



Function spaces on quantum lori

Xiao Xiong

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Xiao Xiong. Function spaces on quantum lori. General Mathematics [math.GM]. Université de Franche-Comté, 2015. English. NNT : 2015BESA2029 . tel-01661517

HAL Id: tel-01661517

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Submitted on 12 Dec 2017

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École Doctorale Carnot-Pasteur

THÈSE DE DOCTORAT

Discipline : Mathématiques

présentée par

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Espaces de fonctions sur les tores quantiques

dirigée par Quanhua XU

Rapporteurs: Éric RICARD et Fedor SUKOCHEV

Soutenue le 2 juillet 2015 devant le jury composé de :

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Remerciements

Cette thèse n'aurait pas pu voir le jour sans la contribution de plusieurs personnes auxquelles j'adresse mes plus sincères remerciements.

Je tiens tout d'abord à remercier mon directeur de thèse, Monsieur Quanhua XU, pour ses conseils précieux et aide indispensable qu'il m'a apportés pendant la réalisation de la thèse. Je le remercie pour le temps qu'il m'a consacré, la multitude des sujets qu'il a proposés ainsi que sa profonde connaissance dont j'ai pu profiter au cours de ces quatre années. Il m'a également encouragé à participer à de nombreuses écoles et conférences, grâce auxquelles j'ai pu enrichir ma culture mathématique, présenter mes résultats et rencontrer d'autres mathématiciens à travers le monde. Je le remercie enfin pour m'avoir donné l'opportunité d'apprendre le français et pour son grand soin sur ma vie en France.

Je souhaite également remercier Messieurs Éric RICARD et Fedor SUKOCHEV qui m'ont fait l'honneur d'être les rapporteurs de ma thèse. Je les remercie pour leur lecture attentive du manuscrit et l'intérêt qu'ils portent à mes recherches. Je voudrais également les remercier aussi pour leurs intuitions et pour avoir répondu à mes questions lors d'innombrables discussions mathématiques.

Je témoigne ma gratitude envers Monsieur Gilles PISIER qui a accepté d'être président du jury. Je tiens aussi à remercier vivement Messieurs Christian LE MERDY et Abdellah YOUSSEFI qui ont accepté de faire partie du jury.

Je voudrais remercier Marius JUNGE pour les discussions stimulantes que nous avons eues. J'exprime de même ma reconnaissance à Tao MEI, pour une observation critique de ma thèse. Mes remerciements se dirigent également vers Zhi YIN, qui est le co-auteur de mon mémoire de recherche. Sans sa contribution, cette thèse n'aurait pas pu arriver à sa forme actuelle. Par ailleurs, l'accueil chaleureux de Guixiang HONG pendant les premiers mois à Besançon m'a beaucoup aidé. Je remercie aussi Yanqi QIU et Mathilde PERRIN, pour les discussions mathématiques stimulantes que nous avons échangées. Je remercie sincèrement Roland SPEICHER pour son accueil à l'Université des Saarlandes (Allemagne).

Cette thèse a été effectuée au sein de l'équipe d'analyse fonctionnelle de Besançon, dont la dynamique et la bonne ambiance ont contribué au bon déroulement de ces quatre années. Je remercie en particulier mes collègues, Mikael DE LA SALLE, Uwe FRANZ, Yulia KUZNETSOVA, Gilles LANCIEN, Christian LE MERDY, Stefan NEUWIRTH et Alexandre NOU pour leurs excellents exposés au groupe de travail dont j'ai beaucoup profité. J'adresse toute ma gratitude aux membres du laboratoire de Mathématiques de Besançon, qui contribuent à en faire un lieu de travail agréable et convivial. Merci en particulier à Emilie DUPRE, Richard FERRERE, Odile HENRY, Romain PACE, Catherine PAGANI, Pascaline SAIRE et Catherine VUILLEMENOT pour leur aide chaleureuse.

Mes remerciements se dirigent également vers Monsieur Hua CHEN, qui était mon

directeur en master et qui m'a permis de decouvrir le monde de l'analyse. C'est lui qui m'a propose de faire une these sous la direction de Quanhua XU. Une partie de cette these a ete effectuee pendant mes sejours a l'Universite de Wuhan (Chine) ou j'ai beneficie des groupes de travail. Je remercie en particulier Zeqian CHEN, Guixiang HONG, Yong JIAO, Tao MA, Tao MEI, Mu SUN, Maofa WANG, Simeng WANG, Xumin WANG et Quanhua XU pour leurs excellents exposes.

Je souhaite remercier la joyeuse bande des doctorants, sans qui la vie au laboratoire n'aurait pas la meme saveur. Merci a Aude, Charlotte, Michaël, Runlian, Simeng, Souleiman, Yahya... Merci aux membres du bureau 401 et ses affiliés, Alexis, Céline, Clément, Cyrille, Emilie, Guillaume, Ibrahim, Karine, Marine, Michel et Pammella pour la bonne ambiance, l'aide sur la vie et le français. En particulier, je remercie sincerement Clément pour l'aide sur la partie française de ma these. Je souhaite remercier aussi tous ceux qui ont contribue aux moments de détente ces dernieres années, lors d'une soirée, une fondue ou un sport. Merci a Bin, Gang, Tianwen, Tingjian, Yazhou, Yinsheng, Yufei...

J'exprime ma profonde reconnaissance et toute ma gratitude a ma famille, mes parents, grand-parents qui m'ont toujours encourage dans mes projets. Le parcours aurait ete beaucoup plus difficile si je n'avais pas ete si bien entoure et epaulé.

Mes derniers remerciements vont bien sur a ma femme Jianqiao, qui m'a apporte le bonheur et a enjolivé ma vie pendant la preparation de cette these.

Résumé

Cette thèse donne une étude systématique des espaces de Sobolev, Besov et Triebel-Lizorkin sur le tore quantique $\mathbb{T}_\theta^d$ pour une matrice $d \times d$ anti-symétrique réelle θ . Ces espaces partagent beaucoup de propriétés avec leurs analogues classiques. Nous prouvons le théorème de réduction pour tous ces espaces et une inégalité de Poincaré pour les espaces de Sobolev. Nous montrons aussi que l'espace de Sobolev $W_\infty^k(\mathbb{T}_\theta^d)$ coïncide avec l'espace de Lipschitz d'ordre k étudié par Weaver. Nous démontrons les inégalités de plongement pour eux, incluant le plongement d'espaces de Besov et d'espaces de Sobolev. Nous obtenons une caractérisation générale à la Littlewood-Paley pour les espaces de Besov et Triebel-Lizorkin, qui implique des caractérisations concrètes par les semigroupes de Poisson et de chaleur ainsi par des différences. Certains d'entre elles sont nouvelles, même dans le cas commutatif; par exemple, celle d'espaces de Besov et Triebel-Lizorkin par le semigroupe de Poisson améliore le résultat classique. En conséquence de la caractérisation d'espaces de Besov par des différences, nous étendons les récents résultats de Bourgain-Brézis-Mironescu et Maz'ya-Shaposhnikova sur les limites de normes de Besov au cadre quantique. La même caractérisation implique que l'espace de Besov $B_{\infty,\infty}^\alpha(\mathbb{T}_\theta^d)$ avec $\alpha > 0$ est l'analogue quantique de l'espace de Zygmund usuel d'ordre α . Nous étudions aussi l'interpolation de ces espaces, et en particulier, déterminons explicitement le K -fonctionnel du couple $(L_p(\mathbb{T}_\theta^d), W_p^k(\mathbb{T}_\theta^d))$, ce qui est l'analogue quantique du résultat classique de Johnen et Scherer. Enfin, nous montrons que les multiplicateurs de Fourier complètement bornés sur tous ces espaces coïncident avec ceux sur les espaces correspondants sur le tore usuel. Nous prouvons également que les multiplicateurs de Fourier sur les espaces de Besov sont complètement déterminés par ceux sur les sous-espaces L_p associés à leurs composantes dans la décomposition de Littlewood-Paley.

Mots-clefs

Tore quantique, espaces L_p non commutatifs, potentiels de Bessel et de Riesz, espaces de Sobolev (potentiels), espaces de Besov, espaces de Triebel-Lizorkin, espaces de Hardy, caractérisation, semigroupes de Poisson et de chaleur, inégalités de plongement, interpolation, multiplicateurs de Fourier (complètement bornés).

Abstract

This thesis gives a systematic study of Sobolev, Besov and Triebel-Lizorkin spaces on a noncommutative d -torus $\mathbb{T}_\theta^d$ for a skew symmetric real $d \times d$ -matrix θ . These spaces share many properties with their classical counterparts. We prove, among other basic properties, the lifting theorem for all these spaces and a Poincaré type inequality for Sobolev spaces. We also show that the Sobolev space $W_\infty^k(\mathbb{T}_\theta^d)$ coincides with the Lipschitz space of order k , already studied by Weaver. We establish the embedding inequalities of all these spaces, including the Besov and Sobolev embedding theorems. We obtain Littlewood-Paley type characterizations for Besov and Triebel-Lizorkin spaces in a general way, as well as the concrete ones in terms of the Poisson, heat semigroups and differences. Some of them are new even in the commutative case, for instance, our Poisson semigroup characterization of Besov and Triebel-Lizorkin spaces improves the classical ones. As a consequence of the characterization of the Besov spaces by differences, we extend to the quantum setting the recent results of Bourgain-Brézis-Mironescu and Maz'ya-Shaposhnikova on the limits of Besov norms. The same characterization implies that the Besov space $B_{\infty,\infty}^\alpha(\mathbb{T}_\theta^d)$ with $\alpha > 0$ is the quantum analogue of the usual Zygmund class of order α . We investigate the interpolation of all these spaces, in particular, determine explicitly the K-functional of the couple $(L_p(\mathbb{T}_\theta^d), W_p^k(\mathbb{T}_\theta^d))$, which is the quantum analogue of a classical result due to John and Scherer. Finally, we show that the completely bounded Fourier multipliers on all these spaces do not depend on the matrix θ , so coincide with those on the corresponding spaces on the usual d -torus. We also give a quite simple description of (completely) bounded Fourier multipliers on the Besov spaces in terms of their behavior on the L_p -components in the Littlewood-Paley decomposition.

Keywords

Quantum tori, noncommutative L_p -spaces, Bessel and Riesz potentials, (potential) Sobolev spaces, Besov spaces, Triebel-Lizorkin spaces, Hardy spaces, characterizations, Poisson and heat semigroups, embedding inequalities, interpolation, (completely bounded) Fourier multipliers.

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Introduction

Cette thèse étudie l'analyse sur le tore quantique. Le tore quantique est un exemple fondamental dans la théorie d'algèbre d'opérateur. Elle est probablement la classe la plus intéressante et la plus accessible aux objets d'étude dans la théorie de géométrie non commutative ([18]), puisqu'elle possède aussi une structure différentielle. Il existe de nombreux travaux étendus dans ce cadre (voir, par exemple, la série d'articles par Rieffel [63]). Cependant, très peu de travaux avaient été faits sur son aspect analytique. Cette lacune est probablement due aux nombreuses difficultés que l'on peut rencontrer en travaillant avec les espaces L_p non commutatifs, et qui apparaissent inévitablement si l'on veut faire de l'analyse.

L'article [17] est le premier travail systématique sur l'analyse harmonique sur le tore quantique. Il étudie certains aspects importants de l'analyse harmonique dans ce cadre, comprenant les inégalités maximales, les convergences en moyenne et ponctuelles correspondantes des séries de Fourier quantiques, les multiplicateurs de Fourier complètement bornés sur les espaces L_p et la théorie des espaces de Hardy et BMO . Il s'inspirait directement de récents progrès de l'analyse harmonique non commutative. [17] devint possible par le développement des inégalités de martingales et ergodiques non commutatives et la théorie de Littlewood-Paley-Stein pour les semi-groupes Markoviens quantiques, qui ont été réalisés grâce aux efforts de nombreux chercheurs; nous renvoyons le lecteur à [56, 31, 36, 37, 61, 62, 52], et [32, 44, 45, 33, 34].

Cette thèse vise à étudier les espaces de fonctions sur le tore quantique. Elle est donc la continuation naturelle de [17]. Les espaces à étudier sont les espaces de Sobolev, Besov et Triebel-Lizorkin. Dans le cadre classique, ces espaces sont fondamentaux dans beaucoup de branches des mathématiques, telles que l'analyse harmonique, les EDP, l'analyse fonctionnelle et la théorie d'approximation. Nos références pour la théorie classique sont [1, 42, 49, 53, 73, 74]. À notre connaissance, ils n'avaient jusqu'à présent jamais été étudiés dans le cas quantique, sauf dans deux cas particuliers. Les espaces de Sobolev avec la norme L_2 ont été étudiés par Spera [65], en vue d'applications à la théorie de Yang-Mills pour le tore quantique [66] (ou [26, 39, 58, 64] pour les travaux associés). D'autre part, inspirés par la géométrie non commutative de Connes [18], ou plus précisément, la partie sur les espaces métriques non commutatifs, Weaver [78, 79] a développé les classes de Lipschitz d'ordre α pour $0 < \alpha \leq 1$ sur le tore quantique. Le fait que seuls ces deux cas aient été étudiés montre une fois de plus les difficultés associées à la non commutativité qui ont été évoquées précédemment.

Une de ces difficultés est à souligner : elle est pertinente pour cette thèse, et est le manque d'un analogue non commutatif de la fonction maximale ponctuelle usuelle. Cependant, la technique de fonction maximale joue un rôle primordial dans la théorie classique des espaces de Besov et Triebel-Lizorkin (ainsi que dans la théorie des espaces de Hardy). Elle n'est plus disponible dans le cadre quantique. Nous nous efforçons d'inventer de nouveaux outils, comme dans les travaux cités précédemment sur les inégalités de

martingales non commutatives et la théorie quantique de Littlewood-Paley-Stein, où la même difficulté apparaissait déjà.

Un puissant outil utilisé dans [17] est la méthode de transference. Elle transfère les problèmes sur le tore quantique aux problèmes correspondants dans le cas des fonctions à valeurs opérateurs sur le tore usuel, de manière à utiliser des résultats existants dans ce dernier cas ou à adapter les arguments classiques. Cette méthode est efficace pour plusieurs problèmes étudiés dans [17], incluant les inégalités maximales et les espaces de Hardy. Elle est encore utile dans certaines parties de cette thèse, par exemple, pour les espaces de Besov qui peuvent être étudiés à travers les espaces de Besov à valeurs vectorielles classiques par des moyens de transference, les espaces de Banach concernés étant des espaces L_p non commutatifs sur le tore quantique. Mais elle est inefficace pour les autres. Par exemple, les inégalités de plongement de Sobolev ne peuvent pas être prouvées par transference. D'autre part, si nous voulons étudier les espaces de Triebel-Lizorkin sur le tore quantique via transference, nous devons d'abord développer la théorie des espaces de Triebel-Lizorkin à valeurs opérateurs sur le tore classique. Cette dernière théorie est aussi difficile que la première. Contrairement à [17], la méthode de transference va jouer un rôle très restreint dans cette thèse. Cependant, nous allons utiliser les multiplicateurs de Fourier de manière cruciale. Nous développons une analyse différentielle intrinsèque sur le tore quantique, sans se référer fréquemment au tore usuel via transference. C'est un grand avantage de la méthode présente par rapport à celle de [17]. Nous espérons que l'étude menée ici ouvre de nouvelles perspectives d'applications, et motive d'autres travaux de recherche sur le tore quantique ou dans des circonstances similaires.

Cette thèse est constituée de sept chapitres, rédigés en anglais. Elle se base sur un travail en commun avec Quanhua Xu et Zhi Yin. Dans cette introduction, nous rappellerons dans un premier temps les connaissances fondamentales et les notations (voir les sections respectives suivantes pour plus de détails), et donnerons les définitions des espaces à étudier. Ensuite nous décrirons les principaux résultats prouvés dans la thèse en les classifiant en cinq familles.

0.1 Notations et définitions

Soient $d \geq 2$ et $\theta = (\theta_{kj})$ une matrice $d \times d$ anti-symétrique réelle. Rappelons que le tore non commutatif $\mathcal{A}_\theta$ de d générateurs est la C^* -algèbre universelle, engendrée par d opérateurs unitaires $U_1, \dots, U_d$ vérifiant la relation de commutation suivante

$$U_k U_j = e^{2\pi i \theta_{kj}} U_j U_k, \quad j, k = 1, \dots, d.$$

Soit $U = (U_1, \dots, U_d)$. Pour $m = (m_1, \dots, m_d) \in \mathbb{Z}^d$, notons

$$U^m = U_1^{m_1} \dots U_d^{m_d}.$$

Un polynôme en U est une somme finie :

$$x = \sum_{m \in \mathbb{Z}^d} \alpha_m U^m, \quad \alpha_m \in \mathbb{C}.$$

Pour un tel polynôme x , nous définissons $\tau(x) = \alpha_0$. Alors, τ s'étend en un état tracial fidèle sur $\mathcal{A}_\theta$. Soit $\mathbb{T}_\theta^d$ l'algèbre de von Neumann obtenue par la représentation GNS de τ . On dit que $\mathbb{T}_\theta^d$ est le d-tore quantique associé à θ . Remarquons que si $\theta = 0$, alors $\mathcal{A}_\theta = C(\mathbb{T}^d)$ et $\mathbb{T}_\theta^d = L_\infty(\mathbb{T}^d)$, où $\mathbb{T}^d$ est le d-tore habituel. En conséquence, un d-tore

quantique est une déformation du d-tore habituel muni de la mesure de Haar normalisée. Il est donc naturel d'espérer que $\mathbb{T}_\theta^d$ partage beaucoup de propriétés avec $\mathbb{T}^d$. Les espaces L_p non commutatifs associés à $\mathbb{T}_\theta^d$ sont notés $L_p(\mathbb{T}_\theta^d)$. La transformation de Fourier de $x \in L_1(\mathbb{T}_\theta^d)$ est définie par

$$\widehat{x}(m) = \tau((U^m)^* x), \quad m \in \mathbb{Z}^d.$$

L'opérateur x est, bien sûr, uniquement déterminé par sa série de Fourier:

$$x \sim \sum_{m \in \mathbb{Z}^d} \widehat{x}(m) U^m.$$

La structure différentielle de $\mathbb{T}_\theta^d$ est modélisée sur celle de $\mathbb{T}^d$. Soit

$$\mathcal{S}(\mathbb{T}_\theta^d) = \left\{ \sum_{m \in \mathbb{Z}^d} a_m U^m : \{a_m\}_{m \in \mathbb{Z}^d} \text{ rapidement décroissant} \right\}.$$

C'est la déformation de l'espace des fonctions infiniment différentiables sur $\mathbb{T}^d$. Comme dans le cas commutatif, $\mathcal{S}(\mathbb{T}_\theta^d)$ a une topologie localement convexe naturelle. Son espace dual topologique $\mathcal{S}'(\mathbb{T}_\theta^d)$ est l'espace de distributions sur $\mathbb{T}_\theta^d$. Les dérivées partielles sur $\mathcal{S}(\mathbb{T}_\theta^d)$ sont déterminées par

$$\partial_j(U_j) = 2\pi i U_j \text{ et } \partial_j(U_k) = 0, \quad k \neq j, 1 \leq j, k \leq d.$$

Soit $\mathbb{N}_0$ l'ensemble des entiers positifs. La dérivée partielle D^m associée à $m = (m_1, \dots, m_d) \in \mathbb{N}_0^d$ est définie par $\partial_1^{m_1} \dots \partial_d^{m_d}$. L'ordre de D^m est $|m|_1 = m_1 + \dots + m_d$. $\Delta = \partial_1^2 + \dots + \partial_d^2$ est le Laplacien sur $\mathbb{T}_\theta^d$. Par dualité, les dérivées et la transformation de Fourier transfèrent aussi sur $\mathcal{S}'(\mathbb{T}_\theta^d)$.

Pour une fonction $\phi : \mathbb{Z}^d \rightarrow \mathbb{C}$, notons M_ϕ le multiplicateur de Fourier associé sur $\mathbb{T}^d$:

$$\widehat{M_\phi f}(m) = \phi(m) \widehat{f}(m)$$

pour tous les polynômes trigonométriques f sur $\mathbb{T}^d$. On dit que ϕ est un multiplicateur sur $L_p(\mathbb{T}^d)$ si M_ϕ s'étend en une application bornée sur $L_p(\mathbb{T}^d)$. Les multiplicateurs de Fourier sur $\mathbb{T}_\theta^d$ sont définis de la même manière. Nous utilisons le même symbole M_ϕ pour le multiplicateur correspondant sur $\mathbb{T}_\theta^d$. Pour une fonction ϕ sur $\mathbb{R}^d$, nous notons ϕ plutôt que $\phi|_{\mathbb{Z}^d}$ un multiplicateur de Fourier sur $L_p(\mathbb{T}^d)$ ou $L_p(\mathbb{T}_\theta^d)$.

Deux multiplicateurs spéciaux vont apparaître fréquemment : le potentiel de Bessel $J^\alpha = (1 - (2\pi)^{-2}\Delta)^{\frac{\alpha}{2}}$ et celui de Riesz $I^\alpha = (-(2\pi)^{-2}\Delta)^{\frac{\alpha}{2}}$. Ils sont définis par

$$J^\alpha(x) = \sum_{m \in \mathbb{Z}^d} (1 + |m|^2)^{\frac{\alpha}{2}} \widehat{x}(m) U^m \text{ et } I^\alpha(x) = \sum_{m \in \mathbb{Z}^d} |m|^\alpha \widehat{x}(m) U^m,$$

pour toute distribution x sur $\mathbb{T}_\theta^d$. Ici le potentiel de Riesz est défini pour les distributions x avec $\widehat{x}(0) = 0$. Notons aussi $J_\alpha(\xi) = (1 + |\xi|^2)^{\frac{\alpha}{2}}$ sur $\mathbb{R}^d$ et $I_\alpha(\xi) = |\xi|^\alpha$ sur $\mathbb{R}^d \setminus \{0\}$. Alors $J_\alpha(\xi)$ et $I_\alpha(\xi)$ sont les symboles de J^α et I^α respectivement.

Fixons une fonction de Schwartz φ sur $\mathbb{R}^d$ vérifiant la propriété de décomposition usuelle de Littlewood-Paley :

$$\begin{cases} \text{supp } \varphi \subset \{\xi : 2^{-1} \leq |\xi| \leq 2\}, \\ \varphi > 0 \text{ sur } \{\xi : 2^{-1} < |\xi| < 2\}, \\ \sum_{k \in \mathbb{Z}} \varphi(2^{-k}\xi) = 1, \quad \xi \neq 0. \end{cases}$$

Pour tout $k \geq 0$, notons φ_k la fonction dont la transformation de Fourier est égale à $\varphi(2^{-k}\cdot)$. Pour une distribution x sur $\mathbb{T}_\theta^d$, définissons

$$\tilde{\varphi}_k * x = \sum_{m \in \mathbb{Z}^d} \varphi(2^{-k}m) \hat{x}(m) U^m.$$

Alors $x \mapsto \tilde{\varphi}_k * x$ est le multiplicateur de Fourier avec le symbole $\varphi(2^{-k}\cdot)$.

Nous pouvons maintenant définir les quatre familles d'espaces de fonction sur $\mathbb{T}_\theta^d$ à étudier. Soient $1 \leq p, q \leq \infty$ et $k \in \mathbb{N}, \alpha \in \mathbb{R}$.

- *Les espaces de Sobolev :*

$$W_p^k(\mathbb{T}_\theta^d) = \{x \in \mathcal{S}'(\mathbb{T}_\theta^d) : D^m x \in L_p(\mathbb{T}_\theta^d) \text{ pour tout } m \in \mathbb{N}_0^d \text{ avec } |m|_1 \leq k\}.$$

- *Les espaces de Sobolev Potentiels :*

$$H_p^\alpha(\mathbb{T}_\theta^d) = \{x \in \mathcal{S}'(\mathbb{T}_\theta^d) : J^\alpha x \in L_p(\mathbb{T}_\theta^d)\}.$$

- *Les espaces de Besov :*

$$B_{p,q}^\alpha(\mathbb{T}_\theta^d) = \{x \in \mathcal{S}'(\mathbb{T}_\theta^d) : (|\hat{x}(0)|^q + \sum_{k \geq 0} 2^{qk\alpha} \|\tilde{\varphi}_k * x\|_p^q)^{\frac{1}{q}} < \infty\}.$$

- *Les espaces de Triebel-Lizorkin pour $p < \infty$:*

$$F_p^{\alpha,c}(\mathbb{T}_\theta^d) = \{x \in \mathcal{S}'(\mathbb{T}_\theta^d) : \|(|\hat{x}(0)|^2 + \sum_{k \geq 0} 2^{2k\alpha} |\tilde{\varphi}_k * x|^2)^{\frac{1}{2}}\|_p < \infty\}.$$

Munis de leurs normes naturelles, tous ces espaces deviennent des espaces de Banach. Notons que $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ est l'espace de Triebel-Lizorkin colonne, qui a des versions ligne et de mélange: $F_p^{\alpha,r}(\mathbb{T}_\theta^d)$ consiste de tous les x tels que $x^* \in F_p^{\alpha,c}(\mathbb{T}_\theta^d)$, muni de la norme $\|x\|_{F_p^{\alpha,r}} = \|x^*\|_{F_p^{\alpha,c}}$; l'espace de mélange $F_p^\alpha(\mathbb{T}_\theta^d)$ est défini par

$$F_p^\alpha(\mathbb{T}_\theta^d) = \begin{cases} F_p^{\alpha,c}(\mathbb{T}_\theta^d) + F_p^{\alpha,r}(\mathbb{T}_\theta^d) & \text{si } 1 \leq p < 2, \\ F_p^{\alpha,c}(\mathbb{T}_\theta^d) \cap F_p^{\alpha,r}(\mathbb{T}_\theta^d) & \text{si } 2 \leq p < \infty, \end{cases}$$

muni de la norme

$$\|x\|_{F_p^\alpha} = \begin{cases} \inf \{ \|y\|_{F_p^{\alpha,c}} + \|z\|_{F_p^{\alpha,r}} : x = y + z \} & \text{si } 1 \leq p < 2, \\ \max(\|x\|_{F_p^{\alpha,c}}, \|x\|_{F_p^{\alpha,r}}) & \text{si } 2 \leq p < \infty. \end{cases}$$

Nous rappelons également les espaces de Hardy sur le tore quantique, définis dans [17]. Pour un élément x dans $L_1(\mathbb{T}_\theta^d)$, son intégrale de Poisson est définie par (notons $|\cdot|$ la norme euclidienne de $\mathbb{R}^d$)

$$\mathbb{P}_r(x) = \sum_{m \in \mathbb{Z}^d} \hat{x}(m) r^{|m|} U^m, \quad 0 \leq r < 1.$$

Sa g -fonction de Littlewood-Paley associée est

$$s^c(x) = \left(\int_0^1 \left| \frac{\partial}{\partial r} \mathbb{P}_r(x) \right|^2 (1-r) dr \right)^{\frac{1}{2}}.$$

Pour $1 \leq p < \infty$, posons

$$\|x\|_{\mathcal{H}_p^c} = |\widehat{x}(0)| + \|s^c(x)\|_{L_p(\mathbb{T}_\theta^d)}.$$

L'espace de Hardy colonne $\mathcal{H}_p^c(\mathbb{T}_\theta^d)$ est alors défini par

$$\mathcal{H}_p^c(\mathbb{T}_\theta^d) = \{x \in L_1(\mathbb{T}_\theta^d) : \|x\|_{\mathcal{H}_p^c} < \infty\}.$$

L'espace de Hardy ligne $\mathcal{H}_p^r(\mathbb{T}_\theta^d)$ est défini comme l'espace des x tels que $x^* \in \mathcal{H}_p^c(\mathbb{T}_\theta^d)$ muni de la norme naturelle. Les espaces de Hardy mélange sont définis de la façon suivante : si $1 \leq p < 2$,

$$\mathcal{H}_p(\mathbb{T}_\theta^d) = \mathcal{H}_p^c(\mathbb{T}_\theta^d) + \mathcal{H}_p^r(\mathbb{T}_\theta^d)$$

muni de la norme

$$\|x\|_{\mathcal{H}_p} = \inf\{\|y\|_{\mathcal{H}_p^c} + \|z\|_{\mathcal{H}_p^r} : x = y + z, y \in \mathcal{H}_p^c(\mathbb{T}_\theta^d), z \in \mathcal{H}_p^r(\mathbb{T}_\theta^d)\},$$

et si $2 \leq p < \infty$,

$$\mathcal{H}_p(\mathbb{T}_\theta^d) = \mathcal{H}_p^c(\mathbb{T}_\theta^d) \cap \mathcal{H}_p^r(\mathbb{T}_\theta^d)$$

muni de la norme d'intersection

$$\|x\|_{\mathcal{H}_p} = \max\{\|x\|_{\mathcal{H}_p^c}, \|x\|_{\mathcal{H}_p^r}\}.$$

Rappelons aussi que $\text{BMO}^c(\mathbb{T}_\theta^d)$ (dans [17]) est la dualité de $\mathcal{H}_1^c(\mathbb{T}_\theta^d)$ avec la norme

$$\max\{|\widehat{x}(0)|, \sup_{0 \leq r < 1} \|\mathbb{P}_r(|x - \mathbb{P}_r(x)|^2)\|_{\mathbb{T}_\theta^d}^{\frac{1}{2}}\},$$

et que $\text{BMO}^r(\mathbb{T}_\theta^d)$ est l'ensemble de x tel que $x^* \in \text{BMO}^c(\mathbb{T}_\theta^d)$, avec

$$\|x\|_{\text{BMO}^r} = \|x^*\|_{\text{BMO}^c}.$$

Alors l'espace mélange $\text{BMO}(\mathbb{T}_\theta^d)$ est l'intersection :

$$\text{BMO}(\mathbb{T}_\theta^d) = \text{BMO}^c(\mathbb{T}_\theta^d) \cap \text{BMO}^r(\mathbb{T}_\theta^d),$$

muni de la norm

$$\|x\|_{\text{BMO}} = \max(\|x\|_{\text{BMO}^c}, \|x\|_{\text{BMO}^r}).$$

Nous allons utiliser fréquemment la notation $A \lesssim B$, qui représente l'inégalité $A \leq cB$ pour une constante c . Les constantes correspondantes dans toutes telles inégalités dépendent possiblement de la dimension d , les fonctions de test φ ou ψ , les indices α , p ou q , etc. mais ne dépendent jamais des fonctions f ou distributions x en considération. $A \approx B$ signifie alors $A \lesssim B$ et $A \gtrsim B$ au même temps.

