C for 50 minutes in air as shown in Table 23. All three samples have kept their n-type conductivity character and show a better electrical conductivity after the heating. The electrical resistivities of the treated samples are found to be 1.1,

1.4 and 1.7  $\Omega\cdot\text{cm}$ , and the corresponding charge carrier concentrations are measured to be  $1.1 \times 10^{19}$ ,  $5 \times 10^{18}$  and  $5.3 \times 10^{18} \text{ cm}^{-3}$ , for  $\text{O}_2/\text{Ar}$  flow rate of 0.055, 0.065 and 0.075, respectively.

The heating, in presence of air, which might have reduced the moisture adsorption in the  $\text{SnO}_x$  thin films, has likely increased the oxidation state of the  $\text{SnO}_x$  towards  $\text{SnO}_2$  which led to a clear improvement in terms of mobility and charge carrier concentration. This improvement is reflected in the electrical resistivity which has decreased significantly.

Charge carrier concentration ( $\text{cm}^{-3}$ )

$\text{O}_2/\text{Ar}$	As deposited	Annealed at 600°C 50 minutes	Annealed at 600°C for 50 min Aged for 7 days	Annealed at 600°C for 50 min Aged for 7 days heating at 150°C
0.055	N/A*	4.50E+17	-2.36E+18	-1.11E+19
0.065	N/A*	4.40E+16	-1.72E+18	-4.97E+18
0.075	N/A*	-2.50E+16	-1.46E+18	-5.29E+18

Resistivity ( $\Omega\cdot\text{cm}$ )

$\text{O}_2/\text{Ar}$	As deposited	Annealed at 600°C 50 minutes	Annealed at 600°C for 50 min Aged for 7 days	Annealed at 600°C for 50 min Aged for 7 days heating at 150°C
0.055	N/A*	1.40E+02	1.67E+01	1.12E+00
0.065	N/A*	1.10E+02	2.25E+01	1.36E+00
0.075	N/A*	2.00E+02	2.67E+01	1.68E+00

Mobility ( $\text{cm}^2/\text{V}\cdot\text{s}$ )

$\text{O}_2/\text{Ar}$	As deposited	Annealed at 600°C - 50 minutes	Annealed at 600°C for 50 min Aged for 7 days	Annealed at 600°C for 50 min Aged for 7 days heating at 150°C
0.055	N/A*	1.30E-01	1.61E-01	5.95E-01
0.065	N/A*	1.30E+00	1.66E-01	9.23E-01
0.075	N/A*	1.30E+00	1.65E-01	7.12E-01

**Table 21:** Hall effect measurement for as deposited, thermally annealed, aged, and heated  $\text{SnO}_x$  samples.

\* Values with orange background represent n-type conductive samples and blue represent p-type conductive samples. Values with white background represent the samples where the conductivity type could not be determined. The samples with the mark. N/A\*: high resistivity  $\text{SnO}_x$ .

#### 2.5.4. Summary

The study of thermal annealing under Ar atmosphere has shown that the higher temperature of 500 and 600°C has enabled p-type conductive samples while for treatment at 400°C no p-type conductive

samples. This has confirmed the importance of both temperature and ambient during the annealing. It is also confirmed that the p-type character is correlated to the presence of SnO phase thanks to the transmission electron microscopy diffraction and to the optical measurements that revealed lower values of the optical bandgap. The annealing time experiments under Ar ambient has shown that p-type films are obtained for shorter time of annealing while extending the annealing duration to two hours conducted to the reverse effect where all samples become n-type.

On the other hand, the ageing effect under air ambient was assessed after seven days. We show that treatment have converted the SnO<sub>x</sub> films from p-type to n-type conductivity and their resistivity has decreased. This could be explained by the further oxidation of SnO<sub>x</sub> films from SnO towards SnO<sub>2</sub> and the increasing formation of the intrinsic n-type charge carriers which compensated the p-type charge carriers and increases the n-type carriers which reduced the resistivity. Finally, the effect of presence of moisture on tin oxide after seven days of the ageing has been assessed using heating process at 150°C in the presence of air. This process did not recover the p-type conductivity and instead it has led to further oxidation of SnO<sub>x</sub> film and further decrease in resistivity by forming more n-type charge carriers.

## **2.6. General summary**

This chapter show that it is possible to tune the structural, optical, and electrical properties of sputtered SnO<sub>x</sub> films by optimising the growth conditions. This work has led to the synthesis of high conductive SnO<sub>2</sub> thin films with high optical transmittance which has enabled achieving a high optoelectronic performance as revealed by FoM. Thermal annealing for SnO<sub>x</sub> thin films under Ar at a relatively high temperature around 600°C has shown a potential ability to convert n-type conductivity to p-type. Material properties of SnO<sub>x</sub> before and after the annealing have shown that this conductivity is related to the formation of tin monoxide (SnO) phase.

In the next chapter, the performance assessment of perovskite solar cells using SnO<sub>2</sub> as Electron Transport Layer will be explored.

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# **Chapter 3: Study of Tin Oxide (SnO<sub>2</sub>) Material as Electron Transport Layer (ETL) for Perovskite Solar Cells**

### 3. Study of Tin Oxide (SnO<sub>2</sub>) Material as Electron Transport Layer (ETL) for Perovskite Solar Cells

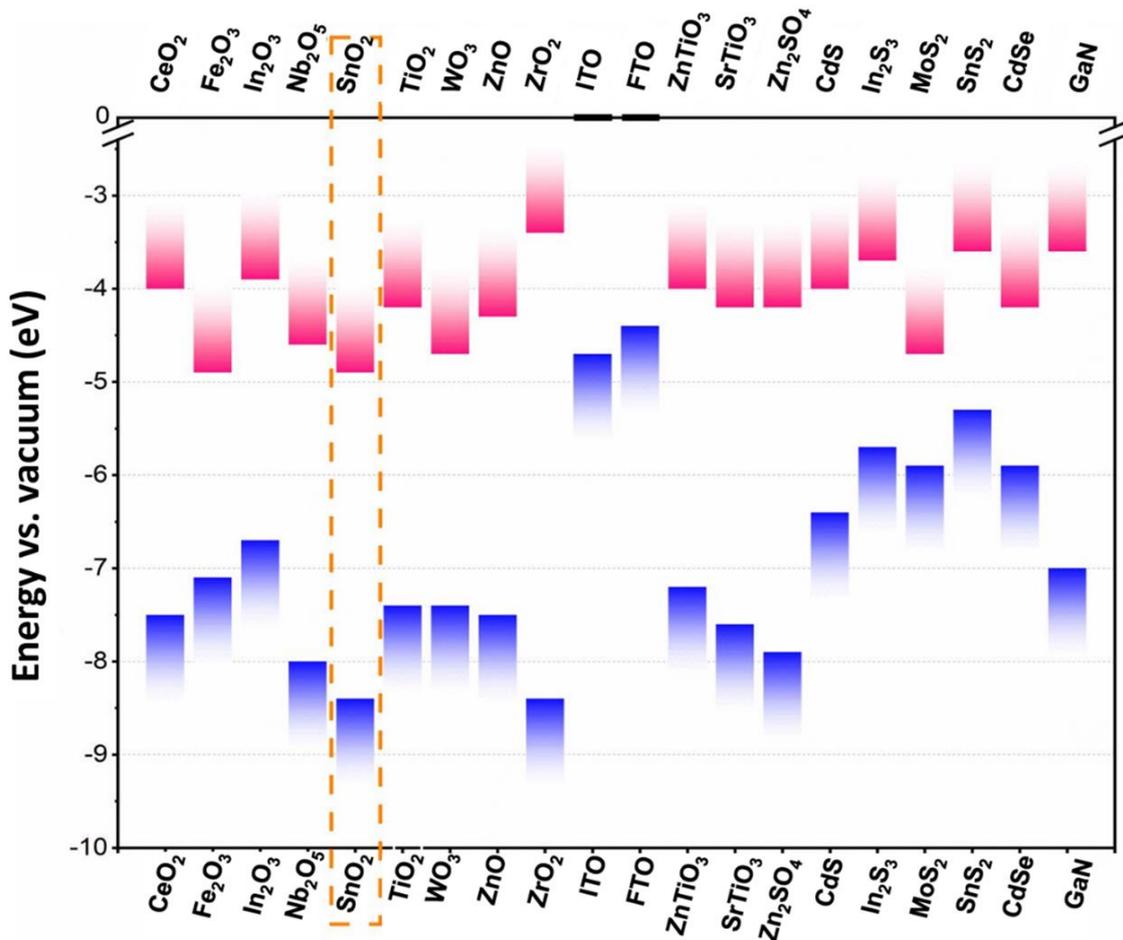
#### 3.1. Introduction

The performance and cost-effectiveness fabrication of the perovskite solar cells (PSCs) are the two main assets which are increasingly attracting academic and industrial attention. Certified Power Conversion Efficiency (PCE) for the best solar cell efficiency has shown a 25.7% for PSCs as achieved by UNIST in 2021<sup>[1]</sup>. Focus is put nowadays on the PSCs commercialization <sup>[2]</sup>, and this aim is still facing two main challenges, namely a descent device operational-stability and the fabrication scalability. The stability of the PSCs has been the cornerstone of extensive research and development over the last years. Nevertheless, this research effort has been found to be one of the most complex physico-chemical issues that involves multiple factors and various physical phenomena. These issues are also a subject of the device configuration and materials' characteristics. In fact, the device stability can directly be affected by the electrode material and its characteristics (work function, dimensions, etc.) <sup>[3]</sup>, electron transport layer (ETL) and hole transport layer (HTL) properties <sup>[4, 5]</sup>, the nature of the interface between the absorber-perovskite layer and the charge transport materials <sup>[6]</sup>, and indeed, the stability of the perovskite material itself <sup>[7]</sup>. In 2016, Ahn et al. <sup>[8]</sup> proposed that the ETL based on TiO<sub>2</sub> is among the most responsible factors for the light-induced degradation in PSCs. This suggestion was also supported by the research outcome of Qiu and coworkers in 2018 <sup>[9]</sup>. On the other hand, SnO<sub>2</sub> as ETL has demonstrated its capability to replace the conventional TiO<sub>2</sub> due to the fact that a power conversion efficiency (PCE) of more than 21% has been already achieved using SnO<sub>2</sub> ETL <sup>[10]</sup>. SnO<sub>2</sub> shows several benefits over TiO<sub>2</sub>, including a higher electron mobility and an excellent energy level matching <sup>[11]</sup>. More importantly, SnO<sub>2</sub> as ETL is highly efficient against the perovskite solar cells degradation, which is induced by TiO<sub>2</sub> ETL, thereby considerably improving the device operational lifetime under continuous light illumination at the maximum power point. In this context, Christians et al. <sup>[12]</sup> have recently demonstrated a much longer lifetime with un-encapsulated perovskite solar cells based on SnO<sub>2</sub> as ETL compared to TiO<sub>2</sub>. On the other hand, the second big challenge deals with the scalability of the PSC fabrication, to reach the module scale (i.e. perovskite solar modules (PSMs)), while maintaining performance similar to PSCs of small areas <sup>[2]</sup>. As large-scale thin film growth processes for the PSCs fabrication have been introduced, the number of reports related to PSMs has drastically increased <sup>[13]</sup>. For instance, M. Green et al. have reported a PCE of 16 % with an aperture area (AA) of 16.29 cm<sup>2</sup> <sup>[14]</sup> and H. Chen et al. have achieved a certified PCE of 12.1% with a larger AA of 36.1 cm<sup>2</sup> <sup>[15]</sup>. Other key parameters are related to the cost-effectiveness and large-scale deposition processes of ETL <sup>[2]</sup>. Currently, the majority of PSMs are based on TiO<sub>2</sub> as ETL, which requires high processing temperature. TiO<sub>2</sub> is also the origin of many instability issues <sup>[16]</sup> due to its relatively higher resistance

and a costly laser patterning method which is often used to remove  $\text{TiO}_2$  coating from the interconnection paths between sub-cells <sup>[17]</sup>. This is required to avoid the rise in the series resistance value, thereby decreasing the overall PSM performance <sup>[18]</sup>.

Unlike  $\text{TiO}_2$  material,  $\text{SnO}_2$  can be processed at much lower temperatures using different deposition technologies, including solution processes <sup>[11, 19, 20]</sup>, electrodeposition, <sup>[21]</sup>, electron-beam <sup>[22]</sup>, atomic layer deposition <sup>[23]</sup> and magnetron sputtering <sup>[24]</sup>. It is worth noting that the majority of reports so far related to  $\text{SnO}_2$  as ETL are only for small area devices <sup>[16]</sup>. Among all these thin film fabrication methods, magnetron sputtering (MS) is one of the most promising technologies due to its advantages related to cost-effectiveness and uniform large-scale  $\text{SnO}_2$  thin films. Thus far, there are only a few reports about  $\text{SnO}_2$  thin films, deposited by MS, as ETL for PSCs. For instance, F. Ali et al. have reported a PCE of 14% for an area of  $0.09 \text{ cm}^2$  <sup>[25]</sup>. Furthermore, the film uniformity across large area MS is demonstrated as well as the superior electrical conductivity and electron mobility of  $\text{SnO}_2$  versus that of  $\text{TiO}_2$  which are also demonstrated as an advantage to improve the interconnection quality between different sub-cells in PSMs. On the other hand, the state-of-the-art related to PCE of PSC based on  $\text{SnO}_2$  ETL is achieved by a very thin layer ( $\sim 25 \text{ nm}$ ) of  $\text{SnO}_2$  deposited by spin coating technique <sup>[26]</sup>. However, this deposition method associated with the very thin layer can lead to fringe effects, pinholes, and thickness non-uniformity related to large area thin films, especially for upscaling small devices into large PV modules.

Figure 1 displays the conduction band minimum (CBM) and valence band maximum (VBM) of commonly implemented inorganic metal materials as electron transport materials (ETLs) in perovskite solar cells (PSCs) where metal oxides, metal sulfide, CdSe and GaN were included. To deliver an efficient and reliable PSC, it is essential to meet the following key characteristics: (i) good optical transmittance; (ii) a low photon-energy loss; (iii) an appropriate bandgap-matching/alignment; (iv) high electrical conductivity and electron mobility; (d) cost-effectiveness; and an acceptable rate of reproducibility (i.e. stability), <sup>[27, 28]</sup>.



**Figure 1:** Schematic drawing showing the Conduction band (CB) and valence band (VB) of commonly employed inorganic materials as ETLs in PSCs. <sup>[29-39]</sup>.

As a replacement to  $\text{TiO}_2$ , various materials, which offer particular optoelectronic properties, have already been studied. This has included binary metal oxides ( $\text{ZnO}$ ,  $\text{In}_2\text{O}_3$ ,  $\text{Nb}_2\text{O}_5$ ,  $\text{WO}_3$ ,  $\text{Fe}_2\text{O}_3$  and  $\text{CeO}_2$ ) <sup>[32]</sup>, ternary metal oxides ( $\text{Zn}_2\text{SnO}_4$ ,  $\text{BaSnO}_3$  and  $\text{SrTiO}_4$ ), <sup>[30]</sup> metal sulfides ( $\text{MoS}_2$ ,  $\text{CdS}$ ,  $\text{In}_2\text{S}_3$ ,  $\text{SnS}_2$  and  $\text{Bi}_2\text{S}_3$ ), <sup>[29]</sup> as well as  $\text{GaN}$ ,  $\text{CdSe}$ , and  $\text{InGaZnO}_4$  <sup>[31]</sup>. Nevertheless, many pending issues are still to be addressed, including the low interfacial contact quality.

In 2015, the innovating work of Dai's group demonstrated -for the first time- a PSC based  $\text{SnO}_2$  ETL with a PCE of 6.5 %. This was followed, in the same year, by the work of Wan et al. who reached a PCE of more than 15 % <sup>[20, 40]</sup>. The power conversion efficiency of PSC based on  $\text{SnO}_2$  as electron transport material has recently achieved values above 20% <sup>[10]</sup>. This has demonstrated the  $\text{SnO}_2$  as an attractive and promising ETL material from various point of views, especially for perovskite solar cell, and as an excellent substitute to the conventional  $\text{TiO}_2$ , owing to the following factors: (i) an optimized energy-level matching <sup>[41]</sup>; (ii) a significantly higher electron mobility than  $\text{TiO}_2$ ; (iii) a high electrical conductivity <sup>[42, 43]</sup>; (iv) a large bandgap (3.6–4.5 eV); (v) a high transparency (i.e. optical transmittance) <sup>[44, 45]</sup>; (vi) a large flexibility in terms of processing temperature from relatively low temperature down to room

temperature; (vii) a high stability under light (i.e. low photoactivity), and (viii) a high stability under heat and humidity <sup>[46, 47]</sup>.

In addition, spin coating or spray methods are usually used for synthesis of TiO<sub>2</sub> ETL in PSC. This requires an elevated post-treatment temperature, generally above 450 °C, to enable the formation of the mesoporous layer with a dense structure, a crystalline structure, and a good electrical conductivity. Conversely, SnO<sub>2</sub> is routinely grown at much lower temperatures ( $\leq 250$  °C) and in some cases it is grown at room temperature when the crystalline structure is not required. This advantage is very appealing for large-scale industrial applications.

### 3.2. SnO<sub>2</sub> thin film as ETL

To prepare metal-oxide thin films with a high quality, either in the lab- and/or at an industrial-scale, magnetron sputtering is demonstrated to be a reliable and mature deposition process, also offering the possibility to use low-cost materials' targets. The SnO material is sputtered through high-energy argon-plasma ions, it reacts with oxygen and then it is deposited on the top of FTO layer under a continuous process. An accurate control of the film thickness and density can be achieved with MS via the deposition rate. It's a relatively cost-effective process, with a very low rate of waste. The thin film growth process is usually operated in a high vacuum chamber, which enables a high level of reproducibility <sup>[49]</sup>.

Qiu et al. studied in a systematic way the physical and chemical properties of the sputtered SnO<sub>2</sub> <sup>[26]</sup>. The fabricated perovskite solar cells implementing SnO<sub>2</sub> as ETL were found to exhibit a PCE of about 20% and a stability of about 625 hours as measured under T80 standard, thereby demonstrating the enhanced electrical conductivity and stability thanks to SnO<sub>2</sub> <sup>[26]</sup>. On the other hand, Bai and co-workers highlighted the impact of the Ar/O<sub>2</sub> gas ratio on the structural and morphological characteristics of the sputtered SnO<sub>2</sub> films <sup>[50]</sup>. The level of the trap states and dynamics of the carrier transit were also investigated in their PSC devices, which demonstrated a champion PCE up to 18% <sup>[50]</sup>. Moreover, Otoufi and co-workers experimented the bilayer architecture made with sputtered SnO<sub>2</sub> on TiO<sub>2</sub> layer, which was found to improve charge collection capacity, which has led to a PCE of about 12%, which is 4 % higher than that obtained with only TiO<sub>2</sub> as ETL ( $\sim 8\%$ ) <sup>[51]</sup>. It is worth noting at this level that the flowing gas during the deposition process plays a key role in controlling the oxygen vacancies <sup>[52]</sup>. In addition, defects states within the SnO<sub>2</sub> bandgap, which may originate from the amorphous and/or nano-crystallinity present in the films, might be effectively suppressed by a post-thermal annealing process under air, thereby leading to the passivation of the interface with the perovskite material.

In this work, we have achieved the implementation of moderate temperature radio frequency magnetron sputtered SnO<sub>2</sub> as electron transport layer for spin-coated triple-cation based perovskite

solar cells. We demonstrated that both procedures of mesoporous-scaffold and high-temperature processing are not essential in order to achieve high PSC device performance. Furthermore, no passivation process has been performed and no encapsulation has been used. Nevertheless, a PSC of more than 17% PCE has been demonstrated.

We have explored the material's properties of SnO<sub>2</sub>, namely structural, morphological, electrical and optical properties, as well as its chemical states. We have also studied the temperature dependency of the perovskite layer performance through PL measurement under various temperature, which has suggested an interplay of different physical phenomena, including charge transfer dynamics and charge recombination. Our developed MS SnO<sub>2</sub> ETL is demonstrated to achieve a good device photovoltaic performance and a relatively good lifetime, which also could help for further development and integration of SnO<sub>2</sub> films into PSMs.

### 3.3. Experiments and Characterization

**Electron Transport Layer (ETL):** The glass coated FTO substrates with 7Ω/□ sheet resistance were first cleaned with 1 wt.% Na dodecylsulfate aqueous, deionized water, acetone, and isopropyl alcohol prior to their utilization. SnO<sub>2</sub> thin films were grown on these Glass/FTO, using magnetron sputtering, manufactured by Torr Int. (described in the appendix in section 1.3.), at 200 °C with a high purity SnO target (purity 99.99%, 2" diameter, Codex International™), under O<sub>2</sub>:Ar mixture at 4:200 sccm, and a power of 50 W for 10 min, where the aim for these conditions is to improve the electron mobility without compromising the charge carrier concentration. A base pressure of 6.66 ×10<sup>-3</sup> Pa was first reached and then a deposition pressure which depends on the variable oxygen flow rate, was set around 0.66 Pa. As a sufficient oxygen flow was provided to assure the highest chemical state of Sn oxide in the grown film which are Sn(IV) oxide, the argon is necessary to form and maintain the magnetron plasma. The film thickness was about 40 nm. A deposition rate of about 4 nm min<sup>-1</sup> was deliberately chosen for a better film' coverage. Films were then treated with UV-ozone for 15 min and then one SnO<sub>2</sub> sample was post annealed at 250 °C in air, for 30 min. Samples were then placed into a nitrogen-filled glovebox for the perovskite layer deposition. In order to study the material's properties, some ETL - SnO<sub>2</sub> thin films were deposited directly on glass substrates (as-deposited SnO<sub>2</sub>) and some were air-annealed at 250 °C for 30 minutes.

**Perovskite Absorber Layer:** "The (FA, MA, and Cs: FMC) triple-cation mixed-halide perovskite was prepared inside glovebox conditions to maintain oxygen and H<sub>2</sub>O levels under 1 ppm. The one-step and 'antisolvent' processes were employed to obtain FAPbI<sub>3</sub> and FMC based films. Cs<sub>0.05</sub>MA<sub>0.10</sub>FA<sub>0.85</sub>Pb(I<sub>0.85</sub>Br<sub>0.15</sub>)<sub>3</sub> was prepared by dissolving the precursors in a 1:4 DMSO: DMF mixed solvent. A 1.4 M precursor solution was obtained by mixing CsI (0.07 M), FAI (1.13 M), PbI<sub>2</sub> (1.19 M),

MABr (0.2 M), and PbBr<sub>2</sub> (0.2 M) in the DMSO: DMF solvent with a 5% excess of PbI<sub>2</sub>. The perovskite films were deposited onto the 40-nm SnO<sub>2</sub> layer by spin-coating. This process started at 1000 rpm, for 10 s with a first-step acceleration rate of 200 rpm/s, followed by a second step acceleration at 1500 rpm/s until reaching a final speed of 4000 rpm, and then maintained for 35 s. During the last 20 s of spinning, 120 μL of chlorobenzene was dropped onto the substrate. The obtained film was finally annealed at 100 °C for 60 min”, this process is reported previously by A. Manekkathodi, B. Aïssa et al. [48].

**Hole-Transport Layer (HTL):** “The doped Spiro-OMeTAD was prepared by dissolving 65 mg of Spiro-OMeTAD in 1 mL of chlorobenzene. Then, 17.5 μL/mL of 26 mM solution of Li-dopant (bis(trifluoromethane)-sulfonimide lithium salt solution in acetonitrile) in acetonitrile, 21.9 μL/mL of 26 mM solution of cobalt-dopant (FK209) solution in acetonitrile, and 20 μL/mL of tert-butylpyridine were added as an additive. This solution was spin-coated at 4000 rpm for 20 s on the perovskite films immediately after their growth. After the Spiro-OMeTAD deposition, samples were taken out from the glovebox for oxidation purposes”, this process is also reported previously by A. Manekkathodi, B. Aïssa et al. [48].

**Back-Contact Electrode:** Back-contact electrode of 100 nm Au film was thermally evaporated under vacuum to complete the device. Film thicknesses were measured by a stylus profiler (Bruker Dektak).

**Optical Properties:** UV and visible light Transmittance and Absorbance spectra were determined by UV-Visible spectrometry (Jasco V670).

**Grazing Incidence X-ray Diffraction (GIXRD)** has been performed for both as-deposited SnO<sub>2</sub> and the annealed samples. The incident angle of the x-ray source was fixed at 0.55° while the 2θ was scanned from 15 to 65° for both diffractograms. The step size was fixed at 0.02° and the scan speed was kept at 2°/min for both diffractograms. The X-ray source is Cu K-alpha, and its wavelength is 1.54 Å. GIXRD analysis was performed using Rigaku Smartlab (Japan).

**X-ray photoelectron spectroscopy (XPS)** has been conducted on both as-deposited and annealed SnO<sub>2</sub> thin film samples. High resolution spectra for Sn3d and O1s were performed using 20 eV pass energy, 0.1 eV Step size and 5 periods (average of 5 spectra) while survey spectra were achieved using 100 eV pass energy, 1 eV step size and 1 period (no averaging). Prior to the XPS analysis, both samples were first cleaned using low energy/atom argon cluster source to remove the surface contamination without altering the surface chemistry of the SnO<sub>2</sub> thin film samples, which is confirmed by reducing drastically the carbon contamination while keeping the Sn3d spectra intact. It is worth noting that XPS equipment is calibrated using triple pure standard samples of Au, Ag and Cu. During the measurements, all spectra

were referenced using C1s to correct all the surface charging related shift. XPS analysis was performed using Thermo Fisher Scientific– Escalab 250Xi (United Kingdom).

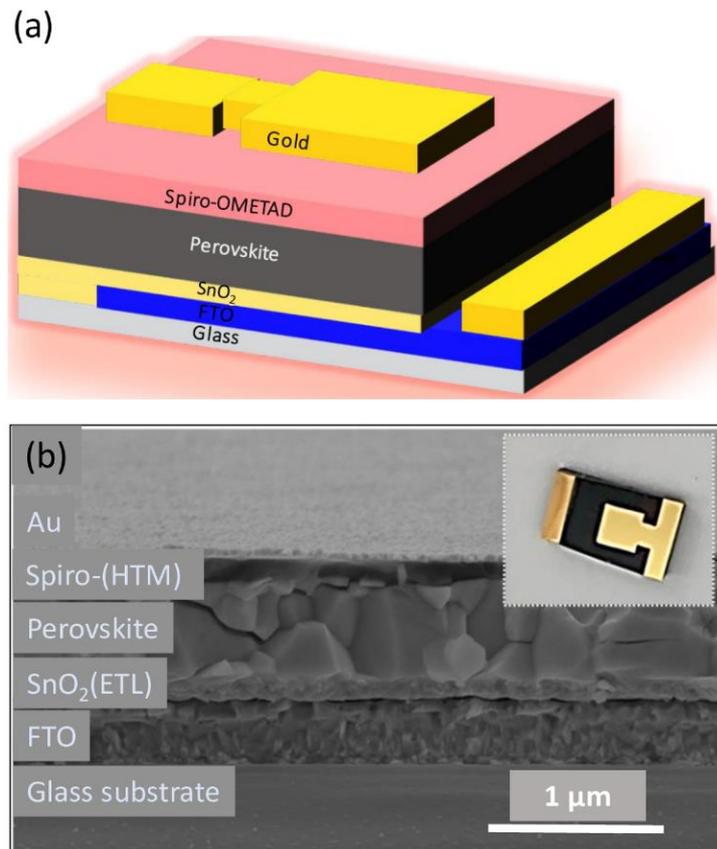
**Temperature-Dependent PL:** An iHR320 fluorescence spectrometer (Horiba, Jobin Yvon iHR320) equipped with CCD camera was used for the photoluminescence (PL) spectra recording. To decouple the response of the absorbing material from any other influence, only perovskite films were grown on glass substrate following the same procedure depicted above. A green PL excitation source of 532 nm wavelength was used and based on a diode-pumped solid-state laser (from Laser Quantum).

The aim is to probe the temperature dependency of the PL response of the absorber films. To do so, the temperature was varied from RT to 75 °C with a ramping rate of 5 °C/min. At each temperature, we stabilize the sample's temperature for 5 min prior to its measurement. Laser was turned off during the heating and stabilization intervals to avoid any photo-degradation of the samples. For all the temperatures, the PL response was recorded from the same spot on the sample and with the same exposure and integration-time.

**Solar Cell Device Performance:** A standard AM1.5G AAA solar simulator with 100 mW/cm<sup>2</sup> and a certified reference Si solar cell was employed to analyze the PV performance (Newport, Inc.). Measurements were performed in ambient air. An appropriate mask designed an active area of 0.16 cm<sup>2</sup>. I–V plots were recorded under different temperatures between 25 and 75 °C upon illumination (Keithley Model 2400 source meter). The temperature ramp/cooling rate of 3–5 °C/min was adopted through a solid-state heating/cooling system (ThermoCube). The PSCs were stabilized at each temperature setting for 10 min prior to launch the subsequent measurement. All the cooling and/or heating steps were performed under dark conditions.

A load resistance of 180 Ω was connected to the device for its stability. Current–voltage (I-V curve) scans were acquired. No filters have been used during the measurements. The stability measurement was performed at RH of 5% in a nitrogen box. The EQE scans were acquired by an IPCE measurement system.

**PSC device configuration and cross-sectional imaging:** The PSC device configuration is based on the conventional architecture where the layers on the glass substrate are fluorine doped tin oxide (FTO), SnO<sub>2</sub>, perovskite, Spiro-OMETAD and Au, respectively as shown in Figure 2. FTO material and SnO<sub>2</sub> are used as anode and ETL, respectively, while the perovskite as the absorber layer. Spiro-OMETAD acts as HTL, and the metallic cathode is made of Au contact.

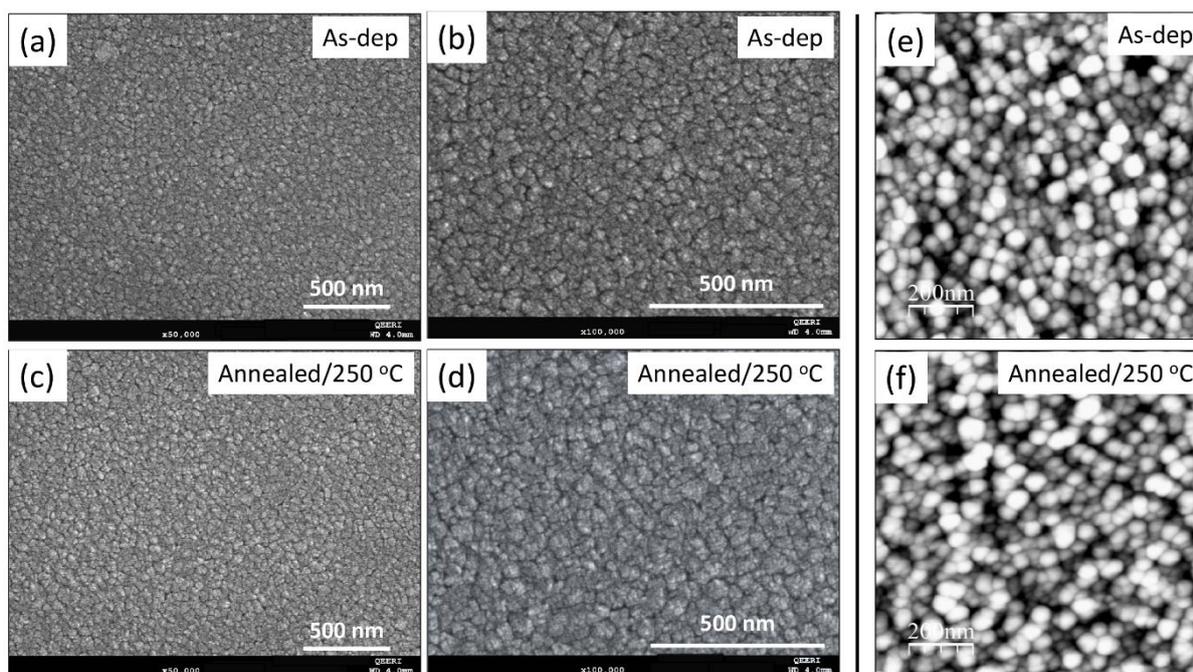


**Figure 2:** (a) Schematic diagram of the perovskite solar cell (PSC) used in this study. (b) Cross-sectional SEM image of Cs<sub>0.05</sub>MA<sub>0.10</sub>FA<sub>0.85</sub>Pb(I<sub>0.85</sub>Br<sub>0.15</sub>)<sub>3</sub> PSC. The inset shows a photograph of a PSC that is representative of those used in our experiments.

### 3.4. Results

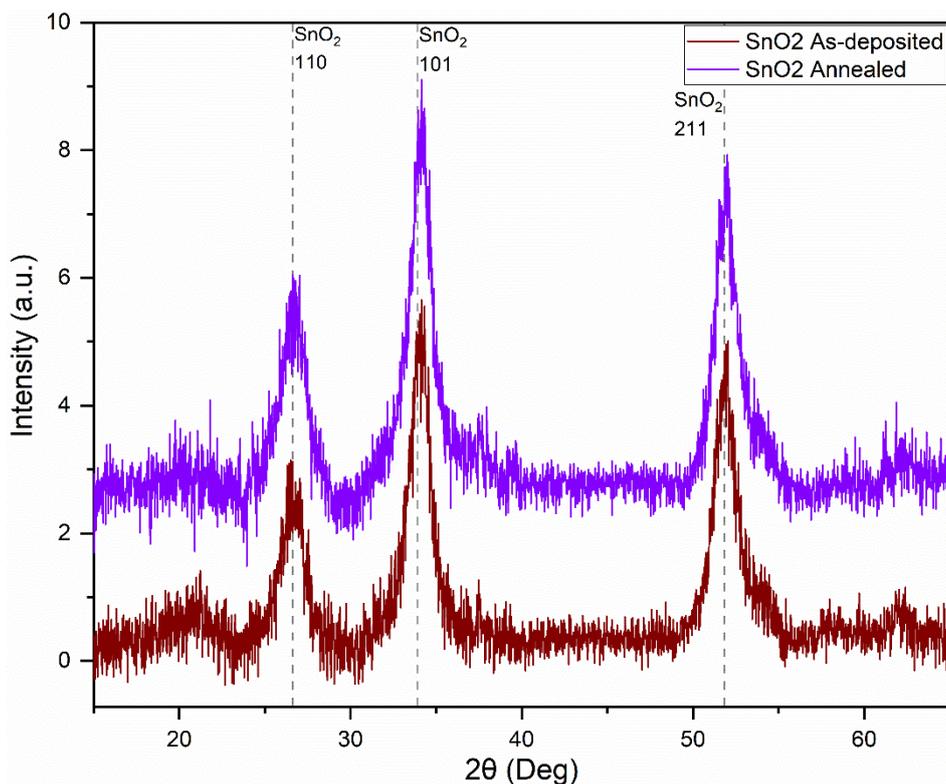
#### 3.4.1. Properties of the ETL-SnO<sub>2</sub> film

Figures 3a, 3b, 3c and 3d show representative top-view SEM images of as-deposited and thermally annealed SnO<sub>2</sub> films deposited onto glass substrate. SnO<sub>2</sub> samples demonstrate a uniform smooth surface morphology, with a large grain size, pinhole- and crack-free films. The thermal annealing treatment was found to have little or no impact on the morphology of the SnO<sub>2</sub> thin films. Figures 3e and 3f show the associated AFM images of these SnO<sub>2</sub> films.



**Figure 3:** SEM (a-d) and AFM images (e, f): Top-view SEM images showing the SnO<sub>2</sub> films (a) as deposited magnified at x50k, (b) as deposited magnified at x100k, (c) annealed at 250 °C magnified at x50k, (d) annealed at 250 °C magnified at x100k; and AFM images of the SnO<sub>2</sub> films (e) as-deposited, (f) annealed at 250 °C.

The as-deposited SnO<sub>2</sub> thin films have shown a clear crystalline microstructure of the SnO<sub>2</sub> phase by displaying clear triple peaks of (110), (101) and (211) orientation planes, as shown in Figure 4. There is no presence of any secondary phases which reveals the high crystalline quality of this SnO<sub>2</sub> thin film. After the air annealing the SnO<sub>2</sub> thin film samples have kept their SnO<sub>2</sub> microstructural phase as revealed by the same observed triple peaks of (110), (101) and (211) (Figure 4). Furthermore, it is noticed that the crystallinity has enhanced after the air annealing as demonstrated by the slight decrease of the full width at half maximum (FWHM) related to the SnO<sub>2</sub> phase peak (101). By using the Scherrer equation, the calculated crystallite size of SnO<sub>2</sub> phase using (101) peak has increased from 60.7 Å for the as-deposited SnO<sub>2</sub> sample to 69.2 Å for the annealed SnO<sub>2</sub> sample. These results are shown in Table 1. As per the GIXRD analysis, it is clear that the growth conditions favored the formation of SnO<sub>2</sub> phase without any secondary phases. Furthermore, the post air annealing process has kept and enhanced the crystallinity of the SnO<sub>2</sub> phase without promoting any secondary phases. These results are matching with previously reported findings<sup>[52]</sup> where improved crystallinity of the SnO<sub>2</sub> films was demonstrated with respect to an air-thermal-annealing treatment, as more oxygen incorporated in the film increased its crystallinity.



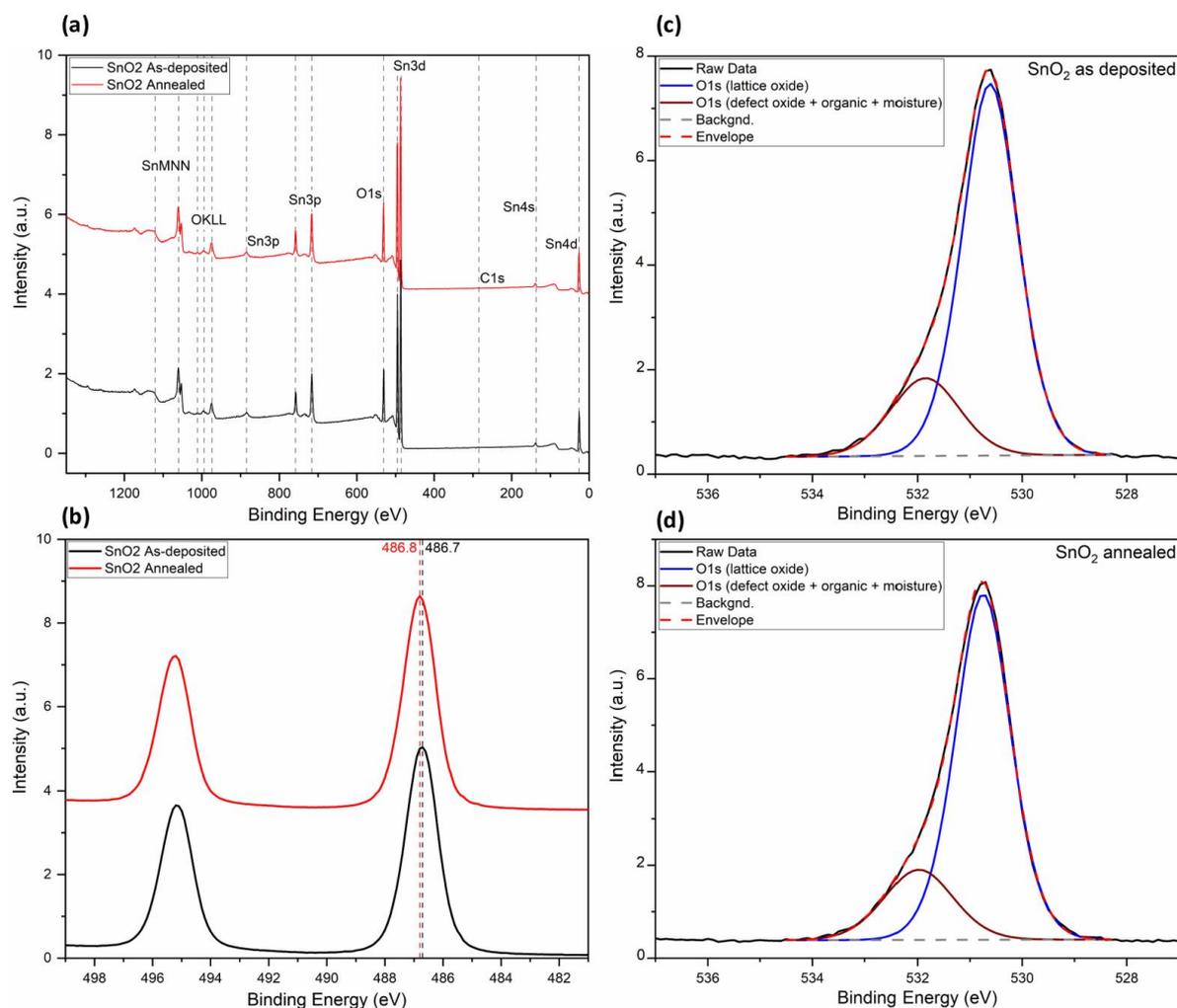
**Figure 4:** GIXRD for as-deposited SnO<sub>2</sub> and annealed SnO<sub>2</sub> thin film samples.

Sample	FWHM (°)	Crystallite size (Å)
SnO <sub>2</sub> as deposited	1.36	60.7
SnO <sub>2</sub> annealed	1.34	69.2

**Table 1:** Crystallite size calculation using SnO<sub>2</sub> phase peak (101) from Scherrer equation.

Survey XPS spectra has revealed the high purity of the SnO<sub>2</sub> thin film samples owing to the presence of only O and Sn as well as the low-to-absent content of C, particularly after the surface cleaning. All the peaks are related to Sn and O photoelectrons and Auger electrons as shown in Figure 5a. For the as deposited SnO<sub>2</sub> samples, Sn3d<sub>5/2</sub> peak is positioned on a higher binding energy at 486.7 eV, which indicates the presence of higher oxidation state related to Sn(IV). After the air annealing process, the peak position of Sn3d<sub>5/2</sub> has slightly shifted towards the higher binding energy at 486.8 eV, which indicates a slight increase in oxidation state after the annealing treatment as displayed by Figure 5b. These results match well with the crystallite size increase as discussed above. XPS related O 1s spectra have shown two main component peaks, the first peak is related to the photoelectrons originating from the oxygen atoms in the SnO<sub>2</sub> lattice, while the second peak is related to the other oxygen chemical states which might be defects in Sn oxide, or remaining surface organic molecules, and/or surface adsorbed moisture, as displayed in Figure 5c and 5d. For the as-deposited SnO<sub>2</sub> sample, the

peak position of O related to the lattice  $\text{SnO}_2$  is located at 530.6 eV, which confirms the presence of Sn(IV) oxide related to  $\text{SnO}_2$ . This is also corroborating well with the previous GIXRD results. The lattice oxide peak FWHM is of 1.26 eV, which is a reduced value that indicates the low chemical disorder as expected to be achieved by the vacuum-based material deposition.



**Figure 5:** Post cleaned XPS spectra for as-deposited  $\text{SnO}_2$  and annealed  $\text{SnO}_2$  thin film samples: (a) surveys, (b)  $\text{Sn}3d$ , (c)  $\text{O}1s$  for as deposited  $\text{SnO}_2$ , (d)  $\text{O}1s$  for annealed  $\text{SnO}_2$ .

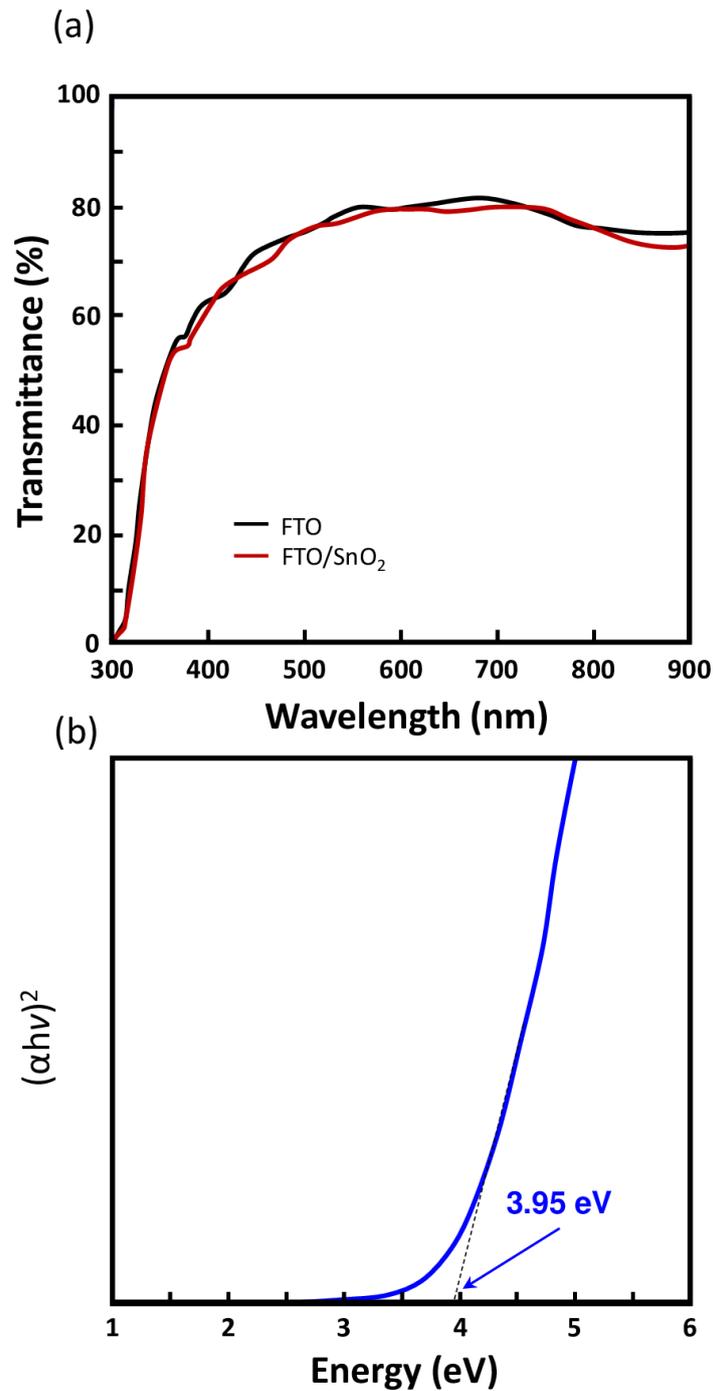
After the annealing, the peak position of O related to the lattice oxide has increased to 530.8 eV, which reveals the higher oxidation state of  $\text{SnO}_2$  thin film after the annealing. Furthermore, the FWHM of O related to lattice oxide has slightly decreased from 1.26 to 1.25 eV, revealing a slight enhancement in the chemical disorder. The percentage of O related to lattice oxide has slightly increased from 79.8 to 80.1%, which also reveals that the thermal annealing process has improved the micro-structuring of the  $\text{SnO}_2$  films as shown in Table 2 and it has enabled a slight reduction of the oxygen vacancies. Overall, this process has improved the surface chemistry of the  $\text{SnO}_2$  thin film [53].

Sample	Binding Energy (eV)	FWHM (eV)	Chemical State	Percentage (%)
SnO <sub>2</sub> as deposited	530.6	1.26	Lattice Oxide	79.8
	531.9	1.54	Defect/Organic/H <sub>2</sub> O	20.2
SnO <sub>2</sub> annealed	530.8	1.25	Lattice Oxide	80.1
	532.0	1.54	Defect/Organic/H <sub>2</sub> O	19.9

**Table 2:** XPS O1s spectra fitting: peak positions, FWHM and percentage of chemical states.

As for the optical properties, the SnO<sub>2</sub> films have shown rather a high transmittance in the visible range which is higher than 80 %. The associated bandgap obtained from Tauc plot was about 3.95 eV, a value that is greater than that of TiO<sub>2</sub> and/or ZnO thin films (Figure 6). It is worth noting that larger bandgap can act as an efficient hole blocking and as a barrier against high-energy photons absorption, thereby decreasing the current losses, which is a requirement for stable halide perovskite solar cells. Indeed, both high optical transmittance and film quality are critical for an effective ETL layer for the *n-i-p* planar structure of PSCs <sup>[54]</sup>.

Furthermore, Kam et al. <sup>[54]</sup> used a semiconductor band structure to calculate the position of the conduction band minimum (CBM) of a sputtered SnO<sub>2</sub> film, which was found to be -4.36 eV, which is even lower than that of TiO<sub>2</sub> and ZnO, where both are around -4.2 eV. More specifically, a deeper conduction band value will improve the electron transfer from the absorber layer, i.e. perovskite film, to the SnO<sub>2</sub> ETL layer. In the same context, Kam et al. calculations also showed the position of valance band maximum (VBM) at -8.08 eV <sup>[54]</sup>, which is clearly deeper than that of TiO<sub>2</sub> and ZnO, which are at -7.4 eV and -7.6 eV, respectively. Here again, a deeper valence band of SnO<sub>2</sub> conjugated to its larger bandgap will enhance the capacity of the perovskite film to block the holes towards the SnO<sub>2</sub> ETL.



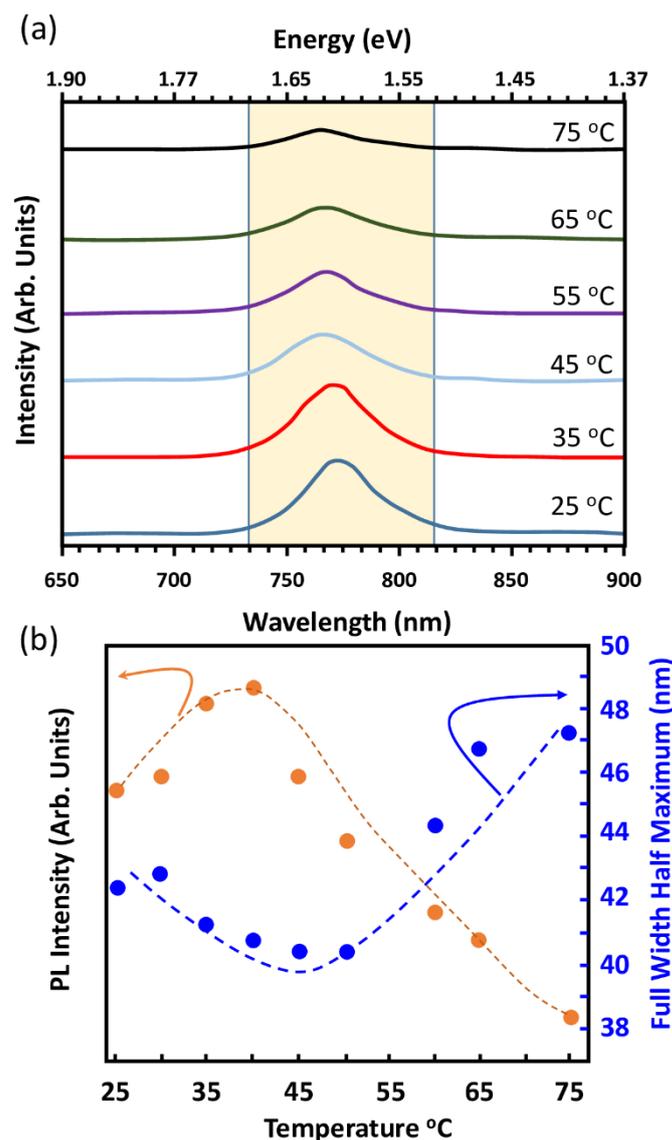
**Figure 6:** (a) Transmittance spectra % of the sputtered SnO<sub>2</sub> film deposited on fluorine doped tin oxide (FTO) glass and FTO on glass only. (b) Associated Tauc plot showing a bandgap of 3.95 eV.

On the other hand, the electrical resistivity of SnO<sub>2</sub> films has decreased from 0.245 Ω·cm (as-deposited) to 0.134 Ω·cm after a thermal annealing in air. The associated electron mobility has increased from 4.38 cm<sup>2</sup>/V·s as-deposited to 11.29 cm<sup>2</sup>/V·s after annealing treatment. However, the electron density has slightly decreased from 5.82 × 10<sup>18</sup> cm<sup>-3</sup> to 3.86 × 10<sup>18</sup> cm<sup>-3</sup>. The as-deposited and the annealed SnO<sub>2</sub> samples have negative charge carrier type (electrons). The thermal annealing process has clearly enhanced the electrical conductivity by significantly increasing the electron mobility and this finding is

matching with the increase in crystallite size after the annealing process. The decrease in charge carrier concentration is related to the oxygen vacancy reduction as the thermal annealing in presence of oxygen in air enables the vacant oxygen site filling. On the other hand, its effect on the surface morphology is not very evident. The effect of the thermal annealing has also been highlighted and discussed above through microstructural study and chemical state analysis as well as the measurements of the average roughness as conducted by AFM (Figure 3e-f). In fact, the root mean square (RMS) roughness value was found to be in the low-level values, and it changed only slightly between the as-deposited and the annealed film (measured in the range of 1.45 to 1.33 nm). Moreover, from morphological point of view, the sputtered SnO<sub>2</sub> film was uniformly deposited, thereby showing rather a low value of surface roughness. This process is also boosted by thermal annealing treatment, and the obtained roughness values are highly accommodating for the solution-deposited perovskite on SnO<sub>2</sub>/FTO/Glass. Therefore, this critical issue of the roughness effect has not been discussed in depth in the relevant literature. It is worth noting that a higher level of roughness of SnO<sub>2</sub> film will lead to inconstant rate for perovskite crystallization and hence leads ultimately to an increase of the carrier recombination probability between the perovskite absorber layer and the SnO<sub>2</sub> ETL.

#### 3.4.2. Perovskite absorber layer

To highlight the impact of real-world conditions on these PSC devices, especially when they are subject to operate in a harsh condition, such as desert environment, a temperature-sensitive Photoluminescence (PL) study was conducted through PL spectroscopy to elucidate the correlation with the charge-carrier and bandgap dynamics. Figure 7a shows the temperature-dependency of the PL measurements of the tri-cation perovskite films deposited on glass substrates. A well-defined PL peak centered at ~773 nm is measured and is the fingerprint of the band-to-band recombination. It shows an associated bandgap of ~1.6 eV, corresponding typically to the tetragonal phase of the tri-cation perovskite material <sup>[13]</sup>. The variation in PL emission peak position (Figure 7a) as well as in PL emission peak broadening and intensity (Figure 7b) as a function of temperature, were investigated and the results are presented in Figure 7. The PL intensity was found to increase with respect to the temperature up to ~40 °C and then start decreasing when the temperature continues increasing from 40 to 75 °C.



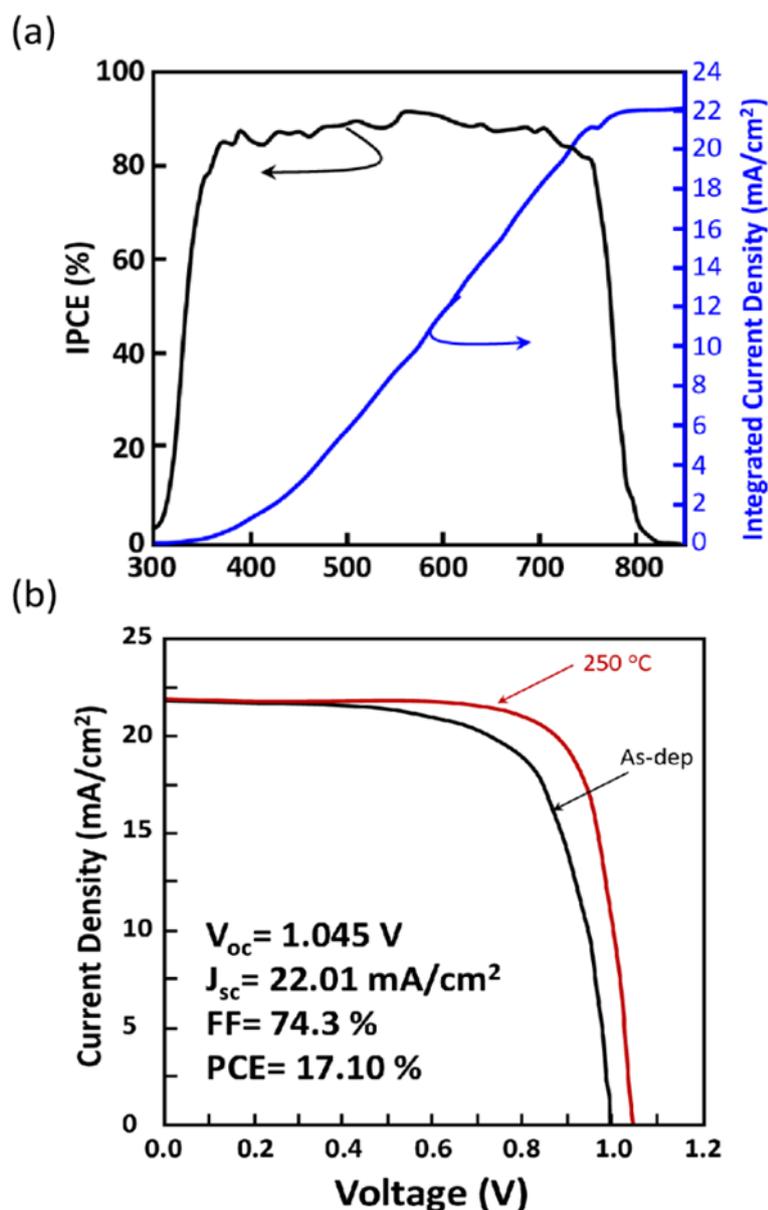
**Figure 7:** (a) Variation of the PL spectrum of the absorber perovskite films deposited on glass substrate. Measurements are performed in the range 25–75 °C. PL  $\lambda_{\text{ext}} = 532$  nm. (b) Associated change in the PL peak intensity at  $\sim 773$  nm and the FWHM as a function of temperature.

A rise in the PL emission intensity is generally the consequence of a decrease of the non-radiative recombination at traps and defects level in the bandgap, which lowers the band-to-band recombination rate and lower the charge–phonon coupling in the perovskite film <sup>[15,55]</sup>, leading thereby to an improvement of the  $V_{\text{oc}}$  and the fill factor (FF) values. Moreover, this decrease in the PL-emission-peak intensity while increasing the temperature above 40 °C is because of the charge trapping due to a boost of the charge–phonon coupling. This is also accompanied by a PL-emission-peak broadening (Figure 7a) <sup>[13]</sup>. The variations in the charge–phonon coupling and the orbital splitting, which are a direct consequence of the perovskite lattice expansion with respect to the temperature increasing, generate a singular bandgap broadening in addition to a blue shift of  $\sim 15$  meV in the PL emission <sup>[13]</sup>. At this range of temperature, both the frequency and the population of the particular involved phonon modes

increase with temperature <sup>[55]</sup> and support the suggestion that the charge trapping due to the electron–phonon coupling is dominant, which might be the reason behind the reduced PSC performance at elevated temperatures. Figure 7b shows an additional highlight of the temperature-dependency of the charge–phonon coupling <sup>[13, 56]</sup>, namely the variation of the FWHM of the PL emission peak with respect to the temperature. The FWHM was found to decrease slightly from RT to 40 °C and then noticeably increase from 40 °C to 75 °C, denoting a higher charge–phonon interaction at this range of temperature. This broadening of the FWHM further supports the diagnosis and conclusion that PSC PV performance decrease for temperatures beyond 40 °C could be attributed to the charge trapping related to higher charge–phonon interactions. Overall, we attribute the PL changes up to 40 °C to the carrier accumulation near the perovskite/glass interface and/or to the diminution in the non-radiative charge traps, whereas higher charge–phonon interactions dominate at higher temperatures.

### 3.4.3. Perovskite solar cell device performance

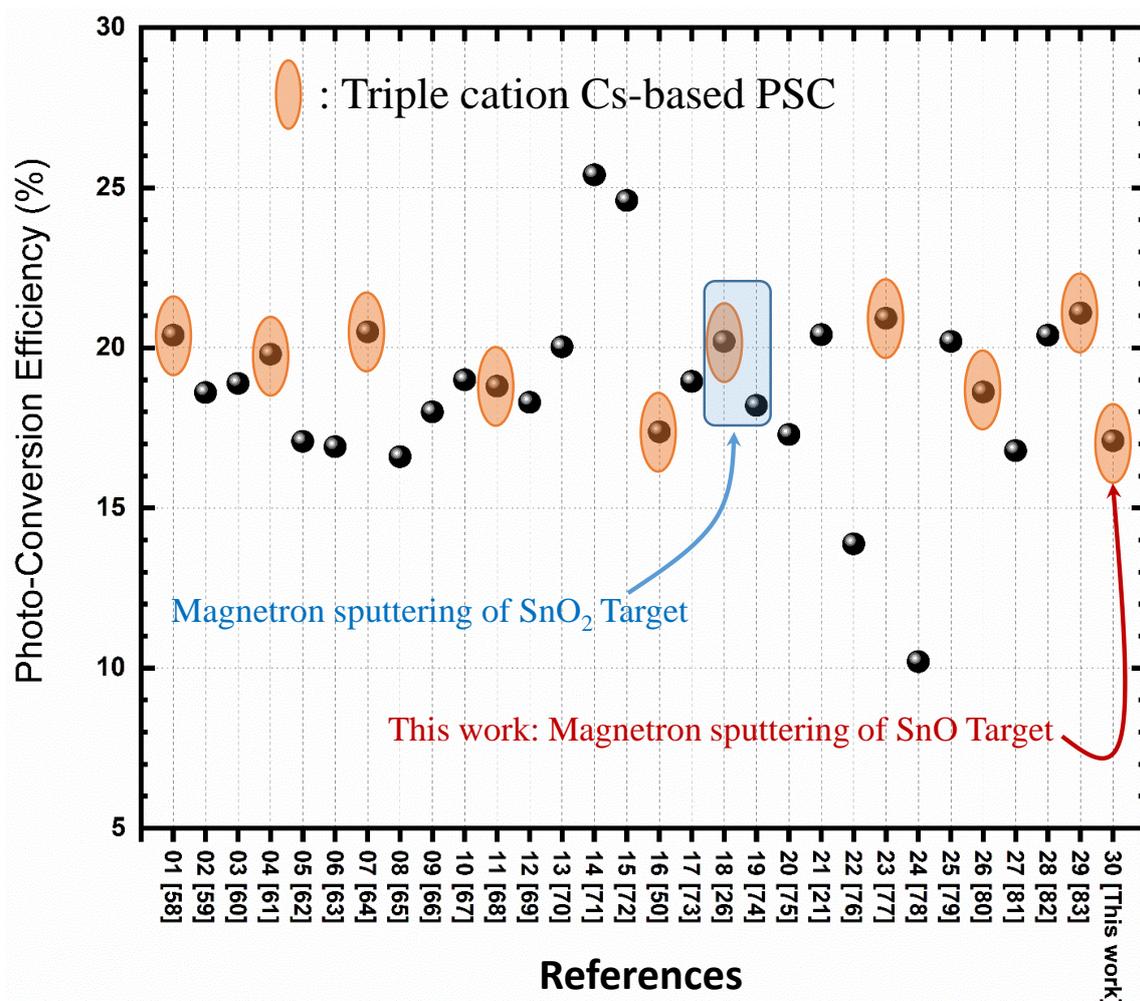
Figure 8 displays the perovskite planar-junction solar-cell performance of the best device based on the optimized sputtered SnO<sub>2</sub> films integrated as ETL. Figure 8a shows the incident-photon-to-current efficiency (IPCE) and associated with the integrated current density. Following the optimization process of SnO<sub>2</sub> thin films using the growth and post-depositions conditions, an optimized SnO<sub>2</sub> ETL could be achieved through improving its material properties related to micro-structure, morphology and surface chemistry as well as the optoelectronic properties. Therefore, a PCE of 17.1% has been reached as shown in Figure 8b after the post-annealing treatment. The improvement of PCE from as-deposited SnO<sub>2</sub> ETL of 15.07% to thermally annealed SnO<sub>2</sub> ETL of 17.1% is a direct result of the improved microstructural and optoelectronic properties, including electrical conductivity and electron mobility. As discussed above, SnO<sub>2</sub> has high transmittance in visible light region and a deep valence band maximum position which has improved the hole-blocking process while minimizing the recombination at the SnO<sub>2</sub>/perovskite interface. The high value related to J<sub>SC</sub> of 22 mA cm<sup>-2</sup> is corroborating well with the incident photon-to-electron conversion efficiency (Figure 8a) integrated J<sub>SC</sub>, which also confirms the high transmittance of the SnO<sub>2</sub> ETL-layer.



**Figure 8:** Perovskite solar cell device performance. a) IPCE and associated integrated current density of the best device based on the sputtered SnO<sub>2</sub> films as ETL. (b) *J*–*V* curves of the perovskite devices based on the SnO<sub>2</sub> films as-deposited and that annealed at 250 °C (device performance results are related to the annealed SnO<sub>2</sub> at 250 °C).

Another advantage of the SnO<sub>2</sub> as ETL for perovskite solar cells compared to TiO<sub>2</sub>, is its longer operational stability and lifetime. Indeed, the operational lifetime of the solar cell based on SnO<sub>2</sub> ETL has been measured under its maximum power point, under a continuous light illumination, at 45 °C (the detailed results are not shown here but will be a subject for a separate report). The best-reached T80 (the time required to decrease to 80% of the initial efficiency) lifetime was 250 h and the average lifetime over 200 h. This protocol is also regarded to be the most reliable and reproducible way to test the operational stability<sup>[57]</sup>.

Finally, Figure 9 summarizes a literature survey of various PCE values recorded for different PSC solar cells based SnO<sub>2</sub> ETL from about thirty references. In this literature survey, SnO<sub>2</sub> ETL was grown by different methods from different material sources. Only two references related to magnetron sputtering were found and show values comparable to our current work. Further optimization might be provided to enhance the PV properties through a systematic study by improving the microstructural and optoelectronic properties via post-deposition thermal annealing.



**Figure 9:** Summary of a literature survey of various PCE values recorded for different PSC solar cells based SnO<sub>2</sub> ETL. SnO<sub>2</sub> was grown by different methods from different sources [21, 26, 50, 58-82].