0.2 Propriétés fondamentales

Une propriété fondamentale commune aux espaces de Sobolev, Besov and Triebel-Lizorkin est le théorème de réduction par le potentiel de Bessel ou Riesz. Celui-ci est évident pour les espaces de Sobolev Potentiels $H_p^\alpha(\mathbb{T}_\theta^d)$ par définition. Le théorème de réduction par le potentiel de Riesz de l'espace de Sobolev Potentiel $H_p^\alpha(\mathbb{T}_\theta^d)$ peut être montré par la caractérisation de $H_p^\alpha(\mathbb{T}_\theta^d)$ suivante :

Théorème 0.1. *Soit $1 \leq p \leq \infty$. Alors*

$$\|x\|_{H_p^\alpha} \approx \left(|\hat{x}(0)|^p + \|I^\alpha(x - \hat{x}(0))\|_p^p \right)^{\frac{1}{p}},$$

où les constantes équivalentes dépendent seulement de α et d .

Pour les espaces de Sobolev, nous obtenons une inégalité de type Poincaré, qui assure le théorème de réduction de $W_p^k(\mathbb{T}_\theta^d)$. Pour $x \in W_p^k(\mathbb{T}_\theta^d)$, posons

$$|x|_{W_p^k} = \left(\sum_{m \in \mathbb{N}_0^d, |m|_1 = k} \|D^m x\|_p^p \right)^{\frac{1}{p}}.$$

Nous prouvons dans le chapitre 2 le résultat suivant. Soulignons que la preuve diffère de l'argument standard dans le cas classique.

Théorème 0.2. *Soit $1 \leq p \leq \infty$. Alors pour $x \in W_p^1(\mathbb{T}_\theta^d)$,*

$$\|x - \hat{x}(0)\|_p \lesssim |x|_{W_p^1}.$$

Plus généralement, si $k \in \mathbb{N}$ et $x \in W_p^k(\mathbb{T}_\theta^d)$ avec $\hat{x}(0) = 0$, alors

$$|x|_{W_p^j} \lesssim |x|_{W_p^k}, \quad \forall 0 \leq j < k.$$

Par conséquent, $|\hat{x}(0)| + |x|_{W_p^k}$ est une norme équivalente sur $W_p^k(\mathbb{T}_\theta^d)$.

Pour les espaces de Besov et Triebel-Lizorkin, le théorème de réduction est aussi valable pour le potentiel de Riesz et les dérivées partielles. De plus, ce théorème de réduction est valable même pour les dérivées partielles d'ordre fractionnel, définies de la façon suivante : Pour $a \in \mathbb{R}_+$, nous définissons $D_{i,a}(\xi) = (2\pi i \xi_i)^a$ pour $\xi \in \mathbb{R}^d$, et D_i^a le multiplicateur de Fourier associé sur $\mathbb{T}_\theta^d$. Posons $D_a = D_{1,a_1} \cdots D_{d,a_d}$ et $D^a = D_1^{a_1} \cdots D_d^{a_d}$ pour tout $a = (a_1, \dots, a_d) \in \mathbb{R}_+^d$. Notons que si a est un nombre entier positif, $D_i^a = \partial_i^a$, donc cette généralisation de dérivées partielles ne pose pas de conflit de notation.

Théorème 0.3. *Soient $1 \leq p, q \leq \infty$, $\alpha \in \mathbb{R}$.*

(i) *Pour tout $\beta \in \mathbb{R}$, J^β et I^β sont des isomorphismes entre $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ et $B_{p,q}^{\alpha-\beta}(\mathbb{T}_\theta^d)$.*

(ii) *Soit $a \in \mathbb{R}_+^d$. Si $x \in B_{p,q}^\alpha(\mathbb{T}_\theta^d)$, alors $D^a x \in B_{p,q}^{\alpha-|a|_1}(\mathbb{T}_\theta^d)$ et*

$$\|D^a x\|_{B_{p,q}^{\alpha-|a|_1}} \lesssim \|x\|_{B_{p,q}^\alpha}.$$

(iii) *Soit $\beta > 0$. Alors $x \in B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ si et seulement si $D_i^\beta x \in B_{p,q}^{\alpha-\beta}(\mathbb{T}_\theta^d)$ pour tout $i = 1, \dots, d$. En outre, dans ce cas,*

$$\|x\|_{B_{p,q}^\alpha} \approx |\hat{x}(0)| + \sum_{i=1}^d \|D_i^\beta x\|_{B_{p,q}^{\alpha-\beta}}.$$

Théorème 0.4. *Soient $1 \leq p < \infty$ et $\alpha \in \mathbb{R}$.*

(i) *Pour tout $\beta \in \mathbb{R}$, J^β et I^β sont des isomorphismes entre $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ et $F_p^{\alpha-\beta,c}(\mathbb{T}_\theta^d)$. En particulier, J^α et I^α sont des isomorphismes entre $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ et $\mathcal{H}_p^c(\mathbb{T}_\theta^d)$.*

(ii) Soit $a \in \mathbb{R}_+^d$. Si $x \in F_p^{\alpha,c}(\mathbb{T}_\theta^d)$, alors $D^a x \in F_p^{\alpha-|a|_1,c}(\mathbb{T}_\theta^d)$ et

$$\|D^a x\|_{F_p^{\alpha-|a|_1,c}} \lesssim \|x\|_{F_p^{\alpha,c}}.$$

(iii) Soit $\beta > 0$. Alors $x \in F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ si et seulement si $D_i^\beta x \in F_p^{\alpha-\beta,c}(\mathbb{T}_\theta^d)$ pour tout $i = 1, \dots, d$. En outre, dans ce cas,

$$\|x\|_{F_p^{\alpha,c}} \approx |\widehat{x}(0)| + \sum_{i=1}^d \|D_i^\beta x\|_{F_p^{\alpha-\beta,c}}.$$

Nous démontrons aussi que $W_\infty^k(\mathbb{T}_\theta^d)$ est l'analogue de la classe de Lipschitz classique d'ordre k pour $\mathbb{T}_\theta^d$. Pour $u \in \mathbb{R}^d$, définissons $\Delta_u x = \pi_z(x) - x$, où $z = (e^{2\pi i u_1}, \dots, e^{2\pi i u_d})$ et π_z est l'automorphisme de $\mathbb{T}_\theta^d$ déterminé par $U_j \mapsto z_j U_j$ pour $1 \leq j \leq d$. Alors pour un entier positif k , Δ_u^k est l'opérateur différentiel d'ordre k sur $\mathbb{T}_\theta^d$, associé à u . Notons que Δ_u^k est aussi le multiplicateur de Fourier de symbole e_u^k , où $e_u(\xi) = e^{2\pi i u \cdot \xi} - 1$. Le module d'ordre k de L_p -lissité de $x \in L_p(\mathbb{T}_\theta^d)$ est défini par

$$\omega_p^k(x, \varepsilon) = \sup_{0 < |u| \leq \varepsilon} \|\Delta_u^k x\|_p.$$

Nous démontrons alors que pour $1 \leq p \leq \infty$ et $k \in \mathbb{N}$,

$$\sup_{\varepsilon > 0} \frac{\omega_p^k(x, \varepsilon)}{\varepsilon^k} \approx \sum_{m \in \mathbb{N}_0^d, |m|_1 = k} \|D^m x\|_p.$$

En particulier, nous retrouvons le résultat de Weaver [78, 79] sur la classe de Lipschitz sur $\mathbb{T}_\theta^d$ quand $p = \infty$ et $k = 1$.

Les autres résultats concernent la relation entre les quatre familles d'espaces de fonction. La majeure partie sont les suivants :

(i) Pour $1 < p < \infty$, $H_p^k(\mathbb{T}_\theta^d) = W_p^k(\mathbb{T}_\theta^d)$ avec des normes équivalentes.

(ii) Pour $1 \leq p \leq \infty$ et $\alpha \in \mathbb{R}^d$, nous avons les inclusions suivantes :

$$\begin{aligned} B_{p,p}^\alpha(\mathbb{T}_\theta^d) &\subset H_p^\alpha(\mathbb{T}_\theta^d) \subset B_{p,2}^\alpha(\mathbb{T}_\theta^d) \quad \text{si } 1 \leq p \leq 2, \\ B_{p,2}^\alpha(\mathbb{T}_\theta^d) &\subset H_p^\alpha(\mathbb{T}_\theta^d) \subset B_{p,p}^\alpha(\mathbb{T}_\theta^d) \quad \text{si } 2 \leq p \leq \infty. \end{aligned}$$

(iii) Pour $1 \leq p < \infty$, $B_{p,\min(p,2)}^\alpha(\mathbb{T}_\theta^d) \subset F_p^{\alpha,c}(\mathbb{T}_\theta^d) \subset B_{p,\max(p,2)}^\alpha(\mathbb{T}_\theta^d)$.

(iv) Pour $1 < p < \infty$ et $\alpha \in \mathbb{R}$, $F_p^\alpha(\mathbb{T}_\theta^d) = H_p^\alpha(\mathbb{T}_\theta^d)$ avec des normes équivalentes.

En conséquence, nous en déduisons que l'espace de Sobolev potentiel $H_p^\alpha(\mathbb{T}_\theta^d)$ admet une caractérisation de Littlewood-Paley pour $1 < p < \infty$.

0.3 Caractérisations

La seconde famille de résultats sont diverses caractérisations des espaces de Besov et Triebel-Lizorkin. C'est la partie la plus difficile et la plus technique dans cette thèse. Dans le cadre classique, toutes les preuves existantes que nous connaissons pour ces caractérisations utilisent des techniques de fonction maximale de façon cruciale. Comme évoqué

précédemment, ces techniques ne sont plus disponibles maintenant. En remplacement, nous utilisons fréquemment les multiplicateurs de Fourier. Soulignons que nos résultats sont meilleurs que ceux existants dans la littérature, même dans le cas classique.

Illustrons cela en énonçant une caractérisation générale des espaces de Besov. Soit h une fonction de Schwartz auxiliaire telle que

$$\text{supp } h \subset \{\xi \in \mathbb{R}^d : |\xi| \leq 4\} \quad \text{et } h = 1 \text{ sur } \{\xi \in \mathbb{R}^d : |\xi| \leq 2\}.$$

Soit $\alpha_0, \alpha_1 \in \mathbb{R}$. Soit ψ une fonction infiniment différentiable sur $\mathbb{R}^d \setminus \{0\}$ vérifiant les conditions suivantes :

$$\begin{cases} |\psi| > 0 \text{ sur } \{\xi : 2^{-1} \leq |\xi| \leq 2\}, \\ \mathcal{F}^{-1}(\psi h I_{-\alpha_1}) \in L_1(\mathbb{R}^d), \\ \sup_{j \in \mathbb{N}_0} 2^{-\alpha_0 j} \|\mathcal{F}^{-1}(\psi(2^j \cdot) \varphi)\|_1 < \infty. \end{cases}$$

La première condition de non annulation sur ψ est une condition Tauberienne. L'intégrabilité de la transformation inverse de Fourier peut être réduite à un critère plus commode grâce à l'espace de Sobolev potentiel $H_2^\sigma(\mathbb{R}^d)$ avec $\sigma > \frac{d}{2}$. Nous utiliserons la même notation pour ψ que pour φ . En particulier, ψ_k est la transformation inverse de Fourier de $\psi(2^{-k} \cdot)$, et $\tilde{\psi}_k$ est le multiplicateur de Fourier sur $\mathbb{T}_\theta^d$ avec le symbole $\psi(2^{-k} \cdot)$. Pour les paramètres continus, nous allons utiliser la notation similaire : étant donné $\varepsilon > 0$, notons ψ_ε la fonction qui a la transformation de Fourier $\psi^{(\varepsilon)} = \psi(\varepsilon \cdot)$, et $\tilde{\psi}_\varepsilon$ le multiplicateur de Fourier sur $\mathbb{T}_\theta^d$ associé à $\psi^{(\varepsilon)}$. Nous avons donc

Théorème 0.5. *Soient $1 \leq p, q \leq \infty$, $\alpha \in \mathbb{R}$ et $\alpha_0 < \alpha < \alpha_1$. Supposons que ψ vérifie l'hypothèse précédente. Alors une distribution x sur $\mathbb{T}_\theta^d$ appartient à $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ si et seulement si*

$$\left(\sum_{k \geq 0} (2^{k\alpha} \|\tilde{\psi}_k * x\|_p)^q \right)^{\frac{1}{q}} < \infty.$$

Dans ce cas, nous avons

$$\|x\|_{B_{p,q}^\alpha} \approx \left(|\hat{x}(0)|^q + \sum_{k \geq 0} (2^{k\alpha} \|\tilde{\psi}_k * x\|_p)^q \right)^{\frac{1}{q}}.$$

De la même façon, pour toute distribution x sur $\mathbb{T}_\theta^d$,

$$\|x\|_{B_{p,q}^\alpha} \approx \left(|\hat{x}(0)|^q + \int_0^1 \varepsilon^{-q\alpha} \|\tilde{\psi}_\varepsilon * x\|_p^q \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}}.$$

Deux exemples précis et importants pour la fonction ψ sont donnés par les semi-groupes de Poisson et de la chaleur. Rappelons que $\mathbb{P}_r(x)$ définit l'intégrale de Poisson circulaire d'une distribution x sur $\mathbb{T}_\theta^d$:

$$\mathbb{P}_r(x) = \sum_{m \in \mathbb{Z}^d} \hat{x}(m) r^{|m|} U^m, \quad 0 \leq r < 1.$$

Nous introduisons aussi le noyau de la chaleur circulaire $\mathbb{W}$ de $\mathbb{T}^d$:

$$\mathbb{W}_r(z) = \sum_{m \in \mathbb{Z}^d} r^{|m|^2} z^m, \quad z \in \mathbb{T}^d, \quad 0 \leq r < 1.$$

Alors pour $x \in \mathcal{S}'(\mathbb{T}_\theta^d)$ nous posons

$$\mathbb{W}_r(x) = \sum_{m \in \mathbb{Z}^d} \widehat{x}(m) r^{|m|^2} U^m, \quad 0 \leq r < 1.$$

Pour une distribution x sur $\mathbb{T}_\theta^d$ et $k \in \mathbb{Z}$, soit

$$\mathcal{J}_r^k \mathbb{P}_r(x) = \sum_{m \in \mathbb{Z}^d} C_{m,k} \widehat{x}(m) r^{|m|-k} U^m, \quad 0 \leq r < 1,$$

où

$$C_{m,k} = |m| \cdots (|m| - k + 1) \quad \text{si } k \geq 0 \quad \text{et} \quad C_{m,k} = \frac{1}{(|m| + 1) \cdots (|m| - k)} \quad \text{si } k < 0.$$

Notons que $\mathcal{J}_r^k$ est la dérivée d'ordre k associée à r si $k \geq 0$, et est l'intégration d'ordre $(-k)$ si $k < 0$. De même, nous pouvons définir $\mathcal{J}_r^k \mathbb{W}_r(x)$.

Théorème 0.6. *Soient $1 \leq p, q \leq \infty$, $\alpha \in \mathbb{R}$ et $k \in \mathbb{Z}$. Pour une distribution x sur $\mathbb{T}_\theta^d$, nous avons :*

(i) *Si $k > \alpha$, alors*

$$\|x\|_{B_{p,q}^\alpha} \approx \left(\max_{|m| < k} |\widehat{x}(m)|^q + \int_0^1 (1-r)^{q(k-\alpha)} \|\mathcal{J}_r^k \mathbb{P}_r(x_k)\|_p^q \frac{dr}{1-r} \right)^{\frac{1}{q}},$$

$$\text{où } x_k = x - \sum_{|m| < k} \widehat{x}(m) U^m.$$

(ii) *Si $k > \frac{\alpha}{2}$, alors*

$$\|x\|_{B_{p,q}^\alpha} \approx \left(\max_{|m|^2 < k} |\widehat{x}(m)|^q + \int_0^1 (1-r)^{q(k-\frac{\alpha}{2})} \|\mathcal{J}_r^k \mathbb{W}_r(x)\|_p^q \frac{dr}{1-r} \right)^{\frac{1}{q}}.$$

L'utilisation des opérateurs intégraux (correspondant à k négatif) dans l'énoncé précédent semble complètement nouveau, même dans le cas $\theta = 0$ (le cas commutatif). Des analogues des Théorèmes 0.5 et 0.6 sont aussi valables pour $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$. Pour les espaces de Triebel-Lizorkin, une autre amélioration de notre caractérisation par rapport à celle classique est l'hypothèse faite sur k : dans le cas classique, k doit être plus grand que $d + \max(\alpha, 0)$, mais ici nous ne requérons que $k > \alpha$.

Théorème 0.7. *Avec la même notation que dans les deux théorèmes précités, toutes les normes suivantes sont équivalentes à la norme de $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$:*

- $|\widehat{x}(0)| + \left\| \left(\sum_{k \geq 0} (2^{2\alpha k} |\widetilde{\psi}_k * x|^2)^{\frac{1}{2}} \right) \right\|_p.$
- $|\widehat{x}(0)| + \left\| \left(\int_0^1 \varepsilon^{-2\alpha} |\widetilde{\psi}_\varepsilon * x|^2 \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{2}} \right\|_p.$
- $\max_{|m| < k} |\widehat{x}(m)| + \left\| \left(\int_0^1 (1-r)^{2(k-\alpha)} |\mathcal{J}_r^k \mathbb{P}_r(x_k)|^2 \frac{dr}{1-r} \right)^{\frac{1}{2}} \right\|_p, \text{ pour } k > \alpha.$
- $\max_{|m|^2 < k} |\widehat{x}(m)| + \left\| \left(\int_0^1 (1-r)^{2(k-\frac{\alpha}{2})} |\mathcal{J}_r^k \mathbb{W}_r(x)|^2 \frac{dr}{1-r} \right)^{\frac{1}{2}} \right\|_p, \text{ pour } k > \frac{\alpha}{2}.$

La caractérisation de l'espace de Besov classique par différences est aussi étendue au cadre quantique. Ce résultat ressemble au précédent en termes de dérivées du semi-groupe de Poisson.

Théorème 0.8. *Pour $1 \leq p, q \leq \infty$ et $\alpha \in \mathbb{R}, k \in \mathbb{N}$ avec $0 < \alpha < k$, nous posons*

$$\|x\|_{B_{p,q}^{\alpha,\omega}} = \left(\int_0^1 \varepsilon^{-\alpha q} \omega_p^k(x, \varepsilon)^q \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}}.$$

Alors $x \in B_{p,q}^{\alpha}(\mathbb{T}_{\theta}^d)$ si et seulement si $\|x\|_{B_{p,q}^{\alpha,\omega}} < \infty$, et $\|x\|_{B_{p,q}^{\alpha}} \approx |\hat{x}(0)| + \|x\|_{B_{p,q}^{\alpha,\omega}}$.

La caractérisation de l'espace de Besov par différences montre que $B_{\infty,\infty}^{\alpha}(\mathbb{T}_{\theta}^d)$ est l'analogue quantique de la classe de Zygmund classique. En particulier, pour $0 < \alpha < 1$, $B_{\infty,\infty}^{\alpha}(\mathbb{T}_{\theta}^d)$ est la classe de Hölder d'ordre α , déjà étudiée par Weaver [79].

Dans le cadre non commutatif, le comportement de la limite de la quantité $\|x\|_{B_{p,q}^{\alpha,\omega}}$ lorsque $\alpha \rightarrow k$ ou $\alpha \rightarrow 0$ fait l'objet d'une série de publications récentes. Elles ont été initiées par Bourgain, Brézis et Mironescu [13, 14] qui considéraient le cas $\alpha \rightarrow 1$ ($k = 1$). Leur travail a été simplifié et étendu par Maz'ya et Shaposhnikova [41]. Ici, nous obtenons l'analogue suivant de leurs résultats pour $\mathbb{T}_{\theta}^d$:

Corollaire 0.9. *Pour $1 \leq p \leq \infty, 1 \leq q < \infty$ et $0 < \alpha < k$ avec $k \in \mathbb{N}$,*

$$\lim_{\alpha \rightarrow k} (k - \alpha)^{\frac{1}{q}} \|x\|_{B_{p,q}^{\alpha,\omega}} \approx q^{-\frac{1}{q}} \sum_{m \in \mathbb{N}_0^d, |m|_1 = k} \|D^m x\|_p,$$

$$\lim_{\alpha \rightarrow 0} \alpha^{\frac{1}{q}} \|x\|_{B_{p,q}^{\alpha,\omega}} \approx q^{-\frac{1}{q}} \|x\|_p$$

avec des constantes dépendant seulement de d et k .

0.4 Interpolation

Notre troisième famille de résultats concernant l'interpolation. Comme dans le cas usuel, l'interpolation des espaces de Besov est vraiment simple, et celle des espaces de Triebel-Lizorkin peut être réduite facilement au problème correspondant des espaces de Hardy. Nous en faisons la liste de la façon suivante :

- $(B_{p,q_0}^{\alpha_0}(\mathbb{T}_{\theta}^d), B_{p,q_1}^{\alpha_1}(\mathbb{T}_{\theta}^d))_{\eta,q} = B_{p,q}^{\alpha}(\mathbb{T}_{\theta}^d), \quad \alpha_0 \neq \alpha_1, \alpha = (1 - \eta)\alpha_0 + \eta\alpha_1;$
- $(B_{p,q_0}^{\alpha}(\mathbb{T}_{\theta}^d), B_{p,q_1}^{\alpha}(\mathbb{T}_{\theta}^d))_{\eta,q} = B_{p,q}^{\alpha}(\mathbb{T}_{\theta}^d), \quad \frac{1}{q} = \frac{1 - \eta}{q_0} + \frac{\eta}{q_1};$
- $(B_{p_0,q_0}^{\alpha_0}(\mathbb{T}_{\theta}^d), B_{p_1,q_1}^{\alpha_1}(\mathbb{T}_{\theta}^d))_{\eta,q} = B_{p,q}^{\alpha}(\mathbb{T}_{\theta}^d), \quad \alpha_0 \neq \alpha_1, \alpha = (1 - \eta)\alpha_0 + \eta\alpha_1,$
 $\frac{1}{p} = \frac{1 - \eta}{p_0} + \frac{\eta}{p_1}, \frac{1}{q} = \frac{1 - \eta}{q_0} + \frac{\eta}{q_1}, p = q;$
- $(B_{p_0,q_0}^{\alpha_0}(\mathbb{T}_{\theta}^d), B_{p_1,q_1}^{\alpha_1}(\mathbb{T}_{\theta}^d))_{\eta} = B_{p,q}^{\alpha}(\mathbb{T}_{\theta}^d), \quad \alpha = (1 - \eta)\alpha_0 + \eta\alpha_1, \frac{1}{p} = \frac{1 - \eta}{p_0} + \frac{\eta}{p_1},$
 $\frac{1}{q} = \frac{1 - \eta}{q_0} + \frac{\eta}{q_1}, q < \infty;$
- $(F_{\infty}^{\alpha,c}(\mathbb{T}_{\theta}^d), F_1^{\alpha,c}(\mathbb{T}_{\theta}^d))_{\frac{1}{p}} = F_p^{\alpha}(\mathbb{T}_{\theta}^d) = (F_{\infty}^{\alpha,c}(\mathbb{T}_{\theta}^d), F_1^{\alpha,c}(\mathbb{T}_{\theta}^d))_{\frac{1}{p},p} \quad 1 < p < \infty;$
- $(F_p^{\alpha_0,c}(\mathbb{T}_{\theta}^d), F_p^{\alpha_1,c}(\mathbb{T}_{\theta}^d))_{\eta,q} = B_{p,q}^{\alpha}(\mathbb{T}_{\theta}^d), \quad 1 \leq p, q \leq \infty, \alpha_0 \neq \alpha_1,$
 $\alpha = (1 - \eta)\alpha_0 + \eta\alpha_1;$

- $(F_{\infty}^{\alpha_0, c}(\mathbb{T}_{\theta}^d), F_1^{\alpha_1, c}(\mathbb{T}_{\theta}^d))_{\frac{1}{p}} = F_p^{\alpha, c}(\mathbb{T}_{\theta}^d), \quad \alpha = (1 - \frac{1}{p})\alpha_0 + \frac{\alpha_1}{p}, 1 < p < \infty.$

Ici l'espace $F_{\infty}^{\alpha, c}(\mathbb{T}_{\theta}^d)$ est défini comme la dualité de $F_1^{-\alpha, c}(\mathbb{T}_{\theta}^d)$. Alors par la propriété de réduction de l'espace de Triebel-Lizorkin, J^{α} est aussi un isomorphisme entre $F_{\infty}^{\alpha, c}(\mathbb{T}_{\theta}^d)$ et $\text{BMO}^c(\mathbb{T}_{\theta}^d)$.

La tâche vraiment difficile concerne l'interpolation des espaces de Sobolev pour laquelle nous avons seulement obtenu des résultats partiels. Le couple le plus intéressant est $(W_1^k(\mathbb{T}_{\theta}^d), W_{\infty}^k(\mathbb{T}_{\theta}^d))$. Rappelons que l'interpolation complexe de ce couple reste toujours ouverte même dans le cas commutatif (un problème ouvert connu de longue date; il a été explicitement formulé par P. Jones dans [27, p. 173]), alors que son interpolation réelle a été déterminée par DeVore et Scherer [21]. Nous ne savons pas, malheureusement, comment prouver l'analogue quantique du théorème de DeVore et Scherer. Cependant, nous pouvons étendre la formule de K -fonctionnel du couple $(L_p(\mathbb{R}^d), W_p^k(\mathbb{R}^d))$ obtenue par Johnen et Scherer [30] au tore quantique. Ce résultat s'écrit de la façon suivante :

Théorème 0.10. *Soient $1 \leq p \leq \infty$ et $k \in \mathbb{N}$. Alors*

$$K(x, \varepsilon^k; L_p(\mathbb{T}_{\theta}^d), W_p^k(\mathbb{T}_{\theta}^d)) \approx \varepsilon^k |\widehat{x}(0)| + \omega_p^k(x, \varepsilon), \quad 0 < \varepsilon \leq 1$$

avec des constantes dépendant seulement de d et k .

En conséquence, nous déterminons les espaces d'interpolation réelle de $(L_p(\mathbb{T}_{\theta}^d), W_p^k(\mathbb{T}_{\theta}^d))$, qui sont les espaces de Besov.

Pour le couple $(H_{p_0}^{\alpha_0}(\mathbb{T}_{\theta}^d), H_{p_1}^{\alpha_1}(\mathbb{T}_{\theta}^d))$, notons que si $\alpha_0 = \alpha_1$, ce qui suit est facile à déduire :

$$(H_{p_0}^{\alpha}(\mathbb{T}_{\theta}^d), H_{p_1}^{\alpha}(\mathbb{T}_{\theta}^d))_{\eta} = H_p^{\alpha}(\mathbb{T}_{\theta}^d) \text{ et } (H_{p_0}^{\alpha}(\mathbb{T}_{\theta}^d), H_{p_1}^{\alpha}(\mathbb{T}_{\theta}^d))_{\eta, p} = H_p^{\alpha}(\mathbb{T}_{\theta}^d),$$

pour $\frac{1}{p} = \frac{1-\eta}{p_0} + \frac{\eta}{p_1}$, parce que J^{α} est une isométrie entre $H_p^{\alpha}(\mathbb{T}_{\theta}^d)$ et $L_p(\mathbb{T}_{\theta}^d)$ pour tout $1 \leq p \leq \infty$. Si $\alpha_0 \neq \alpha_1$, l'interpolation réelle de ce couple est déduite par l'interpolation réelle correspondante d'espaces de Besov. Le résultat est

$$(H_{p_0}^{\alpha_0}(\mathbb{T}_{\theta}^d), H_{p_1}^{\alpha_1}(\mathbb{T}_{\theta}^d))_{\eta, q} = B_{p, q}^{\alpha}(\mathbb{T}_{\theta}^d), \quad \alpha = (1 - \eta)\alpha_0 + \eta\alpha_1.$$

Pour l'interpolation complexe, nous considérons le couple $(H_{\text{BMO}}^{\alpha_0}(\mathbb{T}_{\theta}^d), H_{\mathcal{H}_1}^{\alpha_1}(\mathbb{T}_{\theta}^d))$, où $H_{\mathcal{H}_1}^{\alpha_1}(\mathbb{T}_{\theta}^d)$ est défini par

$$H_{\mathcal{H}_1}^{\alpha}(\mathbb{T}_{\theta}^d) = \{x \in \mathcal{S}'(\mathbb{T}_{\theta}^d) : J^{\alpha}x \in \mathcal{H}_1(\mathbb{T}_{\theta}^d)\} \text{ avec } \|x\|_{H_{\mathcal{H}_1}^{\alpha}} = \|J^{\alpha}x\|_{\mathcal{H}_1},$$

et $H_{\text{BMO}}^{\alpha}(\mathbb{T}_{\theta}^d)$ est défini de la même façon, par la BMO-norme de $J^{\alpha}x$.