### 3.5. Summary

We have successfully employed the radio frequency magnetron sputtered SnO<sub>2</sub> as electron transport layer for spin-coated-(FA, MA, and Cs) triple cation-based perovskite solar cells on FTO substrates. We demonstrated that neither mesoporous scaffold nor high-temperature processing procedures were essential to achieve high device performance. In addition, no passivation process has been performed and no encapsulation has been used. Nevertheless, PSCs of 17.10% PCE have been achieved. Material's

characterization study has demonstrated that the air annealing has enhanced the material's structural and electrical properties, particularly the improvement in crystallite size which led to an improvement in electron mobility. These enhancements have improved the overall performance of the PSC device. Our results on the temperature dependency suggested an interplay of different physical phenomena, including charge transfer dynamics and charge recombination that govern the PV performance at different temperature intervals. We attributed the PL changes up to 40 °C to the possible carrier accumulation near the SnO<sub>2</sub>/perovskite interface and to the diminution in the non-radiative charge traps, whereas higher charge-phonon interactions are believed to be dominating at higher temperatures (i.e. from 40 °C to 75 °C). Our work offers a significant insight into the operational reliability and variability of the PV performance of PSCs, as well as the underlying mechanisms in perovskite materials, at real-world operating temperatures. More importantly, sputtered SnO<sub>2</sub> as ETL was demonstrated to enable a good device photovoltaic performance, stability and lifetime and could serve as a promising route for further development and integration of sputtered SnO<sub>2</sub> films into large-scale and cost-effective perovskite photovoltaics modules.

In the next chapter, nitrogen doping of SnO<sub>x</sub> via a two-step method will be explored aiming at achieving a p-type conductivity where synthesis, conduction, materials characterization will be discussed thoroughly.

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# **Chapter 4: Properties of Nitrogen doped Tin Oxide**

## 4. Properties of Nitrogen doped Tin Oxide

### 4.1. Introduction

As elaborated in Chapter 1, transparent conductive oxides (TCOs) are very appealing materials due to their combined properties of optical transparency, electrical conductivity, and chemical stability. TCOs have been heavily investigated and are crucial materials in various optoelectronic applications such as solar cells, light emitting diodes (LEDs), flat panel display (FPDs)<sup>[1-8]</sup>. The widely studied TCOs are In-, Sn-, Zn-based oxides<sup>[9-15]</sup>. These oxides are known to be intrinsically n-type conductive materials. Nowadays, there is a huge demand for high quality p-type TCOs to achieve the next generation optoelectronic applications including all oxide devices<sup>[3, 16, 17]</sup>. In Chapter 2, we have investigated the growth of undoped SnO<sub>x</sub> films using two different magnetron sputtering methods that differ by the target material (SnO and Sn), deposition parameters and the atmosphere of the annealing process; The optimization of the grown and annealed undoped SnO<sub>x</sub> has enabled a major improvement in optoelectronic properties as well as the ability to control the conductivity, n-type and p-type. In Chapter 3, the performance of SnO<sub>2</sub> thin film as an electron transport layer in perovskite solar cell has been assessed as well as the thermal annealing effect of this layer on the solar cell performance. In this chapter, the extrinsic doping of SnO<sub>x</sub> using nitrogen will be evaluated to achieve the p-type conductivity.

There have been several reports about p-type TCO materials<sup>[18-21]</sup>. SnO<sub>2</sub> is a very attractive material due to its higher earth-abundance, wide bandgap, and higher exciton binding energy of 130 mV<sup>[22, 23]</sup>. P-type doping for SnO<sub>2</sub> have been reported using doping metals with lower valence cations such as In to substitute the Sn atoms and to act as acceptor defects<sup>[24, 25]</sup> or using nitrogen atoms as doping elements to substitute oxygen atoms<sup>[26-28]</sup>. Nitrogen doping is highly desirable due to its ion size, high solubility, electronegativity, earth-abundance, and non-toxicity. Nitrogen ionic radius of 1.71 Å is larger than oxygen ionic radius of 1.40 Å, so the nitrogen doping increases the lattice constant for the host SnO<sub>2</sub> and it requires a relatively higher temperature for the nitrogen incorporation into the oxygen lattice<sup>[22, 23, 26-32]</sup>. Nitrogen doped SnO<sub>2</sub> (N:SnO<sub>2</sub>) has been synthesized previously by different methods. Jie et al.<sup>[31]</sup> reported N:SnO<sub>2</sub> thin films using Chemical Vapor Deposition. Bhawna et al.<sup>[33]</sup> reported N:SnO<sub>2</sub> nanoparticles using a wet chemistry synthesis process. Wang et al.<sup>[34]</sup> reported nanoparticles N:SnO<sub>2</sub> synthesis using a Laser Induced Pyrolysis process. There are several reports about nitrogen doped SnO<sub>2</sub> synthesis using magnetron sputtering and N<sub>2</sub> partial pressure during the growth. However, there are issues related to the incorporation of atomic nitrogen doping within the SnO host material, control of SnO<sub>x</sub>N<sub>y</sub> stoichiometry and formation of Sn nitride phases<sup>[22, 23, 28-30, 32]</sup>. In addition, there is also a previous report about plasma post-treated p-type ZnO by Cao et al.<sup>[35]</sup>. Nitrogen incorporation in SnO<sub>2</sub> can result in different chemical states, particularly atomic nitrogen doping which substitutes

oxygen in the host SnO<sub>2</sub> lattice and makes bonds with Sn(IV). This nitrogen atomic doping is reported to generate acceptor levels close to the valence band within the bandgap<sup>[22, 23, 27-30]</sup>.

As reported in literature, reactive magnetron sputtering of Sn using a mixture of both reactive gases O<sub>2</sub> and N<sub>2</sub> can be challenging in terms of incorporating atomic nitrogen within the SnO<sub>2</sub> host material. In this study, we report on the synthesis of N:SnO<sub>x</sub> using a two-step method. The first step consists of growing SnO<sub>2</sub> films by magnetron sputtering and the second one is a plasma treatment of the films under a high purity nitrogen precursor, namely ammonia (NH<sub>3</sub>), nitrous oxide (N<sub>2</sub>O), or nitrogen gas (N<sub>2</sub>). The plasma treated SnO<sub>x</sub> films were characterized by several techniques to monitor the effects of nitrogen treatment on the structural, composition, electrical and optical properties. We show that within our experimental conditions, the N<sub>2</sub> plasma treatment allowed the conversion from n-type character to p-type one.

## **4.2. Materials and Methods**

### **4.2.1. Material Synthesis**

SnO<sub>x</sub> were synthesized using magnetron sputtering, manufactured by AJA Int. (described in the appendix in section 1.3.) on quartz substrates using Sn metallic target and O<sub>2</sub>/Ar flow rate ratio at 0.25 where O<sub>2</sub>:Ar is 3:12 sccm at a pressure of 0.45 Pa, for a deposition time of 12 minutes. The sputtering power was fixed at 50 W. The substrate temperature was set at 100 °C. SnO<sub>x</sub> films were grown at 100 °C which is a low deposition temperature and using a relatively O<sub>2</sub> poor condition during the growth.

Upon deposition, the SnO<sub>x</sub> thin films were submitted to plasma treatments using an electron cyclotron resonance - plasma enhanced chemical vapor deposition (ECR-PECVD) reactor. Its model is MycroSys 400 System from the manufacturer Roth & Rau as shown in Figure 1. The conditions of plasma treatments are summarized in the Table 1. The nitrogen gas precursors were NH<sub>3</sub>, N<sub>2</sub>O or N<sub>2</sub> with two levels of treatment for each precursor namely, moderate (NH<sub>3</sub>-1, N<sub>2</sub>O-1, N<sub>2</sub>-1) and intense (NH<sub>3</sub>-2, N<sub>2</sub>O-2, N<sub>2</sub>-2). The exposure time and microwave power were fixed at 30 min and 500 W, respectively, for all the treatments.



**Figure 1:** Photo of an electron cyclotron resonance - plasma enhanced chemical vapor deposition (ECR-PECVD) reactor.

Treatment	Flow rate (sccm)	Temperature (°C)	RF power (W)	Acceleration (V)
NH <sub>3</sub> -1	10	100	30	70
NH <sub>3</sub> -2	20	300	50	150
N <sub>2</sub> O-1	10	100	30	70
N <sub>2</sub> O-2	20	300	50	150
N <sub>2</sub> -1	10	100	30	70
N <sub>2</sub> -2	20	300	50	150

**Table 1:** Conditions of plasma treatments applied to SnO<sub>x</sub> sputtered films.

#### 4.2.2. Material Characterization

The structural properties of as-deposited (reference) and plasma treated SnO<sub>x</sub> films were analysed using Grazing Incidence X-ray Diffraction (GIXRD Smartlab - Rigaku) where  $\theta$  was fixed at 0.55° and  $2\theta$  scanned from 15 to 65°. Surface imaging has been conducted using Atomic Force Microscopy in ScanAsyst imaging mode (AFM Dimension Icon - Bruker). Depth profiling was conducted using Time of Flight - Secondary Ion Mass Spectroscopy (ToF.SIMS<sup>5</sup> - IONTOF) with negative polarity, Bi<sup>+</sup> 100 x 100  $\mu\text{m}^2$  as a primary scanning analysis beam and Cs<sup>+</sup> 400 x 400  $\mu\text{m}^2$  as a secondary scanning sputtering beam. The optical properties were analysed using UV-Visible Spectroscopy (UV-Vis Spectrometer

Lambda 19 - Perkin Elmer) scanning from 1000 to 248 nm with 4 nm step size. Finally, the electrical properties were monitored using a Hall Effect probe system (HMS 3000 - Ecopia).

In order to understand the chemistry of SnO<sub>x</sub> films, a detailed study of elemental composition and the chemical states has been conducted using X-ray Photoelectron Spectroscopy (XPS Escalab 250Xi – Thermo Fisher Scientific) with a monochromated Al K alpha source at 1486.7 eV. XPS pass energy was 20 eV for high resolution scans and 100 eV for survey scans. Prior to performing the measurements, surface cleaning was performed using Argon cluster beam to remove the adventitious carbon contamination. The cleaning has enhanced the photoelectron intensities for N 1s, Sn 3d<sub>5/2</sub>, and O 1s. The remaining traces of surface adventitious carbon was used as a reference at 284.8 eV.

Previous reports about XPS analysis for N doped SnO<sub>x</sub> suggest that the N 1s peak position related to nitride Sn-N bonds varies from 396.5 to 397.6 eV depending on the Sn chemical state. N 1s peak position for N incorporated within SnO<sub>x</sub> is between 399 and 400 eV which is denoted as N:SnO<sub>x</sub> where nitrogen atoms are expected to occupy oxygen vacancy sites. The higher binding energy compared to nitride Sn-N is mainly related to the oxygen rich environment with higher electronegativity. The N 1s peak position ranging from 402 to 404 eV is related to -NO<sub>x</sub> adsorbed and/or substitution for oxygen sublattice which is denoted as NO<sub>y</sub>:SnO<sub>x</sub>. If the content of N is relatively small, nitrogen atoms are expected to fill the oxygen vacancies without altering the host SnO<sub>2</sub> chemical state. If nitrogen content becomes high, the chemical state of N is expected to be nitride by forming SnN<sub>x</sub> regions. XPS related O 1s chemical state analysis for as-deposited and plasma treated SnO<sub>x</sub> samples are expected to form mainly five chemical states for O 1s spectra representing: i) O-Sn(II) at 530.1 eV; ii) O-Sn(IV) at 530.6 eV are respectively related to SnO and SnO<sub>2</sub> lattice oxides, iii) O-Sn defect lattice oxide is around 531.7 eV, iv) oxygen related to organic and moisture is around 532.3 eV, and v) O-Si related to SiO<sub>2</sub> O peak is expected around 532.8 eV. Sn 3d<sub>5/2</sub> chemical state analysis for as-deposited and plasma treated SnO<sub>x</sub> samples which reveals mainly three possible chemical states represented by peaks for Sn 3d<sub>5/2</sub> spectra metallic related Sn(0) around 484.6 eV, Sn(II) around 486.4 eV and Sn(IV) around 486.9 eV<sup>[22, 26, 28, 29, 36-41]</sup>.

#### 4.3. Nitrogen doped SnO<sub>x</sub> thin films

All SnO<sub>x</sub> thin films were grown at 100 °C which is a low deposition temperature and using a relatively O<sub>2</sub> poor condition during the growth. These growth conditions (low temperature, poor O<sub>2</sub>) were intentional to form a high concentration of oxygen vacancy defect which is supposed to enhance the incorporation of atomic nitrogen doping into the SnO<sub>x</sub> thin film.

It has been reported that amorphous undoped tin oxide can exhibit a relatively high n-type conductivity due to the presence of high defect concentration related to oxygen vacancies<sup>[3, 5]</sup>. NH<sub>3</sub> treatments are

expected to release nitrogen ions, by breaking N-H bonds, and incorporating these ions within the SnO<sub>x</sub> host thin films. N<sub>2</sub>O treatments are likely to produce nitrogen, oxygen and NO<sub>y</sub> (where y is expected to be mainly 1) ions and to incorporate these ions within the host SnO<sub>x</sub> thin films. NO<sub>y</sub> is normally known by NO<sub>x</sub> but we are denoting it NO<sub>y</sub> to differentiate the x of the SnO<sub>x</sub>. The oxidation state change is expected due to the potential oxygen stoichiometry increase. N<sub>2</sub> treatments are expected to break N-N bonds and deliver only nitrogen ions to be incorporated within the SnO<sub>x</sub> host material.

The optical transmittance spectra in the UV-visible-NIR range of the as-deposited and N<sub>2</sub>O and N<sub>2</sub> plasma treated SnO<sub>x</sub> films will be plotted at each section. The measured average transmittance in the visible region from 500 to 800 nm and the optical bandgap as deduced from the Tauc plot will also be reported. The electrical properties and optoelectronic performance will also be presented for as-deposited and N<sub>2</sub>O and N<sub>2</sub> plasma treated SnO<sub>x</sub> films.

#### 4.3.1. As-deposited SnO<sub>x</sub> thin film (reference)

As elaborated previously in this chapter, the as-deposited SnO<sub>x</sub> thin film (reference SnO<sub>x</sub>) is intentionally grown at lower temperature with oxygen poor condition to be amorphous with a high oxygen vacancy concentration to facilitate the incorporation of nitrogen atoms within the SnO<sub>x</sub> film. GIXRD of the SnO<sub>x</sub> reference sample has not shown any peak confirming the amorphous microstructure character of this film. AFM image for the as-deposited SnO<sub>x</sub> film did not show any clear crystalline morphology. However, this film has revealed a highly smooth surface with a sub-nanometre roughness of 0.5 nm using root mean square calculation method.

XPS spectra of the as-deposited SnO<sub>x</sub> film shows a high quality SnO<sub>x</sub> witnessed by the presence of only Sn and O in the survey spectrum. High resolution spectrum of Sn 3d<sub>5/2</sub> shows that the chemical state is Sn(IV) except a very minor presence of Sn(II). The O 1s has revealed that majority of oxygen are related to SnO<sub>2</sub> lattice oxide and that a small portion are related SnO<sub>x</sub> defect oxide. O-Sn(IV) bonds are about 91% of the total O and a small part related to O-Sn defect oxide around 7 %. There is a minor fraction of less than 1% of total O for O-Sn(II).

As-deposited SnO<sub>x</sub> film has the average transmittance in the visible region from 500 to 800 nm of about 76.5% and its optical bandgap is 4.15 eV. The reference SnO<sub>x</sub> film exhibits a resistivity of 0.25 Ω-cm, mobility of 1.45 cm<sup>2</sup>/Vs, and a charge carrier concentration of 1.7 x 10<sup>19</sup> cm<sup>-3</sup>. The reached low mobility and relatively high charge carrier concentration are mainly due to the amorphous structure and oxygen vacancies, respectively.

These results related to as-deposited undoped SnO<sub>x</sub>, will be presented as reference in the following sections along with the results of plasma treated SnO<sub>x</sub> films.

#### 4.3.2. Ammonia (NH<sub>3</sub>) treatment for SnO<sub>x</sub> thin films

Following a visual check for NH<sub>3</sub> plasma treated SnO<sub>x</sub> films, it was clear that the ammonia treatments has led to a deterioration of the SnO<sub>x</sub> thin films by the absence of yellowish mirror like surface of both SnO<sub>x</sub> NH<sub>3</sub> treated films. In addition, a chemical alteration at the surface of the NH<sub>3</sub>-2 treated SnO<sub>x</sub> film has been observed due to the formation of a powdery material. As the NH<sub>3</sub> treated SnO<sub>x</sub> samples were completely deteriorated, the following characterization was only conducted to understand the chemistry alterations that took place. Therefore, the material characterization for NH<sub>3</sub> plasma treated SnO<sub>x</sub> films were limited to structural properties and chemical state analysis using GIXRD and XPS.

##### a) Structural Properties by GIXRD

Figure 2 plots the GIXRD patterns for the as-deposited and plasma treated SnO<sub>x</sub> films using NH<sub>3</sub> as plasma gas precursor. The as-deposited SnO<sub>x</sub> film has clearly an amorphous structure as witnessed by the large band in the range 25-35°. On the other hand, the SnO<sub>x</sub> thin film subject to moderate treatment using NH<sub>3</sub> (NH<sub>3</sub>-1) did not show any peaks except the typical broad peak related to amorphous quartz substrate around 21°. The intense NH<sub>3</sub>-2 treatment for SnO<sub>x</sub> has led to a formation of a crystalline metallic Sn as shown by the strong metallic Sn peaks features in the GIXRD pattern. This revealed that this plasma treatment has resulted in a reduction of SnO<sub>x</sub> to metallic Sn. Furthermore, the presence of a broad amorphous quartz peak related to the substrate around 21° also confirms the etching to the SnO<sub>x</sub>. The NH<sub>3</sub>-1 sample does not exhibit even the Sn peaks suggesting a potential powerful etching of the SnO<sub>x</sub> film. The presence of an intense SiO<sub>2</sub> broad peak confirm the etching of SnO<sub>x</sub>. It is most likely that NH<sub>3</sub> treatment release a high concentration of atomic hydrogen that reacted with oxygen of the SnO<sub>x</sub> film to form volatile OH species leaving behind unbonded Sn elements and/or resulting in an etching of the entire film.

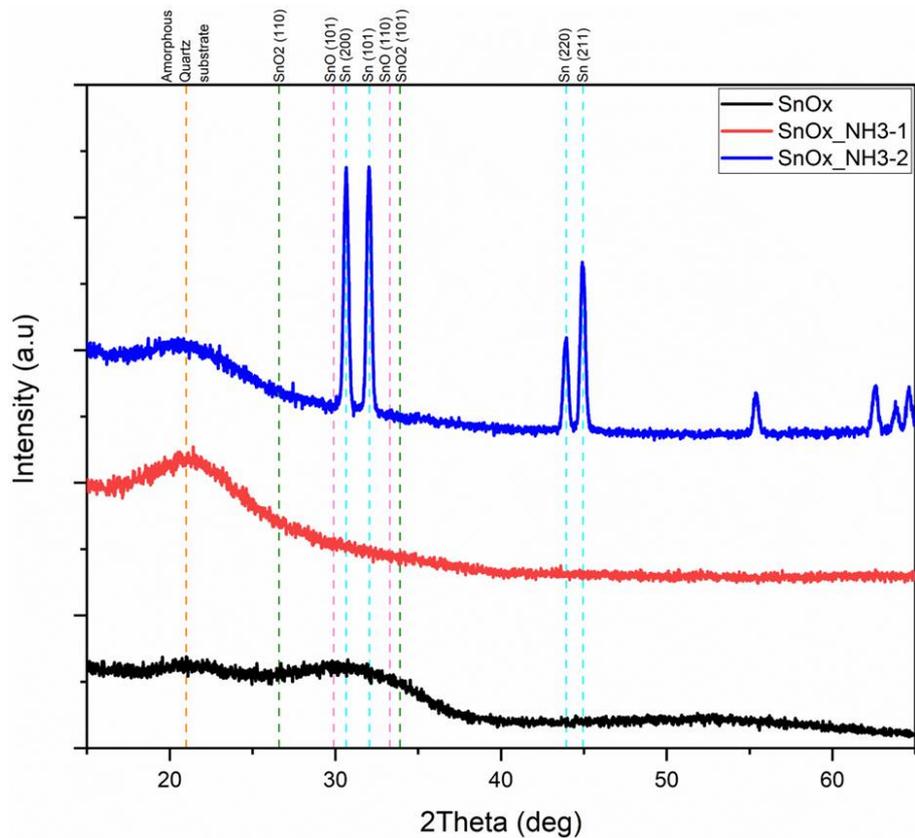


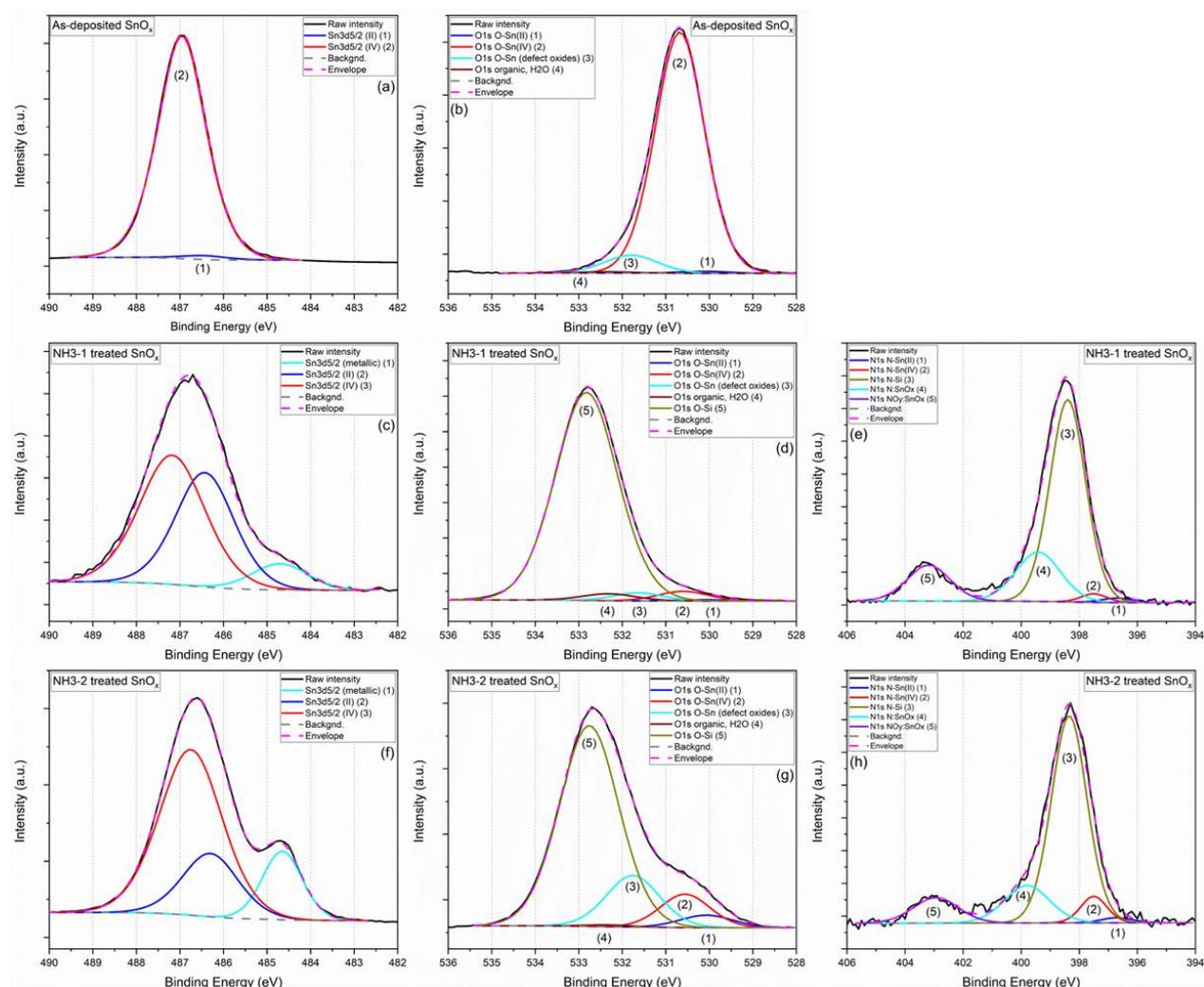
Figure 2: GIXRD patterns for as-deposited and NH<sub>3</sub> plasma treated SnO<sub>x</sub> films.

#### b) Surface chemical state analysis by XPS

Figure 3 plots the XPS data for as-deposited SnO<sub>x</sub> and NH<sub>3</sub> plasma treated SnO<sub>x</sub> thin films. The main five nitrogen related chemical states for the treated SnO<sub>x</sub> samples are represented by five peaks: (1) N 1s peak at 396.7 eV is related to nitride Sn(II)-N; (2) N 1s peak at 397.5 eV is related to nitride Sn(IV)-N; (3) N 1s related to Si-N chemical state that is expected to be around 398.2 eV. This higher binding energy is due to the higher ionicity of Sn-N compared to Si-N<sup>[37]</sup>. (4) The peak in the range from 399 to 400 eV is attributed to nitrogen atoms occupying oxygen vacancy site denoted N:SnO<sub>x</sub><sup>[22]</sup>; (5) N 1s peak around 403 eV is related to NO<sub>y</sub>:SnO<sub>x</sub>. The chemical states related to N-Sn(IV) and N:SnO<sub>2</sub> are reported to generate acceptor sites and contribute to p-type conductivity while NO<sub>y</sub>:SnO<sub>x</sub> and Sn(II)-N chemical states are reported to generate donor sites and contribute to n-type conductivity<sup>[22, 23, 27-30, 41]</sup>.

The survey results displayed in Table 2 show the presence of mainly oxygen and silicon as well as a minor presence of N, Sn and C. The observed Si peak positioned at 103.2 and 103.4 eV for NH<sub>3</sub>-1 and NH<sub>3</sub>-2, respectively, are very close to the characteristic SiO<sub>2</sub> related peak position at 103.5 eV. This reveals that the SnO<sub>x</sub> thin films have been substantially etched by NH<sub>3</sub> and the analysed surface is that of the substrate (quartz). Further analysis using high resolution Sn 3d<sub>5/2</sub> spectra for both treatment conditions show the presence of Sn(IV), Sn(II) and metallic Sn related peaks. Therefore, the NH<sub>3</sub>

treatments have etched and reduced the SnO<sub>x</sub> films, which is in good agreement with XRD analysis. It should be noted that the quartz substrate and metallic Sn were only observed after exposure of SnO<sub>x</sub> samples to NH<sub>3</sub> plasma. The NH<sub>3</sub> treatments has also formed some N-Si bonds at the quartz substrate. The summary of chemical state analysis is shown in Table 2.



**Figure 3:** XPS Sn3d<sub>5/2</sub>, O1s, and N1s spectra fitting for as-deposited and NH<sub>3</sub> plasma treated SnO<sub>x</sub> films.

Sample	Survey spectra		High resolution spectra deconvolution				Total Atomic (%)
	Element	Atomic (%)	Chemical state	Binding Energy (eV)	Area (CPS.eV)	Area (%)	
SnO <sub>x</sub> (as-deposited)	C 1s	0.87	--	--	--	--	0.9
	N 1s	0	--	--	--	--	0.0
	Sn 3d <sub>5/2</sub>	34.04	Sn(II)	486.5	19936	1.7	0.6
			Sn(IV)	487.0	1152286	98.3	33.5

			O-Sn(II)	530.1	2256	0.79	0.5	
	O 1s	65.09	O-Sn(IV)	530.7	262011	91.32	59.4	
			Defect oxide	531.8	21131	7.37	4.8	
			Organic, H <sub>2</sub> O	532.4	1505	0.52	0.3	
Total %							100	
SnO <sub>x</sub> _NH <sub>3</sub> - 1	Si 2p	31.8	--	--	--	--	31.8	
	C 1s	2.31	--	--	--	--	2.3	
	N 1s	6.68	N-Sn(II)	396.7	376	1.23	0.1	
			N-Sn(IV)	397.5	603	2.1	0.1	
			N-Si	398.4	19302	63.86	4.3	
			N:SnO <sub>x</sub>	399.4	5962	19.27	1.3	
			NO <sub>y</sub> :SnO <sub>x</sub>	403.2	4106	13.53	0.9	
	Sn 3d5/2	0.97	Sn(0)	484.5	2795	2.87	0.0	
			Sn(II)	486.4	11526	20.04	0.2	
			Sn(IV)	486.9	44328	77.09	0.7	
	O 1s	58.24	O-Sn(II)	530.1	2073	0.44	0.3	
O-Sn(IV)			530.6	15964	3.39	2.0		
Defect oxide			531.6	12926	2.74	1.6		
Organic, H <sub>2</sub> O			532.3	12148	2.58	1.5		
SiO <sub>2</sub>			532.8	428071	90.85	52.9		
Total %							100	
SnO <sub>x</sub> _NH <sub>3</sub> - 2	Si 2p	26.16	--	--	--	--	26.2	
	C 1s	8.46	--	--	--	--	8.5	
	N 1s	4.53	N-Sn(II)	396.7	215	1.4	0.1	

			N-Sn(IV)	397.5	931	6.3	0.3
			N-Si	398.4	9882	66.6	3.0
			N:SnO <sub>x</sub>	399.8	2270	15.2	0.7
			NO <sub>y</sub> :SnO <sub>x</sub>	402.9	1544	10.4	0.5
Sn 3d5/2	5.44		Sn(0)	484.7	42012	10.2	0.6
			Sn(II)	486.3	54897	40.9	2.2
			Sn(IV)	486.8	161662	49.0	2.7
O 1s	55.35		O-Sn(II)	530.1	11909	3.9	2.2
			O-Sn(IV)	530.6	29772	9.7	5.4
			Defect oxide	531.8	49353	16.1	8.9
			Organic, H <sub>2</sub> O	532.3	1558	0.5	0.3
			SiO <sub>2</sub>	532.8	213494	69.8	38.6
Total %							100

**Table 2:** XPS survey and chemical states quantification for as-deposited and NH<sub>3</sub> plasma treated SnO<sub>x</sub> films.

### c) Summary

The ammonia treatment has led to the etching and the reduction of SnO<sub>x</sub> films likely due to the formation of -OH related species induced by the presence of hydrogen. Therefore, the (moderate and intense) NH<sub>3</sub> plasma treatment with the used conditions is not suitable to form a nitrogen doped SnO<sub>x</sub>.

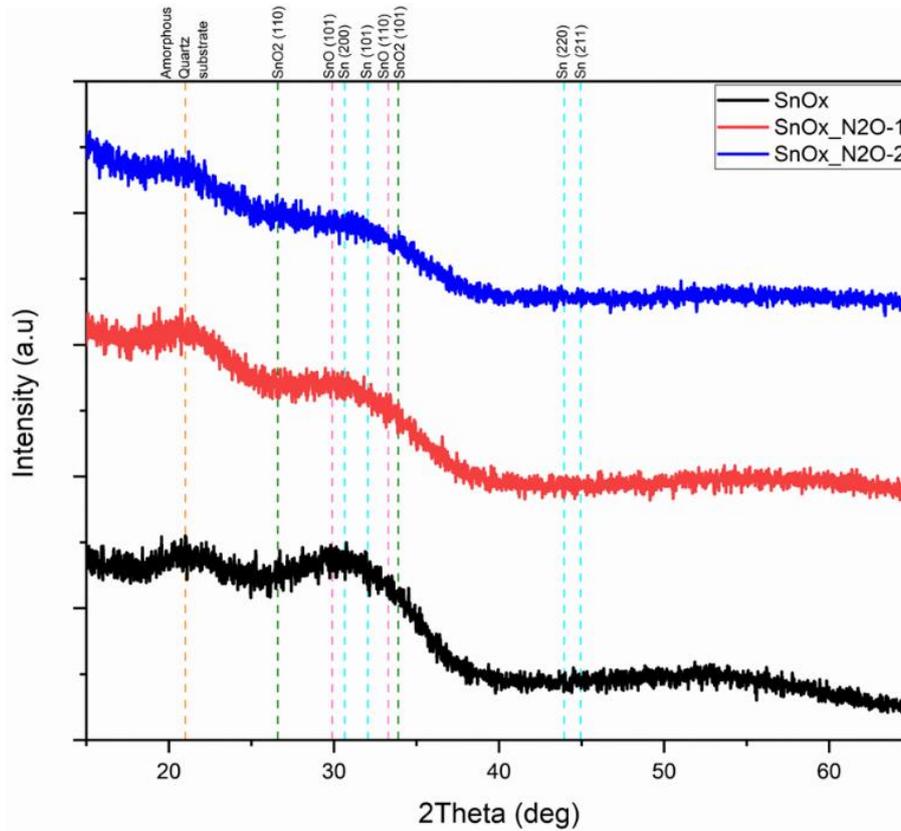
#### 4.3.3. Nitrous oxide (N<sub>2</sub>O) treatment for SnO<sub>x</sub> thin films

Exposure to N<sub>2</sub>O plasma is expected to produce N, O and NO<sub>x</sub> ions to be incorporated within the SnO<sub>x</sub>. Consequently, the oxidation state is supposed to change due to the potential oxygen stoichiometry increase.