Alors nous avons le résultat suivant

Théorème 0.11. *Soient $\alpha_0, \alpha_1 \in \mathbb{R}$ et $1 < p < \infty$. Alors*

$$(H_{\text{BMO}}^{\alpha_0}(\mathbb{T}_{\theta}^d), H_{\mathcal{H}_1}^{\alpha_1}(\mathbb{T}_{\theta}^d))_{\frac{1}{p}} = H_p^{\alpha}(\mathbb{T}_{\theta}^d), \quad \alpha = (1 - \frac{1}{p})\alpha_0 + \frac{\alpha_1}{p}.$$

Par le théorème de réitération, le théorème précédent implique le suivant :

Corollaire 0.12. *Soient $0 < \eta < 1$, $\alpha_0, \alpha_1 \in \mathbb{R}$ et $1 < p_0, p_1 < \infty$. Alors*

$$(H_{p_0}^{\alpha_0}(\mathbb{T}_{\theta}^d), H_{p_1}^{\alpha_1}(\mathbb{T}_{\theta}^d))_{\eta} = H_p^{\alpha}(\mathbb{T}_{\theta}^d), \quad \alpha = (1 - \eta)\alpha_0 + \eta\alpha_1, \quad \frac{1}{p} = \frac{1 - \eta}{p_0} + \frac{\eta}{p_1}.$$

Ce corollaire donne le résultat partiel de l'interpolation complexe de $(W_{p_0}^{k_0}(\mathbb{T}_\theta^d), W_{p_1}^{k_1}(\mathbb{T}_\theta^d))$ pour $1 < p_0 < p_1 < \infty$. Pour $k_0 = k_1$ et $p_0 = 1$, nous avons besoin de considérer les espaces de Hardy-Sobolev $W_{\text{BMO}}^k(\mathbb{T}_\theta^d)$ et $W_{\mathcal{H}_1}^k(\mathbb{T}_\theta^d)$. Ils sont définis par les normes :

$$\sup_{0 \leq |m|_1 \leq k} \|D^m x\|_{\text{BMO}} \quad \text{et} \quad \sum_{0 \leq |m|_1 \leq k} \|D^m x\|_{\mathcal{H}_1}$$

respectivement. Nous pouvons prouver que pour tout $k \in \mathbb{N}$, $W_{\text{BMO}}^k(\mathbb{T}_\theta^d) = H_{\text{BMO}}^k(\mathbb{T}_\theta^d)$ et $W_{\mathcal{H}_1}^k(\mathbb{T}_\theta^d) = H_{\mathcal{H}_1}^k(\mathbb{T}_\theta^d)$. Alors nous obtenons

Théorème 0.13. *Soient $k \in \mathbb{N}$ et $1 < p < \infty$. Alors pour $X = W_{\mathcal{H}_1}^k(\mathbb{T}_\theta^d)$ ou $X = W_1^k(\mathbb{T}_\theta^d)$,*

$$(W_{\text{BMO}}^k(\mathbb{T}_\theta^d), X)_{\frac{1}{p}} = W_p^k(\mathbb{T}_\theta^d) = (W_{\text{BMO}}^k(\mathbb{T}_\theta^d), X)_{\frac{1}{p}, p}.$$

Conséquemment, pour tout $0 < \eta < 1$ et $1 < p_0 < \infty$,

$$(W_{p_0}^k(\mathbb{T}_\theta^d), W_1^k(\mathbb{T}_\theta^d))_\eta = W_p^k(\mathbb{T}_\theta^d) = (W_{p_0}^k(\mathbb{T}_\theta^d), W_1^k(\mathbb{T}_\theta^d))_{\eta, p}, \quad \frac{1}{p} = \frac{1-\eta}{p_0} + \eta.$$

0.5 Plongements

La quatrième famille de résultats concernent le plongement des espaces définis précédemment. Un résultat typique est l'analogie de l'inégalité de plongement de Sobolev classique pour $W_p^k(\mathbb{T}_\theta^d)$.

Théorème 0.14. *Supposons que $1 < p < \infty$. Si $\alpha \in \mathbb{R}^+$, $\alpha p < d$ et $\frac{1}{p_1} = \frac{1}{p} - \frac{\alpha}{d}$, alors nous avons :*

$$H_p^\alpha(\mathbb{T}_\theta^d) \subset L_{p_1}(\mathbb{T}_\theta^d).$$

De la même façon, si $k \in \mathbb{N}$, $kp < d$ et $\frac{1}{p_1} = \frac{1}{p} - \frac{k}{d}$, alors nous avons :

$$W_p^k(\mathbb{T}_\theta^d) \subset L_{p_1}(\mathbb{T}_\theta^d).$$

D'autre part, si $\alpha p < d$ (ou $kp < d$), les espaces de Sobolev sont plongés dans les espaces de Hölder-Zygmund, ce qui coïncide bien avec les théorèmes de Sobolev classiques.

Théorème 0.15. *Supposons que $1 \leq p < \infty$ et $\alpha p > d$. Alors pour $\alpha_1 = \alpha - \frac{d}{p}$, nous avons*

$$H_p^\alpha(\mathbb{T}_\theta^d) \subset B_{\infty, \infty}^{\alpha_1}(\mathbb{T}_\theta^d).$$

De la même façon, si $1 < p < \infty$ et $k \in \mathbb{N}$, $kp > d$, alors nous avons pour $\alpha_1 = k - \frac{d}{p}$,

$$W_p^k(\mathbb{T}_\theta^d) \subset B_{\infty, \infty}^{\alpha_1}(\mathbb{T}_\theta^d).$$

Pour les espaces de Besov, nous avons

Théorème 0.16. *Supposons que $1 \leq p \leq p_1 \leq \infty$, $1 \leq q \leq q_1 \leq \infty$, et $\alpha - \frac{d}{p} = \alpha_1 - \frac{d}{p_1}$. Alors nous avons l'inclusion suivante :*

$$B_{p, q}^\alpha(\mathbb{T}_\theta^d) \subset B_{p_1, q_1}^{\alpha_1}(\mathbb{T}_\theta^d).$$

Des inégalités de plongement similaires sont aussi valables pour les espaces de Triebel-Lizorkin, ainsi que $F_p^\alpha(\mathbb{T}_\theta^d) = H_p^\alpha(\mathbb{T}_\theta^d)$ avec des normes équivalentes pour $1 < p < \infty$. Combinée avec l'interpolation réelle, l'inégalité de plongement de $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ assure l'inégalité de plongement des espaces de Sobolev. Nos preuves des inégalités de ces plongements sont basées sur le procédé célèbre de semi-groupe de Varopoulos [76] pour la théorie de Littlewood-Sobolev, qui était déjà employé par Junge et Mei [34] dans leur étude des espaces de BMO sur les semi-groupes quantiques Markoviens.

De plus, les trois théorèmes précédents ont des versions compactes si on réduit un peu les indices p_1 de $L_{p_1}(\mathbb{T}_\theta^d)$ et $B_{p_1,q_1}^{\alpha_1}(\mathbb{T}_\theta^d)$, ou α_1 de $B_{\infty,\infty}^{\alpha_1}(\mathbb{T}_\theta^d)$. Ils étendent le théorème de compacité de Rellich-Kondrachof dans la théorie classique de plongement Sobolev.

Théorème 0.17. (i) *Soient $1 \leq p < p_1 \leq \infty, 1 \leq q \leq q_1 \leq \infty$ et $\alpha - \frac{d}{p} = \alpha_1 - \frac{d}{p_1}$. Alors le plongement $B_{p,q}^\alpha(\mathbb{T}_\theta^d) \hookrightarrow B_{p^*,q_1}^{\alpha_1}(\mathbb{T}_\theta^d)$ est compact pour $1 \leq p^* < p_1$.*

(ii) *Si $p > 1, \alpha - \frac{d}{p} = \alpha_1 - \frac{d}{p_1}$, alors $H_p^\alpha(\mathbb{T}_\theta^d) \hookrightarrow H_{p^*}^{\alpha_1}(\mathbb{T}_\theta^d)$ est compact pour $1 \leq p^* < p_1$. En particulier, si $\alpha = k$ et $\alpha_1 = k_1$ sont des nombres entiers positifs, alors $W_p^k(\mathbb{T}_\theta^d) \hookrightarrow W_{p^*}^{k_1}(\mathbb{T}_\theta^d)$ est compact.*

(iii) *Si $p > 1, p(\alpha - \alpha_1) > d$ et $\alpha^* < \alpha_1 = \alpha - \frac{d}{p}$, alors $H_p^\alpha(\mathbb{T}_\theta^d) \hookrightarrow B_{\infty,\infty}^{\alpha^*}(\mathbb{T}_\theta^d)$ est compact. En particulier, si $\alpha = k \in \mathbb{N}$, alors $W_p^k(\mathbb{T}_\theta^d) \hookrightarrow B_{\infty,\infty}^{\alpha^*}(\mathbb{T}_\theta^d)$ est compact.*

0.6 Multiplicateur

La dernière famille de résultats de cette thèse décrivent les multiplicateurs de Fourier sur les espaces précédents. Comme dans le cas L_p traité dans [17], nous nous intéressons particulièrement aux multiplicateurs de Fourier complètement bornés. Tous les espaces considérés ont une structure naturelle d'espace d'opérateur au sens de Pisier. Inspiré par le théorème de transférence de Neuwirth et Ricard [48], les multiplicateurs de Fourier sur $L_p(\mathbb{T}_\theta^d)$ sont associés aux multiplicateurs de Schur sur la classe de Schatten S_p dans [17]. Pour une distribution x sur $\mathbb{T}_\theta^d$, nous notons sa matrice en base $(U^m)_{m \in \mathbb{Z}^d}$:

$$[x] = \left(\langle x U^n, U^m \rangle \right)_{m,n \in \mathbb{Z}^d} = \left(\hat{x}(m-n) e^{i n \tilde{\theta} (m-n)^t} \right)_{m,n \in \mathbb{Z}^d}.$$

Ici k^t note la transposée de $k = (k_1, \dots, k_d)$, et $\tilde{\theta}$ est la $d \times d$ -matrice suivante déduite de la matrice anti-symétrique θ :

$$\tilde{\theta} = -2\pi \begin{pmatrix} 0 & \theta_{12} & \theta_{13} & \dots & \theta_{1d} \\ 0 & 0 & \theta_{23} & \dots & \theta_{2d} \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & \dots & \theta_{d-1,d} \\ 0 & 0 & 0 & \dots & 0 \end{pmatrix}.$$

Soient $\phi : \mathbb{Z}^d \rightarrow \mathbb{C}$ et M_ϕ le multiplicateur de Fourier associé sur $\mathbb{T}_\theta^d$. Posons $\mathring{\phi} = (\phi_{m-n})_{m,n \in \mathbb{Z}^d}$. Alors

$$[M_\phi x] = (\phi_{m-n} \hat{x}(m-n) e^{i n \tilde{\theta} (m-n)^t})_{m,n \in \mathbb{Z}^d} = S_{\mathring{\phi}}([x]),$$

où $S_{\mathring{\phi}}$ est le multiplicateur de Schur avec le symbole $\mathring{\phi}$. Si X est un espace de Banach de distributions sur $\mathbb{T}_\theta^d$, notons $\mathbf{M}(X)$ l'espace des multiplicateurs de Fourier bornés sur X ; si

X a de plus une structure d'espaces d'opérateur, $M_{cb}(X)$ est l'espace des multiplicateurs de Fourier complètement bornés sur X . De la même façon, si X est un espace de Schatten L_p (ou un sous-espace fermé), nous notons $M(X)$ l'espace des multiplicateurs de Schur bornés sur X ; avec sa structure naturelle d'espace d'opérateur, nous notons aussi $M_{cb}(X)$ l'espace de multiplicateurs de Schur complètement bornés sur X . Ces espaces de multiplicateurs sont équipés avec leurs normes naturelles. Il est démontré dans [17] que, pour $1 \leq p \leq \infty$, $M_{cb}(L_p(\mathbb{T}_\theta^d)) = M_{cb}(S_p)$ avec des normes égales. Conséquemment, $M_{cb}(L_p(\mathbb{T}_\theta^d))$ est indépendant de θ .

En utilisant de nouveau le transference de Neuwirth-Ricard entre les multiplicateurs de Fourier et de Schur, nous déduisons un résultat similaire pour l'espace de Sobolev $W_p^k(\mathbb{T}_\theta^d)$ et l'espace de Triebel-Lizorkin $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$.

Théorème 0.18. *Soient $1 \leq p \leq \infty$, et $k \in \mathbb{N}$, $\alpha \in \mathbb{R}$. Alors*

$$M_{cb}(W_p^k(\mathbb{T}_\theta^d)) = M_{cb}(W_p^k(\mathbb{T}^d)) \text{ avec des normes égales,}$$

et

$$M_{cb}(F_p^{\alpha,c}(\mathbb{T}_\theta^d)) = M_{cb}(F_p^{\alpha,c}(\mathbb{T}^d)) \text{ avec des normes égales.}$$

Notons que, bien qu'il s'agisse d'un espace de fonction commutatif, l'espace $F_p^\alpha(\mathbb{T}^d)$ (correspondant à $\theta = 0$) est muni de trois structures différentes d'espace d'opérateur, les deux premières définies par les plongements dans $L_p(\mathbb{T}^d; \ell_2^{\alpha,c})$ et $L_p(\mathbb{T}^d; \ell_2^{\alpha,r})$, la troisième étant le mélange des deux. Il en résulte trois espaces d'opérateur différents, notés $F_p^{\alpha,c}(\mathbb{T}^d)$, $F_p^{\alpha,r}(\mathbb{T}^d)$ et $F_p^\alpha(\mathbb{T}^d)$, respectivement. Le théorème précédent donne en particulier que $M_{cb}(\mathcal{H}_p(\mathbb{T}_\theta^d)) = M_{cb}(\mathcal{H}_p^c(\mathbb{T}^d))$ avec des normes égales.

La situation pour les espaces de Besov est très satisfaisante puisqu'il est bien connu que les multiplicateurs de Fourier se comportent bien mieux sur les espaces de Besov que sur les espaces L_p (dans le cas commutatif). Nous prouvons le résultat suivant

Théorème 0.19. *Soient $\alpha \in \mathbb{R}$ et $1 \leq p, q \leq \infty$. Soit $\phi : \mathbb{Z}^d \rightarrow \mathbb{C}$. Alors ϕ est un multiplicateur de Fourier sur $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ si et seulement si les $\phi\varphi^{(k)}$ sont des multiplicateurs de Fourier sur $L_p(\mathbb{T}_\theta^d)$ uniformément en k . Dans ce cas, nous avons*

$$\|\phi\|_{M(B_{p,q}^\alpha(\mathbb{T}_\theta^d))} \approx |\phi(0)| + \sup_{k \geq 0} \|\phi\varphi^{(k)}\|_{M(L_p(\mathbb{T}_\theta^d))}$$

avec des constantes dépendant seulement de α . Une version complètement bornée similaire est également vraie.

En conséquence, les multiplicateurs de Fourier sur $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ sont complètement déterminés par les multiplicateurs de Fourier sur $L_p(\mathbb{T}_\theta^d)$ associés à leurs composantes dans la décomposition de Littlewood-Paley. Donc les multiplicateurs complètement bornés sur $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ dépendent seulement de p . Dans le cas $p = 1$, un multiplicateur est borné sur $B_{1,q}^\alpha(\mathbb{T}_\theta^d)$ si et seulement si il est complètement borné si et seulement si il est la transformation de Fourier d'une fonction dans $B_{1,\infty}^0(\mathbb{T}^d)$. En utilisant un exemple classique de Stein-Zygmund [70], nous démontrons qu'il existe un ϕ qui est un multiplicateur complètement borné sur $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ pour tout p mais non borné sur $L_p(\mathbb{T}_\theta^d)$ pour tout $p \neq 2$.

Introduction

This thesis deals with analysis on quantum tori. Quantum or noncommutative torus are fundamental examples in operator algebras and probably the most accessible interesting class of objects in noncommutative geometry (cf. [18]). Their algebraic and geometric aspects have been sufficiently well understood; there exist extensive works on these; see, for instance, the survey paper by Rieffel [63]. However, very little had been done about their analytic aspect. Presumably, this deficiency is due to numerous difficulties one may encounter when dealing with noncommutative L_p -spaces, since these spaces come up unavoidably if one wishes to do analysis.

[17] is the first systematic work on harmonic analysis on quantum tori. It studies several subjects of harmonic analysis, including maximal inequalities, mean and pointwise convergences of Fourier series, completely bounded Fourier multipliers on L_p -spaces and Hardy spaces. It was directly inspired by the current development on noncommutative harmonic analysis, and was made possible by the recent developments on noncommutative martingale/ergodic inequalities and the Littlewood-Paley-Stein theory for quantum Markovian semigroups, which had been achieved thanks to the efforts of many researchers; see, for instance, [56, 31, 36, 37, 61, 62, 52], and [32, 44, 45, 33, 34].

This thesis intends to study function spaces on quantum tori. Thus it is a natural continuation of [17]. The spaces to be investigated are (potential) Sobolev, Besov and Triebel-Lizorkin spaces. In the classical setting, these spaces are fundamental for many branches of mathematics such as harmonic analysis, PDE, functional analysis and approximation theory. Our references for the classical theory are [1, 42, 49, 53, 73, 74]. However, they have never been investigated so far in the quantum case, except two special cases to our best knowledge. Sobolev spaces with the L_2 -norm were studied by Spera [65] in view of applications to the Yang-Mills theory for quantum tori [66] (see also [26, 39, 58, 64] for some related works). On the other hand, inspired by Connes' noncommutative geometry [18], or more precisely, the part on noncommutative metric spaces, Weaver [78, 79] developed the Lipschitz classes of order α for $0 < \alpha \leq 1$ on quantum tori. The fact that only these two cases have been studied illustrates once more the above mentioned difficulties related to noncommutativity.

Among these difficulties, one is to be emphasized: it is notably relevant to this thesis, and is the lack of a noncommutative analogue of the usual pointwise maximal function. However, maximal function techniques play a paramount role in the classical theory of Besov and Triebel-Lizorkin spaces (as well as in the theory of Hardy spaces). They are no longer available in the quantum setting, which forces us to invent new tools, like in the previously quoted works on noncommutative martingale inequalities and the quantum Littlewood-Paley-Stein theory where the same difficulty has already appeared.

One powerful tool used in [17] is the transference method. It consists in transferring problems on quantum tori to the corresponding ones in the case of operator-valued func-

tions on the usual tori, in order to use existing results in the latter case or adapt classical arguments. This method is efficient for several problems studied in [17], including the maximal inequalities and Hardy spaces. It is still useful for some parts of the present work, for instance, for Besov spaces which can be investigated through the classical vector-valued Besov spaces by means of transference, the relevant Banach spaces being the noncommutative L_p -spaces on quantum tori. However, it becomes inefficient for others. For example, the Sobolev embedding inequalities cannot be proved by transference. On the other hand, if one wishes to study Triebel-Lizorkin spaces on quantum tori via transference, one should first develop the theory of operator-valued Triebel-Lizorkin spaces on the classical tori. The latter is as hard as the former. Contrary to [17], the transference method will play a very limited role in this thesis. Instead, we will use Fourier multipliers in a crucial way, this approach is of interest in its own right. We thus develop an intrinsic differential analysis on quantum tori, without frequently referring to the usual tori via transference as in the previous one. This is a major advantage of the present methods over [17]. We hope that the study carried out here would open new perspectives of applications and motivate more future research works on quantum tori or in similar circumstances.

This thesis consists of seven chapters, written in English. It is based on a work joint with Quanhua Xu and Zhi Yin. In this introduction, we will first recall some necessary definitions and notation (see the respective sections below for more details), and give the definitions of the spaces to be considered. Then we will describe the main results of the thesis by classifying them into five families.

0.1 Notation and definitions

Let $d \geq 2$ and $\theta = (\theta_{kj})$ be a real skew-symmetric $d \times d$ -matrix. The d -dimensional noncommutative torus $\mathcal{A}_\theta$ is the universal C^* -algebra generated by d unitary operators $U_1, \dots, U_d$ satisfying the following commutation relation

$$U_k U_j = e^{2\pi i \theta_{kj}} U_j U_k, \quad 1 \leq j, k \leq d.$$

Let $U = (U_1, \dots, U_d)$. For $m = (m_1, \dots, m_d) \in \mathbb{Z}^d$, set

$$U^m = U_1^{m_1} \dots U_d^{m_d}.$$

A polynomial in U is a finite sum:

$$x = \sum_{m \in \mathbb{Z}^d} \alpha_m U^m, \quad \alpha_m \in \mathbb{C}.$$

For such a polynomial x , we define $\tau(x) = \alpha_0$. Then, τ extends to a faithful tracial state on $\mathcal{A}_\theta$. Let $\mathbb{T}_\theta^d$ be the w^* -closure of $\mathcal{A}_\theta$ in the GNS representation of τ . This is our d -dimensional quantum torus. It is to be viewed as a deformation of the usual d -torus $\mathbb{T}^d$, or more precisely, of the commutative algebra $L_\infty(\mathbb{T}^d)$. The noncommutative L_p -spaces associated to $\mathbb{T}_\theta^d$ are denoted by $L_p(\mathbb{T}_\theta^d)$. The Fourier transform of an element $x \in L_1(\mathbb{T}_\theta^d)$ is defined by

$$\widehat{x}(m) = \tau((U^m)^* x), \quad m \in \mathbb{Z}^d.$$

The formal Fourier series of x is thus given by

$$x \sim \sum_{m \in \mathbb{Z}^d} \widehat{x}(m) U^m.$$

The differential structure of $\mathbb{T}_\theta^d$ is modeled on that of $\mathbb{T}^d$. Let

$$\mathcal{S}(\mathbb{T}_\theta^d) = \left\{ \sum_{m \in \mathbb{Z}^d} a_m U^m : \{a_m\}_{m \in \mathbb{Z}^d} \text{ rapidly decreasing} \right\}.$$

This is the deformation of the space of infinitely differentiable functions on $\mathbb{T}^d$. Like in the commutative case, $\mathcal{S}(\mathbb{T}_\theta^d)$ carries a natural locally convex topology. Its topological dual $\mathcal{S}'(\mathbb{T}_\theta^d)$ is the space of distributions on $\mathbb{T}_\theta^d$. The partial derivations on $\mathcal{S}(\mathbb{T}_\theta^d)$ are determined by

$$\partial_j(U_j) = 2\pi i U_j \quad \text{and} \quad \partial_j(U_k) = 0, \quad k \neq j, \quad 1 \leq j, k \leq d.$$

Denote $\mathbb{N}_0$ the set of nonnegative integers. Given $m = (m_1, \dots, m_d) \in \mathbb{N}_0^d$, the associated partial derivation D^m is defined to be $\partial_1^{m_1} \dots \partial_d^{m_d}$. The order of D^m is $|m|_1 = m_1 + \dots + m_d$. Let $\Delta = \partial_1^2 + \dots + \partial_d^2$ be the Laplacian. By duality, the derivations and Fourier transform transfer to $\mathcal{S}'(\mathbb{T}_\theta^d)$ too.

Given a function $\phi : \mathbb{Z}^d \rightarrow \mathbb{C}$, let M_ϕ denote the associated Fourier multiplier on $\mathbb{T}^d$, namely, $\widehat{M_\phi f}(m) = \phi(m) \widehat{f}(m)$ for any trigonometric polynomial f on $\mathbb{T}^d$. We call ϕ a multiplier on $L_p(\mathbb{T}^d)$ if M_ϕ extends to a bounded map on $L_p(\mathbb{T}^d)$. Fourier multipliers on $\mathbb{T}_\theta^d$ are defined exactly in the same way. We still use the same symbol M_ϕ to denote the corresponding multiplier on $\mathbb{T}_\theta^d$ too. For a function ϕ on $\mathbb{R}^d$, we will call ϕ rather than $\phi|_{\mathbb{Z}^d}$ a Fourier multiplier on $L_p(\mathbb{T}^d)$ or $L_p(\mathbb{T}_\theta^d)$. This should not cause any ambiguity in concrete contexts.

Two special multipliers will be used frequently: the Bessel potential $J^\alpha = (1 - (2\pi)^{-2}\Delta)^{\frac{\alpha}{2}}$ and the Riesz potential $I^\alpha = (- (2\pi)^{-2}\Delta)^{\frac{\alpha}{2}}$. They are defined by

$$J^\alpha(x) = \sum_{m \in \mathbb{Z}^d} (1 + |m|^2)^{\frac{\alpha}{2}} \widehat{x}(m) U^m$$

and

$$I^\alpha(x) = \sum_{m \in \mathbb{Z}^d} |m|^\alpha \widehat{x}(m) U^m$$

for a distribution x on $\mathbb{T}_\theta^d$. Here the Riesz potential is defined for distributions x with $\widehat{x}(0) = 0$. We denote also $J_\alpha(\xi) = (1 + |\xi|^2)^{\frac{\alpha}{2}}$ on $\mathbb{R}^d$ and $I_\alpha(\xi) = |\xi|^\alpha$ on $\mathbb{R}^d \setminus \{0\}$. Then $J_\alpha(\xi)$ and $I_\alpha(\xi)$ are the symbols of the Fourier multipliers J^α and I^α , respectively.

Fix a Schwartz function φ on $\mathbb{R}^d$ satisfying the usual Littlewood-Paley decomposition property:

$$\begin{cases} \text{supp } \varphi \subset \{\xi : 2^{-1} \leq |\xi| \leq 2\}, \\ \varphi > 0 \text{ on } \{\xi : 2^{-1} < |\xi| < 2\}, \\ \sum_{k \in \mathbb{Z}} \varphi(2^{-k}\xi) = 1, \quad \xi \neq 0. \end{cases}$$

For each $k \geq 0$ let φ_k be the function whose Fourier transform is equal to $\varphi(2^{-k}\cdot)$. For a distribution x on $\mathbb{T}_\theta^d$, define

$$\tilde{\varphi}_k * x = \sum_{m \in \mathbb{Z}^d} \varphi(2^{-k}m) \widehat{x}(m) U^m.$$

So $x \mapsto \tilde{\varphi}_k * x$ is the Fourier multiplier with symbol $\varphi(2^{-k}\cdot)$.

We can now define the four families of function spaces on $\mathbb{T}_\theta^d$ to be studied. Let $1 \leq p, q \leq \infty$ and $k \in \mathbb{N}, \alpha \in \mathbb{R}$.

- *Sobolev spaces:*

$$W_p^k(\mathbb{T}_\theta^d) = \{x \in \mathcal{S}'(\mathbb{T}_\theta^d) : D^m x \in L_p(\mathbb{T}_\theta^d) \text{ for each } m \in \mathbb{N}_0^d \text{ with } |m|_1 \leq k\}.$$

- *Potential Sobolev spaces:*

$$H_p^\alpha(\mathbb{T}_\theta^d) = \{x \in \mathcal{S}'(\mathbb{T}_\theta^d) : J^\alpha x \in L_p(\mathbb{T}_\theta^d)\}.$$

- *Besov spaces:*

$$B_{p,q}^\alpha(\mathbb{T}_\theta^d) = \{x \in \mathcal{S}'(\mathbb{T}_\theta^d) : (|\widehat{x}(0)|^q + \sum_{k \geq 0} 2^{qk\alpha} \|\tilde{\varphi}_k * x\|_p^q)^{\frac{1}{q}} < \infty\}.$$

- *Triebel-Lizorkin spaces for $p < \infty$:*

$$F_p^{\alpha,c}(\mathbb{T}_\theta^d) = \{x \in \mathcal{S}'(\mathbb{T}_\theta^d) : \|(|\widehat{x}(0)|^2 + \sum_{k \geq 0} 2^{2k\alpha} |\tilde{\varphi}_k * x|^2)^{\frac{1}{2}}\|_p < \infty\}.$$

Equipped with their natural norms, all these spaces become Banach spaces. Note that $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ is the column Triebel-Lizorkin space, which has the row and mixture versions: $F_p^{\alpha,r}(\mathbb{T}_\theta^d)$ consists of all x such that $x^* \in F_p^{\alpha,c}(\mathbb{T}_\theta^d)$, equipped with the norm $\|x\|_{F_p^{\alpha,r}} = \|x^*\|_{F_p^{\alpha,c}}$; the mixture space $F_p^\alpha(\mathbb{T}_\theta^d)$ is defined by

$$F_p^\alpha(\mathbb{T}_\theta^d) = \begin{cases} F_p^{\alpha,c}(\mathbb{T}_\theta^d) + F_p^{\alpha,r}(\mathbb{T}_\theta^d) & \text{if } 1 \leq p < 2, \\ F_p^{\alpha,c}(\mathbb{T}_\theta^d) \cap F_p^{\alpha,r}(\mathbb{T}_\theta^d) & \text{if } 2 \leq p < \infty, \end{cases}$$

equipped with

$$\|x\|_{F_p^\alpha} = \begin{cases} \inf \{\|y\|_{F_p^{\alpha,c}} + \|z\|_{F_p^{\alpha,r}} : x = y + z\} & \text{if } 1 \leq p < 2, \\ \max(\|x\|_{F_p^{\alpha,c}}, \|x\|_{F_p^{\alpha,r}}) & \text{if } 2 \leq p < \infty. \end{cases}$$

We also recall the Hardy spaces on quantum tori defined in [17]. For an element x in $L_1(\mathbb{T}_\theta^d)$, its Poisson integral is defined by (with $|\cdot|$ denoting the Euclidean norm of $\mathbb{R}^d$)

$$\mathbb{P}_r(x) = \sum_{m \in \mathbb{Z}^d} \widehat{x}(m) r^{|m|} U^m, \quad 0 \leq r < 1.$$

Its associated Littlewood-Paley g -function is

$$s^c(x) = \left(\int_0^1 \left| \frac{\partial}{\partial r} \mathbb{P}_r(x) \right|^2 (1-r) dr \right)^{\frac{1}{2}}.$$

For $1 \leq p < \infty$ let

$$\|x\|_{\mathcal{H}_p^c} = |\widehat{x}(0)| + \|s^c(x)\|_{L_p(\mathbb{T}_\theta^d)}.$$

The column Hardy space $\mathcal{H}_p^c(\mathbb{T}_\theta^d)$ is then defined to be

$$\mathcal{H}_p^c(\mathbb{T}_\theta^d) = \{x \in L_1(\mathbb{T}_\theta^d) : \|x\|_{\mathcal{H}_p^c} < \infty\}.$$

The row Hardy space $\mathcal{H}_p^r(\mathbb{T}_\theta^d)$ is defined to be the space of all x such that $x^* \in \mathcal{H}_p^c(\mathbb{T}_\theta^d)$, equipped with the natural norm. The mixture Hardy space is

$$\mathcal{H}_p(\mathbb{T}_\theta^d) = \begin{cases} \mathcal{H}_p^c(\mathbb{T}_\theta^d) + \mathcal{H}_p^r(\mathbb{T}_\theta^d) & \text{if } 1 \leq p < 2, \\ \mathcal{H}_p^c(\mathbb{T}_\theta^d) \cap \mathcal{H}_p^r(\mathbb{T}_\theta^d) & \text{if } 2 \leq p < \infty, \end{cases}$$

equipped with the sum and intersection norms, respectively:

$$\|x\|_{\mathcal{H}_p} = \begin{cases} \inf \{ \|y\|_{\mathcal{H}_p^c} + \|z\|_{\mathcal{H}_p^r} : x = y + z \} & \text{if } 1 \leq p < 2, \\ \max(\|x\|_{\mathcal{H}_p^c}, \|x\|_{\mathcal{H}_p^r}) & \text{if } 2 \leq p < \infty. \end{cases}$$

Recall also that $\text{BMO}^c(\mathbb{T}_\theta^d)$ (in [17]) is the duality of $\mathcal{H}_1^c(\mathbb{T}_\theta^d)$ with the norm

$$\max \{ |\widehat{x}(0)|, \sup_{0 \leq r < 1} \|\mathbb{P}_r(|x - \mathbb{P}_r(x)|^2)\|_{\mathbb{T}_\theta^d}^{\frac{1}{2}} \},$$

and that row $\text{BMO}^r(\mathbb{T}_\theta^d)$ consists of all x such that $x^* \in \text{BMO}^c(\mathbb{T}_\theta^d)$, equipped with the norm

$$\|x\|_{\text{BMO}^r} = \|x^*\|_{\text{BMO}^c}.$$

Then the mixture space $\text{BMO}(\mathbb{T}_\theta^d)$ is the intersection of the column and row BMO spaces:

$$\text{BMO}(\mathbb{T}_\theta^d) = \text{BMO}^c(\mathbb{T}_\theta^d) \cap \text{BMO}^r(\mathbb{T}_\theta^d),$$

equipped with

$$\|x\|_{\text{BMO}} = \max(\|x\|_{\text{BMO}^c}, \|x\|_{\text{BMO}^r}).$$

We will frequently use the notation $A \lesssim B$, which is an inequality up to a constant: $A \leq cB$ for some constant $c > 0$. The relevant constants in all such inequalities may depend on the dimension d , the test function φ or ψ , the indices α , p or q , etc. but never on the functions f or distributions x in consideration. $A \approx B$ means thus $A \lesssim B$ and $A \gtrsim B$.

0.2 Basic properties

A common basic property of potential Sobolev, Besov and Triebel-Lizorkin spaces is a reduction theorem by the Bessel potential. This is obvious for the potential Sobolev space $H_p^\alpha(\mathbb{T}_\theta^d)$ just by definition. The reduction theorem of $H_p^\alpha(\mathbb{T}_\theta^d)$ for the Riesz potential is ensured by a characterization of $H_p^\alpha(\mathbb{T}_\theta^d)$ norm below:

Theorem 0.1. *Let $1 \leq p \leq \infty$. Then*

$$\|x\|_{H_p^\alpha} \approx \left(|\widehat{x}(0)|^p + \|I^\alpha(x - \widehat{x}(0))\|_p^p \right)^{\frac{1}{p}},$$

where the equivalence constants depend only on α and d .

We obtain a Poincaré type inequality, which yields the reduction theorem of Sobolev spaces. For $x \in W_p^k(\mathbb{T}_\theta^d)$ let

$$|x|_{W_p^k} = \left(\sum_{m \in \mathbb{N}_0^d, |m|_1 = k} \|D^m x\|_p^p \right)^{\frac{1}{p}}.$$

Then we show in chapter 2 the following result.

Theorem 0.2. *Let $1 \leq p \leq \infty$. Then for any $x \in W_p^1(\mathbb{T}_\theta^d)$,*

$$\|x - \widehat{x}(0)\|_p \lesssim |x|_{W_p^1}.$$

More generally, if $k \in \mathbb{N}$ and $x \in W_p^k(\mathbb{T}_\theta^d)$ with $\widehat{x}(0) = 0$, then

$$|x|_{W_p^j} \lesssim |x|_{W_p^k}, \quad \forall 0 \leq j < k.$$

Consequently, $|\widehat{x}(0)| + |x|_{W_p^k}$ is an equivalent norm on $W_p^k(\mathbb{T}_\theta^d)$.

Our proof of this inequality greatly differs with standard arguments for such results in the commutative setting.

For Besov and Triebel-Lizorkin spaces, the reduction theorem also holds for the Riesz and Bessel potentials and partial derivatives. Moreover, this reduction theorem of partial derivatives holds in a generalized sense for the so-called fractional derivatives defined as follows: Given $a \in \mathbb{R}_+$, we define $D_{i,a}(\xi) = (2\pi i \xi_i)^a$ for $\xi \in \mathbb{R}^d$, and D_i^a to be the associated Fourier multiplier on $\mathbb{T}_\theta^d$. We set $D_a = D_{1,a_1} \cdots D_{d,a_d}$ and $D^a = D_1^{a_1} \cdots D_d^{a_d}$ for any $a = (a_1, \dots, a_d) \in \mathbb{R}_+^d$. Note that if a is a positive integer, $D_i^a = \partial_i^a$, so this generalization of partial derivatives does not cause any conflict of notation.

Theorem 0.3. *Let $1 \leq p, q \leq \infty$, $\alpha \in \mathbb{R}$.*

(i) *For any $\beta \in \mathbb{R}$, both J^β and I^β are isomorphisms between $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ and $B_{p,q}^{\alpha-\beta}(\mathbb{T}_\theta^d)$.*

(ii) *Let $a \in \mathbb{R}_+^d$. If $x \in B_{p,q}^\alpha(\mathbb{T}_\theta^d)$, then $D^a x \in B_{p,q}^{\alpha-|a|_1}(\mathbb{T}_\theta^d)$ and*

$$\|D^a x\|_{B_{p,q}^{\alpha-|a|_1}} \lesssim \|x\|_{B_{p,q}^\alpha}.$$

(iii) *Let $\beta > 0$. Then $x \in B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ iff $D_i^\beta x \in B_{p,q}^{\alpha-\beta}(\mathbb{T}_\theta^d)$ for all $i = 1, \dots, d$. Moreover, in this case,*

$$\|x\|_{B_{p,q}^\alpha} \approx |\widehat{x}(0)| + \sum_{i=1}^d \|D_i^\beta x\|_{B_{p,q}^{\alpha-\beta}}.$$

Theorem 0.4. *Let $1 \leq p < \infty$ and $\alpha \in \mathbb{R}$.*

(i) *For any $\beta \in \mathbb{R}$, both J^β and I^β are isomorphisms between $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ and $F_p^{\alpha-\beta,c}(\mathbb{T}_\theta^d)$. In particular, J^α and I^α are isomorphisms between $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ and $\mathcal{H}_p^c(\mathbb{T}_\theta^d)$.*

(ii) *Let $a \in \mathbb{R}_+^d$. If $x \in F_p^{\alpha,c}(\mathbb{T}_\theta^d)$, then $D^a x \in F_p^{\alpha-|a|_1,c}(\mathbb{T}_\theta^d)$ and*

$$\|D^a x\|_{F_p^{\alpha-|a|_1,c}} \lesssim \|x\|_{F_p^{\alpha,c}}.$$

(iii) *Let $\beta > 0$. Then $x \in F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ iff $D_i^\beta x \in F_p^{\alpha-\beta,c}(\mathbb{T}_\theta^d)$ for all $i = 1, \dots, d$. Moreover, in this case,*

$$\|x\|_{F_p^{\alpha,c}} \approx |\widehat{x}(0)| + \sum_{i=1}^d \|D_i^\beta x\|_{F_p^{\alpha-\beta,c}}.$$

We also show that $W_\infty^k(\mathbb{T}_\theta^d)$ is the analogue for $\mathbb{T}_\theta^d$ of the classical Lipschitz class of order k . For $u \in \mathbb{R}^d$, define $\Delta_u x = \pi_z(x) - x$, where $z = (e^{2\pi i u_1}, \dots, e^{2\pi i u_d})$ and π_z is the automorphism of $\mathbb{T}_\theta^d$ determined by $U_j \mapsto z_j U_j$ for $1 \leq j \leq d$. Then for a positive

integer k , Δ_u^k is the k th difference operator on $\mathbb{T}_\theta^d$ associated to u . Note that Δ_u^k is also the Fourier multiplier with symbol e_u^k , where $e_u(\xi) = e^{2\pi i u \cdot \xi} - 1$. The k th order modulus of L_p -smoothness of an $x \in L_p(\mathbb{T}_\theta^d)$ is defined by

$$\omega_p^k(x, \varepsilon) = \sup_{0 < |u| \leq \varepsilon} \|\Delta_u^k x\|_p.$$

We then prove that for any $1 \leq p \leq \infty$ and $k \in \mathbb{N}$,

$$\sup_{\varepsilon > 0} \frac{\omega_p^k(x, \varepsilon)}{\varepsilon^k} \approx \sum_{m \in \mathbb{N}_0^d, |m|_1 = k} \|D^m x\|_p.$$

In particular, we recover Weaver's results [78, 79] on the Lipschitz class on $\mathbb{T}_\theta^d$ when $p = \infty$ and $k = 1$.

Other results concern the links between these spaces. Part of them are listed as follows:

- (i) For $1 < p < \infty$, $H_p^k(\mathbb{T}_\theta^d) = W_p^k(\mathbb{T}_\theta^d)$ with equivalent norms.
- (ii) For $1 \leq p \leq \infty$ and $\alpha \in \mathbb{R}^d$, $B_{p, \min(p, 2)}^\alpha(\mathbb{T}_\theta^d) \subset H_p^\alpha(\mathbb{T}_\theta^d) \subset B_{p, \max(p, 2)}^\alpha(\mathbb{T}_\theta^d)$.
- (iii) For $1 \leq p < \infty$, $B_{p, \min(p, 2)}^\alpha(\mathbb{T}_\theta^d) \subset F_p^{\alpha, c}(\mathbb{T}_\theta^d) \subset B_{p, \max(p, 2)}^\alpha(\mathbb{T}_\theta^d)$.
- (iv) For $1 < p < \infty$ and $\alpha \in \mathbb{R}$, $F_p^\alpha(\mathbb{T}_\theta^d) = H_p^\alpha(\mathbb{T}_\theta^d)$ with equivalent norms.

As a consequence of the last one, we deduce that the potential Sobolev space $H_p^\alpha(\mathbb{T}_\theta^d)$ admits a Littlewood-Paley characterization for $1 < p < \infty$.

0.3 Characterizations

The second family of results are various characterizations of Besov and Triebel-Lizorkin spaces. This is the most difficult and technical part of the thesis. In the classical case, all existing proofs of these characterizations that we know use maximal function techniques in a crucial way. As pointed out earlier, these techniques are unavailable now. Instead, we use frequently Fourier multipliers. We would like to emphasize that our results are better than those in literature even in the commutative case. Let us illustrate this by stating a general characterization of Besov spaces. Let h be an auxiliary Schwartz function such that

$$\text{supp } h \subset \{\xi \in \mathbb{R}^d : |\xi| \leq 4\} \quad \text{and} \quad h = 1 \quad \text{on} \quad \{\xi \in \mathbb{R}^d : |\xi| \leq 2\}.$$

Let $\alpha_0, \alpha_1 \in \mathbb{R}$. Let ψ be an infinitely differentiable function on $\mathbb{R}^d \setminus \{0\}$ satisfying the following conditions

$$\begin{cases} |\psi| > 0 \quad \text{on} \quad \{\xi : 2^{-1} \leq |\xi| \leq 2\}, \\ \mathcal{F}^{-1}(\psi h I_{-\alpha_1}) \in L_1(\mathbb{R}^d), \\ \sup_{j \in \mathbb{N}_0} 2^{-\alpha_0 j} \|\mathcal{F}^{-1}(\psi(2^j \cdot) \varphi)\|_1 < \infty. \end{cases}$$

The integrability of the inverse Fourier transforms can be reduced to a handier criterion by means of the potential Sobolev space $H_2^\sigma(\mathbb{R}^d)$ with $\sigma > \frac{d}{2}$. We will use the same notation for ψ as for φ . In particular, ψ_k is the inverse Fourier transform of $\psi(2^{-k} \cdot)$ and $\tilde{\psi}_k$ is the Fourier multiplier on $\mathbb{T}_\theta^d$ with symbol $\psi(2^{-k} \cdot)$. For continuous parameters, we will use similar notation: given $\varepsilon > 0$, ψ_ε denotes the function with Fourier transform $\psi^{(\varepsilon)} = \psi(\varepsilon \cdot)$, and $\tilde{\psi}_\varepsilon$ denotes the Fourier multiplier on $\mathbb{T}_\theta^d$ associated to $\psi^{(\varepsilon)}$. Then we have

Theorem 0.5. *Let $1 \leq p, q \leq \infty$ and $\alpha \in \mathbb{R}$. Assume $\alpha_0 < \alpha < \alpha_1$. Let ψ satisfy the above assumption. Then a distribution x on $\mathbb{T}_\theta^d$ belongs to $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ iff*

$$\left(\sum_{k \geq 0} (2^{k\alpha} \|\tilde{\psi}_k * x\|_p)^q \right)^{\frac{1}{q}} < \infty.$$

If this is the case, then

$$\|x\|_{B_{p,q}^\alpha} \approx \left(|\hat{x}(0)|^q + \sum_{k \geq 0} (2^{k\alpha} \|\tilde{\psi}_k * x\|_p)^q \right)^{\frac{1}{q}}.$$

Similarly, for any distribution x on $\mathbb{T}_\theta^d$,

$$\|x\|_{B_{p,q}^\alpha} \approx \left(|\hat{x}(0)|^q + \int_0^1 \varepsilon^{-q\alpha} \|\tilde{\psi}_\varepsilon * x\|_p^q \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}}.$$

Two precise and important examples of ψ above are given in terms of the circular Poisson and heat semigroup. Recall that $\mathbb{P}_r(x)$ denote the Poisson integral of a distribution x on $\mathbb{T}_\theta^d$:

$$\mathbb{P}_r(x) = \sum_{m \in \mathbb{Z}^d} \hat{x}(m) r^{|m|} U^m, \quad 0 \leq r < 1.$$

Accordingly, we introduce the circular heat kernel $\mathbb{W}$ of $\mathbb{T}^d$:

$$\mathbb{W}_r(z) = \sum_{m \in \mathbb{Z}^d} r^{|m|^2} z^m, \quad z \in \mathbb{T}^d, \quad 0 \leq r < 1.$$

Then for $x \in \mathcal{S}'(\mathbb{T}_\theta^d)$ we put

$$\mathbb{W}_r(x) = \sum_{m \in \mathbb{Z}^d} \hat{x}(m) r^{|m|^2} U^m, \quad 0 \leq r < 1.$$

Given a distribution x on $\mathbb{T}_\theta^d$ and $k \in \mathbb{Z}$, let

$$\mathcal{J}_r^k \mathbb{P}_r(x) = \sum_{m \in \mathbb{Z}^d} C_{m,k} \hat{x}(m) r^{|m| - k} U^m, \quad 0 \leq r < 1,$$

where

$$C_{m,k} = |m| \cdots (|m| - k + 1) \quad \text{if } k \geq 0 \quad \text{and} \quad C_{m,k} = \frac{1}{(|m| + 1) \cdots (|m| - k)} \quad \text{if } k < 0.$$

Note that $\mathcal{J}_r^k$ is the k -th derivation relative to r if $k \geq 0$, and the $(-k)$ -th integration if $k < 0$. $\mathcal{J}_r^k \mathbb{W}_r(x)$ is defined similarly. Then we have the following

Theorem 0.6. *Let $1 \leq p, q \leq \infty$, $\alpha \in \mathbb{R}$ and $k \in \mathbb{Z}$. Let x be a distribution on $\mathbb{T}_\theta^d$.*

(i) *If $k > \alpha$, then*

$$\|x\|_{B_{p,q}^\alpha} \approx \left(\max_{|m| < k} |\hat{x}(m)|^q + \int_0^1 (1-r)^{q(k-\alpha)} \|\mathcal{J}_r^k \mathbb{P}_r(x_k)\|_p^q \frac{dr}{1-r} \right)^{\frac{1}{q}},$$

$$\text{where } x_k = x - \sum_{|m| < k} \hat{x}(m) U^m.$$

(ii) If $k > \frac{\alpha}{2}$, then

$$\|x\|_{B_{p,q}^\alpha} \approx \left(\max_{|m|^2 < k} |\widehat{x}(m)|^q + \int_0^1 (1-r)^{q(k-\frac{\alpha}{2})} \|\mathcal{J}_r^k \mathbb{W}_r(x)\|_p^q \frac{dr}{1-r} \right)^{\frac{1}{q}}.$$

The use of the integration operators (corresponding to negative k) in the above statement seems completely new even in the case $\theta = 0$ (the commutative case). Counterparts of Theorems 0.5 and 0.6 hold for $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ too. But they are much more difficult than Besov spaces. For Triebel-Lizorkin spaces, another improvement of our characterization over the classical one lies on the assumption on k : in the classical case, k is required to be greater than $d + \max(\alpha, 0)$, while we require only $k > \alpha$. See the third term below.

Theorem 0.7. *With the same assumptions as in the above two theorems, the following norms are all equivalent to the norm of $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$:*

- $|\widehat{x}(0)| + \left\| \left(\sum_{k \geq 0} (2^{2\alpha k} |\widetilde{\psi}_k * x|^2)^{\frac{1}{2}} \right) \right\|_p.$
- $|\widehat{x}(0)| + \left\| \left(\int_0^1 \varepsilon^{-2\varepsilon\alpha} |\widetilde{\psi}_\varepsilon * x|^2 \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{2}} \right\|_p.$
- $\max_{|m|^2 < k} |\widehat{x}(m)| + \left\| \left(\int_0^1 (1-r)^{2(k-\alpha)} |\mathcal{J}_r^k \mathbb{P}_r(x_k)|^2 \frac{dr}{1-r} \right)^{\frac{1}{2}} \right\|_p, \text{ for } k > \alpha.$
- $\max_{|m|^2 < k} |\widehat{x}(m)| + \left\| \left(\int_0^1 (1-r)^{2(k-\frac{\alpha}{2})} |\mathcal{J}_r^k \mathbb{W}_r(x)|^2 \frac{dr}{1-r} \right)^{\frac{1}{2}} \right\|_p, \text{ for } k > \frac{\alpha}{2}.$

In the classic setting, the characterizations of both Besov and Triebel-Lizorkin spaces are strongly based on maximal function techniques, which, as pointed out earlier, are not available in quantum case. Instead, the Fourier multiplier will play an important role. Moreover, since the noncommutative Triebel-Lizorkin spaces are subspaces of Hilbert-valued noncommutative L_p spaces, we will develop Fourier multiplier theory for such spaces. Besides, we will also need variance of square function characterizations of Hardy spaces developed in [81].

The classical characterization of Besov spaces by differences is also extended to the quantum setting. This result resembles the previous one in terms of the derivations of the Poisson semigroup.

Theorem 0.8. *For $1 \leq p, q \leq \infty$ and $\alpha \in \mathbb{R}, k \in \mathbb{N}$ with $0 < \alpha < k$, let*

$$\|x\|_{B_{p,q}^{\alpha,\omega}} = \left(\int_0^1 \varepsilon^{-\alpha q} \omega_p^k(x, \varepsilon)^q \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}}.$$

Then $x \in B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ iff $\|x\|_{B_{p,q}^{\alpha,\omega}} < \infty$, and $\|x\|_{B_{p,q}^\alpha} \approx |\widehat{x}(0)| + \|x\|_{B_{p,q}^{\alpha,\omega}}$.

The difference characterization of Besov spaces shows that $B_{\infty,\infty}^\alpha(\mathbb{T}_\theta^d)$ is the quantum analogue of the classical Zygmund class. In particular, for $0 < \alpha < 1$, $B_{\infty,\infty}^\alpha(\mathbb{T}_\theta^d)$ is the Hölder class of order α , already studied by Weaver's [79].

In the commutative case, the limit behavior of the quantity $\|x\|_{B_{p,q}^{\alpha,\omega}}$ as $\alpha \rightarrow k$ or $\alpha \rightarrow 0$ are object of a recent series of publications. This line of research was initiated by Bourgain, Brézis and Mironescu [13, 14] who considered the case $\alpha \rightarrow 1$ ($k = 1$). Their work was later simplified and extended by Maz'ya and Shaposhnikova [41]. Here, we obtain the following analogue for $\mathbb{T}_\theta^d$ of their results:

Corollary 0.9. For $1 \leq p \leq \infty$, $1 \leq q < \infty$ and $0 < \alpha < k$ with $k \in \mathbb{N}$,

$$\lim_{\alpha \rightarrow k} (k - \alpha)^{\frac{1}{q}} \|x\|_{B_{p,q}^{\alpha,\omega}} \approx q^{-\frac{1}{q}} \sum_{m \in \mathbb{N}_0^d, |m|_1 = k} \|D^m x\|_p,$$

$$\lim_{\alpha \rightarrow 0} \alpha^{\frac{1}{q}} \|x\|_{B_{p,q}^{\alpha,\omega}} \approx q^{-\frac{1}{q}} \|x\|_p$$

with relevant constants depending only on d and k .

0.4 Interpolation

Our third family of results deal with interpolation. Like in the classical setting, the interpolation of Besov spaces is quite simple, and that of Triebel-Lizorkin spaces can be easily reduced to the corresponding problem of Hardy spaces. Here we list them as follows

- $(B_{p,q_0}^{\alpha_0}(\mathbb{T}_\theta^d), B_{p,q_1}^{\alpha_1}(\mathbb{T}_\theta^d))_{\eta,q} = B_{p,q}^\alpha(\mathbb{T}_\theta^d), \quad \alpha_0 \neq \alpha_1, \alpha = (1 - \eta)\alpha_0 + \eta\alpha_1;$
- $(B_{p,q_0}^\alpha(\mathbb{T}_\theta^d), B_{p,q_1}^\alpha(\mathbb{T}_\theta^d))_{\eta,q} = B_{p,q}^\alpha(\mathbb{T}_\theta^d), \quad \frac{1}{q} = \frac{1 - \eta}{q_0} + \frac{\eta}{q_1};$
- $(B_{p_0,q_0}^{\alpha_0}(\mathbb{T}_\theta^d), B_{p_1,q_1}^{\alpha_1}(\mathbb{T}_\theta^d))_{\eta,q} = B_{p,q}^\alpha(\mathbb{T}_\theta^d), \quad \alpha_0 \neq \alpha_1, \alpha = (1 - \eta)\alpha_0 + \eta\alpha_1,$
 $\frac{1}{p} = \frac{1 - \eta}{p_0} + \frac{\eta}{p_1}, \frac{1}{q} = \frac{1 - \eta}{q_0} + \frac{\eta}{q_1}, p = q;$
- $(B_{p_0,q_0}^{\alpha_0}(\mathbb{T}_\theta^d), B_{p_1,q_1}^{\alpha_1}(\mathbb{T}_\theta^d))_\eta = B_{p,q}^\alpha(\mathbb{T}_\theta^d), \quad \alpha = (1 - \eta)\alpha_0 + \eta\alpha_1, \frac{1}{p} = \frac{1 - \eta}{p_0} + \frac{\eta}{p_1},$
 $\frac{1}{q} = \frac{1 - \eta}{q_0} + \frac{\eta}{q_1}, q < \infty;$
- $(F_\infty^{\alpha,c}(\mathbb{T}_\theta^d), F_1^{\alpha,c}(\mathbb{T}_\theta^d))_{\frac{1}{p}} = F_p^\alpha(\mathbb{T}_\theta^d) = (F_\infty^{\alpha,c}(\mathbb{T}_\theta^d), F_1^{\alpha,c}(\mathbb{T}_\theta^d))_{\frac{1}{p},p} \quad 1 < p < \infty;$
- $(F_p^{\alpha_0,c}(\mathbb{T}_\theta^d), F_p^{\alpha_1,c}(\mathbb{T}_\theta^d))_{\eta,q} = B_{p,q}^\alpha(\mathbb{T}_\theta^d), \quad 1 \leq p, q \leq \infty, \alpha_0 \neq \alpha_1,$
 $\alpha = (1 - \eta)\alpha_0 + \eta\alpha_1;$
- $(F_\infty^{\alpha_0,c}(\mathbb{T}_\theta^d), F_1^{\alpha_1,c}(\mathbb{T}_\theta^d))_{\frac{1}{p}} = F_p^{\alpha,c}(\mathbb{T}_\theta^d), \quad \alpha = (1 - \frac{1}{p})\alpha_0 + \frac{\alpha_1}{p}, 1 < p < \infty.$

Here the space $F_\infty^{\alpha,c}(\mathbb{T}_\theta^d)$ is defined to be the duality of $F_1^{-\alpha,c}(\mathbb{T}_\theta^d)$. Then by the reduction property of Triebel-Lizorkin space, we have that J^α is an isomorphism between $F_\infty^{\alpha,c}(\mathbb{T}_\theta^d)$ and $\text{BMO}^c(\mathbb{T}_\theta^d)$.

The really hard task concerns the interpolation of Sobolev spaces for which we have obtained only partial results. The most interesting couple is $(W_1^k(\mathbb{T}_\theta^d), W_\infty^k(\mathbb{T}_\theta^d))$. Recall that the complex interpolation of this couple remains always unsolved even in the commutative case (a well-known longstanding open problem which has been recorded by P. Jones in [27, p. 173]), while its real interpolation spaces were completely determined by DeVore and Scherer [21]. We do not know, unfortunately, how to prove the quantum analogue of DeVore and Scherer's theorem. However, we are able to extend to the quantum tori the K-functional formula of the couple $(L_p(\mathbb{R}^d), W_p^k(\mathbb{R}^d))$ obtained by Johnen and Scherer [30]. This result reads as follows:

Theorem 0.10. Let $1 \leq p \leq \infty$ and $k \in \mathbb{N}$. Then

$$K(x, \varepsilon^k; L_p(\mathbb{T}_\theta^d), W_p^k(\mathbb{T}_\theta^d)) \approx \varepsilon^k |\hat{x}(0)| + \omega_p^k(x, \varepsilon), \quad 0 < \varepsilon \leq 1$$

with relevant constants depending only on d and k .

As a consequence, we determine the real interpolation spaces of $(L_p(\mathbb{T}_\theta^d), W_p^k(\mathbb{T}_\theta^d))$, which are Besov spaces.

For the couple $(H_{p_0}^{\alpha_0}(\mathbb{T}_\theta^d), H_{p_1}^{\alpha_1}(\mathbb{T}_\theta^d))$, note that if $\alpha_0 = \alpha_1$ we easily have

$$(H_{p_0}^\alpha(\mathbb{T}_\theta^d), H_{p_1}^\alpha(\mathbb{T}_\theta^d))_\eta = H_p^\alpha(\mathbb{T}_\theta^d) \quad \text{and} \quad (H_{p_0}^\alpha(\mathbb{T}_\theta^d), H_{p_1}^\alpha(\mathbb{T}_\theta^d))_{\eta,p} = H_p^\alpha(\mathbb{T}_\theta^d),$$

for $\frac{1}{p} = \frac{1-\eta}{p_0} + \frac{\eta}{p_1}$, since J^α is an isometry between $H_p^\alpha(\mathbb{T}_\theta^d)$ and $L_p(\mathbb{T}_\theta^d)$ for all $1 \leq p \leq \infty$. If $\alpha_0 \neq \alpha_1$, the real interpolation of this couple is deduced from the corresponding one of Besov spaces. The result is

$$(H_{p_0}^{\alpha_0}(\mathbb{T}_\theta^d), H_{p_1}^{\alpha_1}(\mathbb{T}_\theta^d))_{\eta,q} = B_{p,q}^\alpha(\mathbb{T}_\theta^d), \quad \alpha = (1-\eta)\alpha_0 + \eta\alpha_1.$$

For the complex interpolation, our method is to consider the couple $(H_{\text{BMO}}^{\alpha_0}(\mathbb{T}_\theta^d), H_{\mathcal{H}_1}^{\alpha_1}(\mathbb{T}_\theta^d))$, where $H_{\mathcal{H}_1}^{\alpha_1}(\mathbb{T}_\theta^d)$ is defined as

$$H_{\mathcal{H}_1}^\alpha(\mathbb{T}_\theta^d) = \{x \in \mathcal{S}'(\mathbb{T}_\theta^d) : J^\alpha x \in \mathcal{H}_1(\mathbb{T}_\theta^d)\} \quad \text{with} \quad \|x\|_{H_{\mathcal{H}_1}^\alpha} = \|J^\alpha x\|_{\mathcal{H}_1},$$

and $H_{\text{BMO}}^\alpha(\mathbb{T}_\theta^d)$ is defined similarly, by the BMO -norm of $J^\alpha x$.

Then we have the following

Theorem 0.11. *Let $\alpha_0, \alpha_1 \in \mathbb{R}$ and $1 < p < \infty$. Then*

$$(H_{\text{BMO}}^{\alpha_0}(\mathbb{T}_\theta^d), H_{\mathcal{H}_1}^{\alpha_1}(\mathbb{T}_\theta^d))_{\frac{1}{p}} = H_p^\alpha(\mathbb{T}_\theta^d), \quad \alpha = (1 - \frac{1}{p})\alpha_0 + \frac{\alpha_1}{p}.$$

By the reiteration theorem, the above theorem implies the following:

Corollary 0.12. *Let $0 < \eta < 1$, $\alpha_0, \alpha_1 \in \mathbb{R}$ and $1 < p_0, p_1 < \infty$. Then*

$$(H_{p_0}^{\alpha_0}(\mathbb{T}_\theta^d), H_{p_1}^{\alpha_1}(\mathbb{T}_\theta^d))_\eta = H_p^\alpha(\mathbb{T}_\theta^d), \quad \alpha = (1-\eta)\alpha_0 + \eta\alpha_1, \quad \frac{1}{p} = \frac{1-\eta}{p_0} + \frac{\eta}{p_1}.$$

This corollary also answers partly the question about the complex interpolation of the couple $(W_{p_0}^{k_0}(\mathbb{T}_\theta^d), W_{p_1}^{k_1}(\mathbb{T}_\theta^d))$ for $1 < p_0 < p_1 < \infty$. To deduce the complex interpolation of the couple $(W_{p_0}^k(\mathbb{T}_\theta^d), W_{p_1}^k(\mathbb{T}_\theta^d))$ for $p_0 = 1$, we need to consider the Hardy Sobolev spaces $W_{\text{BMO}}^k(\mathbb{T}_\theta^d)$ and $W_{\mathcal{H}_1}^k(\mathbb{T}_\theta^d)$ instead of $W_\infty^k(\mathbb{T}_\theta^d)$ and $W_1^k(\mathbb{T}_\theta^d)$. They are defined by the norms:

$$\sup_{0 \leq |m|_1 \leq k} \|D^m x\|_{\text{BMO}} \quad \text{and} \quad \sum_{0 \leq |m|_1 \leq k} \|D^m x\|_{\mathcal{H}_1}$$

respectively. We can prove that for any $k \in \mathbb{N}$, $W_{\text{BMO}}^k(\mathbb{T}_\theta^d) = H_{\text{BMO}}^k(\mathbb{T}_\theta^d)$ and $W_{\mathcal{H}_1}^k(\mathbb{T}_\theta^d) = H_{\mathcal{H}_1}^k(\mathbb{T}_\theta^d)$. Then we obtain

Theorem 0.13. *Let $k \in \mathbb{N}$ and $1 < p < \infty$. Then for $X = W_{\mathcal{H}_1}^k(\mathbb{T}_\theta^d)$ or $X = W_1^k(\mathbb{T}_\theta^d)$,*

$$(W_{\text{BMO}}^k(\mathbb{T}_\theta^d), X)_{\frac{1}{p}} = W_p^k(\mathbb{T}_\theta^d) = (W_{\text{BMO}}^k(\mathbb{T}_\theta^d), X)_{\frac{1}{p},p}.$$

Consequently, for any $0 < \eta < 1$ and $1 < p_0 < \infty$,

$$(W_{p_0}^k(\mathbb{T}_\theta^d), W_1^k(\mathbb{T}_\theta^d))_\eta = W_p^k(\mathbb{T}_\theta^d) = (W_{p_0}^k(\mathbb{T}_\theta^d), W_1^k(\mathbb{T}_\theta^d))_{\eta,p}, \quad \frac{1}{p} = \frac{1-\eta}{p_0} + \eta.$$

0.5 Embeddings

The fourth family of results concern the embedding of the preceding spaces. A typical one is the analogue of the classical Sobolev embedding inequality for $W_p^k(\mathbb{T}_\theta^d)$.