#### a) Structural Properties by GIXRD and surface morphology by AFM

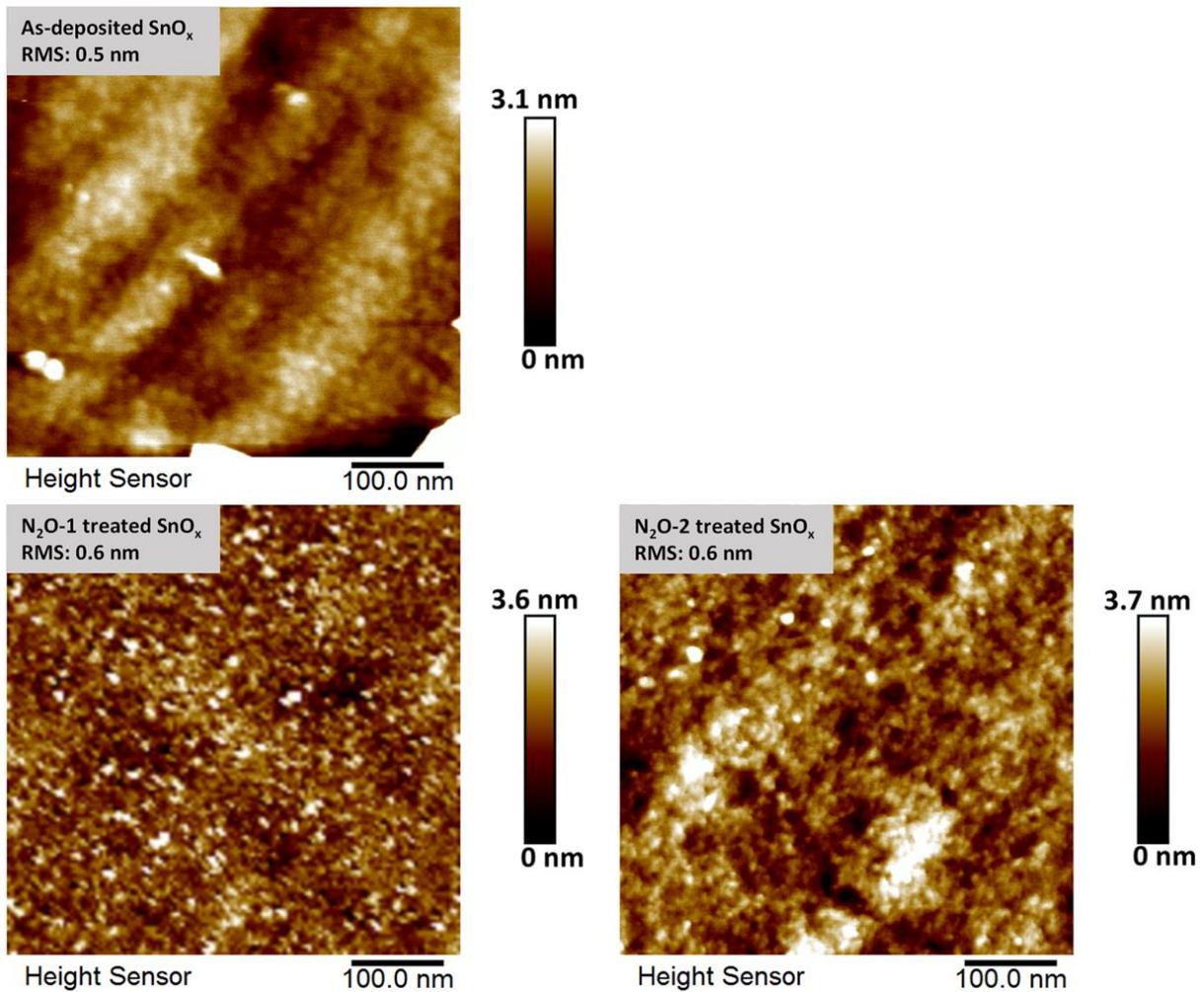
SnO<sub>x</sub> films after the plasma treatments under N<sub>2</sub>O (both moderate N<sub>2</sub>O-1 and intense N<sub>2</sub>O-2) have revealed that the amorphous microstructure character has been kept after both treatments. This amorphous microstructure is revealed by the absence of any peak in GIXRD pattern as shown by Figure

4. Moreover, the structural properties by GIXRD did not show any sign of SnO<sub>x</sub> film etching as the SiO<sub>2</sub> broad and intense peak was not observed.



**Figure 4:** GIXRD patterns for as-deposited and N<sub>2</sub>O plasma treated SnO<sub>x</sub> films.

AFM analysis has further confirmed the results of the amorphous character related to N<sub>2</sub>O plasma treated SnO<sub>x</sub>. Indeed, AFM images after N<sub>2</sub>O treatment have not revealed any crystalline microstructure formation as shown in Figure 5. The RMS roughness has slightly increased to 0.6 nm as compared to 0.5 nm for the reference SnO<sub>x</sub> film.



**Figure 5:** AFM images for as-deposited and N<sub>2</sub>O plasma treated SnO<sub>x</sub> films.

b) Surface chemical state analysis by XPS

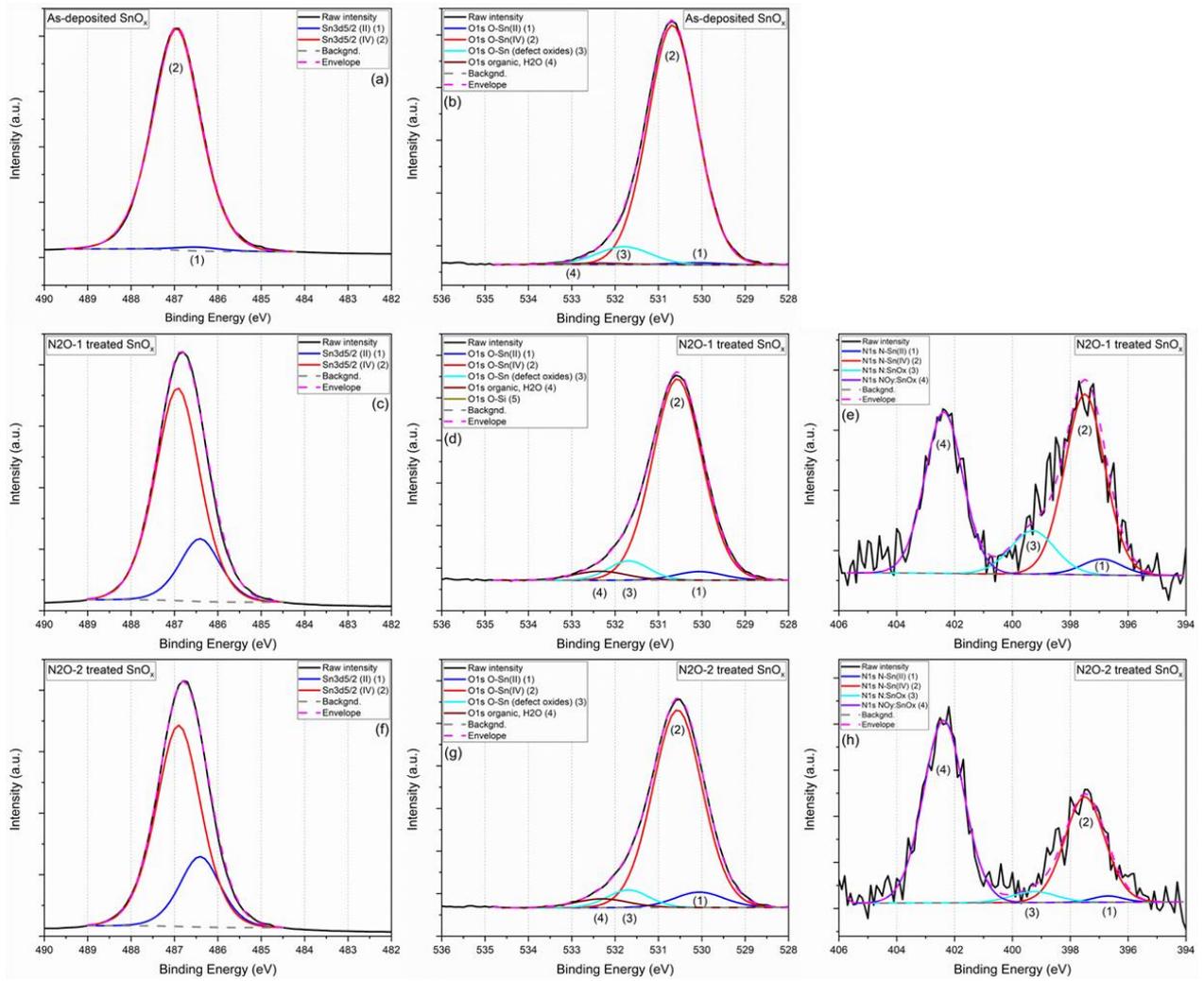
The data analysis from survey spectra of N<sub>2</sub>O treat SnO<sub>x</sub> films has revealed the presence of essentially O and Sn, with a minor atomic concentration of C and N where the survey atomic concentrations related to the moderate treatment N<sub>2</sub>O-1 for O, Sn, C and N are around 63.5, 30.8, 3.5 and 2.3 at.%, respectively. For the intense treatment N<sub>2</sub>O-2, the survey atomic concentrations for O, Sn, C and N are around 64, 30.6, 3.6 and 1.8 at.%, respectively. Figure 6 shows the XPS spectra fitting following a chemical state analysis of Sn 3d<sub>5/2</sub>, O 1s and N 1s spectra for SnO<sub>x</sub> films before and after both N<sub>2</sub>O treatments namely, N<sub>2</sub>O-1 and N<sub>2</sub>O-2.

Sn 3d<sub>5/2</sub> spectra for both N<sub>2</sub>O treatment conditions reveal that the chemical state is mainly Sn(IV) with a smaller presence of Sn(II) and the absence of Sn(0). For the moderate treatment N<sub>2</sub>O-1, the chemical state analysis shows that the presence of Sn(IV) is 92.4% while the presence of Sn(II) is 7.6%. For the intense treatment N<sub>2</sub>O-2, the chemical state analysis shows that the presence of Sn(IV) is 86.5% while the presence of Sn(II) is 13.5%.

As per O 1s spectra, the moderate N<sub>2</sub>O-1 treatment has kept the same chemical states within SnO<sub>x</sub> which are O-Sn(II), O-Sn(IV), defect oxide, and organic/H<sub>2</sub>O and their presence is 5.5%, 85.6%, 6.8%, 4.1%, respectively. Moreover, the moderate N<sub>2</sub>O-1 treatment has led to a minor partial chemical state reduction for O-Sn(IV) which is the main chemical state. Conversely, the states of Sn(II) and defect oxide have increased. For N<sub>2</sub>O-2 treated SnO<sub>x</sub>, there is a similar distribution of same chemical states within SnO<sub>x</sub> compared to N<sub>2</sub>O-1 treated SnO<sub>x</sub> which are O-Sn(II), O-Sn(IV), defect oxide, and organic/H<sub>2</sub>O and their presence is 6.5%, 83.2%, 6.3%, 4%, respectively. The intense N<sub>2</sub>O-2 treatment has also resulted in a minor partial chemical state reduction for O-Sn(IV). Nevertheless, O-Sn(IV) is remained the main chemical state. On the other hand, the states of Sn(II) and defect oxide have increased slightly more compared to the moderate N<sub>2</sub>O-1 treatment.

As per N 1s spectra, the moderate N<sub>2</sub>O-1 treatment has produced mainly molecular NO<sub>y</sub><sup>-</sup> (mainly NO<sup>-</sup>) incorporation in the SnO<sub>2</sub> host film and N-Sn(IV) chemical state. The presence of the following chemical states N-Sn(II), N-Sn(IV), N:SnO<sub>x</sub>, NO<sub>y</sub>:SnO<sub>x</sub> is 4%, 45.4%, 12.2%, 38.4%, respectively. For the intense N<sub>2</sub>O-2 treatment, the chemical state analysis has revealed similar results where the two main chemical states are molecular NO<sub>y</sub><sup>-</sup> (mainly NO<sup>-</sup>) incorporation in the SnO<sub>2</sub> host film and N-Sn(IV) chemical state. However, the presence of molecular NO<sub>y</sub><sup>-</sup> is higher compared to the moderate N<sub>2</sub>O-1 treatment. The other chemical states for this treatment are relatively minor. The presence of the following chemical states N-Sn(II), N-Sn(IV), N:SnO<sub>x</sub>, NO<sub>y</sub>:SnO<sub>x</sub> is 1.6%, 34.4%, 4%, 60%, respectively. All the chemical state analysis results are presented in table 3.

The majority of oxygen (around 84%) is bound to Sn(IV) forming the SnO<sub>2</sub> for both N<sub>2</sub>O treatments. However, there is a small fraction of oxygen which is bound to Sn(II) forming the SnO. There is also a minor part of oxygen related to SnO<sub>x</sub> defects oxide for both N<sub>2</sub>O treatments. The latter include oxygen deficiency defect which are important for n-type electronic conductivity as well as Sn-NO<sub>x</sub> defects [3, 5].



**Figure 6:** XPS Sn3d<sub>5/2</sub>, O1s, and N1s spectra fitting for as-deposited and N<sub>2</sub>O plasma treated SnO<sub>x</sub> films.

The summary of all XPS survey and chemical states quantification, shown in Table 3, reveals that N<sub>2</sub>O treatments have increased oxygen atomic concentration and favoured the incorporation of nitrogen within the SnO<sub>x</sub> thin film at around 2 at. %.

Sample	Survey spectra		High resolution spectra deconvolution				Total Atomic (%)
	Element	Atomic (%)	Chemical state	Binding Energy (eV)	Area (CPS.eV)	Area (%)	
SnO <sub>x</sub> (as-deposited)	C 1s	0.87	--	--	--	--	0.9
	N 1s	0	--	--	--	--	0.0
	Sn 3d <sub>5/2</sub>	34.04	Sn(II)	486.5	19936	1.7	0.6
			Sn(IV)	487.0	1152286	98.3	33.5

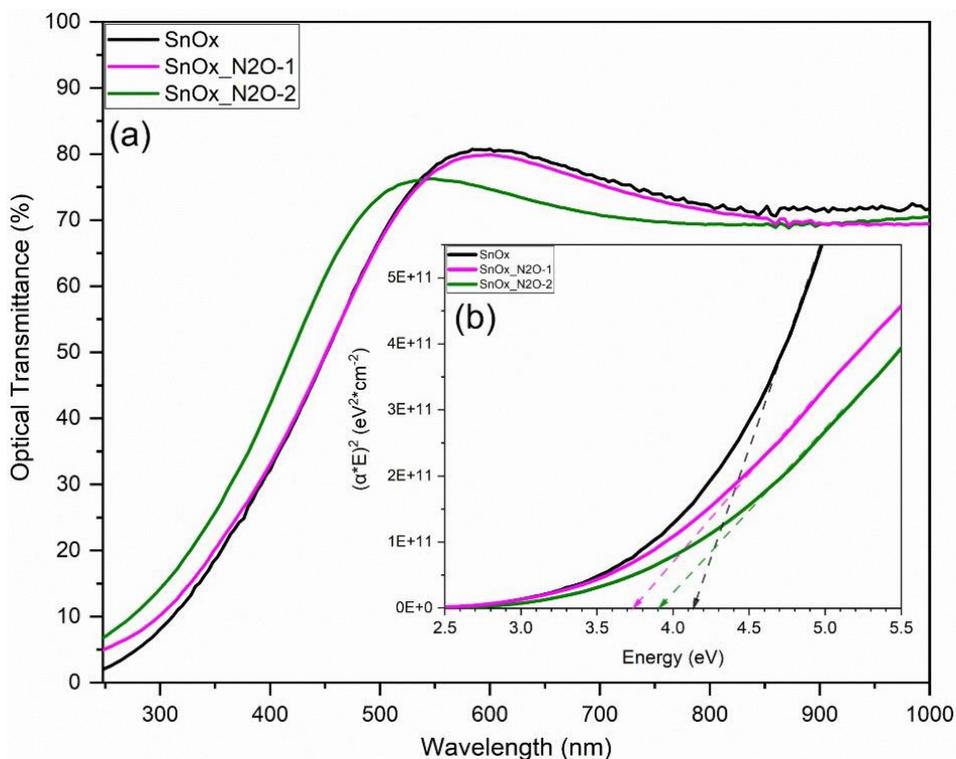
			O-Sn(II)	530.1	2256	0.79	0.5
			O-Sn(IV)	530.7	262011	91.32	59.4
			Defect oxide	531.8	21131	7.37	4.8
			Organic, H <sub>2</sub> O	532.4	1505	0.52	0.3
Total %							100
SnO <sub>x</sub> _N <sub>2</sub> O- 1	C 1s	3.47	--	--	--	--	3.5
	N 1s	2.32	N-Sn(II)	396.9	235	4.02	0.1
			N-Sn(IV)	397.5	2662	45.45	1.1
			N:SnO <sub>x</sub>	399.3	713	12.18	0.3
			NO <sub>y</sub> :SnO <sub>x</sub>	402.4	2243	38.36	0.9
	Sn3d5/2	30.75	Sn(II)	486.4	90496	7.62	2.3
			Sn(IV)	486.9	1097363	92.38	28.4
	O 1s	63.46	O-Sn(II)	530.1	11404	3.53	2.2
			O-Sn(IV)	530.6	275984	85.55	54.3
			Defect oxide	531.7	21843	6.77	4.3
Organic, H <sub>2</sub> O			532.3	13363	4.14	2.6	
Total %							100
SnO <sub>x</sub> _N <sub>2</sub> O- 2	C 1s	3.62	--	--	--	--	3.6
	N 1s	1.76	N-Sn(II)	396.7	93	1.58	0.0
			N-Sn(IV)	397.5	2028	34.36	0.6
			N:SnO <sub>x</sub>	399.3	238	4.03	0.1
			NO <sub>y</sub> :SnO <sub>x</sub>	402.4	3538	60.03	1.1
	Sn3d5/2	30.6	Sn(II)	486.4	156101	13.51	4.1
			Sn(IV)	486.9	999095	86.49	26.5
	O 1s	64.02	O-Sn(II)	530.1	21171	6.5	4.2
			O-Sn(IV)	530.6	270991	83.23	53.3
			Defect oxide	531.7	20390	6.26	4.0
Organic, H <sub>2</sub> O			532.3	13032	4	2.6	
Total %							100

**Table 3:** XPS survey and chemical states quantification for as-deposited and N<sub>2</sub>O plasma treated SnO<sub>x</sub> films.

### c) Optical Properties

The N<sub>2</sub>O-1 treated SnO<sub>x</sub> film exhibits (see Figure 7a) an average optical transmittance of 75.6% which is slightly lower as compared to as-deposited SnO<sub>x</sub> of 76.5%. Nevertheless, N<sub>2</sub>O-2 treated SnO<sub>x</sub> was a subject of a more decrease in average optical transmittance down to 72.6% compared to the as-deposited SnO<sub>x</sub>. This lower optical transmittance is suggested to be related to the incorporation of nitrogen within SnO<sub>x</sub> which led to the decrease in oxygen as shown in chemical state analysis of N<sub>2</sub>O treated SnO<sub>x</sub> (table 3).

N<sub>2</sub>O treatments for SnO<sub>x</sub> has resulted in lowering the optical bandgap as shown in Figure 7b. N<sub>2</sub>O-1 treated SnO<sub>x</sub> has an optical bandgap of 3.75 eV compared to 4.15 eV for as-deposited SnO<sub>x</sub> film. On the other hand, the N<sub>2</sub>O-2 treatment of SnO<sub>x</sub> has an optical bandgap of 3.92 eV compared to 4.15 eV for as-deposited SnO<sub>x</sub> as shown in Table 4. This can be originated from the high disorder within the SnO<sub>x</sub> film due to nitrogen enrichment treatment. Furthermore, the high decrease in optical bandgap following N<sub>2</sub>O-1 treatment at 100°C compared to N<sub>2</sub>O-2 treatment at 300°C is likely due to the lower temperature of the treatment N<sub>2</sub>O-1. These results are in agreement with previous reports in the literature which show that the bandgap energy can be affected by the disorder within the SnO<sub>x</sub> host film due to the nitrogen incorporation that forms energy levels close to the conduction or valence bands and reduces the bandgap energy<sup>[22, 23, 29, 40, 41]</sup>.



**Figure 7:** (a) Optical transmittance and (b) Tauc plot for as-deposited and N<sub>2</sub>O treated SnO<sub>x</sub> films.

Samples	Average optical transmittance (%)	Optical bandgap (eV)
SnO <sub>x</sub> as deposited	76.5	4.15
N <sub>2</sub> O-1 treated SnO <sub>x</sub>	75.6	3.75
N <sub>2</sub> O-2 treated SnO <sub>x</sub>	72.6	3.92

**Table 4:** Average optical transmittance from 500 to 800 nm and optical bandgap energy for as-deposited and N<sub>2</sub>O treated SnO<sub>x</sub> films.

d) Electrical properties

The electrical properties of N<sub>2</sub>O treated SnO<sub>x</sub> films have been investigated using Hall effect measurements. All the results are summarized in table 5.

Upon the N<sub>2</sub>O-1 treatment, the n-type conductivity of the SnO<sub>x</sub> thin film has been enhanced. This improvement is related to the contribution of the mobility which has increased 7 times. The charge carrier concentration has slightly decreased compared to as-deposited SnO<sub>x</sub> film. The behaviour is likely related to the decrease of oxygen vacancies due to the incorporation of oxygen, nitrogen and NO<sub>y</sub> ions. The improved mobility can be attributed to the lower structural disorder due to the treatment temperature and the reduction of oxygen vacancies.

Regarding N<sub>2</sub>O-2 treated SnO<sub>x</sub> film, the conductivity has been also improved while keeping the conductivity of type-n. This improvement is related to the contribution of the mobility which has increased by 11-fold after the N<sub>2</sub>O-2 treatment. The charge carrier concentration is relatively similar for both N<sub>2</sub>O treatments and has slightly decreased compared to as-deposited SnO<sub>x</sub> film. Similar to the previous N<sub>2</sub>O-1, the improved mobility can be attributed to the lower structural disorder due to the treatment temperature and the reduction of oxygen vacancies.

Samples	Thickness (nm)	Resistivity (Ω·cm)	Mobility (cm <sup>2</sup> /V·s)	Charge carrier (cm <sup>-3</sup> )	Conductivity
SnO <sub>x</sub> as deposited	120	2.52 x 10 <sup>-1</sup>	1.45	1.71 x 10 <sup>19</sup>	n-type
N <sub>2</sub> O-1 treated SnO <sub>x</sub>	122	7.52 x 10 <sup>-2</sup>	9.77	8.49 x 10 <sup>18</sup>	n-type
N <sub>2</sub> O-2 treated SnO <sub>x</sub>	119	4.58 x 10 <sup>-2</sup>	16.07	8.47 x 10 <sup>18</sup>	n-type

**Table 5:** Resistivity, mobility, and charge carrier concentration for as-deposited and N<sub>2</sub>O treated SnO<sub>x</sub> thin films.

e) Optoelectronic Performance

The optoelectronic performance for all SnO<sub>x</sub> samples was evaluated using Haacke figure of merit (FoM) which was described in the previous Chapter<sup>[42]</sup>. The results are shown in Table 6. The N<sub>2</sub>O plasma

treated SnO<sub>x</sub> films have a relatively higher figure of merit compared to the as-deposited undoped SnO<sub>x</sub>. This is due to the lower resistivity of N<sub>2</sub>O plasma treated SnO<sub>x</sub> thin films. The highest FoM was recorded for N<sub>2</sub>O-2 plasma treated SnO<sub>x</sub> which is  $1.06 \times 10^{-5} \Omega^{-1}$  for an average transmittance between 500 to 800 nm.

Samples	Figure of merit ( $10^{-3} \Omega^{-1}$ )		
	Average Transmittance 400 – 700 nm	Average Transmittance 500 – 800 nm	Transmittance at 550 nm
As-deposited SnO <sub>x</sub>	$1.04 \times 10^{-03}$	$3.27 \times 10^{-03}$	$4.28 \times 10^{-03}$
N <sub>2</sub> O-1 plasma treated SnO <sub>x</sub>	$3.29 \times 10^{-03}$	$9.89 \times 10^{-03}$	$1.32 \times 10^{-02}$
N <sub>2</sub> O-2 plasma treated SnO <sub>x</sub>	$6.60 \times 10^{-03}$	$1.06 \times 10^{-02}$	$1.70 \times 10^{-02}$

**Table 6:** Figure of Merit for as-deposited and N<sub>2</sub>O plasma treated SnO<sub>x</sub> thin films.

#### f) Summary

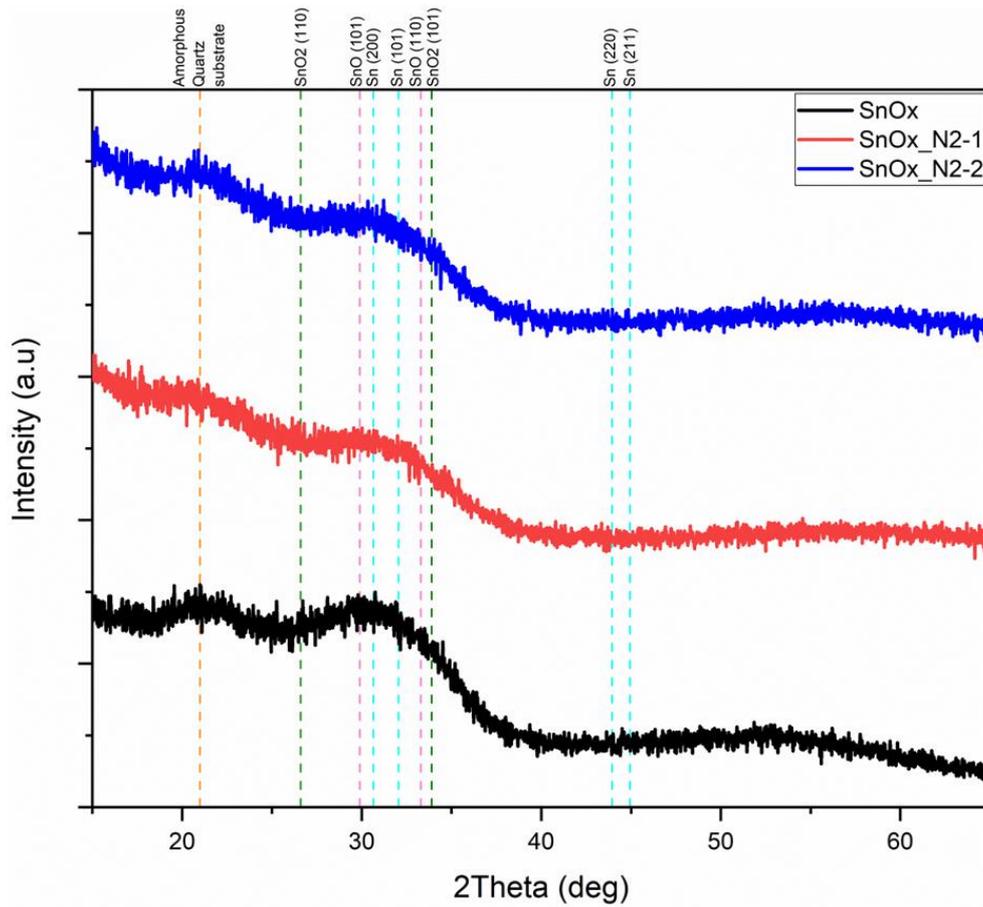
The above study has shown that both N<sub>2</sub>O plasma treatments have enabled the incorporation of nitrogen with the SnO<sub>x</sub> films while keeping the main chemical states of the host material. Furthermore, the host SnO<sub>x</sub> films have kept their amorphous microstructure character and they highly smooth surface morphology. The N<sub>2</sub>O plasma treatments have enabled an improvement in electrical properties of the SnO<sub>x</sub> films via improving the electron mobility. The decrease in electrical resistivity has significantly improved the optoelectronic performance of N<sub>2</sub>O plasma treated SnO<sub>x</sub> films. Nevertheless, these treatments were not able to convert the conductivity from n-type to p-type.

#### 4.3.4. Nitrogen (N<sub>2</sub>) treatment for SnO<sub>x</sub> thin films

N<sub>2</sub> treatments are likely to produce only nitrogen ions which are suggested to be incorporated within the SnO<sub>x</sub>. The oxidation state change is expected due to the potential oxygen stoichiometry decrease following the nitrogen incorporation. N<sub>2</sub> plasma treatments were expected to break N-N bonds and deliver only nitrogen ions which are supposed to be incorporated within the SnO<sub>x</sub> host material.

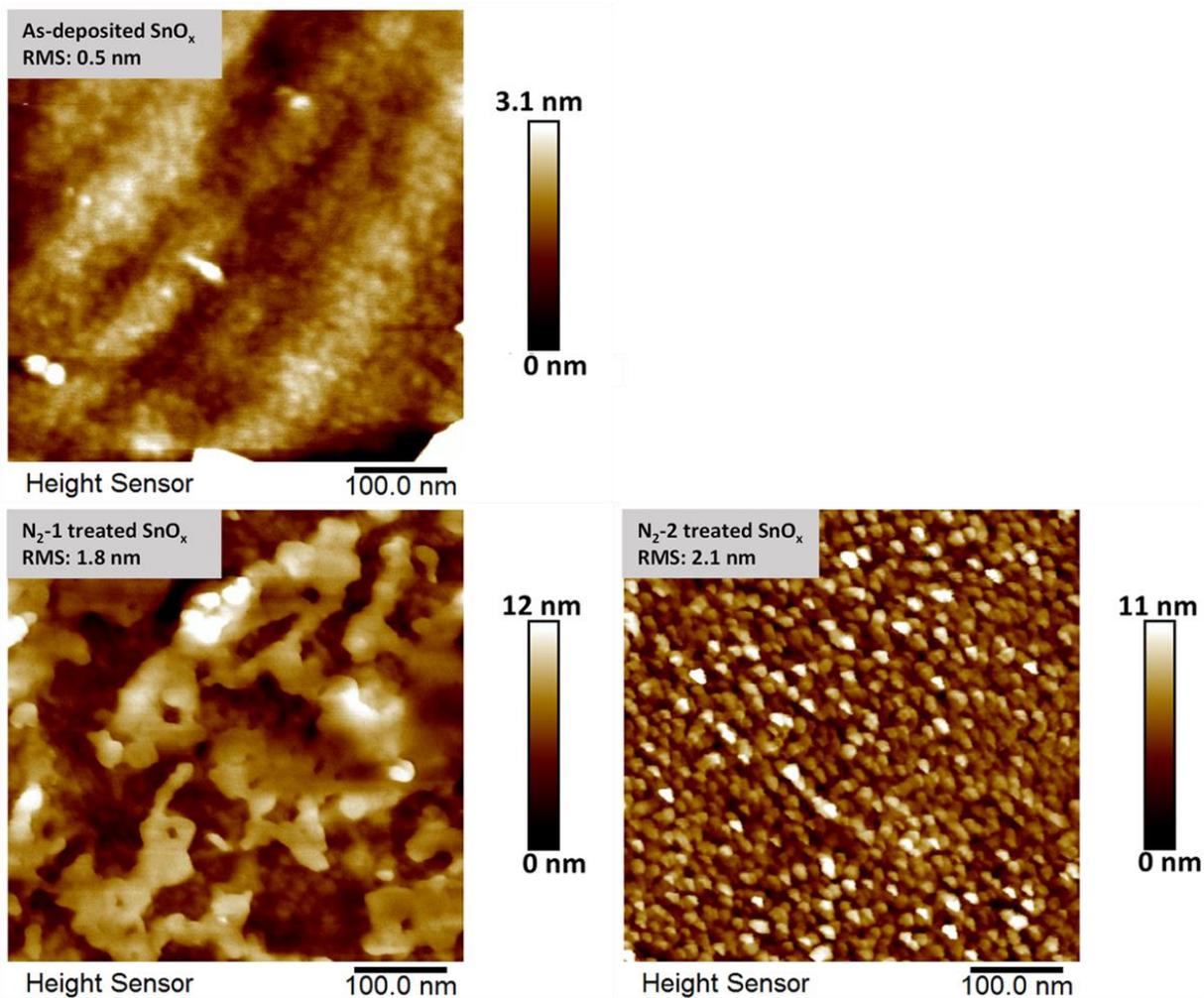
##### a) Structural Properties by GIXRD and surface morphology by AFM

GIXRD pattern for both N<sub>2</sub> plasma treated SnO<sub>x</sub> did not show any peak which indicates that N<sub>2</sub> plasma treatments have resulted in conserving the amorphous microstructure of SnO<sub>x</sub> host films as shown in Figure 8. Moreover, the structural properties by GIXRD did not show any sign of SnO<sub>x</sub> film etching as the SiO<sub>2</sub> broad peak was not observed.



**Figure 8:** GIXRD patterns for as-deposited and N<sub>2</sub> plasma treated SnO<sub>x</sub> films.

AFM analysis has further confirmed the results of the amorphous character related to N<sub>2</sub>O plasma treated SnO<sub>x</sub> films. Indeed, AFM image after N<sub>2</sub>O-1 treatment has shown non-uniform surface morphology without any clear crystalline microstructure. On the other hand, AFM image after N<sub>2</sub>O-2 plasma treatment has revealed a particle-like surface morphology without any clear crystalline microstructure formation as shown in Figure 9. The RMS roughness has increased to 1.8 nm after the N<sub>2</sub>O-1 plasma treatment as compared to 0.5 nm for the reference SnO<sub>x</sub> film, while the RMS roughness has increased to 2.1 nm after the N<sub>2</sub>O-2 plasma treatment.