Theorem 0.14. *Assume that $1 < p < \infty$. If $\alpha \in \mathbb{R}^+$, $\alpha p < d$ and $\frac{1}{p_1} = \frac{1}{p} - \frac{\alpha}{d}$, then we have:*

$$H_p^\alpha(\mathbb{T}_\theta^d) \subset L_{p_1}(\mathbb{T}_\theta^d).$$

Similarly, if $k \in \mathbb{N}$, $kp < d$ and $\frac{1}{p_1} = \frac{1}{p} - \frac{k}{d}$, then we have:

$$W_p^k(\mathbb{T}_\theta^d) \subset L_{p_1}(\mathbb{T}_\theta^d).$$

On the other hand, if $\alpha p > d$ (or $kp > d$), Sobolev spaces are embedded into Hölder-Zygmund spaces, which also coincides well with the classic Sobolev embedding theorems.

Theorem 0.15. *Assume that $1 \leq p < \infty$ and $\alpha p > d$. Then for $\alpha_1 = \alpha - \frac{d}{p}$, we have*

$$H_p^\alpha(\mathbb{T}_\theta^d) \subset B_{\infty,\infty}^{\alpha_1}(\mathbb{T}_\theta^d).$$

Similarly, if $1 < p < \infty$ and $k \in \mathbb{N}$, $kp > d$, then we have for $\alpha_1 = k - \frac{d}{p}$,

$$W_p^k(\mathbb{T}_\theta^d) \subset B_{\infty,\infty}^{\alpha_1}(\mathbb{T}_\theta^d).$$

For Besov spaces, we have

Theorem 0.16. *Assume that $1 \leq p \leq p_1 \leq \infty$, $1 \leq q \leq q_1 \leq \infty$, and $\alpha - \frac{d}{p} = \alpha_1 - \frac{d}{p_1}$. Then we have the following bounded inclusion:*

$$B_{p,q}^\alpha(\mathbb{T}_\theta^d) \subset B_{p_1,q_1}^{\alpha_1}(\mathbb{T}_\theta^d).$$

Similar embedding inequalities hold for Triebel-Lizorkin spaces too, since $F_p^\alpha(\mathbb{T}_\theta^d) = H_p^\alpha(\mathbb{T}_\theta^d)$ with equivalent norms for $1 < p < \infty$. Combined with real interpolation, the embedding inequality of $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ yields that of Sobolev spaces. Our proofs of these embedding inequalities are based on Varopoulos's celebrated semigroup approach [76] to the Littlewood-Sobolev theory, which were already employed by Junge and Mei [34] in their study of BMO spaces on quantum Markovian semigroups.

Moreover, the above three theorems have compact versions if we reduce a little the indices p_1 of $L_{p_1}(\mathbb{T}_\theta^d)$ and $B_{p_1,q_1}^{\alpha_1}(\mathbb{T}_\theta^d)$, or α_1 of $B_{\infty,\infty}^{\alpha_1}(\mathbb{T}_\theta^d)$. The one of Sobolev embedding extends the Rellich-Kondrachof compactness theorem in classic Sobolev embedding theory.

Theorem 0.17. (i) *Assume that $1 \leq p < p_1 \leq \infty$, $1 \leq q \leq q_1 \leq \infty$ and $\alpha - \frac{d}{p} = \alpha_1 - \frac{d}{p_1}$. Then the embedding $B_{p,q}^\alpha(\mathbb{T}_\theta^d) \hookrightarrow B_{p^*,q_1}^{\alpha_1}(\mathbb{T}_\theta^d)$ is compact for $1 \leq p^* < p_1$.*

(ii) *If $p > 1$, $\alpha - \frac{d}{p} = \alpha_1 - \frac{d}{p_1}$, then $H_p^\alpha(\mathbb{T}_\theta^d) \hookrightarrow H_{p^*}^{\alpha_1}(\mathbb{T}_\theta^d)$ is compact for $1 \leq p^* < p_1$. In particular, if additionally $\alpha = k$ and $\alpha_1 = k_1$ are nonnegative integers, then $W_p^k(\mathbb{T}_\theta^d) \hookrightarrow W_{p^*}^{k_1}(\mathbb{T}_\theta^d)$ is compact.*

(iii) *If $p > 1$, $p(\alpha - \alpha_1) > d$ and $\alpha^* < \alpha_1 = \alpha - \frac{d}{p}$, then $H_p^\alpha(\mathbb{T}_\theta^d) \hookrightarrow B_{\infty,\infty}^{\alpha^*}(\mathbb{T}_\theta^d)$ is compact. In particular, if additionally $\alpha = k \in \mathbb{N}$, then $W_p^k(\mathbb{T}_\theta^d) \hookrightarrow B_{\infty,\infty}^{\alpha^*}(\mathbb{T}_\theta^d)$ is compact.*

0.6 Multipliers

The last family of results of the thesis describe Fourier multipliers on the preceding spaces. Like in the L_p case treated in [17], we are mainly concerned with completely bounded Fourier multipliers. All spaces in consideration carry a natural operator space structure in Pisier's sense [54]. Inspired by Neuwirth and Ricard's transference theorem [48], Fourier multipliers on $L_p(\mathbb{T}_\theta^d)$ are related with Schur multipliers on Schatten Class S_p in [17]. Given a distribution x on $\mathbb{T}_\theta^d$, we write its matrix in the basis $(U^m)_{m \in \mathbb{Z}^d}$:

$$[x] = \left(\langle x U^n, U^m \rangle \right)_{m,n \in \mathbb{Z}^d} = \left(\hat{x}(m-n) e^{i n \tilde{\theta} (m-n)^t} \right)_{m,n \in \mathbb{Z}^d}.$$

Here k^t denotes the transpose of $k = (k_1, \dots, k_d)$ and $\tilde{\theta}$ is the following $d \times d$ -matrix deduced from the skew symmetric matrix θ :

$$\tilde{\theta} = -2\pi \begin{pmatrix} 0 & \theta_{12} & \theta_{13} & \dots & \theta_{1d} \\ 0 & 0 & \theta_{23} & \dots & \theta_{2d} \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & \dots & \theta_{d-1,d} \\ 0 & 0 & 0 & \dots & 0 \end{pmatrix}.$$

Now let $\phi : \mathbb{Z}^d \rightarrow \mathbb{C}$ and M_ϕ be the associated Fourier multiplier on $\mathbb{T}_\theta^d$. Set $\check{\phi} = (\phi_{m-n})_{m,n \in \mathbb{Z}^d}$. Then

$$[M_\phi x] = (\phi_{m-n} \hat{x}(m-n) e^{i n \tilde{\theta} (m-n)^t})_{m,n \in \mathbb{Z}^d} = S_{\check{\phi}}([x]), \quad (0.1)$$

where $S_{\check{\phi}}$ the Schur multiplier with symbol $\check{\phi}$. If X is a Banach space of distributions on $\mathbb{T}_\theta^d$, we denote by $\mathbf{M}(X)$ the space of bounded Fourier multipliers on X ; if X is further equipped with an operator space structure, $\mathbf{M}_{\text{cb}}(X)$ is the space of c.b. Fourier multipliers on X . Similarly, if X is a Schatten L_p space (or its closed subspace), we denote by $\mathbf{M}(X)$ the space of bounded Schur multipliers on X ; with its natural operator space structure, we denote also $\mathbf{M}_{\text{cb}}(X)$ the space of c.b. Schur multipliers on X . These spaces of multipliers are endowed with their natural norms. It is shown in [17] that, for $1 \leq p \leq \infty$, $\mathbf{M}_{\text{cb}}(L_p(\mathbb{T}_\theta^d)) = \mathbf{M}_{\text{cb}}(S_p)$ with equal norms, whence $\mathbf{M}_{\text{cb}}(L_p(\mathbb{T}_\theta^d))$ is independent of θ .

Using again Neuwirth-Ricard's transference between Fourier multipliers and Schur multipliers, we deduce similar result for Sobolev space $W_p^k(\mathbb{T}_\theta^d)$ and Triebel-Lizorkin spaces $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$.

Theorem 0.18. *Let $1 \leq p \leq \infty$, and $k \in \mathbb{N}$, $\alpha \in \mathbb{R}$. Then*

$$\mathbf{M}_{\text{cb}}(W_p^k(\mathbb{T}_\theta^d)) = \mathbf{M}_{\text{cb}}(W_p^k(\mathbb{T}^d)) \text{ with equal norms,}$$

and

$$\mathbf{M}_{\text{cb}}(F_p^{\alpha,c}(\mathbb{T}_\theta^d)) = \mathbf{M}_{\text{cb}}(F_p^{\alpha,c}(\mathbb{T}^d)) \text{ with equal norms.}$$

Here we should point out that, even though it is a commutative function space, the space $F_p^\alpha(\mathbb{T}^d)$ (corresponding to $\theta = 0$) is endowed with three different operator space structures, the first two being defined by its embedding into $L_p(\mathbb{T}^d; \ell_2^c)$ and $L_p(\mathbb{T}^d; \ell_2^r)$ and the third one being the mixture of these two, resulting three different operator spaces denoted by $F_p^{\alpha,c}(\mathbb{T}^d)$, $F_p^{\alpha,r}(\mathbb{T}^d)$ and $F_p^\alpha(\mathbb{T}^d)$, respectively. The above theorem gives in particular that $\mathbf{M}_{\text{cb}}(\mathcal{H}_p^c(\mathbb{T}_\theta^d)) = \mathbf{M}_{\text{cb}}(\mathcal{H}_p^c(\mathbb{T}^d))$ with equal norms.

The situation for Besov spaces is very satisfactory since it is well known that Fourier multipliers behave much better on Besov spaces than on L_p -spaces (in the commutative case). We prove the following

Theorem 0.19. *Let $\alpha \in \mathbb{R}$ and $1 \leq p, q \leq \infty$. Let $\phi : \mathbb{Z}^d \rightarrow \mathbb{C}$. Then ϕ is a Fourier multiplier on $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ iff the $\phi\varphi^{(k)}$'s are Fourier multipliers on $L_p(\mathbb{T}_\theta^d)$ uniformly in k . In this case, we have*

$$\|\phi\|_{\mathbf{M}(B_{p,q}^\alpha(\mathbb{T}_\theta^d))} \approx |\phi(0)| + \sup_{k \geq 0} \|\phi\varphi^{(k)}\|_{\mathbf{M}(L_p(\mathbb{T}_\theta^d))}$$

with relevant constants depending only on α . A similar c.b. version holds too.

Consequently, the Fourier multipliers on $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ are completely determined by the Fourier multipliers on $L_p(\mathbb{T}_\theta^d)$ associated to their components in the Littlewood-Paley decomposition. So the completely bounded multipliers on $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ depend solely on p . In the case of $p = 1$, a multiplier is bounded on $B_{1,q}^\alpha(\mathbb{T}_\theta^d)$ iff it is completely bounded iff it is the Fourier transform of an element of $B_{1,\infty}^0(\mathbb{T}^d)$. Using a classical example of Stein-Zygmund [70], we show that there exists a ϕ which is a completely bounded Fourier multiplier on $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ for all p but bounded on $L_p(\mathbb{T}_\theta^d)$ for no $p \neq 2$.

Chapter 1

Preliminaries

This chapter collects the necessary preliminaries for the whole paper. The first two sections present the definitions and some basic facts about noncommutative L_p -spaces and quantum tori which are the central objects of the paper. The third one contains some results on Fourier multipliers that will play a paramount role in the whole paper. The last section gives the definitions and some fundamental results on operator-valued Hardy spaces on the usual and quantum tori. This section will be needed only starting from chapter 4 on Triebel-Lizorkin spaces.

1.1 Noncommutative L_p -spaces

Let $\mathcal{M}$ be a von Neumann algebra equipped with a normal semifinite faithful trace τ and $S_{\mathcal{M}}^+$ be the set of all positive elements x in $\mathcal{M}$ with $\tau(s(x)) < \infty$, where $s(x)$ denotes the support of x , i.e., the smallest projection e such that $exe = x$. Let $S_{\mathcal{M}}$ be the linear span of $S_{\mathcal{M}}^+$. Then every $x \in S_{\mathcal{M}}$ has finite trace, and $S_{\mathcal{M}}$ is a w^* -dense $*$ -subalgebra of $\mathcal{M}$.

Let $0 < p < \infty$. For any $x \in S_{\mathcal{M}}$, the operator $|x|^p$ belongs to $S_{\mathcal{M}}^+$ (recalling $|x| = (x^*x)^{\frac{1}{2}}$). We define

$$\|x\|_p = (\tau(|x|^p))^{\frac{1}{p}}.$$

One can check that $\|\cdot\|_p$ is a norm on $S_{\mathcal{M}}$. The completion of $(S_{\mathcal{M}}, \|\cdot\|_p)$ is denoted by $L_p(\mathcal{M})$, which is the usual noncommutative L_p -space associated to $(\mathcal{M}, \tau)$. For convenience, we set $L_{\infty}(\mathcal{M}) = \mathcal{M}$ equipped with the operator norm $\|\cdot\|_{\mathcal{M}}$. The norm of $L_p(\mathcal{M})$ will be often denoted simply by $\|\cdot\|_p$. But if different L_p -spaces appear in a same context, we will sometimes precise their norms in order to avoid possible ambiguity. The reader is referred to [57] and [83] for more information on noncommutative L_p -spaces.

The elements of $L_p(\mathcal{M})$ can be described as closed densely defined operators on H (H being the Hilbert space on which $\mathcal{M}$ acts). A closed densely defined operator x on H is said to be affiliated with $\mathcal{M}$ if $ux = xu$ for any unitary u in the commutant $\mathcal{M}'$ of $\mathcal{M}$. An operator x affiliated with $\mathcal{M}$ is said to be measurable with respect to $(\mathcal{M}, \tau)$ (or simply measurable) if for any $\delta > 0$ there exists a projection $e \in B(H)$ such that

$$e(H) \subset \text{Dom}(x) \quad \text{and} \quad \tau(e^{\perp}) \leq \delta,$$

where $\text{Dom}(x)$ defines the domain of x . We denote by $L_0(\mathcal{M}, \tau)$, or simply $L_0(\mathcal{M})$ the family of all measurable operators. For such an operator x , we define

$$\lambda_s(x) = \tau(e_s^{\perp}(|x|)), \quad s > 0$$

where $e_s^\perp(x) = \mathbb{1}_{(s,\infty)}(x)$ is the spectrum projection of x corresponding to the interval (s, ∞) , and

$$\mu_t(x) = \inf\{s > 0 : \lambda_s(x) < t\}, \quad t > 0.$$

The function $s \mapsto \lambda_s(x)$ is called the distribution function of x and the $\mu_t(x)$ the generalized singular numbers of x . Similarly to the classical case, for $0 < p < \infty, 0 < q \leq \infty$, the noncommutative Lorentz space $L_{p,q}(\mathcal{M})$ is defined to be the collection of all measurable operators x such that

$$\|x\|_{p,q} = \left(\int_0^\infty (t^{\frac{1}{p}} \mu_t(x))^q \frac{dt}{t} \right)^{\frac{1}{q}} < \infty.$$

Clearly, $L_{p,p}(\mathcal{M}) = L_p(\mathcal{M})$. The space $L_{p,\infty}(\mathcal{M})$ is usually called a weak L_p -space, $0 < p < \infty$, and

$$\|x\|_{p,\infty} = \sup_{s>0} s \lambda_s(x)^{\frac{1}{p}}.$$

Like the classical L_p -spaces, noncommutative L_p -spaces behave well with respect to interpolation. Our reference for interpolation theory is [8]. Let $1 \leq p_0 < p_1 \leq \infty$, $1 \leq q \leq \infty$ and $0 < \eta < 1$. Then

$$(L_{p_0}(\mathcal{M}), L_{p_1}(\mathcal{M}))_\eta = L_p(\mathcal{M}) \quad \text{and} \quad (L_{p_0}(\mathcal{M}), L_{p_1}(\mathcal{M}))_{\eta,q} = L_{p,q}(\mathcal{M}), \quad (1.1)$$

where $\frac{1}{p} = \frac{1-\eta}{p_0} + \frac{\eta}{p_1}$.

Now we introduce noncommutative Hilbert space-valued L_p -spaces $L_p(\mathcal{M}; H^c)$ and $L_p(\mathcal{M}; H^r)$, which are studied at length in [32]. Let H be a Hilbert space and v a norm one element of H . Let p_v be the orthogonal projection onto the one-dimensional subspace generated by v . Then define the following row and column noncommutative L_p -spaces:

$$\begin{aligned} L_p(\mathcal{M}; H^r) &= (p_v \otimes 1_{\mathcal{M}}) L_p(B(H) \overline{\otimes} \mathcal{M}), \\ L_p(\mathcal{M}; H^c) &= L_p(B(H) \overline{\otimes} \mathcal{M}) (p_v \otimes 1_{\mathcal{M}}), \end{aligned}$$

where the tensor product $B(H) \overline{\otimes} \mathcal{M}$ is equipped with the tensor trace while $B(H)$ is equipped with the usual trace. For $f \in L_p(\mathcal{M}; H^c)$,

$$\|f\|_{L_p(\mathcal{M}; H^c)} = \|(f^* f)^{\frac{1}{2}}\|_{L_p(\mathcal{M})}.$$

A similar formula holds for the row space by passing to adjoints: $f \in L_p(\mathcal{M}; H^r)$ iff $f^* \in L_p(\mathcal{M}; H^c)$, and $\|f\|_{L_p(\mathcal{M}; H^r)} = \|f^*\|_{L_p(\mathcal{M}; H^c)}$. It is clear that $L_p(\mathcal{M}; H^c)$ and $L_p(\mathcal{M}; H^r)$ are 1-complemented subspaces of $L_p(B(H) \overline{\otimes} \mathcal{M})$ for any p . Thus they also form interpolation scales with respect to both complex and real interpolation methods: Let $1 \leq p_0, p_1 \leq \infty$ and $0 < \eta < 1$. Then

$$\begin{aligned} (L_{p_0}(\mathcal{M}; H^c), L_{p_1}(\mathcal{M}; H^c))_\eta &= L_p(\mathcal{M}; H^c), \\ (L_{p_0}(\mathcal{M}; H^c), L_{p_1}(\mathcal{M}; H^c))_{\eta,p} &= L_{p,q}(\mathcal{M}; H^c), \end{aligned} \quad (1.2)$$

where $\frac{1}{p} = \frac{1-\eta}{p_0} + \frac{\eta}{p_1}$. The same formulas hold for row spaces too.

1.2 Quantum tori

Let $d \geq 2$ and $\theta = (\theta_{kj})$ be a real skew symmetric $d \times d$ -matrix. The associated d -dimensional noncommutative torus $\mathcal{A}_\theta$ is the universal C^* -algebra generated by d unitary operators $U_1, \dots, U_d$ satisfying the following commutation relation

$$U_k U_j = e^{2\pi i \theta_{kj}} U_j U_k, \quad j, k = 1, \dots, d. \quad (1.3)$$

We will use standard notation from multiple Fourier series. Let $U = (U_1, \dots, U_d)$. For $m = (m_1, \dots, m_d) \in \mathbb{Z}^d$ we define

$$U^m = U_1^{m_1} \dots U_d^{m_d}.$$

A polynomial in U is a finite sum

$$x = \sum_{m \in \mathbb{Z}^d} \alpha_m U^m \quad \text{with} \quad \alpha_m \in \mathbb{C},$$

that is, $\alpha_m = 0$ for all but finite indices $m \in \mathbb{Z}^d$. The involution algebra $\mathcal{P}_\theta$ of all such polynomials is dense in $\mathcal{A}_\theta$. For any polynomial x as above we define

$$\tau(x) = \alpha_0,$$

where $0 = (0, \dots, 0)$. Then, τ extends to a faithful tracial state on $\mathcal{A}_\theta$. Let $\mathbb{T}_\theta^d$ be the w^* -closure of $\mathcal{A}_\theta$ in the GNS representation of τ . This is our d -dimensional quantum torus. The state τ extends to a normal faithful tracial state on $\mathbb{T}_\theta^d$ that will be denoted again by τ . Recall that the von Neumann algebra $\mathbb{T}_\theta^d$ is hyperfinite.

Any $x \in L_1(\mathbb{T}_\theta^d)$ admits a formal Fourier series:

$$x \sim \sum_{m \in \mathbb{Z}^d} \hat{x}(m) U^m,$$

where

$$\hat{x}(m) = \tau((U^m)^* x), \quad m \in \mathbb{Z}^d$$

are the Fourier coefficients of x . The operator x is, of course, uniquely determined by its Fourier series.

We introduced in [17] a transference method to overcome the full noncommutativity of quantum tori and use methods of operator-valued harmonic analysis. Let $\mathbb{T}^d$ be the usual d -torus equipped with normalized Haar measure dz . Let $\mathcal{N}_\theta = L_\infty(\mathbb{T}^d) \otimes \mathbb{T}_\theta^d$, equipped with the tensor trace $\nu = \int dz \otimes \tau$. It is well known that for every $0 < p < \infty$,

$$L_p(\mathcal{N}_\theta, \nu) \cong L_p(\mathbb{T}^d; L_p(\mathbb{T}_\theta^d)).$$

The space on the right-hand side is the space of Bochner p -integrable functions from $\mathbb{T}^d$ to $L_p(\mathbb{T}_\theta^d)$. In general, for any Banach space X and any measure space (Ω, μ) , we use $L_p(\Omega; X)$ to denote the space of Bochner p -integrable functions from Ω to X . For each $z \in \mathbb{T}^d$, define π_z to be the isomorphism of $\mathbb{T}_\theta^d$ determined by

$$\pi_z(U^m) = z^m U^m = z_1^{m_1} \dots z_d^{m_d} U_1^{m_1} \dots U_d^{m_d}. \quad (1.4)$$

Since $\tau(\pi_z(x)) = \tau(x)$ for any $x \in \mathbb{T}_\theta^d$, π_z preserves the trace τ . Thus for every $0 < p < \infty$,

$$\|\pi_z(x)\|_p = \|x\|_p, \quad \forall x \in L_p(\mathbb{T}_\theta^d). \quad (1.5)$$

Now we state the transference method as follows (see [17]).

Lemma 1.1. *For any $x \in L_p(\mathbb{T}_\theta^d)$, the function $\tilde{x} : z \mapsto \pi_z(x)$ is continuous from $\mathbb{T}^d$ to $L_p(\mathbb{T}_\theta^d)$ (with respect to the w^* -topology for $p = \infty$). If $x \in \mathcal{A}_\theta$, it is continuous from $\mathbb{T}^d$ to $\mathcal{A}_\theta$.*

Corollary 1.2. (i) Let $0 < p \leq \infty$. If $x \in L_p(\mathbb{T}_\theta^d)$, then $\tilde{x} \in L_p(\mathcal{N}_\theta)$ and $\|\tilde{x}\|_p = \|x\|_p$, that is, $x \mapsto \tilde{x}$ is an isometric embedding from $L_p(\mathbb{T}_\theta^d)$ into $L_p(\mathcal{N}_\theta)$. Moreover, this map is also an isomorphism from $\mathcal{A}_\theta$ into $C(\mathbb{T}^d; \mathcal{A}_\theta)$.

(ii) Let $\widetilde{\mathbb{T}_\theta^d} = \{\tilde{x} : x \in \mathbb{T}_\theta^d\}$. Then $\widetilde{\mathbb{T}_\theta^d}$ is a von Neumann subalgebra of $\mathcal{N}_\theta$ and the associated conditional expectation is given by

$$\mathbb{E}(f)(z) = \pi_z \left(\int_{\mathbb{T}^d} \pi_{\bar{w}}[f(w)] dw \right), \quad z \in \mathbb{T}^d, f \in \mathcal{N}_\theta.$$

Moreover, $\mathbb{E}$ extends to a contractive projection from $L_p(\mathcal{N}_\theta)$ onto $L_p(\widetilde{\mathbb{T}_\theta^d})$ for $1 \leq p \leq \infty$.

(iii) $L_p(\mathbb{T}_\theta^d)$ is isometric to $L_p(\widetilde{\mathbb{T}_\theta^d})$ for every $0 < p \leq \infty$.

1.3 Fourier multipliers

Fourier multipliers will be the most important tool for the whole work. Now we present some known results on them for later use. Given a function $\phi : \mathbb{Z}^d \rightarrow \mathbb{C}$, let M_ϕ denote the associated Fourier multiplier on $\mathbb{T}^d$, namely, $\widehat{M_\phi f}(m) = \phi(m) \widehat{f}(m)$ for any trigonometric polynomial f on $\mathbb{T}^d$. We call ϕ a multiplier on $L_p(\mathbb{T}^d)$ if M_ϕ extends to a bounded map on $L_p(\mathbb{T}^d)$. Fourier multipliers on $\mathbb{T}_\theta^d$ are defined exactly in the same way, we still use the same symbol M_ϕ to denote the corresponding multiplier on $\mathbb{T}_\theta^d$ too. Note that the isomorphism π_z defined in (1.4) is the Fourier multiplier associated to the function ϕ given by $\phi(m) = z^m$.

It is natural to ask if the boundedness of M_ϕ on $L_p(\mathbb{T}^d)$ is equivalent to that on $L_p(\mathbb{T}_\theta^d)$. This is still an open problem. However, it is proved in [17] that the answer is affirmative if “boundedness” is replaced by “complete boundedness”, a notion from operator space theory for which we refer to [23] and [55]. All noncommutative L_p -spaces are equipped with their natural operator space structure introduced by Pisier [54, 55].

We will use the following fundamental property of completely bounded (c.b. for short) maps due to Pisier [54]. Let E and F be operator spaces. Then a linear map $T : E \rightarrow F$ is c.b. iff $\text{Id}_{S_p} \otimes T : S_p[E] \rightarrow S_p[F]$ is bounded for some $1 \leq p \leq \infty$. In this case,

$$\|T\|_{\text{cb}} = \|\text{Id}_{S_p} \otimes T : S_p[E] \rightarrow S_p[F]\|.$$

Here $S_p[E]$ denotes the E -valued Schatten p -class. In particular, if $E = \mathbb{C}$, $S_p[\mathbb{C}] = S_p$ is the noncommutative L_p -space associated to $B(\ell_2)$, equipped with the usual trace. Applying this criterion to the special case where $E = F = L_p(\mathcal{M})$, we see that a map T on $L_p(\mathcal{M})$ is c.b. iff $\text{Id}_{S_p} \otimes T : L_p(B(\ell_2) \bar{\otimes} \mathcal{M}) \rightarrow L_p(B(\ell_2) \bar{\otimes} \mathcal{M})$ is bounded. The readers unfamiliar with operator space theory can take this property as the definition of c.b. maps between L_p -spaces.

Thus ϕ is a c.b. multiplier on $L_p(\mathbb{T}_\theta^d)$ if M_ϕ is c.b. on $L_p(\mathbb{T}_\theta^d)$, or equivalently, if $\text{Id}_{S_p} \otimes M_\phi$ is bounded on $L_p(B(\ell_2) \bar{\otimes} \mathbb{T}_\theta^d)$. Let $\mathbf{M}(L_p(\mathbb{T}_\theta^d))$ (resp. $\mathbf{M}_{\text{cb}}(L_p(\mathbb{T}_\theta^d))$) denote the space of Fourier multipliers (resp. c.b. Fourier multipliers) on $L_p(\mathbb{T}_\theta^d)$, equipped with the natural norm. When $\theta = 0$, we recover the (c.b.) Fourier multipliers on the usual d -torus $\mathbb{T}^d$. The corresponding multiplier spaces are denoted by $\mathbf{M}(L_p(\mathbb{T}^d))$ and $\mathbf{M}_{\text{cb}}(L_p(\mathbb{T}^d))$. Note that in the latter case ($\theta = 0$), $L_p(B(\ell_2) \bar{\otimes} \mathbb{T}_\theta^d) = L_p(\mathbb{T}^d; S_p)$, thus ϕ is a c.b. multiplier on $L_p(\mathbb{T}^d)$ iff M_ϕ extends to a bounded map on $L_p(\mathbb{T}^d; S_p)$.

The following result is taken from [17].

Lemma 1.3. *Let $1 \leq p \leq \infty$. Then $M_{cb}(L_p(\mathbb{T}_\theta^d)) = M_{cb}(L_p(\mathbb{T}^d))$ with equal norms.*

Remark 1.4. Note that $M_{cb}(L_1(\mathbb{T}^d)) = M_{cb}(L_\infty(\mathbb{T}^d))$ coincides with the space of the Fourier transforms of bounded measures on $\mathbb{T}^d$, and $M_{cb}(L_2(\mathbb{T}^d))$ with the space of bounded functions on $\mathbb{Z}^d$.

The most efficient criterion for Fourier multipliers on $L_p(\mathbb{T}^d)$ for $1 < p < \infty$ is Mikhlin's condition. Although it can be formulated in the periodic case, it is more convenient to state this condition in the case of $\mathbb{R}^d$. On the other hand, the Fourier multipliers on $\mathbb{T}^d$ used later will be the restrictions to $\mathbb{Z}^d$ of continuous Fourier multipliers on $\mathbb{R}^d$. As usual, for $m = (m_1, \dots, m_d) \in \mathbb{N}_0^d$ (recalling that $\mathbb{N}_0$ denotes the set of nonnegative integers), we set

$$D^m = \partial_1^{m_1} \cdots \partial_d^{m_d},$$

where ∂_k denotes the k th partial derivation on $\mathbb{R}^d$. Also put $|m|_1 = m_1 + \cdots + m_d$. The Euclidean norm of $\mathbb{R}^d$ is denoted by $|\cdot|$: $|\xi| = \sqrt{\xi_1^2 + \cdots + \xi_d^2}$.

Definition 1.5. A function $\phi : \mathbb{R}^d \rightarrow \mathbb{C}$ is called a Mikhlin multiplier if it is d -times differentiable on $\mathbb{R}^d \setminus \{0\}$ and satisfies the following condition

$$\|\phi\|_M = \sup \{ |\xi|^{m_1} |D^m \phi(\xi)| : \xi \in \mathbb{R}^d \setminus \{0\}, m \in \mathbb{N}_0^d, |m|_1 \leq d \} < \infty.$$

Note that the usual Mikhlin condition requires only partial derivatives up to order $\lfloor \frac{d}{2} \rfloor + 1$ (see, for instance, [25, section II.6] or [67, Theorem 4.3.2]). Our requirement above up to order d is imposed by the boundedness of these multipliers on UMD spaces. We refer to section 4.1 for the usual Mikhlin condition and more multiplier results on $\mathbb{T}_\theta^d$.

It is a classical result that every Mikhlin multiplier is a Fourier multiplier on $L_p(\mathbb{R}^d)$ for $1 < p < \infty$ (cf. [25, section II.6] or [67, Theorem 4.3.2]), so its restriction $\phi|_{\mathbb{Z}^d}$ is a Fourier multiplier on $L_p(\mathbb{T}^d)$ too. It is, however, less classical that such a multiplier is also c.b. on $L_p(\mathbb{R}^d)$ or $L_p(\mathbb{T}^d)$. This follows from a general result on UMD spaces. Recall that a Banach space X is called a UMD space if the X -valued martingale differences are unconditional in $L_p(\Omega; X)$ for any $1 < p < \infty$ and any probability space (Ω, P) . This is equivalent to the requirement that the Hilbert transform be bounded on $L_p(\mathbb{R}^d; X)$ for $1 < p < \infty$. Any noncommutative L_p -space with $1 < p < \infty$ is a UMD space. We refer to [10, 15, 16] for more information.

Before proceeding further, we make a convention used throughout the paper:

Convention. To simplify the notational system, we will use the same derivation symbols for $\mathbb{R}^d$ and $\mathbb{T}^d$. Thus for a multi-index $m \in \mathbb{N}_0^d$, $D^m = \partial_1^{m_1} \cdots \partial_d^{m_d}$, introduced previously, will also denote the partial derivation of order m on $\mathbb{T}^d$. Similarly, $\Delta = \partial_1^2 + \cdots + \partial_d^2$ will denote the Laplacian on both $\mathbb{R}^d$ and $\mathbb{T}^d$. In the same spirit, for a function ϕ on $\mathbb{R}^d$, we will call ϕ rather than $\phi|_{\mathbb{Z}^d}$ a Fourier multiplier on $L_p(\mathbb{T}^d)$ or $L_p(\mathbb{T}_\theta^d)$. This should not cause any ambiguity in concrete contexts. Considered as a map on $L_p(\mathbb{T}^d)$ or $L_p(\mathbb{T}_\theta^d)$, M_ϕ will be often denoted by $f \mapsto \tilde{\phi} * f$ or $x \mapsto \tilde{\phi} * x$.

Note that the notation $\tilde{\phi} * f$ coincides with the usual convolution when ϕ is good enough. Indeed, let $\tilde{\phi}$ be the 1-periodization of the inverse Fourier transform of ϕ whenever it exists in a reasonable sense:

$$\tilde{\phi}(s) = \sum_{m \in \mathbb{Z}^d} \mathcal{F}^{-1}(\phi)(s + m), \quad s \in \mathbb{R}^d.$$

Viewed as a function on $\mathbb{T}^d$, $\tilde{\phi}$ admits the following Fourier series:

$$\tilde{\phi}(z) = \sum_{m \in \mathbb{Z}^d} \phi(m) z^m.$$

Thus for any trigonometric polynomial f ,

$$\tilde{\phi} * f(z) = \int_{\mathbb{T}^d} \tilde{\phi}(zw^{-1}) f(w) dw, \quad z \in \mathbb{T}^d.$$

The following lemma is proved in [43, 85] (see also [12] for the one-dimensional case).

Lemma 1.6. *Let X be a UMD space and $1 < p < \infty$. Let ϕ be a Mikhlin multiplier. Then ϕ is a Fourier multiplier on $L_p(\mathbb{T}^d; X)$. Moreover, its norm is controlled by $\|\phi\|_M$, p and the UMD constant of X .*

Since S_p is a UMD space for $1 < p < \infty$, combining Lemmas 1.3 and 1.6 and Remark 1.4, we obtain the following result.

Lemma 1.7. *Let ϕ be a function on $\mathbb{R}^d$.*

- (i) *If $\mathcal{F}^{-1}(\phi)$ is integrable on $\mathbb{R}^d$, then ϕ is a c.b. Fourier multiplier on $L_p(\mathbb{T}_\theta^d)$ for $1 \leq p \leq \infty$. Moreover, its c.b. norm is not greater than $\|\mathcal{F}^{-1}(\phi)\|_1$.*
- (ii) *If ϕ is a Mikhlin multiplier, then ϕ is a c.b. Fourier multiplier on $L_p(\mathbb{T}_\theta^d)$ for $1 < p < \infty$. Moreover, its c.b. norm is controlled by $\|\phi\|_M$ and p .*

1.4 Hardy spaces

We now present some preliminaries on operator-valued Hardy spaces on $\mathbb{T}^d$ and Hardy spaces on $\mathbb{T}_\theta^d$. Motivated by the developments of noncommutative martingale inequalities in [56, 36] and quantum Markovian semigroups in [32], Mei [44] developed the theory of operator-valued Hardy spaces on $\mathbb{R}^d$. More recently, Mei's work was extended to the torus case in [17] with the objective of developing the Hardy space theory in the quantum torus setting. We now recall the definitions and results that will be needed later. Throughout this section, $\mathcal{M}$ will denote a von Neumann algebra equipped with a normal faithful tracial state τ and $\mathcal{N} = L_\infty(\mathbb{T}^d) \bar{\otimes} \mathcal{M}$ with the tensor trace. In our future applications, $\mathcal{M}$ will be $\mathbb{T}_\theta^d$.

A cube of $\mathbb{T}^d$ is a product $Q = I_1 \times \cdots \times I_d$, where each I_j is an interval (= arc) of $\mathbb{T}$. As in the Euclidean case, we use $|Q|$ to denote the normalized volume (= measure) of Q . The whole $\mathbb{T}^d$ is now a cube too (of volume 1).

We will often identify $\mathbb{T}^d$ with the unit cube $\mathbb{I}^d = [0, 1)^d$ via $(e^{2\pi i s_1}, \dots, e^{2\pi i s_d}) \leftrightarrow (s_1, \dots, s_d)$. Under this identification, the addition in $\mathbb{I}^d$ is the usual addition modulo 1 coordinatewise; an interval of $\mathbb{I}$ is either a subinterval of $\mathbb{I}$ or a union $[b, 1] \cup [0, a]$ with $0 < a < b < 1$, the latter union being the interval $[b-1, a]$ of $\mathbb{I}$ (modulo 1). So the cubes of $\mathbb{I}^d$ are exactly those of $\mathbb{T}^d$. Accordingly, functions on $\mathbb{T}^d$ and $\mathbb{I}^d$ are identified too; they are considered as 1-periodic functions on $\mathbb{R}^d$. Thus $\mathcal{N} = L_\infty(\mathbb{T}^d) \bar{\otimes} \mathcal{M} = L_\infty(\mathbb{I}^d) \bar{\otimes} \mathcal{M}$.

We define $\text{BMO}^c(\mathbb{T}^d, \mathcal{M})$ to be the space of all $f \in L_2(\mathcal{N})$ such that

$$\|f\|_{\text{BMO}^c} = \max \left\{ \|f_{\mathbb{T}^d}\|_{\mathcal{M}}, \sup_{Q \subset \mathbb{T}^d \text{ cube}} \left\| \frac{1}{|Q|} \int_Q |f(z) - \frac{1}{|Q|} \int_Q f(w) dw|^2 dz \right\|_{\mathcal{M}}^{\frac{1}{2}} \right\} < \infty.$$

The row $\text{BMO}^r(\mathbb{T}^d, \mathcal{M})$ consists of all f such that $f^* \in \text{BMO}^c(\mathbb{T}^d, \mathcal{M})$, equipped with $\|f\|_{\text{BMO}^r} = \|f^*\|_{\text{BMO}^c}$. Finally, we define mixture space $\text{BMO}(\mathbb{T}^d, \mathcal{M})$ as the intersection of the column and row BMO spaces:

$$\text{BMO}(\mathbb{T}^d, \mathcal{M}) = \text{BMO}^c(\mathbb{T}^d, \mathcal{M}) \cap \text{BMO}^r(\mathbb{T}^d, \mathcal{M}),$$

equipped with $\|f\|_{\text{BMO}} = \max(\|f\|_{\text{BMO}^c}, \|f\|_{\text{BMO}^r})$.

As in the Euclidean case, these spaces can be characterized by the circular Poisson semigroup. Let $\mathbb{P}_r$ denote the circular Poisson kernel of $\mathbb{T}^d$:

$$\mathbb{P}_r(z) = \sum_{m \in \mathbb{Z}^d} r^{|m|} z^m, \quad z \in \mathbb{T}^d, \quad 0 \leq r < 1. \quad (1.6)$$

The Poisson integral of $f \in L_1(\mathcal{N})$ is

$$\mathbb{P}_r(f)(z) = \int_{\mathbb{T}^d} \mathbb{P}_r(zw^{-1}) f(w) dw = \sum_{m \in \mathbb{Z}^d} \widehat{f}(m) r^{|m|} z^m.$$

Here $\widehat{f}$ denotes, of course, the Fourier transform of f :

$$\widehat{f}(m) = \int_{\mathbb{T}^d} f(z) z^{-m} dz.$$

It is proved in [17] that

$$\sup_{Q \subset \mathbb{T}^d_{\text{cube}}} \left\| \frac{1}{|Q|} \int_Q |f(z) - \frac{1}{|Q|} \int_Q f(w) dw|^2 dz \right\|_{\mathcal{M}} \approx \sup_{0 \leq r < 1} \|\mathbb{P}_r(|f - \mathbb{P}_r(f)|^2)\|_{\mathcal{N}} \quad (1.7)$$

with relevant constants depending only on d . Thus

$$\|f\|_{\text{BMO}^c} \approx \max \{ \|\widehat{f}(0)\|_{\mathcal{M}}, \sup_{0 \leq r < 1} \|\mathbb{P}_r(|f - \mathbb{P}_r(f)|^2)\|_{\mathcal{N}}^{\frac{1}{2}} \}.$$

Now we turn to the operator-valued Hardy spaces on $\mathbb{T}^d$ which are defined by the Littlewood-Paley functions associated to the circular Poisson kernel. For $f \in L_1(\mathcal{N})$ define

$$s^c(f)(z) = \left(\int_0^1 |\partial_r \mathbb{P}_r(f)(z)|^2 (1-r) dr \right)^{\frac{1}{2}}, \quad z \in \mathbb{T}^d.$$

For $1 \leq p < \infty$, let

$$\mathcal{H}_p^c(\mathbb{T}^d, \mathcal{M}) = \{f \in L_1(\mathcal{N}) : \|f\|_{\mathcal{H}_p^c} < \infty\},$$

where

$$\|f\|_{\mathcal{H}_p^c} = \|\widehat{f}(0)\|_{L_p(\mathcal{M})} + \|s^c(f)\|_{L_p(\mathcal{N})}.$$

The row Hardy space $\mathcal{H}_p^r(\mathbb{T}^d, \mathcal{M})$ is defined to be the space of all f such that $f^* \in \mathcal{H}_p^c(\mathbb{T}^d, \mathcal{M})$, equipped with the natural norm. Then we define

$$\mathcal{H}_p(\mathbb{T}^d, \mathcal{M}) = \begin{cases} \mathcal{H}_p^c(\mathbb{T}^d, \mathcal{M}) + \mathcal{H}_p^r(\mathbb{T}^d, \mathcal{M}) & \text{if } 1 \leq p < 2, \\ \mathcal{H}_p^c(\mathbb{T}^d, \mathcal{M}) \cap \mathcal{H}_p^r(\mathbb{T}^d, \mathcal{M}) & \text{if } 2 \leq p < \infty, \end{cases}$$

equipped with the sum and intersection norms, respectively:

$$\|f\|_{\mathcal{H}_p} = \begin{cases} \inf \{ \|g\|_{\mathcal{H}_p^c} + \|h\|_{\mathcal{H}_p^r} : f = g + h \} & \text{if } 1 \leq p < 2, \\ \max(\|f\|_{\mathcal{H}_p^c}, \|f\|_{\mathcal{H}_p^r}) & \text{if } 2 \leq p < \infty. \end{cases}$$

The following is the main results of [17, Section 8]

Lemma 1.8. (i) Let $1 < p < \infty$. Then $\mathcal{H}_p(\mathbb{T}^d, \mathcal{M}) = L_p(\mathcal{N})$ with equivalent norms.

(ii) The dual space of $\mathcal{H}_1^c(\mathbb{T}^d, \mathcal{M})$ coincides isomorphically with $\text{BMO}^c(\mathbb{T}^d, \mathcal{M})$.

(iii) Let $1 < p < \infty$. Then

$$\begin{aligned} (\text{BMO}^c(\mathbb{T}^d, \mathcal{M}), \mathcal{H}_1^c(\mathbb{T}^d, \mathcal{M}))_{\frac{1}{p}} &= \mathcal{H}_p^c(\mathbb{T}^d, \mathcal{M}) \\ (\text{BMO}^c(\mathbb{T}^d, \mathcal{M}), \mathcal{H}_1^c(\mathbb{T}^d, \mathcal{M}))_{\frac{1}{p}, p} &= \mathcal{H}_p^c(\mathbb{T}^d, \mathcal{M}). \end{aligned}$$

Similar statements hold for the row and mixture spaces too.

By transference, the previous results can be transferred to the quantum torus case. The Poisson integral of an element x in $L_1(\mathbb{T}_\theta^d)$ is defined by

$$\mathbb{P}_r(x) = \sum_{m \in \mathbb{Z}^d} \hat{x}(m) r^{|m|} U^m, \quad 0 \leq r < 1.$$

Its associated Littlewood-Paley g -function is

$$s^c(x) = \left(\int_0^1 |\partial_r \mathbb{P}_r(x)|^2 (1-r) dr \right)^{\frac{1}{2}}.$$

For $1 \leq p < \infty$ let

$$\|x\|_{\mathcal{H}_p^c} = |\hat{x}(0)| + \|s^c(x)\|_{L_p(\mathbb{T}_\theta^d)}.$$

The column Hardy space $\mathcal{H}_p^c(\mathbb{T}_\theta^d)$ is then defined to be

$$\mathcal{H}_p^c(\mathbb{T}_\theta^d) = \{x \in L_1(\mathbb{T}_\theta^d) : \|x\|_{\mathcal{H}_p^c} < \infty\}.$$

On the other hand, inspired by (1.7), we define

$$\text{BMO}^c(\mathbb{T}_\theta^d) = \{x \in L_2(\mathbb{T}_\theta^d) : \sup_{0 \leq r < 1} \|\mathbb{P}_r(|x - \mathbb{P}_r(x)|^2)\|_{\mathbb{T}_\theta^d} < \infty\},$$

equipped with the norm

$$\|x\|_{\text{BMO}^c} = \max \{|\hat{x}(0)|, \sup_{0 \leq r < 1} \|\mathbb{P}_r(|x - \mathbb{P}_r(x)|^2)\|_{\mathbb{T}_\theta^d}^{\frac{1}{2}}\}.$$

The corresponding row and mixture spaces are defined similarly to the preceding torus setting.

Lemma 1.8 admits the following quantum analogue:

Lemma 1.9. (i) Let $1 < p < \infty$. Then $\mathcal{H}_p(\mathbb{T}_\theta^d) = L_p(\mathbb{T}_\theta^d)$ with equivalent norms.

(ii) The dual space of $\mathcal{H}_1^c(\mathbb{T}_\theta^d)$ coincides isomorphically with $\text{BMO}^c(\mathbb{T}_\theta^d)$.

(iii) Let $1 < p < \infty$. Then

$$(\text{BMO}^c(\mathbb{T}_\theta^d), \mathcal{H}_1^c(\mathbb{T}_\theta^d))_{\frac{1}{p}} = \mathcal{H}_p^c(\mathbb{T}_\theta^d) = (\text{BMO}^c(\mathbb{T}_\theta^d), \mathcal{H}_1^c(\mathbb{T}_\theta^d))_{\frac{1}{p}, p}.$$

Similar statements hold for the row and mixture spaces too.

In the above definition of $\mathcal{H}_p^c(\mathbb{T}_\theta^d)$, the Poisson kernel can be replaced by any reasonable smooth function. Let ψ be a Schwartz function on $\mathbb{R}^d$ and ψ_j be the function whose Fourier transform is $\psi(2^{-j}\cdot)$. The map $x \mapsto \tilde{\psi}_j * x$ is the Fourier multiplier on $\mathbb{T}_\theta^d$ associated to ψ_j . Now we define the square function associated to ψ of an element $x \in L_1(\mathbb{T}_\theta^d)$ by

$$s_\psi^c(x) = \left(\sum_{j \geq 0} |\tilde{\psi}_j * x|^2 \right)^{\frac{1}{2}}.$$

The following lemma is the main result of [81]. We will need it essentially in the case of $p = 1$.

Lemma 1.10. *Let $1 \leq p < \infty$, and let ψ be a Schwartz function that does not vanish in $\{\xi : 1 \leq |\xi| < 2\}$. Then $x \in \mathcal{H}_p^c(\mathbb{T}_\theta^d)$ iff $s_\psi^c(x) \in L_p(\mathbb{T}_\theta^d)$. In this case, we have*

$$\|x\|_{\mathcal{H}_p^c} \approx |\hat{x}(0)| + \|s_\psi^c(x)\|_{L_p(\mathbb{T}_\theta^d)},$$

where the equivalence constants depend only on d, p and ψ .

Chapter 2

Sobolev spaces

This chapter starts with a brief introduction to distributions on quantum tori. We then pass to the definitions of Sobolev spaces on $\mathbb{T}_\theta^d$ and give some fundamental properties of them. Two families of Sobolev spaces are studied: $W_p^k(\mathbb{T}_\theta^d)$ and the fractional Sobolev spaces $H_p^\alpha(\mathbb{T}_\theta^d)$. We prove a Poincaré type inequality for $W_p^k(\mathbb{T}_\theta^d)$ for any $1 \leq p \leq \infty$. Our approach to this inequality seems very different from existing proofs for such an inequality in the classical case. We show that $W_\infty^k(\mathbb{T}_\theta^d)$ coincides with the Lipschitz class of order k , studied by Weaver [78, 79]. We conclude the chapter with a section on the link between the quantum Sobolev spaces and the vector-valued Sobolev spaces on the usual d -torus $\mathbb{T}^d$.

2.1 Distributions on quantum tori

In this section we give an outline of the distribution theory on quantum tori. Let

$$\mathcal{S}(\mathbb{T}_\theta^d) = \left\{ \sum_{m \in \mathbb{Z}^d} a_m U^m : \{a_m\}_{m \in \mathbb{Z}^d} \text{ rapidly decreasing} \right\}.$$

This is a w^* -dense $*$ -subalgebra of $\mathbb{T}_\theta^d$ and contains all polynomials. We simply write $\mathcal{S}(\mathbb{T}_\theta^d) = \mathcal{S}(\mathbb{T}^d)$, the algebras of infinitely differentiable functions on $\mathbb{T}^d$. Thus for a general θ , $\mathcal{S}(\mathbb{T}_\theta^d)$ should be viewed as a noncommutative deformation of $\mathcal{S}(\mathbb{T}^d)$. We will need the differential structure on $\mathcal{S}(\mathbb{T}_\theta^d)$, which is similar to that on $\mathcal{S}(\mathbb{T}^d)$.

According to our convention made in section 1.3 and in order to lighten the notational system, we will use the same derivation notation on $\mathbb{T}_\theta^d$ as on $\mathbb{T}^d$. For every $1 \leq j \leq d$, define the following derivations, which are operators on $\mathcal{S}(\mathbb{T}_\theta^d)$:

$$\partial_j(U_j) = 2\pi i U_j \quad \text{and} \quad \partial_j(U_k) = 0 \quad \text{for } k \neq j.$$

These operators ∂_j commute with the adjoint operation $*$, and play the role of the partial derivations in the classical analysis on the usual d -torus. Given $m = (m_1, \dots, m_d) \in \mathbb{N}_0^d$, the associated partial derivation D^m is $\partial_1^{m_1} \cdots \partial_d^{m_d}$. We also use Δ to denote the Laplacian: $\Delta = \partial_1^2 + \cdots + \partial_d^2$. The elementary fact expressed in the following remark will be frequently used later on.

Restricted to $L_2(\mathbb{T}_\theta^d)$, the partial derivation ∂_j is a densely defined closed (unbounded) operator whose adjoint is equal to $-\partial_j$. This is an immediate consequence of the following obvious fact (cf. [64]):

Lemma 2.1. *If $x, y \in \mathcal{S}(\mathbb{T}_\theta^d)$, then $\tau(\partial_j(xy)) = -\tau(x\partial_j(y))$ for $j = 1, \dots, d$.*

Thus $\Delta = -(\partial_1^* \partial_1 + \dots + \partial_d^* \partial_d)$, so $-\Delta$ is a positive operator on $L_2(\mathbb{T}_\theta^d)$ with spectrum equal to $\{4\pi^2|m|^2 : m \in \mathbb{Z}^d\}$.

Remark 2.2. Given $u \in \mathbb{R}^d$ let e_u be the function on $\mathbb{R}^d$ defined by $e_u(\xi) = e^{2\pi i u \cdot \xi}$, where $u \cdot \xi$ denotes the inner product of $\mathbb{R}^d$. The Fourier multiplier on $\mathbb{T}_\theta^d$ associated to e_u coincides with π_z in (1.4) with $z = (e^{2\pi i u_1}, \dots, e^{2\pi i u_d})$. This Fourier multiplier will play an important role in the sequel. By analogy with the classical case, we will call it the translation by u and denote it by T_u : $T_u(x) = \pi_z(x)$ for any $x \in \mathcal{S}(\mathbb{T}_\theta^d)$. Then it is clear that

$$\frac{\partial}{\partial u_j} T_u(x) = T_u(\partial_j x), \text{ so } \frac{\partial}{\partial u_j} T_u(x) \Big|_{u=0} = \partial_j x. \quad (2.1)$$

Following the classical setting as in [22], we now endow $\mathcal{S}(\mathbb{T}_\theta^d)$ with an appropriate topology. For each $k \in \mathbb{N}_0$ define a norm p_k on $\mathcal{S}(\mathbb{T}_\theta^d)$ by

$$p_k(x) = \sup_{0 \leq |m|_1 \leq k} \|D^m x\|_\infty.$$

The sequence $\{p_k\}_{k \geq 0}$ induces a locally convex topology on $\mathcal{S}(\mathbb{T}_\theta^d)$. This topology is metrizable by the following distance:

$$d(x, y) = \sum_{k=0}^{\infty} \frac{2^{-k} p_k(x - y)}{1 + p_k(x - y)}$$

with respect to which $\mathcal{S}(\mathbb{T}_\theta^d)$ is complete, an easily checked fact. The following simple proposition describes the convergence in $\mathcal{S}(\mathbb{T}_\theta^d)$.

Proposition 2.3. *A sequence $\{x_n\} \subset \mathcal{S}(\mathbb{T}_\theta^d)$ converges to $x \in \mathcal{S}(\mathbb{T}_\theta^d)$ if and only if for every $m \in \mathbb{N}_0^d$, $D^m x_n \rightarrow D^m x$ in $\mathbb{T}_\theta^d$.*

Proof. Without loss of generality, we assume $x = 0$. Suppose that $x_n \rightarrow 0$ in $\mathcal{S}(\mathbb{T}_\theta^d)$. Then for $m \in \mathbb{N}_0^d$ and $\varepsilon > 0$, there exists an integer N such that for every $n \geq N$,

$$d(x_n, 0) = \sum_{k=0}^{\infty} \frac{2^{-k} p_k(x_n)}{1 + p_k(x_n)} \leq 2^{-|m|_1} \frac{\varepsilon}{1 + \varepsilon}.$$

Thus, $p_{|m|_1}(x_n) \leq \varepsilon$, so $\|D^m x_n\|_\infty \leq \varepsilon$, which means $D^m x_n \rightarrow 0$ in $\mathbb{T}_\theta^d$.

Conversely, assume that for every $m \in \mathbb{N}_0^d$, $D^m x_n \rightarrow 0$ in $\mathbb{T}_\theta^d$. For $\varepsilon > 0$ let N_0 be an integer such that $\sum_{k > N_0} 2^{-k} < \frac{\varepsilon}{2}$. Since $D^m x_n \rightarrow 0$ for $|m|_1 \leq N_0$, there exists $N \in \mathbb{N}$ such that for $n > N$,

$$\sum_{k=0}^{N_0} \frac{2^{-k} p_k(x_n)}{1 + p_k(x_n)} < \frac{\varepsilon}{2}.$$

Consequently, $d(x_n, 0) < \varepsilon$. □

Definition 2.4. A distribution on $\mathbb{T}_\theta^d$ is a continuous linear functional on $\mathcal{S}(\mathbb{T}_\theta^d)$. $\mathcal{S}'(\mathbb{T}_\theta^d)$ denotes the space of distributions.

As a dual space, $\mathcal{S}'(\mathbb{T}_\theta^d)$ is endowed with the w^* -topology. We will use the bracket $\langle, \rangle$ to denote the duality between $\mathcal{S}(\mathbb{T}_\theta^d)$ and $\mathcal{S}'(\mathbb{T}_\theta^d)$: $\langle F, x \rangle = F(x)$ for $x \in \mathcal{S}(\mathbb{T}_\theta^d)$ and $F \in \mathcal{S}'(\mathbb{T}_\theta^d)$. We list some elementary properties of distributions:

- (1) For $1 \leq p \leq \infty$, the space $L_p(\mathbb{T}_\theta^d)$ naturally embeds into $\mathcal{S}'(\mathbb{T}_\theta^d)$: an element $y \in L_p(\mathbb{T}_\theta^d)$ induces a continuous functional on $\mathcal{S}(\mathbb{T}_\theta^d)$ by $x \mapsto \tau(yx)$.

- (2) $\mathcal{S}(\mathbb{T}_\theta^d)$ acts as a bimodule on $\mathcal{S}'(\mathbb{T}_\theta^d)$: for $a, b \in \mathcal{S}(\mathbb{T}_\theta^d)$ and $F \in \mathcal{S}'(\mathbb{T}_\theta^d)$, aFb is the distribution defined by $x \mapsto \langle F, bxa \rangle$.
- (3) The partial derivations ∂_j extend to $\mathcal{S}'(\mathbb{T}_\theta^d)$ by duality: $\langle \partial_j F, x \rangle = -\langle F, \partial_j x \rangle$. For $m \in \mathbb{N}_0^d$, we use again D^m to denote the associated partial derivation on $\mathcal{S}'(\mathbb{T}_\theta^d)$.
- (4) The Fourier transform extends to $\mathcal{S}'(\mathbb{T}_\theta^d)$: for $F \in \mathcal{S}'(\mathbb{T}_\theta^d)$ and $m \in \mathbb{Z}^d$, $\widehat{F}(m) = \langle F, (U^m)^* \rangle$. The Fourier series of F converges to F according to any (reasonable) summation method:

$$F = \sum_{m \in \mathbb{Z}^d} \widehat{F}(m) U^m.$$

2.2 Definitions and basic properties

We begin this section with a simple observation on Fourier multipliers on $\mathcal{S}(\mathbb{T}_\theta^d)$ and $\mathcal{S}'(\mathbb{T}_\theta^d)$. Let $\phi : \mathbb{Z}^d \rightarrow \mathbb{C}$ be a function of polynomial growth. Then its associated Fourier multiplier M_ϕ is a continuous map on both $\mathcal{S}(\mathbb{T}_\theta^d)$ and $\mathcal{S}'(\mathbb{T}_\theta^d)$. Here and in the sequel, a generic element of $\mathcal{S}'(\mathbb{T}_\theta^d)$ is also denoted by x . Two specific Fourier multipliers will play a key role later: they are the Bessel and Riesz potentials.

Definition 2.5. Let $\alpha \in \mathbb{R}$. Define J_α on $\mathbb{R}^d$ and I_α on $\mathbb{R}^d \setminus \{0\}$ by

$$J_\alpha(\xi) = (1 + |\xi|^2)^{\frac{\alpha}{2}} \quad \text{and} \quad I_\alpha(\xi) = |\xi|^\alpha.$$

Their associated Fourier multipliers are the Bessel and Riesz potentials of order α , denoted by J^α and I^α , respectively.

By the above observation, J^α is a Fourier multiplier on $\mathcal{S}'(\mathbb{T}_\theta^d)$, and I^α is also a Fourier multiplier on the subspace of $\mathcal{S}'(\mathbb{T}_\theta^d)$ of all x such that $\widehat{x}(0) = 0$. Note that

$$J^\alpha = (1 - (2\pi)^{-2} \Delta)^{\frac{\alpha}{2}} \quad \text{and} \quad I^\alpha = (-(2\pi)^{-2} \Delta)^{\frac{\alpha}{2}}.$$

Definition 2.6. Let $1 \leq p \leq \infty$, $k \in \mathbb{N}$ and $\alpha \in \mathbb{R}$.