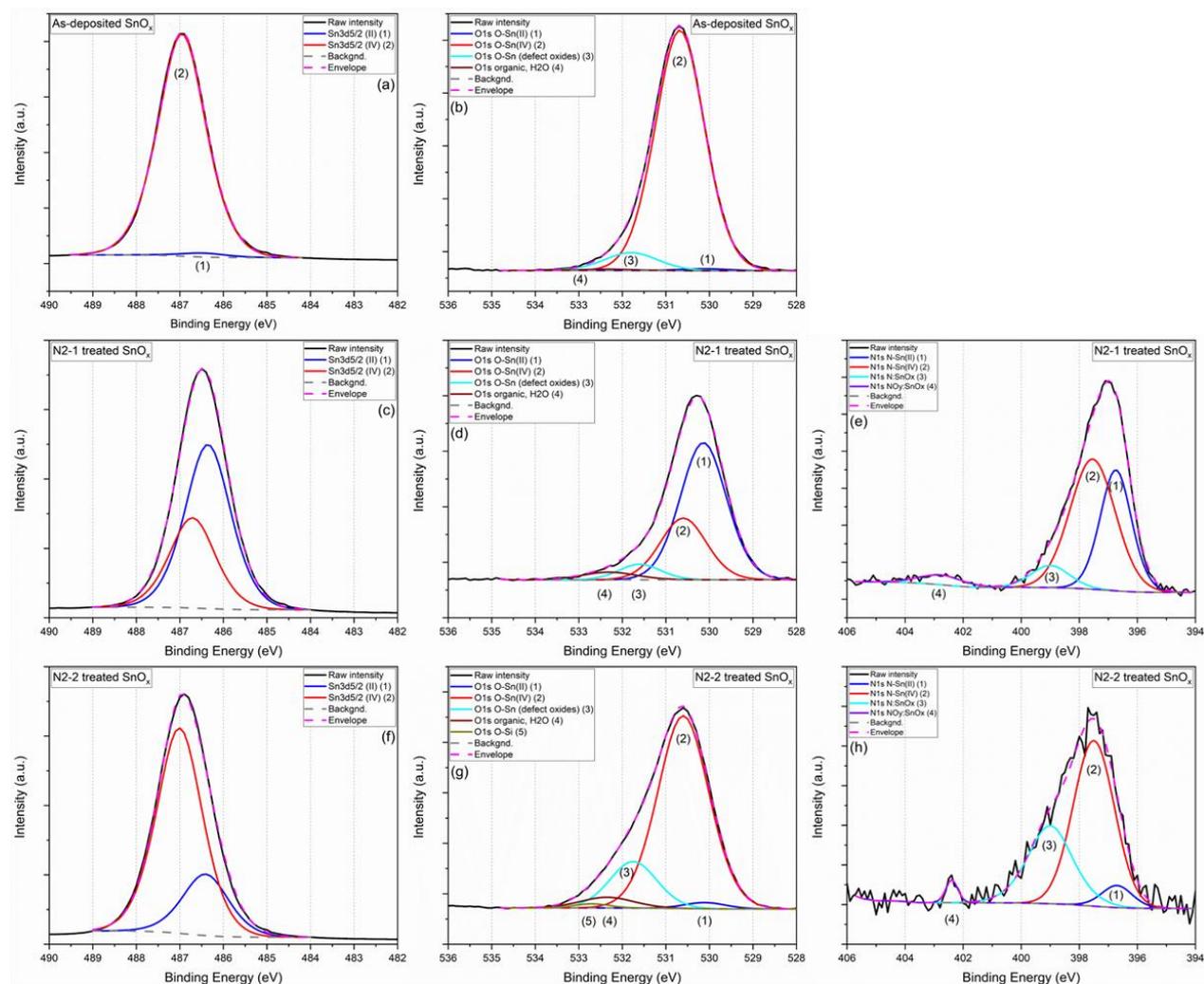


**Figure 9:** AFM images for as-deposited and N<sub>2</sub> plasma treated SnO<sub>x</sub> films.

b) Surface chemical state analysis by XPS

Figure 10 plots the XPS spectra for the N<sub>2</sub> plasma exposed SnO<sub>x</sub> films. The presence of O, Sn, N, C and Si is clearly seen. For N<sub>2</sub>-1 treatment, Sn 3d<sub>5/2</sub> spectrum has revealed that the main chemical state is Sn(II). However, Sn(IV) was found to be a minor chemical state. Nitrogen incorporation has mainly formed N-Sn(II) as well as N-Sn(IV) bonds. N:SnO<sub>x</sub> and NO<sub>y</sub>:SnO<sub>x</sub> are relatively smaller. As per O 1s spectrum, 62% of oxygen formed bonds with Sn(II) and 28% of oxygen have formed bonds with Sn(IV). Oxygen-related SnO<sub>x</sub> defect oxide was found to be 6.5 at. %. For N<sub>2</sub>-2 sample, Sn 3d<sub>5/2</sub> spectra reveal that the main chemical state is Sn(IV) with a smaller fraction of Sn(II). Nitrogen incorporation has mainly formed N-Sn(IV) bonds and partially N:SnO<sub>x</sub>. N-Sn(II) related bonds are minor and there is a small-to-absent fraction of NO<sub>y</sub>:SnO<sub>x</sub>. N<sub>2</sub>-2 treatment has produced the smallest content of molecular NO<sub>y</sub> incorporation and the highest fraction of nitrogen atoms within the SnO<sub>2</sub> and N-Sn(IV) compared to all other treated SnO<sub>x</sub> films. These chemical states are expected to potentially form acceptor sites within the SnO<sub>x</sub> thin film compared to the other samples. On the other hand, the majority of oxygen (around 76%) is bound to Sn(IV) forming the SnO<sub>2</sub> and only a minor fraction of 2% have formed bonds

with Sn(II). There is a relatively small part of oxygen (about 17%) related to SnO<sub>x</sub> defect oxide. The small presence of Si is likely to be a surface contamination which will be discussed in the ToF-SIMS section. The chemical state analysis results are summarized in table 7.



**Figure 10:** XPS Sn3d<sub>5/2</sub>, O1s, and N1s spectra fitting for as-deposited and N<sub>2</sub> plasma treated SnO<sub>x</sub> films.

The summary of all XPS survey and chemical states quantification, shown in Table 7, reveals that N<sub>2</sub>-1 has drastically reduced oxygen atomic concentration and substantially increased nitrogen atomic concentration to 10 at. %. On the other hand, N<sub>2</sub>-2 treatment has reduced the oxygen atomic concentration and favoured the incorporation of nitrogen with similar atomic concentration compared to N<sub>2</sub>O treatments around 2 at. %.

Sample	Survey spectra		High resolution spectra deconvolution				Total Atomic (%)	
	Element	Atomic (%)	Chemical state	Binding Energy (eV)	Area (CPS.eV)	Area (%)		
SnO <sub>x</sub> (as-deposited)	C 1s	0.87	--	--	--	--	0.9	
	N 1s	0	--	--	--	--	0.0	
	Sn3d5/2	34.04	Sn(II)	486.5	19936	1.7	0.6	
			Sn(IV)	487.0	1152286	98.3	33.5	
	O 1s	65.09	O-Sn(II)	530.1	2256	0.79	0.5	
			O-Sn(IV)	530.7	262011	91.32	59.4	
			Defect oxide	531.8	21131	7.37	4.8	
			Organic, H <sub>2</sub> O	532.4	1505	0.52	0.3	
	Total %							100
	SnO <sub>x</sub> _N <sub>2</sub> -1	Si 2p	1.11	--	--	--	--	1.1
C 1s		8.08	--	--	--	--	8.1	
N 1s		10.06	N-Sn(II)	396.7	9048	34.83	3.5	
			N-Sn(IV)	397.5	13915	53.58	5.4	
			N:SnO <sub>x</sub>	399.0	2012	7.75	0.8	
			NO <sub>y</sub> :SnO <sub>x</sub>	402.8	995	3.84	0.4	
Sn3d5/2		31.78	Sn(II)	486.4	841168	84.5	26.9	
			Sn(IV)	486.9	154227	15.5	4.9	
O 1s		48.97	O-Sn(II)	530.1	133793	61.79	30.3	
			O-Sn(IV)	530.6	60207	27.81	13.6	
	Defect oxide		531.6	13723	6.34	3.1		
	Organic, H <sub>2</sub> O		532.3	8711	4.03	2.0		
	SiO <sub>2</sub>		532.8	78	0.04	0.0		

							Total %	100
SnO <sub>x</sub> _N <sub>2</sub> -2	Si 2p	5.16	--	--	--	--	5.2	
	C 1s	0.62	--	--	--	--	0.6	
	N 1s	2.27	N-Sn(II)	396.7	378	6.06	0.1	
			N-Sn(IV)	397.5	3722	59.7	1.4	
			N:SnO <sub>x</sub>	399.0	1984	31.83	0.7	
			NO <sub>y</sub> :SnO <sub>x</sub>	402.4	150	2.41	0.1	
	Sn 3d5/2	30.87	Sn(II)	486.4	53603	5	1.5	
			Sn(IV)	486.9	1019189	95	29.3	
	O 1s	61.08	O-Sn(II)	530.1	6087	2.01	1.2	
			O-Sn(IV)	530.6	231310	76.26	46.6	
			Defect oxide	531.7	46022	15.18	9.3	
			Organic, H <sub>2</sub> O	532.2	391	0.13	0.1	
			SiO <sub>2</sub>	532.6	19485	6.43	3.9	
	Total %							100

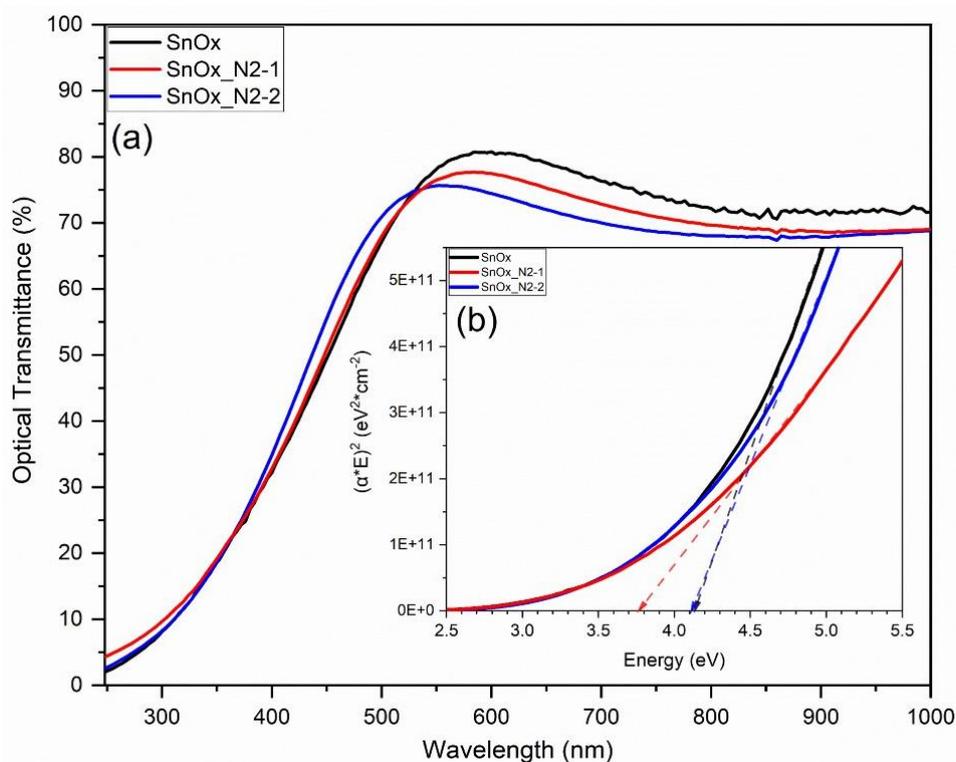
**Table 7:** XPS survey and chemical states quantification for as-deposited SnO<sub>x</sub> and N<sub>2</sub> plasma treated SnO<sub>x</sub> films.

### c) Optical Properties

The N<sub>2</sub>-1 treated SnO<sub>x</sub> film reveals an average optical transmittance of 73.8% which is lower compared to as-deposited SnO<sub>x</sub> of 76.5%. Nevertheless, N<sub>2</sub>-2 treated SnO<sub>x</sub> was a subject of a more decrease in average optical transmittance down to 71.8% compared to the as-deposited SnO<sub>x</sub>. These results are similar to the ones of N<sub>2</sub>O treatments. This lower optical transmittance is suggested to be related to the incorporation of nitrogen within SnO<sub>x</sub> which led to the decrease in oxygen stoichiometry as shown in chemical state analysis of N<sub>2</sub> treated SnO<sub>x</sub> (table 7).

Both N<sub>2</sub> treatments for SnO<sub>x</sub> has led to decrease the optical bandgap as shown in Figure 11. N<sub>2</sub>-1 treated SnO<sub>x</sub> has an optical bandgap of 3.75 eV compared to 4.15 eV for as-deposited SnO<sub>x</sub> film. On the other hand, the N<sub>2</sub>O-2 treatment has led to an optical bandgap of 4.12 eV compared to the optical

bandgap of as-deposited SnO<sub>x</sub> as shown in Table 8. As stated previously, this decrease can be originated from the high disorder within the SnO<sub>x</sub> film due to nitrogen enrichment treatment. However, the high decrease in optical bandgap following N<sub>2</sub>-1 treatment at 100°C compared to slight decrease in N<sub>2</sub>-2 treatment at 300°C is likely due to the lower temperature of the treatment N<sub>2</sub>-1 which did not enable the material restructuring and lowering the disorder within the SnO<sub>x</sub> host film. These results are also in agreement with previous reports in the literature [22, 23, 29, 40, 41].



**Figure 11:** (a) Optical transmittance and (b) Tauc plot for as-deposited and N<sub>2</sub> treated SnO<sub>x</sub> films.

Samples	Average optical transmittance (%)	Optical bandgap (eV)
SnO <sub>x</sub> as deposited	76.5	4.15
N <sub>2</sub> -1 treated SnO <sub>x</sub>	73.8	3.86
N <sub>2</sub> -2 treated SnO <sub>x</sub>	71.8	4.12

**Table 8:** Average optical transmittance from 500 to 800 nm and optical bandgap energy for as-deposited and N<sub>2</sub> treated SnO<sub>x</sub> films.

#### d) Electrical Properties

As for the N<sub>2</sub> plasma treatment, N<sub>2</sub>-1 sample has shown a higher resistivity which might be due to a very high incorporation of nitrogen within the layer. In fact, this treatment has drastically reduced the Sn chemical state from major Sn(IV) to major Sn(II) as revealed by XPS data analysis. This is in

agreement with a previous work which reported that N doping for SnO thin films can increase the resistivity<sup>[41]</sup>.

On the other hand, N<sub>2</sub>-2 sample has shown a moderate resistivity of 4.73 Ω·cm but more importantly it has a p-type character with a hole carrier concentration of about 6.7 x 10<sup>17</sup> cm<sup>-3</sup> (p-type). These results are presented in table 9. As per the literature, this p-type conductivity is related to nitrogen incorporation within the SnO<sub>x</sub>, namely N:SnO<sub>2</sub> and N-Sn(IV). These chemical states of nitrogen form acceptor sites<sup>[27-29]</sup>. The minimal-to-absence of NO<sub>y</sub> formation, which acts as a double donor, has enabled the p-type conductivity. The reduced charge carrier concentration is likely due to the donor compensation<sup>[43]</sup>. This treatment has increased slightly the mobility which can be attributed to the higher treatment temperature at 300 °C compared to the deposition temperature at 100 °C. This p-type conductivity was not permanent and stable over time as the Hall effect measurements have shown a mixture of p-type and n-type after the first measurement. This might be related to the presence of mixed phases as shown previously by surface characterization (XPS) and/or it is related to presence of deep energy level related to nitrogen doping as reported previously<sup>[44]</sup>.

Samples	Thickness (nm)	Resistivity (Ω·cm)	Mobility (cm <sup>2</sup> /V·s)	Charge carrier (cm <sup>-3</sup> )	Conductivity
SnO <sub>x</sub> as deposited	120	2.52 x 10 <sup>-1</sup>	1.45	1.71 x 10 <sup>19</sup>	n-type
N <sub>2</sub> -1 treated SnO <sub>x</sub>	121	N/A*	N/A	N/A	N/A
N <sub>2</sub> -2 treated SnO <sub>x</sub>	113	4.73	1.96	6.71 x 10 <sup>17</sup>	p-type

**Table 9:** Resistivity, mobility, and charge carrier concentration for as-deposited and N<sub>2</sub> plasma treated SnO<sub>x</sub> films.

N/A\* represents a SnO<sub>x</sub> with a very high resistivity.

#### e) Optoelectronic Performance

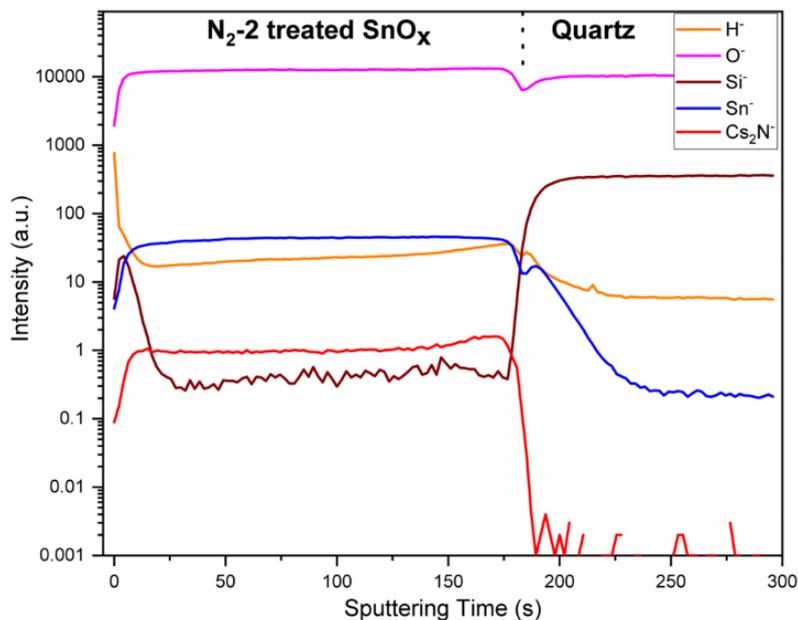
The optoelectronic performance for all SnO<sub>x</sub> samples was evaluated using Haacke figure of merit (FoM)<sup>[42]</sup>. The results are shown in Table 10. The FoM of N<sub>2</sub>-2 plasma treated SnO<sub>x</sub> film could not be determined due to the very high electrical resistivity. The N<sub>2</sub>-2 plasma treated SnO<sub>x</sub> film has resulted in a lower figure of merit around 10<sup>-7</sup> Ω<sup>-1</sup> for an average transmittance between 500 to 800 nm compared to the as-deposited SnO<sub>x</sub>. This is due to the higher resistivity of the N<sub>2</sub>-2 plasma treated SnO<sub>x</sub> film.

Samples	Figure of merit ( $10^{-3} \Omega^{-1}$ )		
	Average Transmittance 400 – 700 nm	Average Transmittance 500 – 800 nm	Transmittance at 550 nm
As-deposited SnO <sub>x</sub>	$1.04 \times 10^{-03}$	$3.27 \times 10^{-03}$	$4.28 \times 10^{-03}$
N <sub>2</sub> -1 plasma treated SnO <sub>x</sub>	N/A	N/A	N/A
N <sub>2</sub> -2 plasma treated SnO <sub>x</sub>	$6.06 \times 10^{-05}$	$9.72 \times 10^{-05}$	$1.56 \times 10^{-04}$

**Table 10:** Figure of Merit for as-deposited and N<sub>2</sub>O plasma treated SnO<sub>x</sub> thin films.

f) Depth Profiling for p-type conductive N doped SnO<sub>x</sub>

To get more insight on the incorporation of nitrogen into the SnO<sub>x</sub> host matrix, ToF-SIMS analysis were carried out on the N<sub>2</sub>-2 sample which revealed a p-type conductivity character. In fact, atomic nitrogen detection can be challenging due to its low sputter yield. On the other hand, Cs<sub>x</sub>N molecular ions, which are generated during the depth profile sputtering at high energy Cs ion beam, are reported to have a higher yield which can be used to probe nitrogen in both positive and negative polarity<sup>[45, 46]</sup>. Among the Cs based ions for this depth profiling, Cs<sub>2</sub>N<sup>-</sup> ions have revealed a much better sensitivity (sputter yield) and provided a stronger intensity. The depth profiles for several elements in the N<sub>2</sub>-2 treated SnO<sub>x</sub> are shown in Figure 12. This nitrogen doped SnO<sub>x</sub> thin film seems to have a good chemical uniformity throughout the whole layer as shown by the flat intensities of O<sup>-</sup> and Sn<sup>-</sup>. The presence of silicon as a surface contaminant is found which confirms the XPS data reported above. The presence of carbon is also confirmed to be a surface contamination related the adventitious carbon. More interestingly, the Cs<sub>2</sub>N<sup>-</sup> distribution is uniform within the N:SnO<sub>x</sub> region. This indirectly confirms that nitrogen atoms have been successfully incorporated uniformly in the host SnO<sub>x</sub> matrix.



**Figure 12:** ToF-SIMS depth profiles for N<sub>2</sub>-2 treated SnO<sub>x</sub> film.

#### g) Summary

For N<sub>2</sub> plasma treatment, the moderate N<sub>2</sub>-1 treatment resulted in a high incorporation of nitrogen as well as a substantial reduction of Sn(IV) to Sn(II). This has led to a large increase in electrical resistivity. The intense N<sub>2</sub>-2 treatment instead resulted in a much smaller incorporation of nitrogen within SnO<sub>x</sub> host film which has led to a relatively smaller increase in electrical resistivity but with a conductivity change from n-type towards p-type. The increase in electrical resistivity has significantly reduced the optoelectronic performance of N<sub>2</sub> plasma treated SnO<sub>x</sub> films compared to as-deposited SnO<sub>x</sub>. XPS studies for the case of intense N<sub>2</sub>-2 plasma treatment revealed the presence mainly of the SnO<sub>2</sub> phase with a minor presence of the SnO phase, and especially the presence of nitrogen in different chemical states (see Figure 10f, g, h). The effective incorporation of elemental nitrogen into the SnO<sub>x</sub> layer was confirmed by ToF-SIMS measurements (Figure 12) which showed that the nitrogen distribution profile is uniform over the entire depth.

#### 4.4. Combined post-annealing and N doping effect

As discussed in our previous studies, both processes of thermal annealing under inert atmosphere namely, Ar and the N incorporation within the SnO<sub>x</sub> have led to a potential ability to convert the conductivity of the SnO<sub>x</sub> thin film from n-type to p-type. Therefore, the aim of this study is to explore the effect of synergy between two potential processes in order to improve the p-type conductivity of the SnO<sub>x</sub> thin films. The applied two step processes are: annealing process which is thermal treatment in Ar atmosphere at 600 °C for 35 minutes and N<sub>2</sub> plasma treatment using ECR-PECVD reactor, which was described previously in this Chapter, whose temperature is 300 °C. For plasma treatment, Microwave power is 500 W, N<sub>2</sub> flow is set at 20 sccm, acceleration voltage at 150 V and treatment time of 30 minutes. The first set of samples were prepared starting by the N<sub>2</sub> plasma treatment and finishing by the annealing (named N-T) while the second set of samples was prepared starting with the annealing and finishing with the N<sub>2</sub> plasma treatment (named T-N).

##### 4.4.1. N<sub>2</sub> intense treatment for SnO<sub>x</sub> followed by thermal annealing under Ar (N-T case)

The SnO<sub>x</sub> film deposited at O<sub>2</sub>/Ar ratio of 0.015 and treated by an intense N<sub>2</sub> plasma, followed by thermal annealing under Ar, has shown a mixture of conductivity type measurements. On the other hand, The SnO<sub>x</sub> film deposited at O<sub>2</sub>/Ar ratio of 0.035, 0.055 and 0.065, and treated by an intense N<sub>2</sub> plasma, followed by thermal annealing under Ar, have all shown a p-type conductivity while the one deposited at O<sub>2</sub>/Ar of 0.075, has shown a n-type conductivity. There is no clear trend for the N<sub>2</sub> plasma treatment followed by thermal annealing compared to thermal annealing alone. However, the SnO<sub>x</sub> deposited at O<sub>2</sub>/Ar of 0.055 has shown a p-type conductivity in both sets of samples. The electrical

resistivity for all samples had no significant change compared to the SnO<sub>x</sub> samples which were thermally annealed only without N<sub>2</sub> plasma treatment.

It is worth noting that SnO<sub>x</sub> deposited at O<sub>2</sub>/Ar of 0.055 exhibit the highest charge carrier concentration of  $9.24 \times 10^{17} \text{ cm}^{-3}$  compared to the other samples which have concentrations in the range of  $10^{16} \text{ cm}^{-3}$ . However, its mobility of  $0.18 \text{ cm}^2/\text{V}\cdot\text{s}$  is low compared to the mobility values of the other samples that range between 1.1 to  $3.9 \text{ cm}^2/\text{V}\cdot\text{s}$ . All the results are presented in Table 11. The step of N<sub>2</sub> plasma treatment prior to the thermal annealing did not significantly improve the p-type conductivity character.

#### 4.4.2. Thermal annealing under Ar for SnO<sub>x</sub> followed by N<sub>2</sub> intense treatment (T-N case)

The SnO<sub>x</sub> film deposited at O<sub>2</sub>/Ar ratio of 0.015 and treated by thermal annealing under Ar, followed by an intense N<sub>2</sub> plasma, has shown a mixture of conductivity type measurements. This result is similar for SnO<sub>x</sub> film deposited at O<sub>2</sub>/Ar ratio of 0.015 with N<sub>2</sub> plasma treatment prior to the thermal annealing. Furthermore, all the other SnO<sub>x</sub> samples have all shown a n-type conductivity. There is a clear trend for the thermal annealing followed by N<sub>2</sub> plasma treatment of converting the p-type conductive SnO<sub>x</sub> samples to n-type conductivity as well as of improving the n-type conductivity for n-type SnO<sub>x</sub> samples.

The n-type charge carrier concentration has significantly improved after the N<sub>2</sub> plasma treatment compared to the SnO<sub>x</sub> samples deposited at O<sub>2</sub>/Ar of 0.045, 0.055 and 0.065 followed only by annealing, their charge carrier concentration increased from  $2.1 \times 10^{16}$ ,  $5.5 \times 10^{15}$  and  $1.4 \times 10^{17} \text{ cm}^{-3}$  to  $4.1 \times 10^{18}$ ,  $3.4 \times 10^{18}$  and  $3.4 \times 10^{18} \text{ cm}^{-3}$ , respectively. For the p-type conductive SnO<sub>x</sub> samples, the additional step of N<sub>2</sub> plasma treatment has changed the charge carrier concentration from relatively low hole (p-type) carrier concentration to high electron (n-type) carrier concentration in the range of  $10^{18} \text{ cm}^{-3}$ . This is due to the charge compensation which was discussed previously in this Chapter.

As the annealing is expected to form SnO phase and form hole charge carriers<sup>[43, 44, 47-51]</sup>, the N<sub>2</sub> plasma treatment for the SnO phase is expected to form n-type charge carriers<sup>[52]</sup>. The electrical resistivity for all samples had a substantial decrease with an n-type conductivity. This is in comparison to the SnO<sub>x</sub> samples which were thermally annealed only, without N<sub>2</sub> plasma treatment, and which have mixed conductivity type. All the electrical measurements' results by Hall effect are summarized in table 11.

Charge carrier concentration ( $\text{cm}^{-3}$ )

O <sub>2</sub> /Ar ratio	As deposited	Annealed at 600°C only - 35 minutes	Treatment → Annealing	Annealing → Treatment
0.015	-1.34E+15	1.51E+16	N/A**	N/A**
0.035	N/A**	-2.13E+16	4.02E+16	-4.10E+18
0.055	N/A*	8.24E+17	9.24E+17	-4.46E+18
0.065	N/A*	-1.37E+17	1.63E+16	-3.35E+18
0.075	N/A*	1.01E+16	-8.21E+16	-3.46E+18

Resistivity ( $\Omega\cdot\text{cm}$ )

O <sub>2</sub> /Ar ratio	As deposited	Annealed at 600°C - 35 minutes	Treatment → Annealing	Annealing → Treatment
0.015	2.00E+02	1.68E+02	1.45E+02**	2.06E+00**
0.035	2.39E+02**	1.91E+02	1.73E+02	3.44E+00
0.055	N/A*	1.66E+02	1.67E+02	3.16E+00
0.065	N/A*	1.31E+02	1.13E+02	3.71E+00
0.075	N/A*	1.86E+02	1.28E+02	4.43E+00

Mobility ( $\text{cm}^2/\text{V}\cdot\text{s}$ )

O <sub>2</sub> /Ar ratio	As deposited	Annealed at 600°C - 35 minutes	Treatment → Annealing	Annealing → Treatment
0.015	2.38E+01	3.27E+00	1.05E+00**	6.07E-01**
0.035	4.48E+00**	1.67E+00	1.18E+00	4.48E-01
0.055	N/A*	6.90E-02	1.76E-01	4.46E-01
0.065	N/A*	3.53E-01	3.91E+00	5.09E-01
0.075	N/A*	3.40E+00	1.07E+00	4.14E-01

**Table 11:** Hall effect measurement for as deposited SnO<sub>x</sub> followed by thermal annealing then nitrogen plasma treatment and as deposited SnO<sub>x</sub> followed by nitrogen plasma treatment then thermal annealing.

\* Values with orange background represent n-type conductive samples and blue represent p-type conductive samples. Values with white background represent the samples where the conductivity type could not be determined. The samples with the mark. N/A\*: high resistivity SnO<sub>x</sub>. \*\*: conductivity type could not be determined

#### 4.4.3. Summary

We have demonstrated that carrying out an intense N<sub>2</sub> plasma treatment followed by a thermal annealing under Ar might lead to an improvement of the p-type conductivity. However, reversing the order of these two processes (annealing under Ar followed by intense N<sub>2</sub> plasma treatment) can lead to the opposite effect that the films retain their n-type conductivity or change from p-type to n-type.

#### 4.5. General Summary

P-type N:SnO<sub>x</sub> has been achieved using a two-step method using magnetron sputtering for SnO<sub>x</sub> deposition and N<sub>2</sub> plasma post treatment for SnO<sub>x</sub> related nitrogen doping. To the best of our knowledge, this is the first detailed material characterization, particularly XPS analysis, for nitrogen-doped Sn oxide using three different nitrogen gas precursors (NH<sub>3</sub>, N<sub>2</sub>O and N<sub>2</sub>) with two levels of treatments for each precursor. Unexpectedly, NH<sub>3</sub> treatments have etched and reduced the SnO<sub>x</sub>. N<sub>2</sub>O treatment has improved the n-type conductivity of SnO<sub>x</sub> as well as the optoelectronic performance. N<sub>2</sub> plasma post-treatment has shown a potential ability to switch the conductivity from n-type to p-type. This conductivity conversion from p-type to n-type is associated with a reduction in optoelectronic performance. XPS N 1s spectra analysis for N<sub>2</sub>-2 plasma treated SnO<sub>x</sub> sample showed that the p-type conductivity was favoured by the atomic doping within the SnO<sub>x</sub> as well as the Sn(IV)-N chemical states. The uniformity of the N:SnO<sub>x</sub> throughout the depth was confirmed by ToF-SIMS analysis.

In other study to assess the synergy between two processes which provides p-type conductivity, we have demonstrated that performing N<sub>2</sub> plasma treatment followed by thermal annealing under Ar has the potential to improve the p-type conductivity. However, reversing the order of these two processes can lead to the opposite effect where all samples kept their n-type conductivity or have been converted from p-type back to n-type. Therefore, it is suggested that the thermal annealing process followed by N<sub>2</sub> plasma treatment has formed nitrogen doping for tin monoxide (N-doped SnO) which was reported in the first chapter to form donor charges.

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# General Conclusion

The aim of this thesis was to conduct a detailed study of undoped tin oxide and nitrogen doped tin oxide as a transparent conducting oxide (TCO) using different growth conditions and exploring the material properties thanks to different characterisation techniques, as well as assessing the performance of tin oxide as electron transport layer (ETL) for perovskite solar cells.

The first chapter of this work has focused on discussing the literature review of transparent conducting oxides (TCOs) in general and their respective applications. In particular, the material synthesis and properties, and the applications of  $\text{In}_2\text{O}_3$  and  $\text{ZnO}$  thin films as TCOs have been discussed thoroughly in this chapter. The state-of-the-art of tin oxide ( $\text{SnO}_x$ ), in both chemical states  $\text{SnO}_2$  and  $\text{SnO}$ , has been addressed in depth in terms of material synthesis and properties as well as their applications.

The work in the second chapter presented different tin oxide thin films synthesised by magnetron sputtering using different conditions namely, sputtering target material, substrate temperature,  $\text{O}_2/\text{Ar}$  flow rate ratio, and post-deposition thermal annealing. The results related to oxidized tin as a target have shown that the temperature substrate plays a major role in the structural properties of the grown tin oxide thin films. Thus, as-deposited samples grown at  $250^\circ\text{C}$  exhibited better crystallinity than the ones deposited at  $100^\circ\text{C}$ . On the other hand, the post-thermal annealing at  $400^\circ\text{C}$  under vacuum has improved the structural properties by improving the crystallinity. Nevertheless, after the annealing, the tin oxide grown at  $250^\circ\text{C}$  have kept a better crystalline structure compared the ones grown at  $100^\circ\text{C}$ . Furthermore, annealing under air has improved the crystalline structure for all samples compared to vacuum annealing. Both  $\text{SnO}_x$  deposited films without adding  $\text{O}_2$  but annealed in air have revealed the presence of both phases  $\text{SnO}$  and  $\text{SnO}_2$ . Vacuum annealing has substantially improved the electrical properties of tin oxide samples compared to the air annealing. On the other hand, air annealing has provided slightly better optical transmittance compared to vacuum annealing. The best conductive sample was deposited at  $250^\circ\text{C}$ , at 0.015  $\text{O}_2/\text{Ar}$  flow rate ratio and annealing under vacuum while the best crystalline sample with highest crystallite size, was deposited at  $250^\circ\text{C}$  without  $\text{O}_2$  flow and annealed in air. It is worth noting that the best conductive sample has also the highest optoelectronic performance as assessed by figure of merit (FoM) of  $5.14 \times 10^{-5} \Omega^{-1}$ . As a result of these investigations, the best operational conditions are deposition at  $250^\circ\text{C}$ , at  $\text{O}_2/\text{Ar}$  flow rate ratio of 0.015 and annealing under vacuum.

The results related to  $\text{SnO}_x$  films produced using a metallic tin as a target have shown that the  $\text{O}_2/\text{Ar}$  flow rate ratio plays a major role in the structural properties of the grown films at relatively lower temperature of  $100^\circ\text{C}$ . At lower  $\text{O}_2/\text{Ar}$  flow rate ratios, the deposited tin oxide thin films have an amorphous structure. At higher  $\text{O}_2/\text{Ar}$  flow rate ratios, the deposited tin oxide thin films have a crystalline structure. The structural properties have influenced the optoelectronic properties of  $\text{SnO}_x$  thin films.  $\text{SnO}_x$  thin films have a lower optical bandgap at lower  $\text{O}_2/\text{Ar}$  flow rate ratios indicating the

presence of mainly SnO phase, and the optical bandgap increase at higher O<sub>2</sub>:Ar flow rate ratios indicating the presence of SnO<sub>2</sub>. Tin oxide thin films deposited at lower O<sub>2</sub>/Ar flow rate ratios have a high resistivity, and the resistivity decreases at higher O<sub>2</sub>/Ar flow rate ratios. Substrate temperature did not significantly influence the structural properties of SnO<sub>x</sub> at a low O<sub>2</sub>/Ar flow rate ratio as grazing incident X-ray diffraction shows a typical SnO phase broad peak around 30° for temperatures between room temperature to 300°C. However, there is small minor peak related to SnO<sub>2</sub> for deposition at 400°C. Optical bandgap was lower for all samples indicating the presence of mainly SnO phase which is due to the poor oxygen condition at the growth. The amorphous structure related to SnO phase have resulted in high resistivity for all samples. **As a result of these investigations, the best operational set of conditions which provided the best conductive n-type SnO<sub>x</sub> film as well as the best optoelectronic performance, is deposition at 100°C using metallic Sn target, at 0.33 O<sub>2</sub>/Ar flow rate ratio without any post-deposition annealing.**

In the previous chapter, the annealing under air has shown a better crystallinity of the SnO<sub>x</sub> films grown by reactive magnetron sputtering with SnO target, with potential ability to form SnO phase. It was also observed that the degree of crystallinity is related mainly due to the pressure of the ambient during the annealing. Using different annealing series, p-type conductivity could be achieved. The first series of annealing under Ar ambient have shown that the higher temperature of 600°C has enabled the formation of more p-type conductive samples, annealing at 500°C have enabled less p-type conductive samples while no p-type conductive samples were produced at 400°C. This has confirmed the importance of high temperature during the annealing to achieve p-type conductivity. Moreover, it is also confirmed that the p-type resistivity is related to SnO phase thanks to the transmission electron microscopy diffraction analysis which revealed the presence of SnO and confirmed by the reduction in the bandgap. On the other hand, it is worth noting the SnO<sub>x</sub> is very sensitive to air, and it oxidises easily in presence of O<sub>2</sub>. The experiments investigating the annealing time under Ar ambient has shown that extending the annealing duration to two hours has conducted to the conversion the conductivity from p to n-type while it provides p-type conductivity for shorter time of annealing. **As a result of these investigations, the best operational set of conditions which provided the best conductive p-type SnO<sub>x</sub> film is deposition at 250°C using SnO target, at 0.065 O<sub>2</sub>/Ar flow rate ratio with post-deposition annealing under Ar at 600°C for 50 minutes.**

The ageing effect in air was assessed after seven days, and as a result tin oxide thin films have been converted from p-type to n-type conductivity and the resistivity has decreased. This could be explained by the further oxidation of SnO<sub>x</sub> from SnO towards SnO<sub>2</sub> and the increasing formation of the intrinsic n-type charge carriers which compensated the p-type charge carriers and increased the n-type carriers reducing consequently the resistivity. On the other hand, the effect of presence of moisture on tin

oxide after seven days of the ageing has been assessed using heating process at 150°C in the presence of air. This process did not recover the p-type conductivity and instead it has led to further oxidation of SnO<sub>x</sub> film and further decrease in resistivity by forming more n-type charge carriers.

In the third, the performance of SnO<sub>2</sub> material, as electron transport layer (ETL) in perovskite solar cells (PSCs) without mesoporous scaffold structure, was assessed using radio frequency magnetron sputtered SnO<sub>2</sub>. Power conversion efficiency (PCE) of PSC using as-deposited SnO<sub>2</sub> thin film applied as ETL, is 15.07% while the use of **annealed SnO<sub>2</sub> thin film increases the PCE up to 17.1%**. It can be concluded that a deposition of SnO<sub>2</sub> ETL and its subsequent air annealing process have enabled a better solar cell performance by improving the optoelectronic properties, particularly the electron mobility.

In the last chapter, the nitrogen doping for SnO<sub>x</sub> was studied using three nitrogen gas precursors (NH<sub>3</sub>, N<sub>2</sub>O and N<sub>2</sub>) with two levels of treatments for each precursor. HN<sub>3</sub> treatments have resulted in etching and damaged the SnO<sub>x</sub> thin films. N<sub>2</sub>O treatments have rather improved the n-type conductivity. One treatment under N<sub>2</sub> plasma has increased the nitrogen atoms' incorporation and enhanced the electrical resistivity of the film while the other **N<sub>2</sub> treatment has moderately increased the electrical resistivity while converting the n-type conductivity to p-type**. Applying thermal treatment to the last N<sub>2</sub> treatment condition has been assessed in this chapter, and it was found that such double processes have a potential to improve the p-type conductivity. However, reversing the order of these two steps lead to the opposite effect where all samples kept their n-type conductivity or have been converted from p-type back to n-type. Therefore, the thermal annealing process followed by N<sub>2</sub> treatment has formed nitrogen doped tin monoxide (N-doped SnO) films.

# Perspectives

This thesis work has enabled us to explore SnO<sub>x</sub> thin films as transparent conducting oxide which can be employed in different applications, particularly in the optoelectronic field. We have demonstrated the possibility to customise the SnO<sub>x</sub> material properties by tuning the growth condition and the post-deposition treatments. Yet, there is a room for further investigation such as, the ageing effect which has shown that can alter seriously the material properties of SnO<sub>x</sub>. This is crucial for the grown thin film reliability particularly for the performance of the device where this film will be implemented. Such study is already planned to be conducted in the near future with real world conditions, accelerated conditions using weathering chamber which is available in the laboratory as well as the effect of radiation on SnO<sub>x</sub> thin films which can simulate the space conditions. Another study which was not possible to be included in this thesis due to time limitation and infrastructure availability, is to develop potential sensing and actuator applications beyond the present state-of-the-art SnO<sub>x</sub>.

SnO<sub>x</sub> has shown the ability to achieve p-type conductivity particularly by forming mainly SnO phase thin film. However, there is still a room for improvement regarding the performance and the stability of this p-type conductivity. One rational way to tackle this issue is to stabilise the SnO phase and enhance the p-type conductivity by introducing doping elements which inhibit the further oxidation of SnO to SnO<sub>2</sub>. The selection of the doping element and its concentration is another study that can further investigate the possibility to achieve a high performance and stable SnO thin film.

# Appendix

### 1.1. Literature review of SnOx thin film as Transparent Conducting Oxide

FoM Ref.	Material	Synthesis method	Doping	Eg (eV)	resistivity ( $\Omega.cm$ )	Sheet resistivity ( $\Omega/\square$ )	T (%)	FoM ( $10^{-3} \Omega^{-1}$ )
1 <sup>[53]</sup>	SnO <sub>2</sub> :Sb	Spray Pyrolysis	Yes	3.82	$5.8 \times 10^{-5}$	0.61	88	47.2
2 <sup>[54]</sup>	SnO <sub>2</sub> :Ba	Spray Pyrolysis	Yes	3.85	$6.95 \times 10^{-4}$	69 (calculated)	90	5.02
3 <sup>[55]</sup>	SnO <sub>2</sub> :Sb	Spray Pyrolysis	Yes	3.77	$1.16 \times 10^{-3}$	141 (calculated)	85	1.39
4 <sup>[56]</sup>	SnO <sub>2</sub> :Sb	Spray Pyrolysis	Yes	3.81	$7.89 \times 10^{-4}$	24.9	89	12.5
5 <sup>[57]</sup>	SnO <sub>2</sub> :Ta	Spray Pyrolysis	Yes	3.77	$4.55 \times 10^{-4}$	25.1	81	6.15
6 <sup>[58]</sup>	SnO <sub>2</sub>	Spin Coating	No	3.92	$2.63 \times 10^{-3}$	63	85	3.3
7 <sup>[59]</sup>	SnO <sub>2</sub> :Ti	Sol Gel Spin Coating	Yes	3.83	$1.77 \times 10^{-3}$	42.3	83	3.66
8 <sup>[60]</sup>	SnO <sub>2</sub> :F	Spray Pyrolysis	Yes	NA	$8 \times 10^{-4}$	40	89	2
9 <sup>[61]</sup>	SnO <sub>2</sub> :F	Sol Gel Dip Coating	Yes	3.91	$7 \times 10^{-4}$	14.5	90	24.3
10 <sup>[62]</sup>	SnO <sub>2</sub> :F	Chemical Vapour Deposition / Pulsed Laser Annealed	Yes	NA	NA	8.8	85	23.2
11 <sup>[63]</sup>	SnO <sub>2</sub> :Te	Pulsed Laser Deposition	Yes	3.5	0.2	$2.22 \times 10^4$	87	0.012
12 <sup>[64]</sup>	SnO <sub>2</sub> :La	Spray Pyrolysis	Yes	3.86	NA	2.1	80	20
13 <sup>[65]</sup>	SnO <sub>2</sub> :Li	Spray Pyrolysis	Yes	3.89	1.19	2.05	80	0.02
14 <sup>[66]</sup>	SnO <sub>2</sub> :F	Atmospheric Pressure Chemical Vapour Deposition	Yes	4.15	$8.4 \times 10^{-4}$	10.5	84	16.7
15 <sup>[67]</sup>	SnO <sub>2</sub> :F	Spray Pyrolysis	Yes	NA	$4 \times 10^{-4}$	3.71	87	61.8
16 <sup>[68]</sup>	SnO <sub>2</sub> :Ta	Spray Pyrolysis	Yes	4.01	$4.36 \times 10^{-4}$	17.96	85	47.3

17 <sup>[69]</sup>	SnO <sub>2</sub> :F	Spray Pyrolysis	Yes	NA	1.14 x 10 <sup>-3</sup> (calculated)	7.48	65	1.93
18 <sup>[70]</sup>	SnO <sub>2</sub> :GO	Spray Pyrolysis	Yes	3.69	NA	26	77	2.83
19 <sup>[71]</sup>	SnO <sub>2</sub> :Gd	Spray Pyrolysis	Yes	3.63	1.02 x 10 <sup>-3</sup>	27	86	8.2
20 <sup>[72]</sup>	SnO <sub>2</sub>	Sputtering	No	3.86	4.45 x 10 <sup>-3</sup>	287 (calculated)	84	0.57
21 <sup>[73]</sup>	SnO <sub>2</sub> :Sb	Spray Pyrolysis	Yes	NA	1.22 x 10 <sup>-3</sup>	15.42	71	2.11
22 <sup>[74]</sup>	SnO <sub>2</sub> :As	Chemical Vapour Deposition	Yes	3.87	1.5 x 10 <sup>-4</sup>	15	85	13.1
23 <sup>[75]</sup>	SnO <sub>2</sub> :Sb	Spray Pyrolysis	Yes	3.88	8 x 10 <sup>-4</sup> (calculated)	40	88	6.9
24 <sup>[75]</sup>	SnO <sub>2</sub> :F	Spray Pyrolysis	Yes	4.07	3.8 x 10 <sup>-4</sup> (calculated)	20	93	24.19
Present Work	SnO <sub>2</sub>	Sputtering	No	4.4	7.2 x 10 <sup>-2</sup>	4.8 x 10 <sup>4</sup>	87	0.05

**Table 1:** Summary of the FoM values and SnO<sub>x</sub> electrical and optical properties selected from relevant literature [53-75].

## 1.2. Perovskite Solar Cell and SnO<sub>2</sub> as Electron Transport Layer

### a) Solar Cells overview

Nowadays, the world is constantly seeking renewable sources of energy due to the increasing energy demand. Predictions are already projecting that the global energy demand will exceed 30 TW by 2050 as shown in Figure 1 and while this demand is quite immense, a major part of it must be renewable to sustain the energy supply for the long term. As the sun is one of the most sustainable sources of energy, solar cells are playing an essential role in supporting the current and future energy demand. The spectral irradiance and contained power of various solar spectra are shown in Figure 2. There is still a lot of effort to employ more solar cells and solar energy in general at regional, national, and international levels. To tackle this issue, scientists, engineers, experts, and policy makers have to work on different aspects particularly, science and technology, economics, and policies. In terms of science and technology, solar cells need to be more efficient and reliable to compete with other sources of

energy. However, the most important factor is to be more cost-effective in order to reach the grid parity.

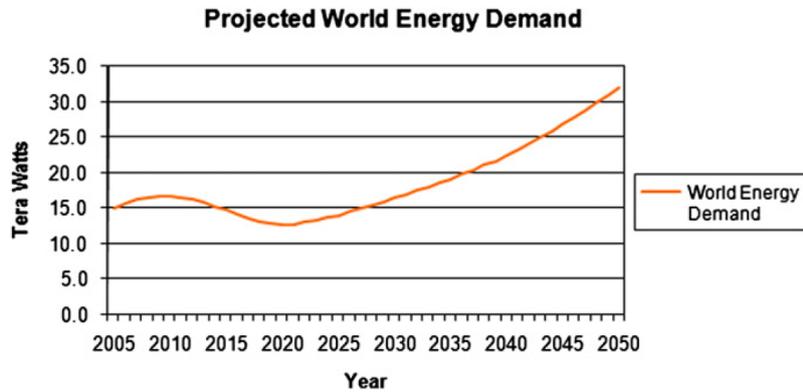


Figure 1: World energy demand projection<sup>[24]</sup>.

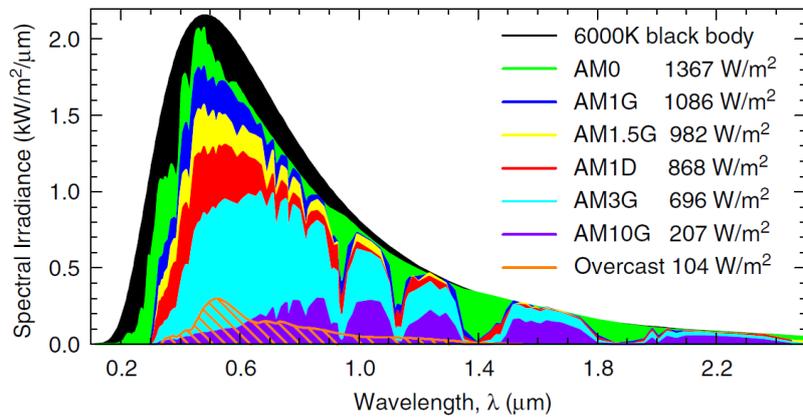


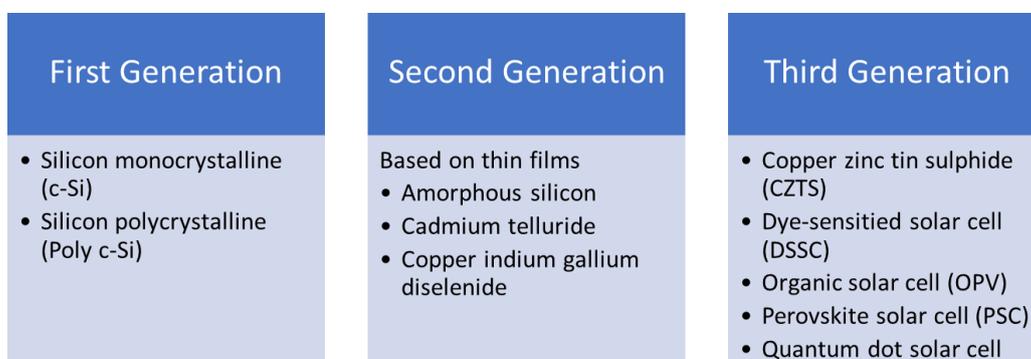
Figure 2: Spectral irradiance and contained power of various solar spectra<sup>[25]</sup>.

#### b) Introduction to Solar Cells

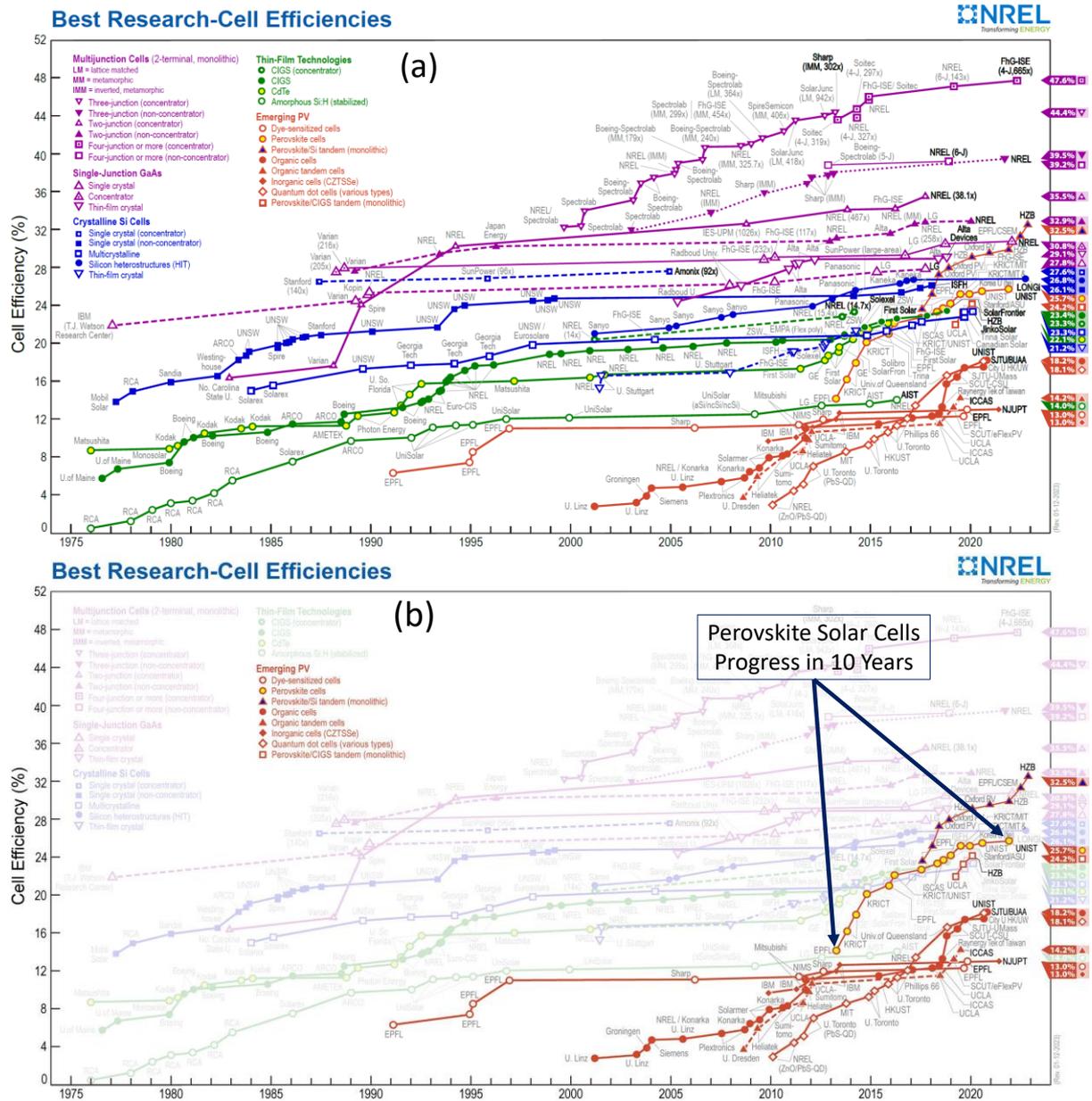
It is important to highlight the properties of the solar spectral irradiance and its related power as it is the source of an enormous energy provided by the sun. The spectral irradiance has a behaviour of black body at 5800 K and its maximum irradiance is at 500 nm. Air mass (AM) coefficient which defines the length of the optical path across the atmosphere. AM0 is the solar spectrum outside the atmosphere with an average solar irradiance power of 1367 W/m<sup>2</sup> and AM1.5 is related to 1.5 times atmosphere thickness where the solar zenith angle is  $z = 48.2^\circ$  with an average solar irradiance power of 982 W/m<sup>2</sup> (nearly 1 kW/m<sup>2</sup>). In solar related technologies, AM1.5 is used a standard for solar irradiance on earth for solar cells and panels.

Solar cells, which are also technically known as photovoltaic cells, are solar light harvesting electronic devices based on semiconducting materials where they convert light from solar spectral irradiance into electricity directly and without any intermediate form of energy. The very basic principle of

photovoltaic energy conversion, which is the main principle of solar cells, consists of three stages. The first is to absorb a photon which has a sufficient energy  $h\nu$  by a semiconducting material (absorber) to generate a pair of hole/electron. The energy of the photon must be equal or higher than the bandgap of the absorber. The second phase is to separate these charge carriers towards the two different contacts, electrons towards the negative contact and holes towards the positive contact. The third is to extract these charges to an external electrical circuit. This whole process enables the generation of electrical current. The basic structure of the solar cell device is based on a p-n junction diode made of two semiconducting materials one is p-type and the other is n-type. There are many technological-based types of solar cells such as silicon solar cells, thin film solar cells and perovskite solar cells as shown in Figure 3 and their performance (power conversion efficiency) progress is illustrated in NREL chart in Figure 4<sup>[24, 26]</sup>.



**Figure 3:** Solar cells classification.



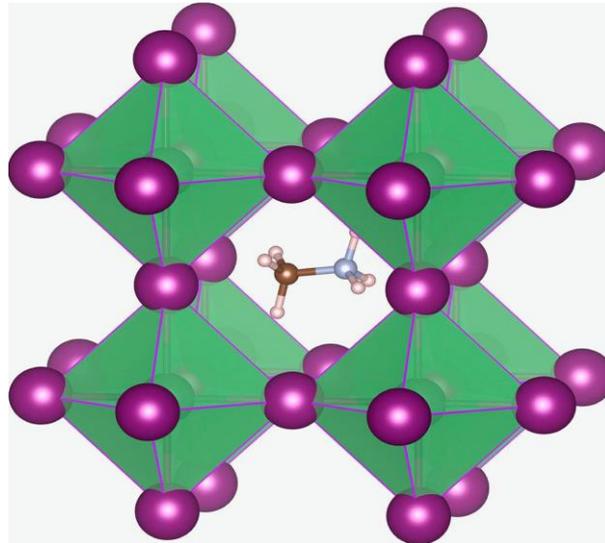
**Figure 4:** (a) NREL chart for best research-cell efficiencies, (b) NREL chart highlighting Perovskite Solar Cells progress in 10 years<sup>[27]</sup>.

c) Perovskite Solar Cell

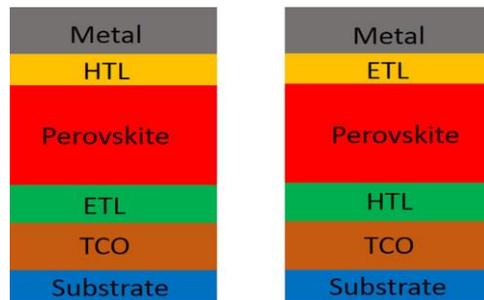
There is a remarkable improvement in various technologies of solar cells in the last decades as shown in NREL chart in Figure 4(a). However, the improvement in performance of hybrid organic-inorganic Perovskite Solar Cells (PSCs) is substantial in the last 15 years. The power conversion efficiency (PCE) has improved from 3.8% in 2009 achieved by researchers from The University of Tokyo to PCE of 25.6% in 2021 achieved by researchers from Ulsan National Institute of Science and Technology (UNIST) as shown in NREL chart in Figure 4(b). Organic-inorganic perovskite is very attractive material as an absorber layer due to its appealing physical properties such as high light absorption coefficient, the ability to tolerate high concentration of defect, long carrier diffusion. Perovskite materials have the

structure of  $ABX_3$  and the first perovskite material, which was discovered a long time back in 1839, was calcium titanium oxide  $CaTiO_3$ . Organic-inorganic perovskite also follow the same structure of  $ABX_3$  where A can be an organic, inorganic, or mixed cation such as  $CH_3NH_3^+$ ,  $CH_3CH_2NH_3^+$ ,  $Cs^+$ . B is a metallic cation namely,  $Pb^{2+}$ ,  $Sn^{2+}$ ,  $Ge^{2+}$ , and X is related to halogen anion namely,  $I^-$ ,  $Br^-$ ,  $Cl^-$ ,  $F^-$ . It is worth noting that the cation A is surrounded by  $BX_6$  octahedra as shown in Figure 5.

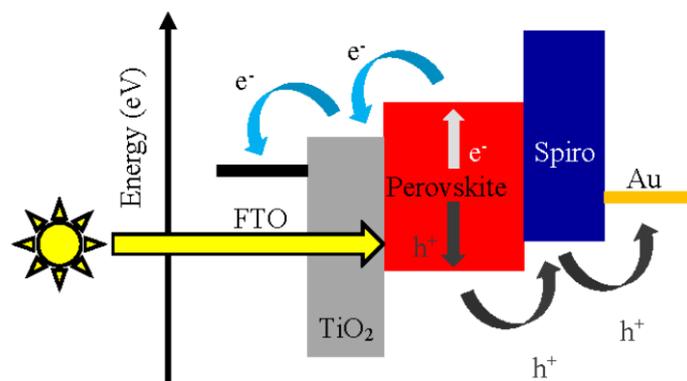
There are various types of architecture for perovskite solar cells. However, there are some layers which are essential for all PSC devices. These layers are the perovskite absorber layer, the electron transport layer (ETL), the hole transport layer (HTL), the transparent conducting oxide layer (anode), the metal contact (cathode), and a transparent substrate which is usually glass. There are two basic architectures of perovskite solar cells, the first is called conventional PSC n-i-p structure which is illustrated in Figure 6(a) and the second is called inverted and has p-i-n structure which is in Figure 6(b). The notation n-i-p is referring to ETL(n)-Absorber(i)-HTL(p) which means ETL on the TCO/Substrate first. The notation p-i-n means the HTL on TCO/Substrate first. The mechanism of PSC is that the perovskite absorber layer with a defined bandgap would absorb a photon, of equal or higher energy than the bandgap, and convert it into a pair of electron and hole. The electrons are transported through ETL while the holes are transported through HTL. These charges (holes and electrons) are collected via contacts anode and cathode and utilised by an external circuit. ETL, HTL and perovskite layer must be selected by their bandgaps and the position of highest occupied molecular orbital (HOMO) and lowest unoccupied molecular orbital (LUMO) to assure a proper band alignment to allow the electrons to move spontaneously towards and through the ETL and vice versa for the holes towards and through the HTL as shown in Figure 7, where FTO ( $F:SnO_2$ ) is TCO,  $TiO_2$  is the ETL, Spiro is HTL and Au is the metallic contact [27-32].



**Figure 5:** Cubic crystal structure ( $ABX_3$ ) of perovskite: A is methylammonium ( $CH_3NH_3^+$ ) surrounded by  $PbX_6$  octahedra <sup>[30]</sup>.



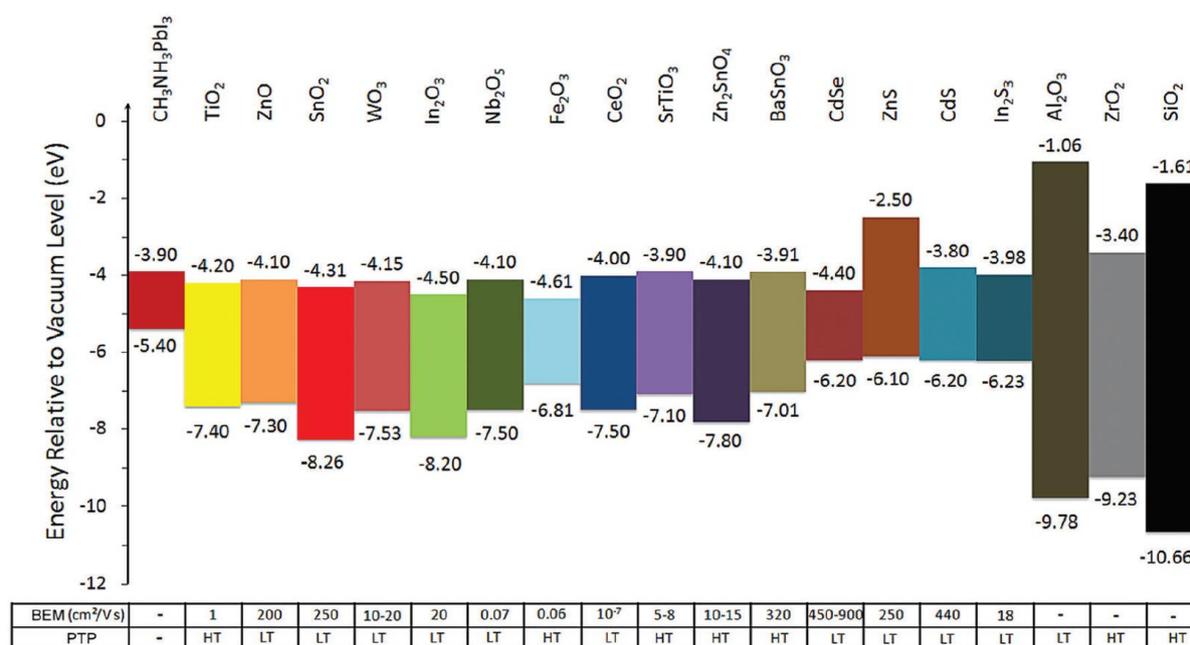
**Figure 6:** Perovskite solar cell device: (a) conventional architecture, (b) inverted architecture <sup>[30]</sup>.



**Figure 7:** Operating mechanism of perovskite solar cell in a conventional architecture <sup>[32]</sup>.

d) ETL in Perovskite Solar Cell

Electron Transport Layer (ETL) is one of the essential layers in perovskite solar cells. An efficient ETL must be aligned successfully with the perovskite absorber layer as described in the previous section. Therefore, both ETL LUMO and HOMO should be higher than the ones of perovskite layer. ETL layer should be a wide bandgap with high optical transmittance in the UV spectra. The generated excitons through the absorption of light within the perovskite layer have to be dissociated prior the collection by ETL or HTL. This dissociation can take place at the interfaces of the absorber perovskite layer. Several metal oxides have been heavily explored at ETL or as a scaffold material such as, TiO<sub>2</sub>, ZnO and SnO<sub>2</sub>. Each of these materials have certain advantages to improve the overall performance of PSC. Figure 8 shows that Energy levels, bulk electron mobility (BEM) and the preferable temperature processing (PTP) of electron transport material in perovskite solar cells. Therefore, there are several metal oxides than can be aligned with perovskite in PSC. Among all these metal oxides, the most appealing ones which have a suitable band alignment with perovskite, requires low temperature processing, and can potentially achieve a high mobility are ZnO, SnO<sub>2</sub> and CdSe. It is worth to note that low processing temperature is an important advantage for industrial applications. Although, ETL made of CdSe can achieve a very high mobility which is the highest among all these materials, the presence of Cd is highly undesirable particularly in industrial applications due to its high toxicity. On the other hand, SnO<sub>2</sub> is the most appealing material as ETL as it has a high mobility and can be process even at room temperature<sup>[33-35]</sup>.



**Figure 8:** Energy levels, bulk electron mobility (BEM) and the preferable temperature processing (PTP) of electron transport material in perovskite solar cells<sup>[35]</sup>.

The survey of PSC performance in terms of Power Conversion Efficiency (PCE) for atomic layer deposition (ALD) grown layers under perovskite absorber has been summarised in the Table 2. There

are many inorganic ETL materials that can fulfil the solar cell structure requirement and provide a high performance in PSC. As also reported previously, this survey shows that SnO<sub>2</sub> processing temperature is relatively lower compared to other materials and the use of it as ETL can enable a very high performance in different device configuration compared to other ETL materials.