- (i) The Sobolev space of order k on $\mathbb{T}_\theta^d$ is defined to be

$$W_p^k(\mathbb{T}_\theta^d) = \{x \in \mathcal{S}'(\mathbb{T}_\theta^d) : D^m x \in L_p(\mathbb{T}_\theta^d) \text{ for each } m \in \mathbb{N}_0^d \text{ with } |m|_1 \leq k\},$$

equipped with the norm

$$\|x\|_{W_p^k} = \left(\sum_{0 \leq |m|_1 \leq k} \|D^m x\|_p^p \right)^{\frac{1}{p}}.$$

- (ii) The potential (or fractional) Sobolev space of order α is defined to be

$$H_p^\alpha(\mathbb{T}_\theta^d) = \{x \in \mathcal{S}'(\mathbb{T}_\theta^d) : J^\alpha x \in L_p(\mathbb{T}_\theta^d)\},$$

equipped with the norm

$$\|x\|_{H_p^\alpha} = \|J^\alpha x\|_p.$$

In the above definition of $\|x\|_{W_p^k}$, we have followed the usual convention for $p = \infty$ that the right-hand side is replaced by the corresponding supremum. This convention will be always made in the sequel. We collect some basic properties of these spaces in the following:

Proposition 2.7. *Let $1 \leq p \leq \infty$, $k \in \mathbb{N}$ and $\alpha \in \mathbb{R}$.*

- (i) $W_p^k(\mathbb{T}_\theta^d)$ and $H_p^\alpha(\mathbb{T}_\theta^d)$ are Banach spaces.
- (ii) The polynomial subalgebra $\mathcal{P}_\theta$ of $\mathbb{T}_\theta^d$ is dense in $W_p^k(\mathbb{T}_\theta^d)$ and $H_p^\alpha(\mathbb{T}_\theta^d)$ for $1 \leq p < \infty$. Consequently, $\mathcal{S}(\mathbb{T}_\theta^d)$ is dense in $W_p^k(\mathbb{T}_\theta^d)$ and $H_p^\alpha(\mathbb{T}_\theta^d)$.
- (iii) For any $\beta \in \mathbb{R}$, J^β is an isometry from $H_p^\alpha(\mathbb{T}_\theta^d)$ onto $H_p^{\alpha-\beta}(\mathbb{T}_\theta^d)$. In particular, J^α is an isometry from $H_p^\alpha(\mathbb{T}_\theta^d)$ onto $L_p(\mathbb{T}_\theta^d)$.
- (iv) $H_p^\alpha(\mathbb{T}_\theta^d) \subset H_p^\beta(\mathbb{T}_\theta^d)$ continuously whenever $\beta < \alpha$.

Proof. (iii) is obvious. It implies (i) for $H_p^\alpha(\mathbb{T}_\theta^d)$.

(i) It suffices to show that $W_p^k(\mathbb{T}_\theta^d)$ is complete. Assume that $\{x_n\} \subset W_p^k(\mathbb{T}_\theta^d)$ is a Cauchy sequence. Then for every $|m|_1 \leq k$, $\{D^m x_n\}$ is a Cauchy sequence in $L_p(\mathbb{T}_\theta^d)$, so $D^m x_n \rightarrow y_m$ in $L_p(\mathbb{T}_\theta^d)$. Particularly, $\{D^m x_n\}_n$ converges to y_m in $\mathcal{S}'(\mathbb{T}_\theta^d)$. On the other hand, since $x_n \rightarrow y_0$ in $L_p(\mathbb{T}_\theta^d)$, for every $x \in \mathcal{S}(\mathbb{T}_\theta^d)$ we have $\tau(x_n D^m x) \rightarrow \tau(y_0 D^m x)$. Thus $\{D^m x_n\}_n$ converges to $D^m y_0$ in $\mathcal{S}'(\mathbb{T}_\theta^d)$. Consequently, $D^m y_0 = y_m$ for $|m|_1 \leq k$. Hence, $y_0 \in W_p^k(\mathbb{T}_\theta^d)$ and $x_n \rightarrow y_0$ in $W_p^k(\mathbb{T}_\theta^d)$.

(ii) Consider the square Fejér mean

$$F_N(x) = \sum_{m \in \mathbb{Z}^d, \max_j |m_j| \leq N} \left(1 - \frac{|m_1|}{N+1}\right) \cdots \left(1 - \frac{|m_d|}{N+1}\right) \hat{x}(m) U^m.$$

By [17, Proposition 3.1], $\lim_{N \rightarrow \infty} F_N(x) = x$ in $L_p(\mathbb{T}_\theta^d)$. On the other hand, F_N commutes with D^m : $F_N(D^m x) = D^m F_N(x)$. We then deduce that $\lim_{N \rightarrow \infty} F_N(x) = x$ in $W_p^k(\mathbb{T}_\theta^d)$ for every $x \in W_p^k(\mathbb{T}_\theta^d)$. Thus $\mathcal{P}_\theta$ is dense in $W_p^k(\mathbb{T}_\theta^d)$. On the other hand, F_N and J^α commute; so by (iii), the density of $\mathcal{P}_\theta$ in $L_p(\mathbb{T}_\theta^d)$ implies its density in $H_p^\alpha(\mathbb{T}_\theta^d)$.

(iv) It is well known that if $\gamma < 0$, the inverse Fourier transform of J_γ is an integrable function on $\mathbb{R}^d$ (see [67, Proposition V.3]). Thus, Lemma 1.7 implies that $J^{\beta-\alpha}$ is a bounded map on $L_p(\mathbb{T}_\theta^d)$ with norm majorized by $\|\mathcal{F}^{-1}(J_{\beta-\alpha})\|_{L_1(\mathbb{R}^d)}$. This is the desired assertion. $\square$

The following shows that the potential Sobolev spaces can be also characterized by the Riesz potential.

Theorem 2.8. *Let $1 \leq p \leq \infty$. Then*

$$\|x\|_{H_p^\alpha} \approx \left(|\hat{x}(0)|^p + \|I^\alpha(x - \hat{x}(0))\|_p^p \right)^{\frac{1}{p}},$$

where the equivalence constants depend only on α and d .

Proof. By changing α to $-\alpha$, we can assume $\alpha > 0$. It suffices to show $\|I^\alpha x\|_p \approx \|J^\alpha x\|_p$ for $\hat{x}(0) = 0$. By [67, Lemma V.3.2], $\frac{I_\alpha}{J_\alpha}$ is the Fourier transform of a bounded measure on $\mathbb{R}^d$, which, together with Lemma 1.7, yields $\|I^\alpha x\|_p \lesssim \|J^\alpha x\|_p$.

To show the converse inequality, let η be an infinite differentiable function on $\mathbb{R}^d$ such that $\eta(\xi) = 0$ for $|\xi| \leq \frac{1}{2}$ and $\eta(\xi) = 1$ for $|\xi| \geq 1$, and let $\phi = J_\alpha I_{-\alpha} \eta$. Then the Fourier multiplier with symbol ϕI_α coincides with J^α on the subspace of distributions on $\mathbb{T}_\theta^d$ with vanishing Fourier coefficients at the origin. Thus we are reduced to proving $\mathcal{F}^{-1}(\phi) \in L_1(\mathbb{R}^d)$. To that end, first observe that for any $m \in \mathbb{N}_0^d$,

$$|D^m \phi(\xi)| \lesssim \frac{1}{|\xi|^{|m|_1+2}}.$$

Consider first the case $d \geq 3$. Choose positive integers ℓ and k such that $\frac{d}{2} - 2 < \ell < \frac{d}{2}$ and $k > \frac{d}{2}$. Then by the Cauchy-Schwarz inequality and the Plancherel theorem,

$$\begin{aligned} \left(\int_{|s|<1} |\mathcal{F}^{-1}\phi(s)| ds \right)^2 &\leq \int_{|s|<1} |s|^{-2\ell} ds \int_{|s|<1} |s|^{2\ell} |\mathcal{F}^{-1}\phi(s)|^2 ds \\ &\lesssim \sum_{m \in \mathbb{N}_0^d, |m|_1 = \ell} \int_{\mathbb{R}^d} |D^m \phi(\xi)|^2 d\xi \\ &\lesssim \int_{|\xi| \geq \frac{1}{2}} \frac{1}{|\xi|^{2(\ell+2)}} d\xi \lesssim 1. \end{aligned}$$

On the other hand,

$$\begin{aligned} \left(\int_{|s| \geq 1} |\mathcal{F}^{-1}\phi(s)| ds \right)^2 &\leq \int_{|s| \geq 1} |s|^{-2k} ds \int_{|s| \geq 1} |s|^{2k} |\mathcal{F}^{-1}\phi(s)|^2 ds \\ &\lesssim \sum_{m \in \mathbb{N}_0^d, |m|_1 = k} \int_{\mathbb{R}^d} |D^m \phi(\xi)|^2 d\xi \lesssim 1. \end{aligned}$$

Thus $\mathcal{F}^{-1}(\phi)$ is integrable for $d \geq 3$.

If $d \leq 2$, the second part above remains valid, while the first one should be modified since the required positive integer ℓ does not exist for $d \leq 2$. We will consider $d = 2$ and $d = 1$ separately. For $d = 2$, choosing $0 < \varepsilon < \frac{1}{2}$, we have

$$\int_{|s|<1} |\mathcal{F}^{-1}\phi(s)| ds \leq \left(\int_{|s|<1} |s|^{-2\varepsilon} ds \right)^{\frac{1}{2}} \left(\int_{\mathbb{R}^d} |s|^{2\varepsilon} |\mathcal{F}^{-1}\phi(s)|^2 ds \right)^{\frac{1}{2}} \lesssim \|I^\varepsilon \phi\|_2.$$

Writing $I^\varepsilon = I^{\varepsilon-1} I^1$ and using the classical Hardy-Littlewood-Sobolev inequality (see [67, Theorem V.1]) and the Riesz transform, we obtain

$$\|I^\varepsilon \phi\|_2 \lesssim \|I^1 \phi\|_q \approx \|\nabla \phi\|_q \lesssim \left(\int_{|\xi| \geq \frac{1}{2}} \frac{1}{|\xi|^{3q}} d\xi \right)^{\frac{1}{q}} \lesssim 1,$$

where $\frac{1}{q} = 1 - \frac{\varepsilon}{2}$ (so $1 < q < \infty$). Thus we are done in the case $d = 2$.

It remains to deal with the one-dimensional case. Write

$$\phi(\xi) = (1 + \xi^{-2})^{\frac{\alpha}{2}} \eta(\xi) = \eta(\xi) + \rho(\xi) \eta(\xi), \quad \xi \in \mathbb{R} \setminus \{0\},$$

where $\rho(\xi) = O(\xi^{-2})$ as $|\xi| \rightarrow \infty$. Since $\eta - 1$ is infinitely differentiable and supported by $[-1, 1]$, its inverse Fourier transform is integrable. So η is the Fourier transform of a finite measure on $\mathbb{R}$. On the other hand, as $\rho\eta \in L_1(\mathbb{R}^d)$, $\mathcal{F}^{-1}(\rho\eta)$ is a bounded continuous function, so it is integrable on $[-1, 1]$. On the other hand, by the second part of the preceding argument for $d \geq 3$, we see that $\mathcal{F}^{-1}(\rho\eta)$ is integrable outside $[-1, 1]$ too, whence $\mathcal{F}^{-1}(\rho\eta) \in L_1(\mathbb{R}^d)$. We thus deduce that ϕ is the Fourier transform of a finite measure on $\mathbb{R}$. Hence the assertion is completely proved. $\square$

Theorem 2.9. *Let $1 < p < \infty$. Then $H_p^k(\mathbb{T}_\theta^d) = W_p^k(\mathbb{T}_\theta^d)$ with equivalent norms.*

Proof. This proof is based on Fourier multipliers by virtue of Lemma 1.7. For any $m \in \mathbb{N}_0^d$ with $|m|_1 \leq k$, the function ϕ , defined by $\phi(\xi) = (2\pi i)^{|m|_1} \xi^m (1 + |\xi|^2)^{-\frac{k}{2}}$, is clearly a Mihlin multiplier. Then for any $x \in \mathcal{S}'(\mathbb{T}_\theta^d)$,

$$\|D^m x\|_p = \|M_\phi(J^k x)\|_p \lesssim \|J^k x\|_p,$$

whence $\|x\|_{W_p^k} \lesssim \|x\|_{H_p^k}$. To prove the converse inequality, choose an infinite differentiable function χ on $\mathbb{R}$ such that $\chi = 0$ on $\{\xi : |\xi| \leq 4^{-1}\}$ and $\chi = 1$ on $\{\xi : |\xi| \geq 2^{-1}\}$. Let

$$\phi(\xi) = \frac{(1 + |\xi|^2)^{\frac{k}{2}}}{1 + \chi(\xi_1)|\xi_1|^k + \cdots + \chi(\xi_d)|\xi_d|^k} \quad \text{and} \quad \phi_j(\xi) = \frac{\chi(\xi_j)|\xi_j|^k}{(2\pi i \xi_j)^k}, \quad 1 \leq j \leq d.$$

These are Mikhlin multipliers too, and

$$J^k x = M_\phi(x + M_{\phi_1} \partial_1^k x + \cdots + M_{\phi_d} \partial_d^k x).$$

It then follows that

$$\|x\|_{H_p^k} \lesssim \left(\|x\|_p^p + \sum_{j=1}^d \|\partial_j^k x\|_p^p \right)^{\frac{1}{p}} \leq \|x\|_{W_p^k}.$$

The assertion is thus proved. $\square$

Remark 2.10. Incidentally, the above proof shows that if $1 < p < \infty$, then

$$\|x\|_{W_p^k} \approx \left(\|x\|_p^p + \sum_{j=1}^d \|\partial_j^k x\|_p^p \right)^{\frac{1}{p}}$$

with relevant constants depending only on p and d .

However, if one allows the above sum to run over all partial derivations of order k , then p can be equal to 1 or ∞ . Namely, for any $1 \leq p \leq \infty$,

$$\|x\|_{W_p^k} \approx \left(\|x\|_p^p + \sum_{m \in \mathbb{N}_0^d, |m|_1 = k} \|D^m x\|_p^p \right)^{\frac{1}{p}}$$

with relevant constants depending only on d . This equivalence can be proved by standard arguments (see Lemma 2.15 below and its proof). In fact, we have a nicer result, a Poincaré-type inequality:

$$\|x\|_p \lesssim \sum_{j=1}^d \|\partial_j x\|_p$$

for any $x \in W_p^1(\mathbb{T}_\theta^d)$ with $\hat{x}(0) = 0$. So $\|x\|_p$ can be removed from the right-hand side of the above equivalence. This inequality will be proved in the next section.

We conclude this section with an easy description of the dual space of $W_p^k(\mathbb{T}_\theta^d)$. Let $N = N(d, k) = \sum_{m \in \mathbb{N}_0^d, 0 \leq |m|_1 \leq k} 1$ and

$$L_p^N = \prod_{j=1}^N L_p(\mathbb{T}_\theta^d) \quad \text{equipped with the norm} \quad \|x\|_{L_p^N} = \left(\sum_{j=1}^N \|x_j\|_p^p \right)^{\frac{1}{p}}.$$

The map $x \mapsto (D^m x)_{0 \leq |m|_1 \leq k}$ establishes an isometry from $W_p^k(\mathbb{T}_\theta^d)$ into L_p^N . Therefore, the dual of $W_p^k(\mathbb{T}_\theta^d)$ with $1 \leq p < \infty$ is identified with a quotient of $L_{p'}^N$, where p' is the conjugate index of p . More precisely, for every $\ell \in (W_p^k(\mathbb{T}_\theta^d))^*$ there exists an element $y = (y_m)_{m \in \mathbb{N}_0^d, 0 \leq |m|_1 \leq k} \in L_{p'}^N$ such that

$$\ell(x) = \sum_{0 \leq |m|_1 \leq k} \tau(y_m D^m x), \quad \forall x \in W_p^k(\mathbb{T}_\theta^d), \quad (2.2)$$

and

$$\|\ell\|_{(W_p^k)^*} = \inf \|y\|_{L_{p'}^N},$$

the infimum running over all $y \in L_{p'}^N$ as above.

$(W_p^k(\mathbb{T}_\theta^d))^*$ can be described as a space of distributions. Indeed, let $\ell \in (W_p^k(\mathbb{T}_\theta^d))^*$ and $y \in L_{p'}^N$ be a representative of ℓ as in (2.2). Define $\ell_y \in \mathcal{S}'(\mathbb{T}_\theta^d)$ by

$$\ell_y = \sum_{0 \leq |m|_1 \leq k} (-1)^{|m|_1} D^m y_m. \quad (2.3)$$

Then

$$\ell_y(x) = \sum_{0 \leq |m|_1 \leq k} \tau(y_m D^m x) = \ell(x), \quad x \in \mathcal{S}(\mathbb{T}_\theta^d).$$

So ℓ is an extension of ℓ_y ; moreover,

$$\|\ell\|_{(W_p^k)^*} = \min\{\|y\|_{L_{p'}^N} : \ell \text{ extends } \ell_y \text{ given by (2.3)}\}.$$

Conversely, suppose ℓ is an element of $\mathcal{S}'(\mathbb{T}_\theta^d)$ of the above form ℓ_y for some $y \in L_{p'}^N$. Then by the density of $\mathcal{S}(\mathbb{T}_\theta^d)$ in $W_p^k(\mathbb{T}_\theta^d)$, ℓ extends uniquely to a continuous functional on $W_p^k(\mathbb{T}_\theta^d)$. Thus we have proved the following

Proposition 2.11. *Let $1 \leq p < \infty$ and $W_{p'}^{-k}(\mathbb{T}_\theta^d)$ be the space of those distributions ℓ which admit a representative ℓ_y as above, equipped with the norm $\inf\{\|y\|_{L_{p'}^N} : y \text{ as in (2.3)}\}$. Then $(W_p^k(\mathbb{T}_\theta^d))^*$ is isometric to $W_{p'}^{-k}(\mathbb{T}_\theta^d)$.*

Note that the duality problem for the potential Sobolev spaces is trivial. Since J^α is an isometry between $H_p^\alpha(\mathbb{T}_\theta^d)$ and $L_p(\mathbb{T}_\theta^d)$, we see that for $1 \leq p < \infty$ and $\alpha \in \mathbb{R}$, the dual space of $H_p^\alpha(\mathbb{T}_\theta^d)$ coincides with $H_{p'}^{-\alpha}(\mathbb{T}_\theta^d)$ isometrically.

2.3 A Poincaré-type inequality

For $x \in W_p^k(\mathbb{T}_\theta^d)$ let

$$|x|_{W_p^k} = \left(\sum_{m \in \mathbb{N}_0^d, |m|_1 = k} \|D^m x\|_p^p \right)^{\frac{1}{p}}.$$

Theorem 2.12. *Let $1 \leq p \leq \infty$. Then for any $x \in W_p^1(\mathbb{T}_\theta^d)$,*

$$\|x - \hat{x}(0)\|_p \lesssim |x|_{W_p^1}.$$

More generally, if $k \in \mathbb{N}$ and $x \in W_p^k(\mathbb{T}_\theta^d)$ with $\hat{x}(0) = 0$, then

$$|x|_{W_p^j} \lesssim |x|_{W_p^k}, \quad \forall 0 \leq j < k.$$

Consequently, $|\hat{x}(0)| + |x|_{W_p^k}$ is an equivalent norm on $W_p^k(\mathbb{T}_\theta^d)$.

The proof given below is quite different from standard approaches to the Poincaré inequality. We will divide it into several lemmas, each of which might be interesting in its own right. We start with the following definition which will be frequently used later. Note that the function e_u and the translation operator T_u have been defined in Remark 2.1.

Definition 2.13. Given $u \in \mathbb{R}^d$ let $d_u = e_u - 1$. The Fourier multiplier on $\mathbb{T}_\theta^d$ defined by d_u is called the difference operator by u and denoted by Δ_u .

Remark 2.14. Note that e_u is the Fourier transform of the Dirac measure δ_u at u . Thus T_u is an isometry and Δ_u is of norm 2 on $L_p(\mathbb{T}_\theta^d)$ for any $1 \leq p \leq \infty$.

Lemma 2.15. Let $1 \leq p \leq \infty$, and $j, k \in \mathbb{N}$ with $j < k$. Then for any $x \in W_p^k(\mathbb{T}_\theta^d)$,

$$|x|_{W_p^j} \lesssim \|x\|_p^{1-\frac{j}{k}} |x|_{W_p^k}^{\frac{j}{k}}.$$

Proof. Fix $x \in W_p^k(\mathbb{T}_\theta^d)$ with $\hat{x}(0) = 0$. For any $u, \xi \in \mathbb{R}^d$ we have

$$\begin{aligned} d_u(\xi) - \frac{\partial}{\partial r} d_{ru}(\xi)|_{r=0} &= \int_0^1 \left(\frac{\partial}{\partial r} d_{ru}(\xi) - \frac{\partial}{\partial r} d_{ru}(\xi)|_{r=0} \right) dr \\ &= \int_0^1 \int_0^r \frac{\partial^2}{\partial s^2} d_{su}(\xi) ds dr. \end{aligned}$$

Since

$$\frac{\partial}{\partial r} d_{ru}(\xi) = e_{ru}(\xi)(2\pi i u \cdot \xi) \quad \text{and} \quad \frac{\partial^2}{\partial s^2} d_{su}(\xi) = e_{su}(\xi)(2\pi i u \cdot \xi)^2,$$

letting $u = te_j$ with $t > 0$ and e_j the j th canonical basic vector of $\mathbb{R}^d$, we deduce

$$\Delta_u x - t \partial_j x = \int_0^1 \int_0^r T_{su}(t^2 \partial_j^2 x) ds dr.$$

Thus

$$t \|\partial_j x\|_p \leq \|\Delta_u x\|_p + t^2 \int_0^1 \int_0^r \|T_{su}(\partial_j^2 x)\|_p ds dr \leq 2\|x\|_p + \frac{t^2}{2} \|\partial_j^2 x\|_p.$$

Dividing by t and taking the infimum over all $t > 0$, we get

$$\|\partial_j x\|_p \leq 2\sqrt{\|x\|_p \|\partial_j^2 x\|_p}. \quad (2.4)$$

This gives the assertion for the case $j = 1$ and $k = 2$. An iteration argument yields the general case. $\square$

Lemma 2.16. Let $j \in \{1, \dots, d\}$ and $x \in W_p^2(\mathbb{T}_\theta^d)$ such that $m_j \neq 0$ whenever $\hat{x}(m) \neq 0$ for $m \in \mathbb{Z}^d$. Then

$$\|x\|_p \leq c \|\partial_j^2 x\|_p,$$

where c is a universal constant. More generally, for any $x \in W_p^2(\mathbb{T}_\theta^d)$ with $\hat{x}(0) = 0$

$$\|x\|_p \leq c \sum_{j=1}^d \|\partial_j^2 x\|_p.$$

Proof. Assume $j = 1$. Consider the one-dimensional heat kernel of $\mathbb{T}$:

$$\mathbb{W}_r^1(z_1) = \sum_{m_1 \in \mathbb{Z}} r^{m_1^2} z_1^{m_1}, \quad z_1 \in \mathbb{T}, \quad 0 \leq r < 1.$$

$\mathbb{W}_r^1$ is also viewed as a function on $\mathbb{T}^d$, independent of $(z_2, \dots, z_d)$. So its associated Fourier multiplier on $\mathcal{S}'(\mathbb{T}_\theta^d)$ is given by

$$\mathbb{W}_r^1(x) = \sum_{m \in \mathbb{Z}^d} \hat{x}(m) r^{m_1^2} U^m.$$

As usual, $\mathbb{W}_r^1(x)$ can be expressed as a convolution:

$$\mathbb{W}_r^1(x) = \int_{\mathbb{T}} \mathbb{W}_r^1(z_1 w_1^{-1}) \pi_w(x) dw_1,$$

where we have identified $\pi_z(\mathbb{W}_r^1(x))$ with $\mathbb{W}_r^1(x)$ via the transference in (1.4). On the other hand, as a Fourier multiplier on $\mathbb{T}_\theta^d$, $\mathbb{W}_r$ admits the following function as symbol

$$\widehat{\mathbb{W}}_\varepsilon^1(\xi_1) = e^{-4\pi^2 \varepsilon \xi_1^2} \quad \text{with } r = e^{-4\pi^2 \varepsilon},$$

where $\mathbb{W}_\varepsilon$ is the one-dimensional heat kernel of $\mathbb{R}$:

$$\mathbb{W}_\varepsilon^1(s_1) = \frac{1}{\sqrt{4\pi\varepsilon}} e^{-\frac{s_1^2}{4\varepsilon}}, \quad s_1 \in \mathbb{R}.$$

By the classical Poisson summation formula, $\mathbb{W}_r^1$ is the 1-periodization of $\mathbb{W}_\varepsilon^1$:

$$\mathbb{W}_r^1(z_1) = \sum_{m_1 \in \mathbb{Z}} \mathbb{W}_\varepsilon^1(s_1 + m_1) \quad \text{with } z_1 = e^{2\pi i s_1} \quad \text{and } r = e^{-4\pi^2 \varepsilon}, \quad (2.5)$$

Thus

$$\|\mathbb{W}_r^1\|_{L_1(\mathbb{T})} = 1, \quad 0 \leq r < 1.$$

It follows that $\mathbb{W}_r^1$ is a contraction on $L_p(\mathbb{T}_\theta^d)$ for any $1 \leq p \leq \infty$ and $0 \leq r < 1$.

Now let $x \in W_p^2(\mathbb{T}_\theta^d)$ such that $m_1 \neq 0$ whenever $\widehat{x}(m) \neq 0$ for $m \in \mathbb{Z}^d$. Let $r_0 \in (0, 1)$ to be determined later. Then

$$x - \mathbb{W}_{r_0}^1(x) = \int_{r_0}^1 \frac{\partial}{\partial r} \mathbb{W}_r^1(x) dr = -\frac{1}{4\pi^2} \int_{r_0}^1 \mathbb{W}_r(\partial_1^2 x) \frac{dr}{r}.$$

It follows that

$$\|x\|_p \leq \|\mathbb{W}_{r_0}(x)\|_p - \frac{\log r_0}{4\pi^2} \|\partial_1^2 x\|_p.$$

We will show that

$$\|\mathbb{W}_{r_0}(x)\|_p \leq \frac{1}{2} \|x\|_p \quad (2.6)$$

for sufficiently small r_0 . Admitting this inequality, we then deduce

$$\|x\|_p \leq -\frac{\log r_0}{2\pi^2} \|\partial_1^2 x\|_p,$$

proving thus the first inequality of the lemma.

Let us show (2.6). By the assumption on x ,

$$\mathbb{W}_{r_0}^1(x) = \int_{\mathbb{T}} \mathbb{W}_{r_0}^1(z_1 w_1^{-1}) \pi_w(x) dw_1 = \int_{\mathbb{T}} (\mathbb{W}_{r_0}^1(z_1 w_1^{-1}) - 1) \pi_w(x) dw_1.$$

Hence,

$$\|\mathbb{W}_{r_0}(x)\|_p \leq \|\mathbb{W}_{r_0} - 1\|_{L_1(\mathbb{T})} \|x\|_p.$$

We must estimate the L_1 -norm of $\mathbb{W}_{r_0} - 1$. To this end, by (2.5) and with $r_0 = e^{-4\pi^2 \varepsilon_0}$, we have

$$\int_{\mathbb{T}} \mathbb{W}_{r_0}(z_1) dz_1 = \int_{-\frac{1}{2}}^{\frac{1}{2}} \mathbb{W}_{r_0}(e^{2\pi i s_1}) ds_1 = \sum_{m_1 \in \mathbb{Z}} \int_{-\frac{1}{2}}^{\frac{1}{2}} \mathbb{W}_{\varepsilon_0}^1(s_1 + m_1) ds_1 = 1.$$

Thus

$$\begin{aligned}
\|\mathbb{W}_{r_0} - 1\|_{L^1(\mathbb{T})} &\leq \sum_{m_1 \in \mathbb{Z}} \int_{-\frac{1}{2}}^{\frac{1}{2}} \left| W_{\varepsilon_0}^1(s_1 + m_1) - \int_{-\frac{1}{2}}^{\frac{1}{2}} W_{\varepsilon_0}^1(t_1 + m_1) dt_1 \right| ds_1 \\
&\leq \frac{4}{\sqrt{\pi}} \int_0^{\frac{1}{4\sqrt{\varepsilon_0}}} e^{-s_1^2} ds_1 \\
&\quad + \sum_{m_1 \in \mathbb{Z} \setminus \{0\}} \int_{-\frac{1}{2}}^{\frac{1}{2}} \int_{-\frac{1}{2}}^{\frac{1}{2}} |W_{\varepsilon_0}^1(s_1 + m_1) - W_{\varepsilon_0}^1(t_1 + m_1)| ds_1 dt_1 \\
&\leq \frac{1}{\sqrt{\pi\varepsilon_0}} + \sum_{m_1 \in \mathbb{Z} \setminus \{0\}} \frac{1}{4\sqrt{\pi\varepsilon_0}} e^{-\frac{m_1^2}{16\varepsilon_0}} \frac{|m_1|}{\varepsilon_0} \leq \frac{c}{\sqrt{\varepsilon_0}},
\end{aligned}$$

where c is a universal constant. Therefore, (2.6) follows if ε_0 is sufficiently large, or equivalently, if r_0 is sufficiently small. So the first inequality of the lemma is proved.

The second one is an immediate consequence of the first. Indeed, let $\mathbb{E}_{U_1, \dots, U_{d-1}}$ be the trace preserving conditional expectation from $\mathbb{T}_\theta^d$ onto the subalgebra generated by $(U_1, \dots, U_{d-1})$. Let $x' = \mathbb{E}_{U_1, \dots, U_{d-1}}(x)$ and $x_d = x - x'$. Then $m_d \neq 0$ whenever $\hat{x}_d(m) \neq 0$ for $m \in \mathbb{Z}^d$. Thus

$$\|x_d\|_p \leq c \|\partial_d^2 x_d\|_p = c \|\partial_d^2 x\|_p.$$

Since x' depends only on $(U_1, \dots, U_{d-1})$, an induction argument then yields the desired inequality. $\square$

Lemma 2.17. *The sequence $\{|x|_{W_p^k}\}_{k \geq 1}$ is increasing, up to constants. More precisely, there exists a constant $c_{d,k}$ such that*

$$|x|_{W_p^k} \leq c_{d,k} |x|_{W_p^{k+1}}, \quad \forall k \geq 1.$$

Proof. The proof is done easily by induction with the help of the previous two lemmas. Indeed, we have (assuming $\hat{x}(0) = 0$)

$$|x|_{W_p^1} \lesssim \sqrt{\|x\|_p |x|_{W_p^2}} \lesssim |x|_{W_p^2}.$$

Thus the assertion is proved for $k = 1$. Then induction gives the general case. $\square$

Proof of Theorem 2.12. By the preceding lemma, it remains to show $\|x\|_p \lesssim |x|_{W_p^1}$ for any $x \in W_p^1(\mathbb{T}_\theta^d)$ with $\hat{x}(0) = 0$. By approximation, we can assume that x is a polynomial. We proceed by induction on the dimension d . Consider first the case $d = 1$. Then

$$x = \sum_{m_1 \in \mathbb{Z} \setminus \{0\}} \hat{x}(m_1) U_1^{m_1}.$$

Define

$$y = \sum_{m_1 \in \mathbb{Z} \setminus \{0\}} \frac{1}{2\pi i m_1} \hat{x}(m_1) U_1^{m_1}.$$

Then $\partial_1 y = x$ and $\partial_1^2 y = \partial_1 x$. Thus Lemma 2.16 implies $\|x\|_p \lesssim \|\partial_1 x\|_p$.