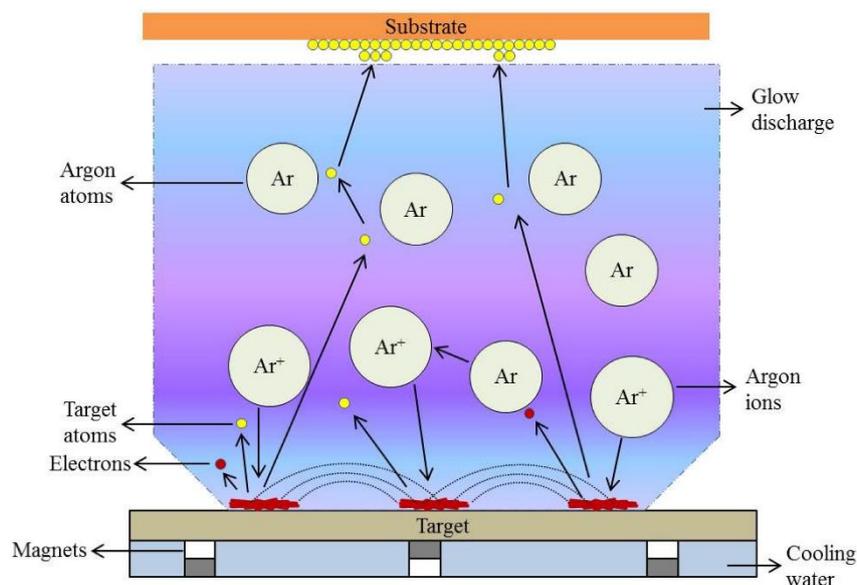
Material	Precursors	Temp. (°C)	Application/ Structure	Device Stack	J <sub>sc</sub> (mA/cm <sup>2</sup> )	V <sub>OC</sub> (V)	FF (%)	η (%)
SnO <sub>2</sub>	TDMA <sub>2</sub> Sn + O <sub>3</sub>	118	ETL/ <i>n-i-p</i>	FTO/SnO <sub>2</sub> (15 nm)/FA <sub>0.85</sub> MA <sub>0.15</sub> Pb(I <sub>0.85</sub> Br <sub>0.15</sub> ) <sub>3</sub> /spiro-OMeTAD/Au	21.3	1.14	74.0	18.4
SnO <sub>2</sub>	TDMA <sub>2</sub> Sn + O <sub>3</sub>	100–120	ETL/ <i>n-i-p</i>	FTO/ <i>d</i> -TiO <sub>2</sub> /SnO <sub>2</sub> /FA <sub>0.85</sub> MA <sub>0.15</sub> Pb(I <sub>0.85</sub> Br <sub>0.15</sub> ) <sub>3</sub> /PTAA/Au	22.7	1.13	78.0	20.0
SnO <sub>2</sub>	TDMA <sub>2</sub> Sn + O <sub>3</sub>	118	ETL/ <i>n-i-p</i>	FTO/SnO <sub>2</sub> (15 nm)/RbI(FA <sub>0.83</sub> MA <sub>0.17</sub> ) <sub>99</sub> Pb(I <sub>0.83</sub> Br <sub>0.17</sub> ) <sub>3</sub> /spiro-OMeTAD/Au	23.0	1.17	71.0	20.0
SnO <sub>2</sub>	TDMA <sub>2</sub> Sn + O <sub>2</sub> Plasma	100	ETL/ <i>n-i-p</i>	FTO/SnO <sub>2</sub> /100 °C/C <sub>60</sub> -SAM/FA <sub>0.30</sub> MA <sub>0.70</sub> PbI <sub>3</sub> /spiro-OMeTAD/Au	22.6	1.13	80.0	20.4
SnO <sub>2</sub>	TDMA <sub>2</sub> Sn + O <sub>2</sub> Plasma	100	ETL/ <i>n-i-p</i>	PET/ITO/SnO <sub>2</sub> /C <sub>60</sub> -SAM/FA <sub>0.30</sub> MA <sub>0.70</sub> PbI <sub>3</sub> /spiro-OMeTAD/Au	22.1	1.10	75.4	18.4
SnO <sub>2</sub>	TDMA <sub>2</sub> Sn + O <sub>2</sub> Plasma	100	ETL/ <i>n-i-p</i>	FTO/SnO <sub>2</sub> /C <sub>60</sub> -SAM/FA <sub>0.30</sub> MA <sub>0.70</sub> PbI <sub>3</sub> /spiro-OMeTAD/Au	21.6	1.13	78.1	19.0
TiO <sub>2</sub>	TiCl <sub>4</sub> + H <sub>2</sub> O	150	ETL/ <i>n-i-p</i>	FTO/TiO <sub>2</sub> (17 nm)/ <i>mp</i> -TiO <sub>2</sub> /MAPbI <sub>3</sub> /Graphene Oxide/spiro-OMeTAD/Au	20.2	1.04	73.0	15.1
TiO <sub>2</sub>	TDMAT + H <sub>2</sub> O	120	ETL/ <i>n-i-p</i>	FTO/TiO <sub>2</sub> (4 nm)/ <i>mp</i> -TiO <sub>2</sub> /MAPbI <sub>3</sub> /spiro-OMeTAD/Au	23.1	1.08	73.4	18.3
TiO <sub>2</sub>	TDMAT + H <sub>2</sub> O	150	ETL/ <i>n-i-p</i>	ITO/TiO <sub>2</sub> (10 nm)/ <i>np</i> -SnO <sub>2</sub> /PC <sub>61</sub> BM/FA <sub>0.30</sub> MA <sub>0.70</sub> Pb(I <sub>1-x</sub> Cl <sub>x</sub> ) <sub>3</sub> /spiro-OMeTAD/Au	23.0	1.08	78.2	19.5
TiO <sub>2</sub>	TDMAT + H <sub>2</sub> O	225	ETL/ <i>n-i-p</i>	FTO/TiO <sub>2</sub> (11 nm)/ <i>mp</i> -TiO <sub>2</sub> /MAPbI <sub>3</sub>	22.3	1.11	74.0	18.4
TiO <sub>2</sub>	TDMAT + H <sub>2</sub> O		ETL/ <i>n-i-p</i>	FTO/ <i>mp</i> -Sb:SnO <sub>2</sub> /TiO <sub>2</sub> (10 nm)/MAPbI <sub>3</sub> /PTAA/Au	23.8	1.10	77.0	20.1
TiO <sub>2</sub>	TDMAT + H <sub>2</sub> O	120	ETL/ Passivation/ <i>n-i-p</i>	FTO/ <i>mp</i> -TiO <sub>2</sub> /TiO <sub>2</sub> (2 nm)/MAPbI <sub>3</sub> /spiro-OMeTAD/Au	17.6	0.97	67.0	11.5
TiO <sub>2</sub>	TTIP + O <sub>2</sub> Plasma	130	ETL/ <i>n-i-p</i>	ITO/CF <sub>4</sub> plasma TiO <sub>2</sub> (20 nm)/MAPbI <sub>3</sub> /spiro-OMeTAD/Au	20.3	1.03	75.5	15.8
TiO <sub>2</sub>	Ti(CpMe)(NMe <sub>2</sub> ) <sub>3</sub> + H <sub>2</sub> O	150	ETL/ Passivation/ <i>n-i-p</i>	ITO/ZnO (80 nm)/TiO <sub>2</sub> (<3 nm)/Cs <sub>0.15</sub> FA <sub>0.75</sub> MA <sub>0.10</sub> PbI <sub>2.9</sub> Br <sub>0.1</sub> /spiro-OMeTAD/MoO <sub>3</sub> /Au	22.5	1.03	74.0	17.1
TiO <sub>2</sub>	TTIP + H <sub>2</sub> O	250	Passivation/ <i>n-i-p</i>	FTO/ <i>c</i> -TiO <sub>2</sub> / <i>NR</i> -TiO <sub>2</sub> /TiO <sub>2</sub> (4 nm)/MAPbI <sub>3</sub> /spiro-OMeTAD/Au	19.8	0.95	72.0	13.5
TiN	TiCl <sub>4</sub> + NH <sub>3</sub>	350	ETL/ Passivation/ <i>n-i-p</i>	FTO/ <i>c</i> -TiO <sub>2</sub> / <i>mp</i> -TiO <sub>2</sub> /TiN (<2 nm)/FA <sub>0.83</sub> MA <sub>0.17</sub> Pb(I <sub>0.83</sub> Br <sub>0.17</sub> ) <sub>3</sub> /PTAA/Au	22.5	1.14	75.0	19.0
TiO <sub>2</sub> -IrO <sub>x</sub>	TDMAT + H <sub>2</sub> O (EtCp)Ir(CHD) + O <sub>3</sub>	175	HTL/ <i>p-i-n</i>	ITO/TiO <sub>2</sub> -IrO <sub>x</sub> (10 nm)/Cs <sub>0.17</sub> FA <sub>0.83</sub> Pb(I <sub>0.83</sub> Br <sub>0.17</sub> ) <sub>3</sub> /C <sub>60</sub> /BCP/Ag	19.6	1.01	80.0	15.8
GaN	TEG + Ar/N <sub>2</sub> /H <sub>2</sub> plasma	280	ETL/ <i>n-i-p</i>	FTO/GaN (5 nm)/FA <sub>0.85</sub> MA <sub>0.15</sub> Pb(I <sub>0.85</sub> Br <sub>0.15</sub> ) <sub>3</sub> /spiro-OMeTAD/Au	22.6	0.98	68.9	15.2
HfO <sub>2</sub>	TEMAHf + H <sub>2</sub> O	90	Passivation/ <i>n-i-p</i>	PEN/ITO/HfO <sub>2</sub> (<1 nm)/SnO <sub>2</sub> /Cs <sub>0.05</sub> (FA <sub>0.83</sub> MA <sub>0.17</sub> ) <sub>0.95</sub> Pb(I <sub>0.83</sub> Br <sub>0.17</sub> ) <sub>3</sub> + RbI + KI/spiro-OMeTAD/Au	21.2	1.14	79.2	19.1
Nb <sub>2</sub> O <sub>5</sub>	( <i>tert</i> -butylimido) bis(diethylamino) niobium + O <sub>3</sub>	170	ETL/ <i>n-i-p</i>	FTO/Nb <sub>2</sub> O <sub>5</sub> (15 nm)/FA <sub>0.85</sub> MA <sub>0.15</sub> Pb(I <sub>0.85</sub> Br <sub>0.15</sub> ) <sub>3</sub> /spiro-OMeTAD/Au				Very low
NiO	Ni(MeCp) <sub>2</sub> + O <sub>2</sub>	350	HTL/ <i>p-i-n</i>	ITO/ <i>s</i> -ALD NiO <sub>x</sub> /FA <sub>0.2</sub> MA <sub>0.8</sub> PbI <sub>3</sub> /PC <sub>61</sub> BM/Al	23.0	1.08	81.0	17.1
NiO	Ni(MeCp) <sub>2</sub> + O <sub>2</sub> plasma	150	HTL/ <i>p-i-n</i>	ITO/NiO (10 nm)/Cs <sub>0.05</sub> (FA <sub>0.83</sub> MA <sub>0.17</sub> ) <sub>0.95</sub> Pb(I <sub>0.83</sub> Br <sub>0.17</sub> ) <sub>3</sub> /C <sub>60</sub> /BCP/Cu	21.8	1.07	73.4	17.1
NiO, AZO, Al <sub>2</sub> O <sub>3</sub>	Ni(dmamb) <sub>2</sub> + O <sub>3</sub> , TMA/DEZ + H <sub>2</sub> O	200, 100, 100	ETL/ <i>p-i-n</i>	FTO/NiO (6 nm)/Cs <sub>0.05</sub> MA <sub>0.95</sub> PbI <sub>3</sub> /PCBM/BCP/AZO (40 nm)/Ag/Al <sub>2</sub> O <sub>3</sub> (50 nm)	22.5	1.03	80.8	18.8
VO <sub>x</sub>	V(dma) <sub>4</sub> + H <sub>2</sub> O	50	HTL/ <i>p-i-n</i>	ITO/VO <sub>x</sub> (1 nm)/MAPbI <sub>3</sub> /PC <sub>61</sub> BM/BCP/Ag	17.9	0.90	71.2	11.5
ZnO/Al <sub>2</sub> O <sub>3</sub>	DEZ + H <sub>2</sub> O	150	ETL/ <i>n-i-p</i>	FTO/ZnO (50 nm)/Al <sub>2</sub> O <sub>3</sub> (<1 nm)/ <i>mp</i> -TiO <sub>2</sub> /MAPbI <sub>3</sub> /spiro-OMeTAD/Au	18.9	1.01	62.0	15.6
ZnS	DEZ + H <sub>2</sub> S	150	Passivation/ <i>n-i-p</i>	FTO/ <i>c</i> -TiO <sub>2</sub> / <i>mp</i> -TiO <sub>2</sub> /ZnS (<2 nm)/FA <sub>0.85</sub> MA <sub>0.15</sub> Pb(I <sub>0.85</sub> Br <sub>0.15</sub> ) <sub>3</sub> /PTAA/Au	22.5	1.13	75.0	18.8

**Table 2:** PSC performance in terms of Power Conversion Efficiency (PCE) for atomic layer deposition (ALD) grown layers under perovskite absorber<sup>[34]</sup>.

### 1.3. Material synthesis and characterization

#### a) Magnetron Sputtering (MS)

Magnetron sputtering technique is one of the most employed Physical Vapour Deposition (PVD) techniques. The vacuum based growth and the high purity material sources assure the high quality of the deposited films. One of the advantages of sputtering is the use of inorganic sources without the need of organic precursors. Thin film sputtering is a mature technology, and it can be used whether at a lab scale or for industrial thin film application. Sputtering has the scalable capability which allows thin films deposition on large areas. In general, the basic principle of sputtering in physics is the bombardment of a solid material surface by high energy plasma or gas which results the ejection of microscopic particles. In magnetron sputtering, the magnetic field, generated by the presence of magnets, is used to restrict the movement of secondary electrons close to surface of the target as shown in Figure 1. The design and strength of the magnetic field can improve the ionisation efficiency and the sputtering deposition rate. Reactive magnetron sputtering refers to the process of magnetron sputtering growth in the presence of reactive gas such as  $O_2$  or  $N_2$ .



**Figure 9:** Basic principle of magnetron sputtering. (Maurya, D.K.; Sardarinejad, A.; Alameh, K. Recent Developments in R.F. Magnetron Sputtered Thin Films for pH Sensing Applications—An Overview. *Coatings* 2014, 4, 756-771 <https://doi.org/10.3390/coatings4040756>)

#### i. Equipment 1 - Torr International

The first equipment, which was used for this thesis, is MagSput™ from Torr International. It has four sputtering 2" diameter target sources where two can be used for radiofrequency (RF) magnetron sputtering and two for direct current (DC) magnetron sputtering. It has a 4" diameter view port with a manual shutter to monitor the growth. This system is equipped with Turbo Molecular vacuum pumping

supported by primary roughing pumping system to assure high vacuum in the deposition chamber while the growth. The magnetron sputtering is not equipped with an entry chamber (load-lock), so the deposition chamber, which is relatively large, is always vented to the atmosphere. The sputtering target to growth  $\text{SnO}_x$  using this equipment is the oxidised SnO (2" diameter) target.



**Figure 10:** Photo of magnetron sputtering made by Torr International.

ii. Equipment 2 - AJA International

The second equipment, which was also used for this thesis, is Orion 5 from AJA International. It has the ability to fit up to (5) 3" or (8) 2" diameter target sources to be used for radiofrequency (RF) magnetron sputtering or for direct current (DC) magnetron sputtering. This system is equipped with Turbo Molecular vacuum pumping supported by primary roughing pumping system to assure high vacuum in the deposition chamber while the growth. The magnetron sputtering is equipped with a high vacuum entry chamber (load-lock), so the deposition chamber, which is relatively smaller compared to the Torr International one, is always kept under high vacuum. The sputtering target to growth  $\text{SnO}_x$  using this equipment is the metallic Sn (2" diameter) target.

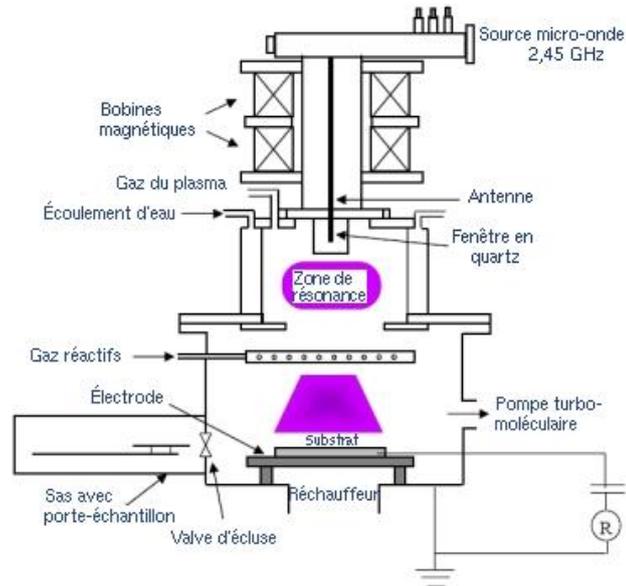


**Figure 11:** Photo of magnetron sputtering made by AJA International.

b) Electron Cyclotron Resonance - Plasma-Enhanced Chemical Vapor Deposition (ECR-PECVD)

Plasma-Enhanced Chemical Vapor Deposition is a thin film deposition technique through gaseous phase. The gas precursors are injected and decomposed inside the growth chamber where the deposition takes place through chemical reaction in presence of a plasma of reacting gases. Electron Cyclotron Resonance (ECR) system, which is also known as a High Density Plasma (HDP) system, is the capability related to microwave source of 2.45 GHz incorporated within the reaction chamber and provides higher plasma ionization and dissociation.

The equipment used for this thesis is made by Roth & Rau and the model is MycroSys 400 PECVD system, which has the capability to deposit thin films from a few nanometres to a few micrometres for areas up to four inches. The temperature range is from room temperature to 500°C and the possible precursor/process gases are SiH<sub>4</sub>, Ar, H<sub>2</sub>, N<sub>2</sub>O, N<sub>2</sub> and NH<sub>3</sub>. In this thesis, this equipment was used only for post deposition treatment where the treatment gases were N<sub>2</sub>O, N<sub>2</sub> and NH<sub>3</sub>.



**Figure 12:** Schematics of Electron Cyclotron Resonance - Plasma-Enhanced Chemical Vapor Deposition

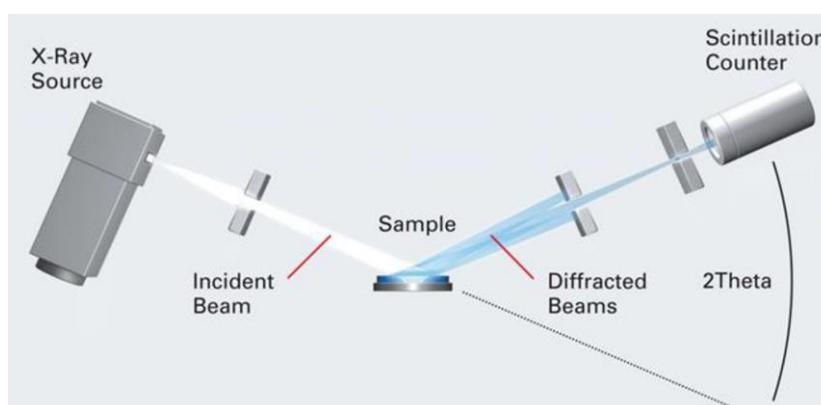


**Figure 13:** Photo of Electron Cyclotron Resonance - Plasma-Enhanced Chemical Vapor Deposition

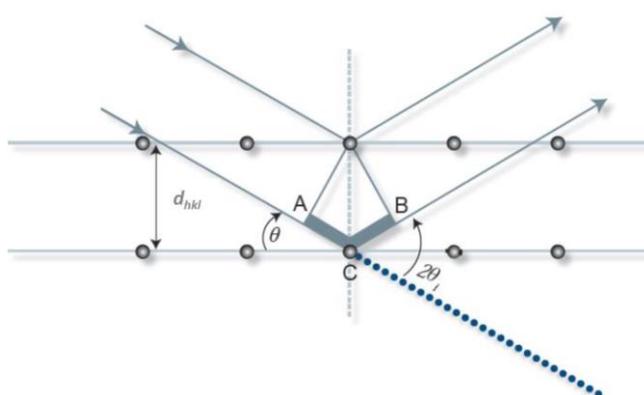
### c) X-ray Diffraction

X-Ray diffraction (XRD) is widely considered as one of the most important and useful materials characterization techniques. It has the advantage of giving detailed information about material crystalline properties particularly the crystalline structure, lattice constant, dislocation density and lattice strain. There are many techniques/configurations based on XRD such as Theta/2Theta or Theta/Theta Bragg Brentano measurements for bulk materials, four circles configuration which is mainly related to single crystals, and grazing incidence XRD (GIXRD) for surface sensitive measurements mainly for thin films where the incident X-ray beam is kept fixed at a very low angle (mainly below  $1^\circ$ )

and the scan is only by varying the detector angle. One of the other advantages of XRD is the non-destructive aspect of this technique. XRD is mainly consisting of X-ray source, sample manipulator, goniometer (which controls the angles) and detector as shown in Figure 14. The physical principle of XRD is based on Bragg's law which relates the x-rays wavelength, the lattice spacing of the crystal and the diffraction angle; so, at Bragg conditions when x-rays scattered through the crystalline material, they interfere out of certain crystalline planes in a certain direction which is the diffraction angle (angle between the incident x-rays and the scattered ones) as shown in Figure 15. In this thesis, only GIXRD configuration was employed.



**Figure 14:** X-ray diffraction schematics (Bruker Corporation).



**Figure 15:** X-ray diffraction principle (Bruker Corporation).

#### d) Surface Analysis

Surface science is a branch of science that focuses on the physical and chemical phenomena that takes place at the surface of materials and the interfaces between two layers. It is involved in many aspects of our lives namely battery electrodes, airplane engine blade, microchips, hetero-junction solar cells, fuel cells, super-hydrophobic, self-cleaning coatings, etc. However, it is still a debate to how define the “Surface” thickness used in surface science. Generally, a bulk material has all dimensions above 1  $\mu\text{m}$ ,

Thin film is a material that has a thickness between 10 and 1000 nm, while materials with a thickness above 1000 nm are generally called films. The ultra-thin film is between 1 and 10 nm and the “surface” is few atomic layers thickness which is up to 1nm. Surface analysis and characterization play a major role in surface science and most of them require sophisticated sources, ultra-high vacuum (UHV) and specialized and sensitive spectrometers, which make them very expensive equipment that are available only in central laboratories. There are many surface analysis techniques, few of them are relatively common, such as X-ray Photoelectron Spectroscopy (XPS) or sometimes called Electron Spectroscopy for Chemical Analysis (ESCA), Time of Flight Secondary Ion Mass Spectroscopy (TOF-SIMS), Auger Electron Spectroscopy (AES). All these techniques are involved in the surface science and provide complementary information. The diagram in Figure 16 (from Evans group) shows different characteristic of these techniques such as the detection range, the analytical spot for many techniques including XPS, SEM, AFM and TOF-SIMS.

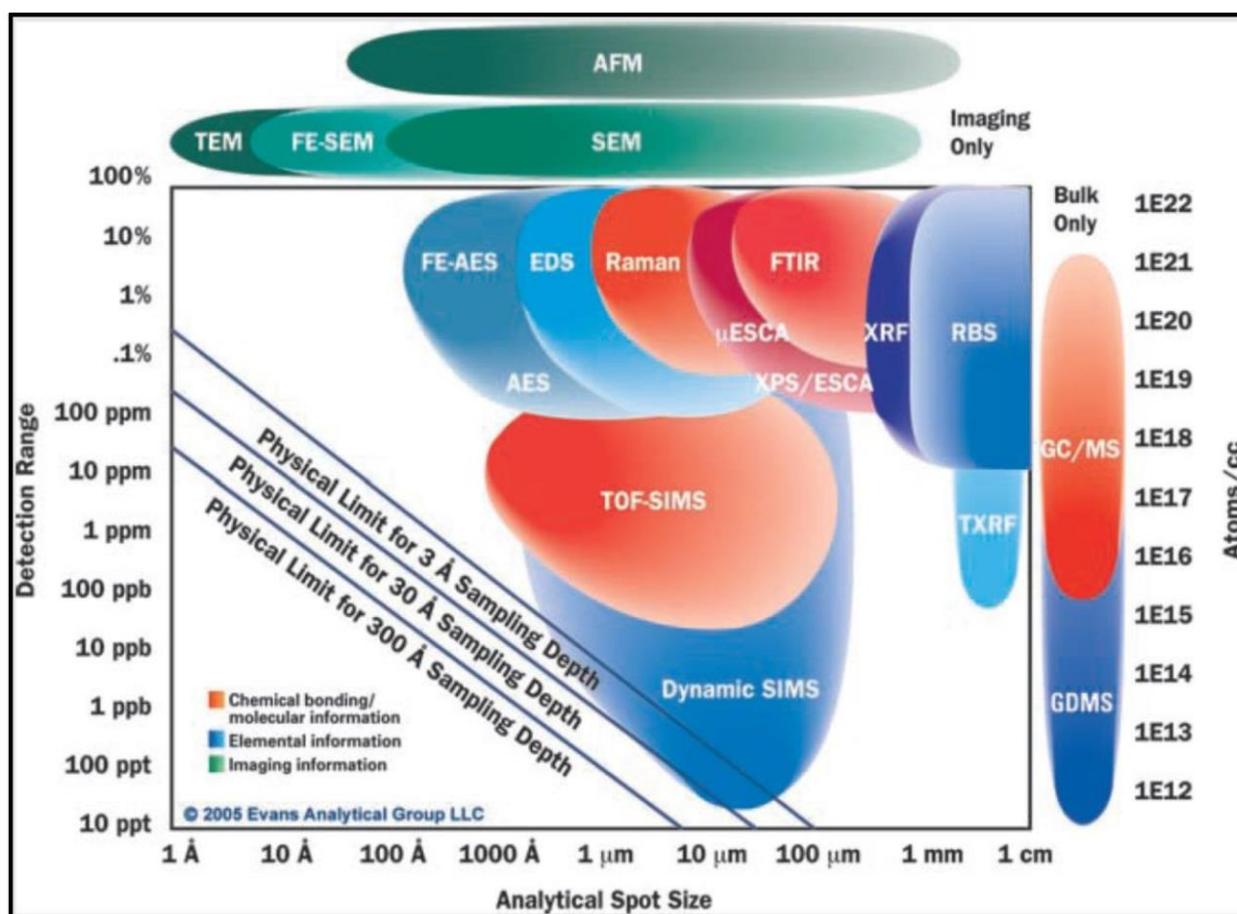
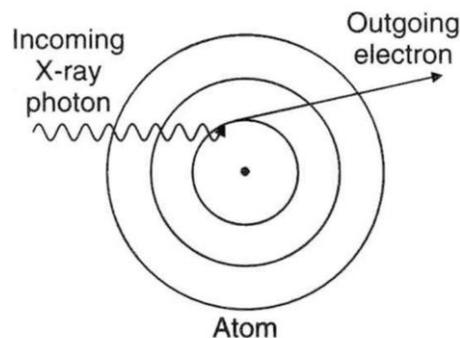


Figure 16: Analytical resolution & detection limit for the most common characterization tools (Evans Analytical Group LLC website).

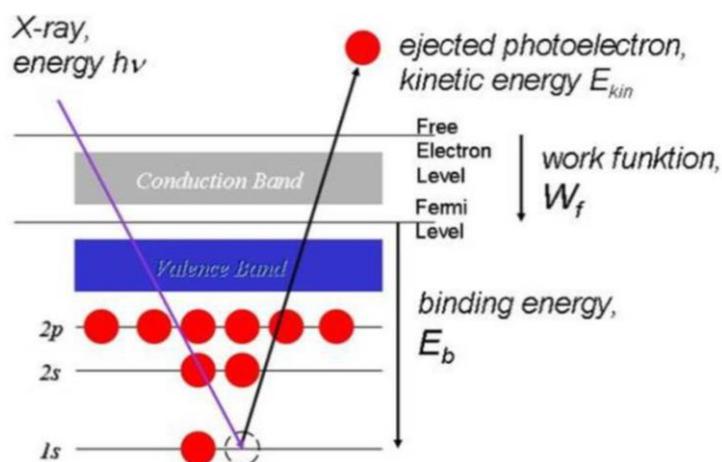
i. X-ray Photoelectron Spectroscopy

XPS is a non-destructive analysis technique and gives information about the existing elements on the surface as well as their chemical state. It also provides the quantitative analysis and other information about the material surface combined with different sources or techniques such as UV Photoelectron Spectroscopy (UPS). XPS is surface sensitive technique which can detect effectively signals only up to approximately 5 nm depending mostly on photoelectron energy. Depth profiling technique using Argon source for surface etching allows performing XPS measurement through depth, this technique is highly important for characterization of hetero junction thin films and interfaces.

As per Evans diagram in figure 16, XPS detection range for a reliable quantitative analysis is between 0.5 to 100% with an analysis area (spot size) from 3  $\mu\text{m}$  up to hundreds of micrometers in diameter using monochromated x-ray source and up to millimeters using a direct source. The basic principle of XPS relies on the Photoelectric effect as shown in figure 17. The incident x-ray interacts with a core level electron; the energy of the x-ray photon is transferred to the electron. Part of this energy will be used to move the electron from its energy level to the Fermi level, in other words to leave the atom, this required energy is known as the binding energy. The second part of the energy is used to carry the electron from the Fermi level to the vacuum level; this energy is required for the electron to leave the material surface which is known as the work function. The rest of the energy will be converted into kinetic energy for the electron as illustrated in figure 18.



**Figure 17:** Schematic of photoelectric effect (University of Hang Kong 2011).



**Figure 18:** Schematic of electrons energy levels in photoelectric process (Leibniz Institute for Solid State and Materials Research Dresden).

ii. Time of Flight – Secondary Ion Mass Spectroscopy (TOF-SIMS)

The TOF-SIMS provides the mass spectroscopy of the existing elements and compounds (atoms and molecules) on the surface, it has high sensitivity suitable for surface traces analysis. It has the capability of detecting the majority of elements and compounds from ppm to ppb range. It has also a very high mass resolution that can easily give information about the isotopes and its mass range is from 0 to 10000 amu (atomic mass unit) including hydrogen. It is extremely surface sensitive because the analysed ions are released from the first two atomic layers. To probe through the material depth, TOFSIMS is usually combined with a sputtering source to enable the depth profiling capability to investigate elements and compounds concentration as a function of depth. This capability is highly required in semiconductor research and development to track the doping and defects as well as secondary phases at the interfaces. Although the TOFSIMS is highly sensitive, the quantification is quite challenging and requires specialized standard samples.

As per Evans diagram, TOF-SIMS analytical spot size (primary beam spot size) varies from 100 nm to 100  $\mu\text{m}$  and the detection limit which is related to the sensitivity is approximately from 100 ppb to 0.1 %. Its basic principle relies on the bombardment of the material surface by a primary ion beam source of high energy up to 50 keV, these ions collide with the surface of material to generate secondary ions as illustrated in figure 19, the primary ions are energetic enough to generate these secondary ions. However, their energy should not exceed certain level to avoid the fragmentation of large molecules.

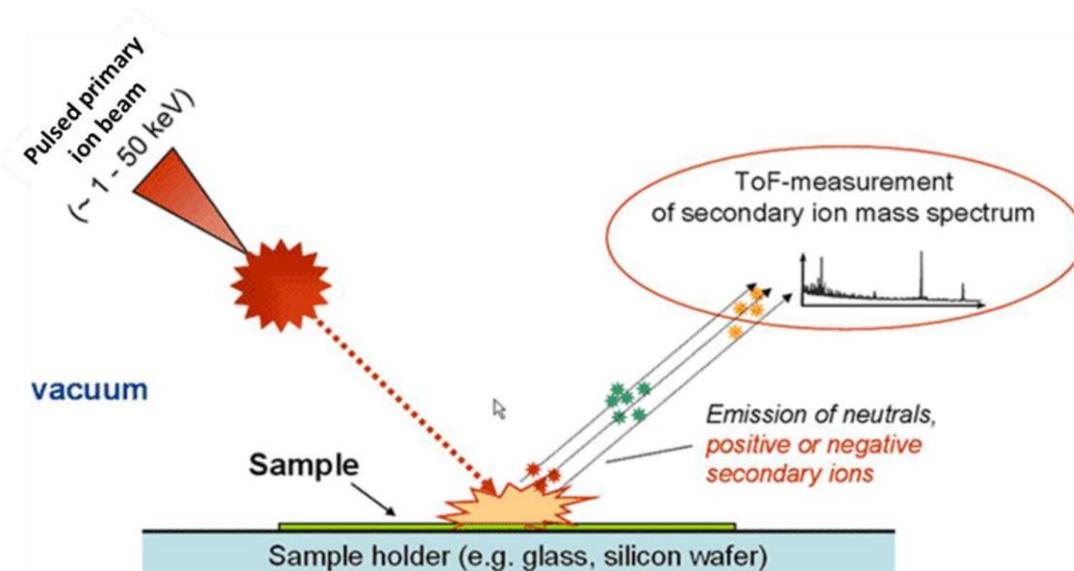


Figure 19: TOF-SIMS basic principle (University of Göttingen).

- e) Surface Imaging
  - i. Scanning Electron Microscopy (SEM)

The SEM technique gives a variety of information about the material surface. It provides mainly imaging and surface topography using a focused electron beam which interacts with the surface and generates a magnified image with resolution in the nanometre range. Other information such as the elemental composition and crystal orientation can be obtained when SEM is combined with certain accessories. The spot size is in the range of a few nanometres (sometimes in sub-nanometre range) which provides a high lateral resolution.

As per duality wave-particle concept, the particles such as electrons can be described as waves and particles. Thus, electrons behave as waves just like photons (light), however, their associated wavelength is extremely small (approximately 100,000 time shorter than light) as shown in Figure 20. This very short wavelength of high energy electrons makes them a very good source of “illumination” to “see” very small features.

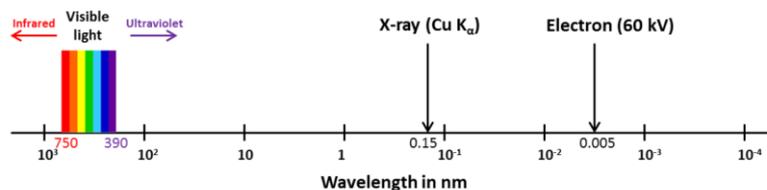
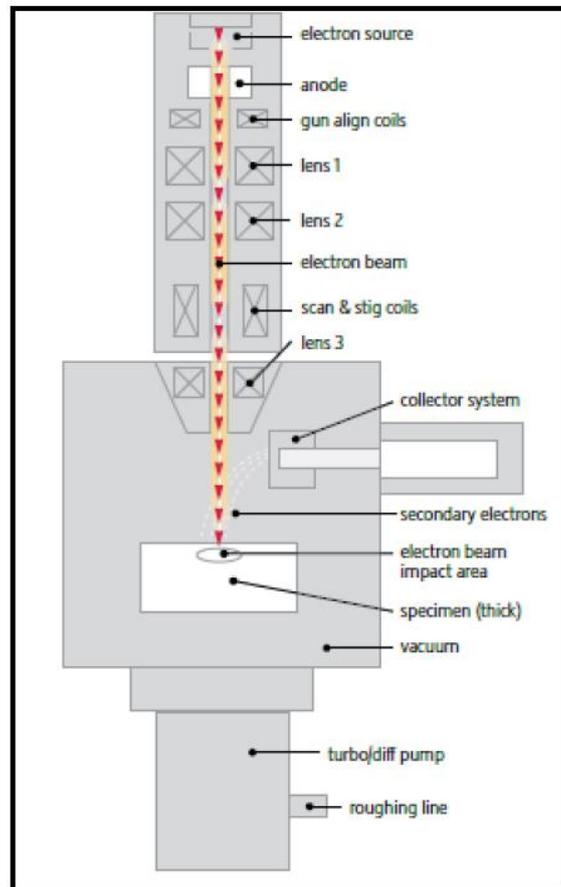


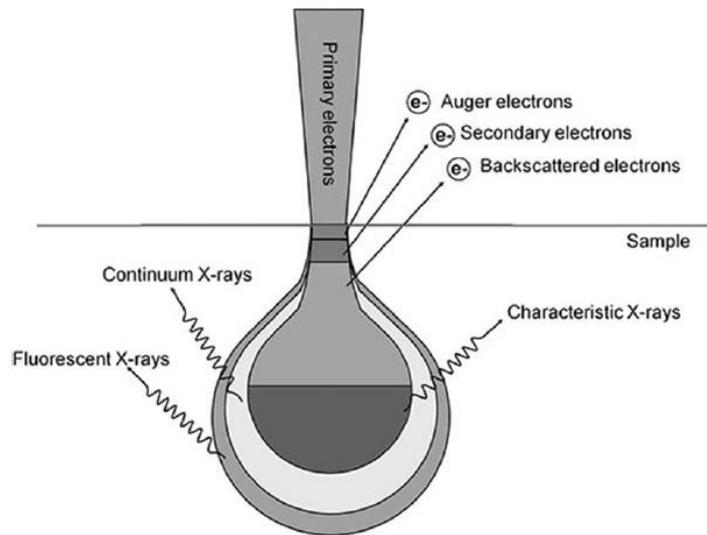
Figure 20: Comparison between the wavelength of a 60 kV electron and visible light.

The electron beam is produced at the electron gun and then through the column it is focused and directed toward the sample surface. The electromagnetic lenses and the electronics control the beam focus and raster to scan the selected area as shown in the schematics in Figure 21.



**Figure 21:** Basic components of Scanning Electron Microscope (FEI part of Thermo Fisher Scientific).

The electrons will interact differently with the surface due mainly to chemical composition and morphology. The electron beam interactions with the surface will produce emissions as shown in Figure 22. Measuring these emissions, namely, secondary electrons provide the imaging capability.



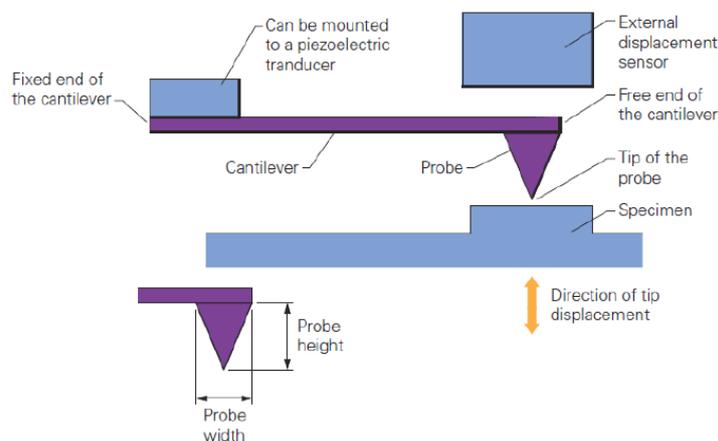
**Figure 22:** Scanning electron microscopy interaction zones and related emissions (Nanotechnologies for the Life Sciences 2007).

ii. Atomic Force Microscopy (AFM)

AFM is one of the most multipurpose imaging microscopes at the nanometer scale, it has three main capabilities which are imaging, measuring, and manipulating.

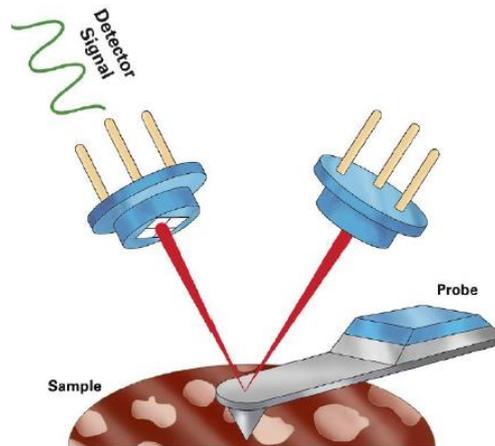
Firstly, imaging is used to provide three-dimensional images to give idea about the surface morphology. Also the scanning probes collect different localized measured information such as electrical properties, mechanical properties within a small area of the sample surface. Finally, AFM has some manipulation applications such as nanolithography using anodic oxidation by a special tip, this technique can provide a small oxide layers for nanofabrication with high precision.

AFM has very high resolution and can provide imaging with atomic lateral resolution and sub-nanometer height resolution. AFM probe consists of an extremely sharp tip (less than 10 nm tip radius) on the edge of the cantilever as shown in Figure 23.



**Figure 23:** AFM cantilever and tip schematic (Bruker Corporation).

AFM basic principle relies on approaching the tip to the sample surface and the interactions with the surface are quantified by sensing the small deflection of the cantilever using a laser source and detector as shown in the Figure 24. In the scanning operation, when the probe is deflected, the laser light spot is deviated from the photo-detector center to a (x,y) position. This change in the signal is monitored and processed to provide a three-dimensional (3D) image.



**Figure 24:** Probe deflection sensing schematic (Bruker Corporation).

There are different modes used in AFM, the main modes are Contact Mode and Tapping Mode. In Contact Mode, the sharp tip is continuously in contact with the surface and the deflection of the laser spot is directly related to change in height. In Tapping Mode, the probe is in continuous mechanical oscillation and its frequency is close or at its resonance frequency. In this case, the detector measures the cantilever oscillation amplitude and with a feedback mode it adjusts the probe height based on the amplitude signal to maintain the same amplitude set-point. The change in height is measured and processed into a 3D image.

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## Extended Abstract in French (Résumé étendu)

Titre de la thèse : **Etude des couches minces d'oxyde d'étain en vue de leurs applications dans les cellules photovoltaïques.**

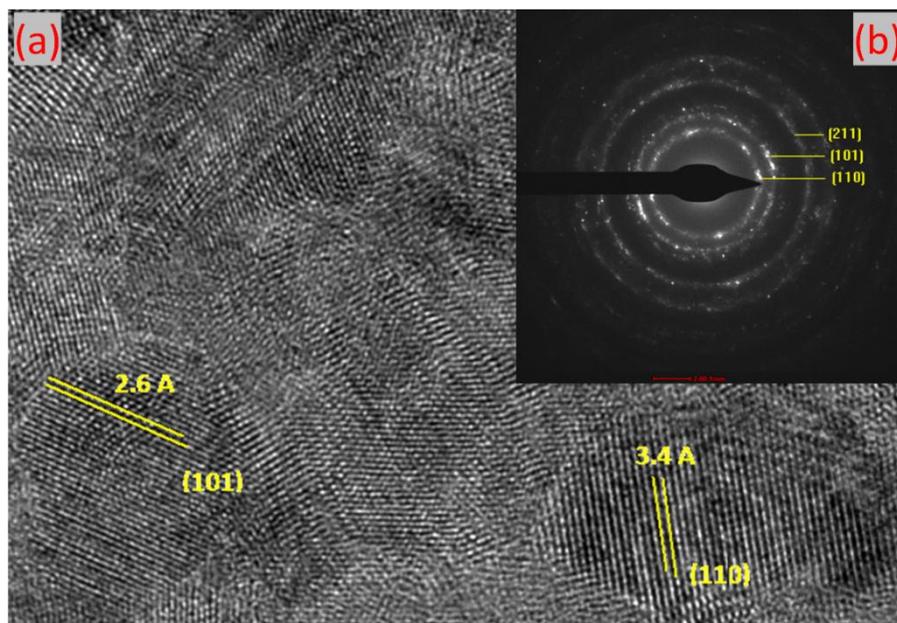
La mise en œuvre de matériaux à base d'éléments abondants et peu coûteux pour des applications dans les cellules solaires photovoltaïques est l'un des aspects clés pour réduire le coût de fabrication et respecter l'environnement. Pour répondre à cet objectif, il y a un effort constant de recherche depuis des décennies pour mettre en œuvre des oxydes conducteurs transparents (OCT) performants très utilisés dans cette technologie. En particulier, l'oxyde d'indium et d'étain, InSnO (ITO), qui est l'un des OCT les plus utilisés dans plusieurs applications liées à l'optoélectronique, est déjà confronté à des problèmes de ressources en raison du coût élevé de l'indium ainsi que de son abondance limitée sur la croûte terrestre par rapport à la plupart des éléments métalliques de transition. D'autre part, l'étain (Sn) est un élément plus accessible et plus abondant que l'indium. Les couches minces d'oxyde d'étain ( $\text{SnO}_x$ ), basées sur l'élément métallique Sn, peuvent atteindre une conductivité et une mobilité électrique élevées en optimisant les conditions de croissance du film  $\text{SnO}_x$ . L'objectif de cette thèse était de mener une étude détaillée de l'oxyde d'étain non dopé et de l'oxyde d'étain enrichi en azote ( $\text{N-SnO}_x$ ) en tant qu'oxyde conducteur transparent (OCT) en utilisant différentes conditions de croissance et d'explorer les propriétés de ces matériaux à l'aide de différentes techniques de caractérisation. Il a été également question d'évaluer les performances de l'oxyde d'étain comme couche de transport d'électrons (CTE) pour les cellules solaires de type pérovskite.

Le **premier chapitre** présente une revue non-exhaustive de littérature sur les oxydes transparents conducteurs (OTC) et leurs applications respectives. En particulier, la synthèse et les propriétés des matériaux, ainsi que les applications des couches minces  $\text{In}_2\text{O}_3$  et  $\text{ZnO}$  en tant que OTC sont abordées dans ce chapitre. Plus particulièrement, l'état de l'art de l'oxyde d'étain ( $\text{SnO}_x$ ), dans les deux états chimiques  $\text{SnO}_2$  et  $\text{SnO}$ , est discuté en termes de synthèse et de propriétés de ces matériaux ainsi que de leurs applications potentielles.

Le **deuxième chapitre** est consacré à la synthèse et à la caractérisation complète de couches minces d'oxyde d'étain élaborées par deux méthodes de pulvérisation magnétron qui diffèrent principalement par le matériau cible lors de la pulvérisation et en faisant varier plusieurs conditions expérimentales, notamment la température du substrat lors du dépôt, le rapport de débit des gaz  $\text{O}_2$  et Ar, et le recuit thermique post-dépôt.

*La première voie de recherche a consisté à préparer les films de  $\text{SnO}_x$  à partir d'une cible d'étain oxydé.* Nous avons montré que la température du substrat joue un rôle majeur dans les propriétés structurales des couches minces d'oxyde d'étain élaborées. Ainsi, les films déposés à  $250^\circ\text{C}$  présentent une

meilleure cristallinité que ceux déposés à 100°C. D'autre part, le post-recuit sous vide à 400°C a amélioré les propriétés structurales en rehaussant la cristallinité. Néanmoins, après le recuit, l'oxyde d'étain développé à 250°C a conservé une meilleure structure cristalline par rapport à celui fabriqué à 100°C. Le recuit sous ambiance air a amélioré la structure cristalline de tous les films par rapport au recuit sous vide. Les deux couches SnO<sub>x</sub> déposées sans O<sub>2</sub> et recuits à l'air ont révélé la présence de deux phases SnO et SnO<sub>2</sub>. Le recuit sous vide a considérablement amélioré les propriétés électriques des échantillons d'oxyde d'étain par rapport au recuit à l'air. Pourtant, le recuit sous air ambiant a permis d'atteindre une transmission optique légèrement meilleure par rapport à la condition de recuit sous vide. Le meilleur film conducteur a été déposé à 250°C sous un rapport de débit O<sub>2</sub>/Ar de 0.015 et suite à un recuit sous vide. Figure 1 montre l'image et la diffraction de microscopie électronique à transmission (MET) de ce film, où la distance interréticulaire et les cercles de diffraction confirment la présence de la phase SnO<sub>2</sub>. Cependant, le film cristallin avec la taille de cristallite la plus élevée a été déposé à 100°C sans flux d'O<sub>2</sub> et avec recuit à l'air. Il est à noter que le meilleur film conducteur a les performances optoélectroniques les plus élevées, évaluées par le facteur de mérite (FdM).



**Figure 1:** Image MET de coupe transversale de lamelle pour le meilleur film conducteur (a) image MET à haute résolution, (b) diagramme de diffraction MET . Ce film SnO<sub>x</sub> a été déposé à 250°C, sous un rapport de débit O<sub>2</sub>:Ar de 0.015 puis recuit sous vide.

La deuxième voie étudiée dans cette thèse est la croissance PVD de SnO<sub>x</sub> utilisant une cible d'étain métallique. Nous avons montré que le rapport de débit de gaz O<sub>2</sub>/Ar joue un rôle majeur dans les propriétés structurales des couches minces d'oxyde d'étain déposées à une température inférieure de 100°C. Pour des rapports de débit O<sub>2</sub>/Ar faibles, les couches de SnO<sub>x</sub> déposées ont une structure amorphe. Par contre, pour les rapports de débit O<sub>2</sub>/Ar plus élevés, les films ont une structure cristalline. Comme attendu, les propriétés structurales ont un effet important sur les propriétés

optoélectroniques des couches minces de  $\text{SnO}_x$  où la bande interdite optique est plus faible pour des rapports de débit  $\text{O}_2/\text{Ar}$  faibles. Cela indique la présence principale de la phase  $\text{SnO}$ . En revanche, pour des rapports de débit  $\text{O}_2/\text{Ar}$  plus élevés, la bande interdite optique augmente, indiquant la présence de  $\text{SnO}_2$ . Nous avons aussi démontré que la résistivité de ces derniers est bien plus faible.

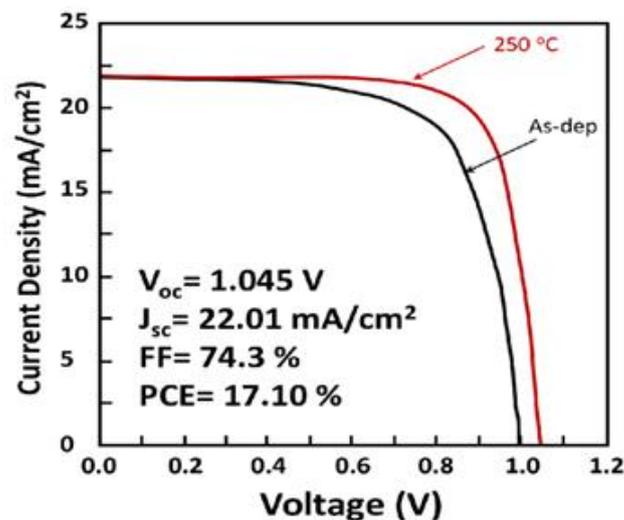
Dans les films  $\text{SnO}_x$  obtenus en utilisant une cible d'étain métallique, nous avons constaté que la température du substrat n'influçait pas de manière significative les propriétés structurales de  $\text{SnO}_x$  produits avec un faible rapport de débit  $\text{O}_2/\text{Ar}$ . Ceci a été confirmé par la diffraction des rayons X - incidente rasante en montrant la présence d'une large bande autour de  $30^\circ$  typique de la phase  $\text{SnO}$  pour des températures du substrat comprises entre la température ambiante et  $400^\circ\text{C}$ . Cependant, nous avons constaté une singularité mineure lié à la présence de  $\text{SnO}_2$  pour le dépôt à  $400^\circ\text{C}$ . La bande interdite optique a légèrement augmenté de 3.5 à 3.7 eV en augmentant la température du substrat à  $400^\circ\text{C}$ , indiquant la présence principalement de la phase  $\text{SnO}$ . Ceci est due à l'appauvrissement de l'atmosphère en oxygène pendant la croissance. On note aussi que la structure amorphe liée à la phase  $\text{SnO}$  a engendrée une résistivité élevée.

Les couches  $\text{SnO}_x$  de type p étant fortement recherchées pour plusieurs applications, certains paramètres expérimentaux ont été variées pour atteindre cet objectif en utilisant des films  $\text{SnO}_x$  obtenus par cible oxydé et soumis à un traitement thermique de recuit. Ainsi, une conductivité de type p a été démontrée dans certains cas spécifiques. La première série d'essais de recuit postérieur au dépôt en variant la température de recuit sous atmosphère d'Argon, a montré que la température de recuit de  $600^\circ\text{C}$  a permis d'obtenir plus de films conducteurs de type p tandis que le recuit à  $500^\circ\text{C}$  a permis de réaliser moins de films conducteurs de type p. Par contre, le recuit à  $400^\circ\text{C}$  n'a pas permis de convertir la conductivité des films  $\text{SnO}_x$  de type n vers type p. Ceci a confirmé l'importance de la température du recuit. Il est également montré que la conductivité de type p est liée à la présence de la phase  $\text{SnO}$ , grâce à la diffraction par microscopie électronique à transmission in-situ et aux mesures de valeurs faibles de la bande interdite. La durée de recuit sous Ar a montré que l'extension du temps de recuit à deux heures a fourni l'effet inverse où tous les films deviennent conducteurs de type n alors qu'ils présentaient une conductivité de type p pour des temps de recuit court.

Il convient de noter que les couches  $\text{SnO}_x$  sont très sensible à l'air et qu'ils s'oxydent facilement en présence d' $\text{O}_2$ . L'effet de vieillissement sous air a également été évalué après sept jours d'exposition, et il s'avère que la conductivité des couches  $\text{SnO}_x$  a évolué du type p au type n et que la résistivité a diminué. Cela pourrait s'expliquer par la continuation de l'oxydation de  $\text{SnO}_x$  par l'air ambiant vers  $\text{SnO}_2$  et en conséquence la formation croissante de porteurs de charge intrinsèques de type n qui compensent les porteurs de charge de type p. L'effet de la présence d'humidité sur l'oxyde d'étain après sept jours de vieillissement a été étudié en utilisant un procédé de chauffage à  $150^\circ\text{C}$  en présence

d'air. Ce processus n'a montré aucune influence de l'humidité et a conduit à une oxydation additionnelle de  $\text{SnO}_x$  et à une diminution de la résistivité en formant davantage de porteurs de charge de type n.

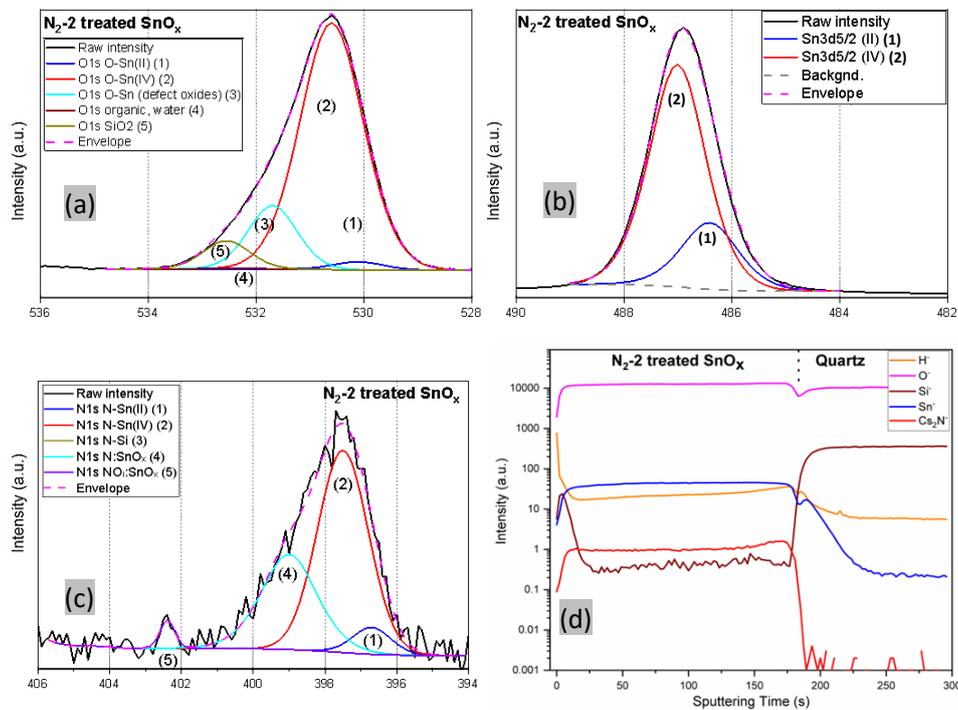
Dans le **troisième chapitre**, les performances du matériau  $\text{SnO}_2$  en tant que semiconducteur de type n, obtenu par pulvérisation, ont été évaluées comme couche de transport d'électrons (CTE) et implémenté dans les cellules solaires pérovskite (CSP) sans structure mésoporeuse. Le rendement de conversion de puissance (RCP) de la cellule CSP, utilisant une couche mince de  $\text{SnO}_2$  (telle que déposée) comme CTE, a révélé un rendement de 15,07 %. En revanche, le RCP du CSP avec une couche mince de  $\text{SnO}_2$  recuite est augmenté de 2 points, à 17,1 %. La Figure 2 montre la caractéristique J-V des deux cellules avec du  $\text{SnO}_2$  tel que déposé et du  $\text{SnO}_2$  recuit en tant que CTE. Le procédé de recuit post-dépôt sous air pour  $\text{SnO}_2$  en tant que CTE a donc permis une meilleure performance des cellules solaires pérovskite en améliorant les propriétés optoélectroniques du  $\text{SnO}_2$ .



**Figure 2:** Caractéristiques densité de courant-tension sous illumination de cellules solaires en pérovskite basées sur les couches minces de  $\text{SnO}_2$  (telle que déposée et recuite à 250°C) utilisés comme couche de transport d'électrons.

Dans le **quatrième et dernier chapitre**, le dopage à l'azote de films  $\text{SnO}_x$  a été étudié en utilisant trois précurseurs d'azote, à savoir des plasmas de  $\text{NH}_3$ ,  $\text{N}_2\text{O}$  et  $\text{N}_2$ , avec deux niveaux de traitements modéré et intense pour chaque précurseur. L'objectif encore une fois est de tenter de former des films  $\text{SnO}_x$  de type p. Les traitements sous plasma  $\text{HN}_3$  ont gravé et endommagé les couches minces de  $\text{SnO}_x$  tandis que les traitements sous plasma  $\text{N}_2\text{O}$  ont augmenté la conductivité de type n. Pour le dopage sous plasma d'azote ( $\text{N}_2$ ), le traitement modéré a abouti à une augmentation importante de la résistivité et le traitement intense a plutôt abouti à une augmentation relativement plus faible de la résistivité mais avec un changement de conductivité du type n vers le type p. Les études par XPS pour le cas de traitement plasma  $\text{N}_2$  intense ont révélé la présence principalement de la phase  $\text{SnO}_2$  avec

une présence mineure de la phase SnO, et surtout la présence de l'azote sous différents états chimiques (voir Figures 3a,b,c). L'incorporation effective d'azote élémentaire dans la couche SnO<sub>x</sub>, a été confirmée par des mesures de ToF-SIMS (Figure 2d) qui ont montré que le profil de distribution de l'azote est uniforme sur toute la profondeur.



**Figure 3:** Analyse de surface pour la couche SnO<sub>x</sub> traités au plasma N<sub>2</sub> intense à l'aide de spectres XPS déconvolués pour (a) O 1s (b) Sn 3d5/2 (c) N 1s et (d) profils de profondeur ToF-SIMS.

Par ailleurs, nous avons également démontré que la réalisation d'un traitement plasma N<sub>2</sub> intense suivi d'un recuit thermique sous Ar conduit à l'amélioration de la conductivité de type p. Cependant, l'inversion de l'ordre de ces deux procédés (recuit sous Ar suivi par traitement plasma N<sub>2</sub> intense) peut conduire à l'effet opposé à savoir que les films conservent leur conductivité de type n ou changent du type p au type n.

Pour résumer les principaux résultats, ce travail de thèse a permis d'ajuster les propriétés structurales, optiques et électriques des couches de SnO<sub>x</sub>, déposées à l'aide de deux configurations différentes de pulvérisation magnétron, en optimisant les conditions de croissance. Pour les couches minces de SnO<sub>x</sub> déposées en utilisant SnO comme cible de pulvérisation, le recuit thermique sous Ar et à une température relativement élevée (autour de 600°C) a montré la possibilité de convertir la conductivité de type n en type p. Les propriétés du matériau SnO<sub>x</sub> avant et après le recuit ont montré que cette conductivité de type p est liée à la formation de la phase de monoxyde d'étain (SnO). D'autre part, nous avons aussi testé l'emploi des couches minces de SnO<sub>2</sub> comme couche de transport d'électrons

dans des cellules solaires à pérovskite. Le recuit thermique de SnO<sub>2</sub> post- dépôt a permis une amélioration nette du rendement de conversion photovoltaïque des cellules testées. L'ensemble de ces résultats montre l'importance des films SnO<sub>x</sub> à utiliser pour la prochaine génération de cellules photovoltaïques (perovskite, tandem perovskite/perovskite, tandem perovskite/Si ou CIGS).

Dans la dernière partie de ce travail, nous avons exploré l'incorporation d'azote, par traitement sous plasma N<sub>2</sub> intense, dans les couches minces de SnO<sub>x</sub>, cette incorporation a indiqué la capacité de convertir la conductivité de type n en type p. Cette conversion est suggérée comme étant le résultat d'éléments azotés occupant les lacunes d'oxygène.

## Publications

1. Y. Zakaria, B. Aïssa, T. Fix, S. Ahzi, A. Samara, S. Mansour, A. Slaoui - Study of wide bandgap SnO<sub>x</sub> thin films grown by a reactive magnetron sputtering via a two-step method. Scientific Reports 12, 15294 (2022). <https://doi.org/10.1038/s41598-022-19270-w>
2. Y. Zakaria, B. Aïssa, T. Fix, S. Ahzi, S. Mansour, A. Slaoui - Moderate temperature deposition of RF magnetron sputtered SnO<sub>2</sub>-based electron transporting layer for triple cation perovskite solar cells – Scientific Reports 13, 9100 (2023). <https://doi.org/10.1038/s41598-023-35651-1>
3. Y. Zakaria, B. Aïssa, T. Fix, S. Roques, G. Ferblantier, S. Ahzi, S. Mansour, A. Slaoui - P-type nitrogen doped SnO<sub>x</sub> synthesis and characterization using magnetron sputtering and plasma post-treatment (In preparation)
4. Y. Zakaria, B. Aïssa, S. Ahzi, A. Slaoui - SnO<sub>x</sub> based thin films as transparent conductive oxides: review (In preparation)

## Communications

1. Y. Zakaria, S. Ahzi, S. Mansour, A. Samara, A. Slaoui - SnO<sub>x</sub> Thin Film RF Sputtering Deposition and Post Annealing (Poster Presentation) - 5th Ed. Smart Materials and Surfaces (2019)
2. Y. Zakaria, A. Slaoui, S. Ahzi, B. Aïssa, T. Fix, G. Ferblantier, A. Samara, S. Mansour- P-type N doped SnO<sub>x</sub> using Magnetron Sputtering and Post Plasma Treatment (Poster Presentation) - European Materials Research Society (Spring 2021)
3. Y. Zakaria, B. Aïssa, T. Fix, S. Ahzi, S. Mansour, A. Slaoui - Synthesis and Characterization of Magnetron Sputtered SnO<sub>2</sub> and its application as Electron Transport Layer (Poster Presentation) - European Materials Research Society (Spring 2023)
4. Y. Zakaria, B. Aïssa, T. Fix, S. Ahzi, S. Mansour, A. Slaoui - Magnetron Sputtered SnO<sub>2</sub> as Electron Transport Layer for Perovskite Solar Cells (Poster Presentation) - European Materials Research Society (Spring 2023)

# Study of Thin Layers of Tin Oxide for Applications in Photovoltaic Solar Cells

## Résumé

L'oxyde d'étain est un oxyde conducteur transparent prometteur qui peut être utilisé dans diverses applications, notamment dans le photovoltaïque. Dans ce travail, nous avons étudié l'effet des conditions de croissance de  $\text{SnO}_x$  en utilisant la pulvérisation magnétron, à savoir le rapport de débit  $\text{O}_2/\text{Ar}$  et la température de croissance, sur les propriétés structurales, optiques et électriques. Nous avons également exploré l'effet du recuit thermique post-dépôt sous vide, air et argon sur les propriétés structurales et optoélectroniques des couches minces de  $\text{SnO}_x$ . De plus, nous avons évalué le film mince  $\text{SnO}_x$  en tant que couche de transport d'électrons pour les cellules solaires à pérovskite et nous avons constaté que l'efficacité de conversion de puissance était passée de 15,07 % à 17,10 % après recuit thermique sous air. Enfin, nous avons étudié la possibilité de doper  $\text{SnO}_x$  avec de l'azote en utilisant trois gaz précurseurs d'azote, à savoir  $\text{NH}_3$ ,  $\text{N}_2\text{O}$  et  $\text{N}_2$ . Nous avons constaté que les traitements  $\text{NH}_3$  ont gravé et endommagé les couches  $\text{SnO}_x$ , les traitements  $\text{N}_2\text{O}$  ont amélioré la conductivité de type n et les traitements  $\text{N}_2$  ont montré la capacité de convertir la conductivité de type n en type p.

**Mots clés :** Oxydes, Pulvérisation Magnétron, Recuit Thermique, Cellule Solaire,  $\text{SnO}_x$

## Résumé en anglais

Tin oxide is a promising transparent conducting oxide which can be employed in various application, particularly in photovoltaic. In this work, we studied the effect of the growth conditions of  $\text{SnO}_x$  using magnetron sputtering namely,  $\text{O}_2/\text{Ar}$  flow rate ratio and growth temperature, on the structural, optical and electrical properties. We also explored the effect of post-deposition thermal annealing under vacuum, air and argon on the structural and optoelectronic properties of  $\text{SnO}_x$  thin films. Furthermore, we assessed the  $\text{SnO}_x$  thin film as electron transport layer for perovskite solar cells and we found that power conversion efficiency has been increased from 15.07% to 17.10% after thermal annealing in air. Finally, we investigated the possibility to dope  $\text{SnO}_x$  with nitrogen using three nitrogen precursor gases namely,  $\text{NH}_3$ ,  $\text{N}_2\text{O}$  and  $\text{N}_2$ . We found that  $\text{NH}_3$  treatments have etched and damaged the  $\text{SnO}_x$  films,  $\text{N}_2\text{O}$  treatments have enhanced the n-type conductivity and  $\text{N}_2$  treatments have shown the ability to convert the conductivity from n-type to p-type.

**Keywords:** Oxides, Magnetron Sputtering, Thermal Annealing, Solar Cell,  $\text{SnO}_x$