Now consider a polynomial x in $(U_1, \dots, U_d)$. As in the proof of Lemma 2.16, let $x' = \mathbb{E}_{U_1, \dots, U_{d-1}}(x)$ and $x_d = x - x'$. The induction hypothesis implies

$$\|x'\|_p \lesssim |x'|_{W_p^1} = \left(\sum_{j=1}^{d-1} \|\mathbb{E}_{U_1, \dots, U_{d-1}}(\partial_j x)\|_p^p \right)^{\frac{1}{p}} \leq |x|_{W_p^1},$$

where we have used the commutation between $\mathbb{E}_{U_1, \dots, U_{d-1}}$ and the partial derivations.

To handle the term x_d , recalling that $m_d \neq 0$ whenever $\hat{x}_d(m) \neq 0$ for $m \in \mathbb{Z}^d$, we introduce

$$y_d = \sum_{m \in \mathbb{Z}^d} \frac{1}{2\pi i m_d} \hat{x}_d(m) U^m.$$

Then $\partial_d y_d = x_d$. So by (2.4) and Lemma 2.16,

$$\|x_d\|_p \lesssim \sqrt{\|y_d\|_p \|\partial_d^2 y_d\|_p} \lesssim \|\partial_d^2 y_d\|_p = \|\partial_d x_d\|_p.$$

Thus we are done, so the theorem is proved. $\square$

2.4 Lipschitz classes

In this section we show that $W_\infty^k(\mathbb{T}_\theta^d)$ is the quantum analogue of the classical Lipschitz class of order k . We will use the translation and difference operators introduced in Remark 2.1 and Definition 2.13. Note that for any positive integer k , $T_u^k = T_{ku}$ and Δ_u^k is the k th difference operator by $u \in \mathbb{R}^d$.

Definition 2.18. Let k be a positive integer and $x \in L_p(\mathbb{T}_\theta^d)$ with $1 \leq p \leq \infty$. The k th order modulus of L_p -smoothness of x is defined by

$$\omega_p^k(x, \varepsilon) = \sup_{0 < |u| \leq \varepsilon} \|\Delta_u^k x\|_p.$$

Remark 2.19. It is clear that $\omega_p^k(x, \varepsilon) \leq 2^k \|x\|_p$ and $\omega_p^k(x, \varepsilon)$ is nondecreasing in ε . On the other hand, $\omega_p^1(x, \varepsilon)$ is subadditive in ε ; for $k \geq 2$, $\omega_p^k(x, \varepsilon)$ is quasi subadditive in the sense that there exists a constant $c = c_k$ such that $\omega_p^k(x, \varepsilon + \eta) \leq c(\omega_p^k(x, \varepsilon) + \omega_p^k(x, \eta))$.

The following is the main result of this section. It shows that $W_\infty^k(\mathbb{T}_\theta^d)$ is the Lipschitz class of order k . We thus recover a result of Weaver [78, 79] for $k = 1$.

Theorem 2.20. For any $x \in W_p^k(\mathbb{T}_\theta^d)$, we have

$$\sup_{\varepsilon > 0} \frac{\omega_p^k(x, \varepsilon)}{\varepsilon^k} \approx |x|_{W_p^k},$$

where the equivalence constants depend only on d and k .

We require the following lemma for the proof.

Lemma 2.21. For any $x \in L_p(\mathbb{T}_\theta^d)$,

$$\lim_{\varepsilon \rightarrow 0} \frac{\omega_p^k(x, \varepsilon)}{\varepsilon^k} = \sup_{0 < \varepsilon \leq 1} \frac{\omega_p^k(x, \varepsilon)}{\varepsilon^k}.$$

Proof. The assertion for $k = 1$ is a common property of increasing and subadditive functions (in ε), and easy to check. Indeed, for any $0 < \varepsilon, \delta < 1$, choose $n \in \mathbb{N}$ such that $n\delta \leq \varepsilon < (n+1)\delta$. Then

$$\frac{\omega_p^1(x, \varepsilon)}{\varepsilon} \leq \frac{n+1}{n} \frac{\omega_p^1(x, \delta)}{\delta},$$

whence the result for $k = 1$. The general case is treated in the same way. Instead of being subadditive, $\omega_p^k(x, \varepsilon)$ is quasi subadditive in the sense that $\omega_p^k(x, n\varepsilon) \leq n^k \omega_p^k(x, \varepsilon)$ for any $n \in \mathbb{N}$. The latter follows immediately from

$$d_{nu}^k = \left(\sum_{j=0}^{n-1} e_{ju} \right)^k d_u^k, \quad \text{so} \quad \Delta_{nu}^k = \left(\sum_{j=0}^{n-1} T_{ju} \right)^k \Delta_u^k.$$

Thus the lemma is proved. $\square$

Proof of Theorem 2.20. If the assertion is proved for all $p < \infty$ with constants independent of p , the case $p = \infty$ will follow by letting $p \rightarrow \infty$. So we will assume $p < \infty$.

We first consider the case $k = 1$ whose proof contains all main ideas. As in the proof of Lemma 2.15, for any $u \in \mathbb{R}^d$, we have

$$d_u(\xi) = \int_0^1 \frac{\partial}{\partial t} d_{tu}(\xi) dt = \int_0^1 e_{tu}(\xi) (2\pi i u \cdot \xi) dt, \quad \xi \in \mathbb{R}^d.$$

In terms of Fourier multipliers, this yields

$$\Delta_u x = \int_0^1 T_{tu}(u \cdot \nabla x) dt,$$

where $u \cdot \nabla x = u_1 \partial_1 x + \cdots + u_d \partial_d x$. Since the translation T_{tu} is isometric, it then follows that

$$\|\Delta_u x\|_p \leq |u| \left\| (|\partial_1 x|^2 + \cdots + |\partial_d x|^2)^{\frac{1}{2}} \right\|_p \stackrel{\text{def}}{=} |u| \|\nabla x\|_p, \quad (2.7)$$

whence

$$\lim_{\varepsilon \rightarrow 0} \frac{\omega_p^1(x, \varepsilon)}{\varepsilon} \leq \|\nabla x\|_p.$$

To show the converse inequality, by the density of $\mathcal{P}_\theta$ in $W_p^k(\mathbb{T}_\theta^d)$ (see Proposition 2.7), we can assume that x is a polynomial. Given $u \in \mathbb{R}^d$ define ϕ on $\mathbb{R}^d$ by

$$\phi(\xi) = d_u(\xi) - \frac{\partial}{\partial t} d_{tu}(\xi) \Big|_{t=0}.$$

Then the Fourier multiplier on $\mathbb{T}_\theta^d$ associated to ϕ is

$$\tilde{\phi} * x = \Delta_u x - u \cdot \nabla x.$$

Thus if $|u| = \varepsilon$,

$$\|u \cdot \nabla x\|_p \leq \omega_p(x, \varepsilon) + \sup_{|u|=\varepsilon} \|\Delta_u x - u \cdot \nabla x\|_p.$$

Since x is a polynomial,

$$\lim_{\varepsilon \rightarrow 0} \sup_{|u|=\varepsilon} \frac{\|\Delta_u x - u \cdot \nabla x\|_p}{\varepsilon} = 0.$$

For $u = (\varepsilon, 0, \dots, 0)$, we then deduce

$$\|\partial_1 x\|_p \leq \lim_{\varepsilon \rightarrow 0} \frac{\omega_p(x, \varepsilon)}{\varepsilon}.$$

Hence the desired assertion for $k = 1$ is proved.

Now we consider the case $k > 1$. (2.7) can be easily iterated as follows:

$$\begin{aligned} \|\Delta_u^k x\|_p &\leq |u| \sum_{j=1}^d \|\partial_j \Delta_u^{k-1} x\|_p = |u| \sum_{j=1}^d \|\Delta_u^{k-1} \partial_j x\|_p \\ &\leq |u|^k \sum_{|m|_1=k} \|\Delta^m x\|_p \approx |u|^k |x|_{W_p^k}. \end{aligned}$$

So

$$\sup_{\varepsilon > 0} \frac{\omega_p^k(x, \varepsilon)}{\varepsilon^k} \lesssim |x|_{W_p^k}.$$

The converse inequality is proved similarly to the case $k = 1$. Let $m \in \mathbb{N}_0^d$ with $|m|_1 = k$. For each j with $m_j > 0$, using the Taylor formula of the function $d_{\varepsilon e_j}$ (recalling that $(e_1, \dots, e_d)$ denotes the canonical basis of $\mathbb{R}^d$), we get

$$\Delta_{\varepsilon e_j}^{m_j} x = \varepsilon^{m_j} \partial_j^{m_j} x + o(\varepsilon^{m_j}),$$

which implies

$$\prod_{j=0}^d \Delta_{\varepsilon e_j}^{m_j} x = \varepsilon^k D^m x + o(\varepsilon^k) \quad \text{as } \varepsilon \rightarrow 0.$$

Thus by the next lemma, we deduce

$$\|D^m x\|_p \leq \varepsilon^{-k} \left\| \prod_{j=0}^d \Delta_{\varepsilon e_j}^{m_j} x \right\|_p + o(1) \lesssim \varepsilon^{-k} \omega_p^k(x, \varepsilon) + o(1),$$

whence the desired converse inequality by letting $\varepsilon \rightarrow 0$. So the theorem is proved modulo the next lemma. $\square$

Lemma 2.22. *Let $u_1, \dots, u_k \in \mathbb{R}^d$. Then*

$$\Delta_{u_1} \cdots \Delta_{u_k} = \sum_{D \subset \{1, \dots, k\}} (-1)^{|D|} T_{\bar{u}_D} \Delta_{u_D}^k,$$

where the sum runs over all subsets of $\{1, \dots, k\}$, and where

$$\bar{u}_D = \sum_{j \in D} u_j, \quad u_D = \sum_{j \in D} \frac{1}{j} u_j.$$

Consequently, for $\varepsilon > 0$ and $x \in L_p(\mathbb{T}_\theta^d)$,

$$\sup_{|u_1| \leq \varepsilon, \dots, |u_k| \leq \varepsilon} \|\Delta_{u_1} \cdots \Delta_{u_k} x\|_p \approx \omega_p^k(x, \varepsilon).$$

Proof. This is a well-known lemma in the classical setting (see [6, Lemma 5.4.11]). We outline its proof for the convenience of the reader. The above equality is equivalent to the corresponding one with Δ_u and T_u replaced by d_u and e_u , respectively. Setting

$$v = \sum_{\ell \in D} \ell u_\ell \quad \text{and} \quad w = - \sum_{\ell \in D} u_\ell,$$

for each $0 \leq j \leq k$, we have

$$\begin{aligned} \prod_{\ell=1}^k d_{(\ell-j)u_\ell} &= \sum_{D \subset \{1, \dots, k\}} (-1)^{k-|D|} \prod_{\ell \in D} e_{(\ell-j)u_\ell} \\ &= \sum_{D \subset \{1, \dots, k\}} (-1)^{k-|D|} e_v (e_w)^j. \end{aligned}$$

The left hand side is nonzero only for $j = 0$. Multiply by $(-1)^{k-j} \binom{k}{j}$ and sum over $0 \leq j \leq k$; then replacing u_ℓ by $\frac{u_\ell}{\ell}$ gives the desired equality. $\square$

Remark 2.23. It might be interesting to note that in the commutative case, the proof of Theorem 2.20 shows

$$\sup_{0 < \varepsilon \leq 1} \frac{\omega_p(x, \varepsilon)}{\varepsilon} = \lim_{\varepsilon \rightarrow 0} \frac{\omega_p(x, \varepsilon)}{\varepsilon} = \|\nabla x\|_p.$$

So Lemma 2.21 is not needed in this case.

2.5 The link with the classical Sobolev spaces

The transference enables us to establish a strong link between the quantum Sobolev spaces defined previously and the vector-valued Sobolev spaces on $\mathbb{T}^d$. Note that the theory of vector-valued Sobolev spaces is well-known and can be found in many books on the subject (see, for instance, [2]). Here we just recall some basic notions. In the sequel, X will always denote a (complex) Banach space.

Let $\mathcal{S}(\mathbb{T}^d; X)$ be the space of X -valued infinitely differentiable functions on $\mathbb{T}^d$ with the standard Fréchet topology, and $\mathcal{S}'(\mathbb{T}^d; X)$ be the space of continuous linear maps from $\mathcal{S}(\mathbb{T}^d)$ to X . All operations on $\mathcal{S}(\mathbb{T}^d)$ such as derivations, convolution and Fourier transform transfer to $\mathcal{S}'(\mathbb{T}^d; X)$ in the usual way. On the other hand, $L_p(\mathbb{T}^d; X)$ naturally embeds into $\mathcal{S}'(\mathbb{T}^d; X)$ for $1 \leq p \leq \infty$, where $L_p(\mathbb{T}^d; X)$ stands for the space of strongly p -integrable functions from $\mathbb{T}^d$ to X .

Definition 2.24. Let $1 \leq p \leq \infty$, $k \in \mathbb{N}$ and $\alpha \in \mathbb{R}$.

(i) The X -valued Sobolev space of order k is

$$W_p^k(\mathbb{T}^d; X) = \{f \in \mathcal{S}'(\mathbb{T}^d; X) : D^m f \in L_p(\mathbb{T}^d; X) \text{ for each } m \in \mathbb{N}_0^d \text{ with } |m|_1 \leq k\}$$

equipped with the norm

$$\|f\|_{W_p^k} = \left(\sum_{0 \leq |m|_1 \leq k} \|D^m f\|_{L_p(\mathbb{T}^d; X)}^p \right)^{\frac{1}{p}}.$$

(ii) The X -valued potential Sobolev space of order α is

$$H_p^\alpha(\mathbb{T}^d; X) = \{f \in \mathcal{S}'(\mathbb{T}^d; X) : J^\alpha f \in L_p(\mathbb{T}^d; X)\}$$

equipped with the norm

$$\|f\|_{H_p^\alpha} = \|J^\alpha f\|_{L_p(\mathbb{T}^d; X)}.$$

Remark 2.25. There exists a parallel theory of vector-valued Sobolev spaces on $\mathbb{R}^d$. In fact, a majority of the literature on the subject is devoted to the case of $\mathbb{R}^d$ which is simpler from the pointview of treatment. The corresponding spaces are $W_p^k(\mathbb{R}^d; X)$ and $H_p^\alpha(\mathbb{R}^d; X)$. They are subspaces of $\mathcal{S}'(\mathbb{R}^d; X)$. The latter is the space of X -valued distributions on $\mathbb{R}^d$, that is, the space of continuous linear maps from $\mathcal{S}(\mathbb{R}^d)$ to X . We will sometimes use the space $\mathcal{S}(\mathbb{R}^d; X)$ of X -valued Schwartz functions on $\mathbb{R}^d$. We set $W_p^k(\mathbb{R}^d) = W_p^k(\mathbb{R}^d; \mathbb{C})$ and $H_p^k(\mathbb{R}^d) = H_p^k(\mathbb{R}^d; \mathbb{C})$.

The properties of the Sobolev spaces on $\mathbb{T}_\theta^d$ in the previous sections also hold for the present setting. For instance, the proof of Proposition 2.9 and Lemma 1.6 give the following well-known result:

Remark 2.26. Let X be a UMD space. Then $W_p^k(\mathbb{T}^d; X) = H_p^k(\mathbb{T}^d; X)$ with equivalent norms for $1 < p < \infty$.

Let us also mention that Theorem 2.12, the Poincaré inequality, transfers to the vector-valued case too. It seems that this result does not appear in literature but it should be known to experts. We record it explicitly here.

Theorem 2.27. *Let X be a Banach space and $1 \leq p \leq \infty$, $k \in \mathbb{N}$. Then*

$$\left(\|\widehat{f}(0)\|_X^p + \sum_{m \in \mathbb{N}_0^d, |m|_1=k} \|D^m f\|_{L_p(\mathbb{T}^d; X)}^p \right)^{\frac{1}{p}}$$

is an equivalent norm on $W_p^k(\mathbb{T}^d; X)$ with relevant constants depending only on d and k .

Now we use the transference method in Corollary 1.2. It is clear that the map $x \mapsto \tilde{x}$ there commutes with ∂_j , that is, $\partial_j \tilde{x} = \widetilde{\partial_j x}$ (noting that the ∂_j on the left-hand side is the j th partial derivation on $\mathbb{T}^d$ and that on the right-hand side is the one on $\mathbb{T}_\theta^d$). On the other hand, the expectation in that corollary commutes with ∂_j too. We then deduce the following:

Proposition 2.28. *Let $1 \leq p \leq \infty$. The map $x \mapsto \tilde{x}$ is an isometric embedding from $W_p^k(\mathbb{T}_\theta^d)$ and $H_p^\alpha(\mathbb{T}_\theta^d)$ into $W_p^k(\mathbb{T}^d; L_p(\mathbb{T}_\theta^d))$ and $H_p^\alpha(\mathbb{T}^d; L_p(\mathbb{T}_\theta^d))$, respectively. Moreover, the ranges of these embeddings are 1-complemented in their respective spaces.*

This result allows us to reduce many problems about $W_p^k(\mathbb{T}_\theta^d)$ to the corresponding ones about $W_p^k(\mathbb{T}^d; L_p(\mathbb{T}_\theta^d))$. For example, we could deduce the properties of $W_p^k(\mathbb{T}_\theta^d)$ in the preceding sections from those of $W_p^k(\mathbb{T}^d; L_p(\mathbb{T}_\theta^d))$. But we have chosen to work directly on $\mathbb{T}_\theta^d$ for the following reasons. It is more desirable to develop an intrinsic quantum theory, so we work directly on $\mathbb{T}_\theta^d$ whenever possible. On the other hand, the existing literature on vector-valued Sobolev spaces often concerns the case of $\mathbb{R}^d$, for instance, there exist few publications on periodic Besov or Triebel-Lizorkin spaces. In order to use existing results, we have to transfer them from $\mathbb{R}^d$ to $\mathbb{T}^d$. However, although it is often easy, this transfer may not be obvious at all, which is the case for Hardy spaces treated in [17] and recalled in section 1.4. This difficulty will appear again later for Besov and Triebel-Lizorkin spaces.

Remark 2.29. The preceding discussion on vector-valued Sobolev spaces on $\mathbb{T}^d$ can be also transferred to the quantum case. We have seen in section 1.3 that all noncommutative L_p -spaces are equipped with their natural operator space structure. Thus $W_p^k(\mathbb{T}_\theta^d)$ and $H_p^\alpha(\mathbb{T}_\theta^d)$ becomes operator spaces in the natural way. More generally, given an operator space E , following Pisier [54], we define the E -valued noncommutative L_p -space $L_p(\mathbb{T}_\theta^d; E)$ (recalling that $\mathbb{T}_\theta^d$ is an injective von Neumann algebra). Similarly, we define the E -valued distribution space $\mathcal{S}'(\mathbb{T}_\theta^d; E)$ that consists of continuous linear maps from $\mathcal{S}(\mathbb{T}_\theta^d)$ to E . Then as in Definition 2.6, we define the corresponding Sobolev spaces $W_p^k(\mathbb{T}_\theta^d; E)$ and $H_p^\alpha(\mathbb{T}_\theta^d; E)$. Almost all previous results remain valid in this vector-valued setting since all Fourier multipliers used in their proofs are c.b. maps. For instance, Theorem 2.9 (or Remark 2.26) now becomes $H_p^k(\mathbb{T}_\theta^d; E) = W_p^k(\mathbb{T}_\theta^d; E)$ for any $1 < p < \infty$ and any OUMD space E (OUMD is the operator space version of UMD; see [54]). Note that we recover $W_p^k(\mathbb{T}^d; E)$ and $H_p^\alpha(\mathbb{T}^d; E)$ if $\theta = 0$ and if E is equipped with its minimal operator space structure.

Chapter 3

Besov spaces

We study Besov spaces on $\mathbb{T}_\theta^d$ in this chapter. The first section presents the relevant definitions and some basic properties of these spaces. The second one shows a general characterization of them. This is the most technical part of the chapter. The formulation of our characterization is very similar to Triebel's classical theorem. Although modeled on Triebel's pattern, our proof is harder than his. The main difficulty is due to the unavailability in the noncommutative setting of the usual techniques involving maximal functions which play an important role in the study of the classical Besov and Triebel-Lizorkin spaces. Like for the Sobolev spaces in the previous chapter, Fourier multipliers are our main tool. We then concretize this general characterization by means of Poisson, heat kernels and differences. We would like to point out that when $\theta = 0$ (the commutative case), these characterizations (except that by differences) improve the corresponding ones in the classical case. Using the characterization by differences, we extend a recent result of Bourgain, Brézis and Mironescu on the limits of Besov norms to the quantum setting. The chapter ends with a section on vector-valued Besov spaces on $\mathbb{T}^d$.

3.1 Definitions and basic properties

We will use Littlewood-Paley decompositions as in the classical theory. Let φ be a Schwartz function on $\mathbb{R}^d$ such that

$$\begin{cases} \text{supp } \varphi \subset \{\xi : 2^{-1} \leq |\xi| \leq 2\}, \\ \varphi > 0 \text{ on } \{\xi : 2^{-1} < |\xi| < 2\}, \\ \sum_{k \in \mathbb{Z}} \varphi(2^{-k}\xi) = 1, \xi \neq 0. \end{cases} \quad (3.1)$$

Note that if $m \in \mathbb{Z}^d$ with $m \neq 0$, $\varphi(2^{-k}m) = 0$ for all $k < 0$, so

$$\sum_{k \geq 0} \varphi(2^{-k}m) = 1, \quad m \in \mathbb{Z}^d \setminus \{0\}.$$

On the other hand, the support of the function $\varphi(2^{-k}\cdot)$ is equal to $\{\xi : 2^{k-1} \leq |\xi| \leq 2^{k+1}\}$, thus $\text{supp } \varphi(2^{-k}\cdot) \cap \text{supp } \varphi(2^{-j}\cdot) = \emptyset$ whenever $|j - k| \geq 2$; consequently,

$$\varphi(2^{-k}\cdot) = \varphi(2^{-k}\cdot) \sum_{j=k-1}^{k+1} \varphi(2^{-j}\cdot), \quad k \geq 0. \quad (3.2)$$

Therefore, the sequence $\{\varphi(2^{-k}\cdot)\}_{k \geq 0}$ is a Littlewood-Paley decomposition of $\mathbb{T}^d$, modulo constant functions.

The Fourier multiplier on $\mathcal{S}'(\mathbb{T}_\theta^d)$ with symbol $\varphi(2^{-k}\cdot)$ is denoted by $x \mapsto \tilde{\varphi}_k * x$:

$$\tilde{\varphi}_k * x = \sum_{m \in \mathbb{Z}^d} \varphi(2^{-k}m) \hat{x}(m) U^m.$$

As noted in section 1.3, the convolution here has the usual sense. Indeed, let φ_k be the function whose Fourier transform is equal to $\varphi(2^{-k}\cdot)$, and let $\tilde{\varphi}_k$ be its 1-periodization, that is,

$$\tilde{\varphi}_k(s) = \sum_{m \in \mathbb{Z}^d} \varphi_k(s + m).$$

Viewed as a function on $\mathbb{T}^d$ by our convention, $\tilde{\varphi}_k$ admits the following Fourier series:

$$\tilde{\varphi}_k(z) = \sum_{m \in \mathbb{Z}^d} \varphi(2^{-k}m) z^m.$$

Thus for any distribution f on $\mathbb{T}^d$,

$$\tilde{\varphi}_k * f(z) = \sum_{m \in \mathbb{Z}^d} \varphi(2^{-k}m) \hat{f}(m) z^m.$$

We will maintain the above notation throughout the remainder of the paper.

We can now start our study of quantum Besov spaces.

Definition 3.1. Let $1 \leq p, q \leq \infty$ and $\alpha \in \mathbb{R}$. The associated Besov space on $\mathbb{T}_\theta^d$ is defined by

$$B_{p,q}^\alpha(\mathbb{T}_\theta^d) = \{x \in \mathcal{S}'(\mathbb{T}_\theta^d) : \|x\|_{B_{p,q}^\alpha} < \infty\},$$

where

$$\|x\|_{B_{p,q}^\alpha} = \left(|\hat{x}(0)|^q + \sum_{k \geq 0} 2^{qk\alpha} \|\tilde{\varphi}_k * x\|_p^q \right)^{\frac{1}{q}}.$$

Let $B_{p,c_0}^\alpha(\mathbb{T}_\theta^d)$ be the subspace of $B_{p,\infty}^\alpha(\mathbb{T}_\theta^d)$ consisting of all x such that $2^{k\alpha} \|\tilde{\varphi}_k * x\|_p \rightarrow 0$ as $k \rightarrow \infty$.

Remark 3.2. The Besov spaces defined above are independent of the choice of the function φ , up to equivalent norms. More generally, let $\{\psi^{(k)}\}_{k \geq 0}$ be a sequence of Schwartz functions such that

$$\begin{cases} \text{supp } \psi^{(k)} \subset \{\xi : 2^{k-1} \leq |\xi| \leq 2^{k+1}\}, \\ \sup_{k \geq 0} \|\mathcal{F}^{-1}(\psi^{(k)})\|_1 < \infty, \\ \sum_{k \geq 0} \psi^{(k)}(m) = 1, \forall m \in \mathbb{Z}^d \setminus \{0\}. \end{cases}$$

Let $\psi_k = \mathcal{F}^{-1}(\psi^{(k)})$ and $\tilde{\psi}_k$ be the periodization of ψ_k . Then

$$\|x\|_{B_{p,q}^\alpha} \approx \left(|\hat{x}(0)|^q + \sum_{k \geq 0} 2^{qk\alpha} \|\tilde{\psi}_k * x\|_p^q \right)^{\frac{1}{q}}.$$

Let us justify this remark. By the discussion leading to (3.2), we have (with $\tilde{\varphi}_{-1} = 0$)

$$\tilde{\psi}_k * x = \sum_{j=k-1}^{k+1} \tilde{\psi}_k * \tilde{\varphi}_j * x.$$

By Lemma 1.7,

$$\|\tilde{\psi}_k * x\|_p \lesssim \sum_{j=k-1}^{k+1} \|\tilde{\varphi}_j * x\|_p.$$

It then follows that

$$\left(|\hat{x}(0)|^q + \sum_{k \geq 0} 2^{qk\alpha} \|\tilde{\psi}_k * x\|_p^q\right)^{\frac{1}{q}} \lesssim \left(|\hat{x}(0)|^q + \sum_{k \geq 0} 2^{qk\alpha} \|\tilde{\varphi}_k * x\|_p^q\right)^{\frac{1}{q}}.$$

The reverse inequality is proved similarly.

Proposition 3.3. *Let $1 \leq p, q \leq \infty$ and $\alpha \in \mathbb{R}$.*

- (i) $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ is a Banach space.
- (ii) $B_{p,q}^\alpha(\mathbb{T}_\theta^d) \subset B_{p,r}^\alpha(\mathbb{T}_\theta^d)$ for $r > q$ and $B_{p,q}^\alpha(\mathbb{T}_\theta^d) \subset B_{p,r}^\beta(\mathbb{T}_\theta^d)$ for $\beta < \alpha$ and $1 \leq r \leq \infty$.
- (iii) $\mathcal{P}_\theta$ is dense in $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ and $B_{p,c_0}^\alpha(\mathbb{T}_\theta^d)$ for $1 \leq p \leq \infty$ and $1 \leq q < \infty$.
- (iv) The dual space of $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ coincides isomorphically with $B_{p',q'}^{-\alpha}(\mathbb{T}_\theta^d)$ for $1 \leq p \leq \infty$ and $1 \leq q < \infty$, where p' denotes the conjugate index of p . Moreover, the dual space of $B_{p,c_0}^\alpha(\mathbb{T}_\theta^d)$ coincides isomorphically with $B_{p',1}^{-\alpha}(\mathbb{T}_\theta^d)$.

Proof. (i) To show the completeness of $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$, let $\{x_n\}$ be a Cauchy sequence in $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$. Then $\{\hat{x}_n(0)\}$ converges to some $\hat{y}(0)$ in $\mathbb{C}$, and for every $k \geq 0$, $\{\tilde{\varphi}_k * x_n\}$ converges to some y_k in $L_p(\mathbb{T}_\theta^d)$. Let

$$y = \hat{y}(0) + \sum_{k \geq 0} y_k.$$

Since $\hat{y}_k$ is supported in $\{m \in \mathbb{Z}^d : 2^{k-1} < |m| < 2^{k+1}\}$, the above series converges in $\mathcal{S}'(\mathbb{T}_\theta^d)$. On the other hand, by (3.2), as $n \rightarrow \infty$, we have

$$\tilde{\varphi}_k * x_n = \sum_{j=k-1}^{k+1} \tilde{\varphi}_k * \tilde{\varphi}_j * x_n \rightarrow \sum_{j=k-1}^{k+1} \tilde{\varphi}_k * y_j = \tilde{\varphi}_k * y.$$

We then deduce that $y \in B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ and $x_n \rightarrow y$ in $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$.

(ii) is obvious.

(iii) We only show the density of $\mathcal{P}_\theta$ in $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ for finite q . For $N \in \mathbb{N}$ let

$$x_N = \hat{x}(0) + \sum_{j=0}^N \tilde{\varphi}_j * x.$$

Then by (3.1), $\tilde{\varphi}_k * (x - x_N) = 0$ for $k \leq N-1$, $\tilde{\varphi}_k * (x - x_N) = \tilde{\varphi}_k * x$ for $k > N+1$, and $\tilde{\varphi}_N * (x - x_N) = \tilde{\varphi}_N * x - \tilde{\varphi}_N * \tilde{\varphi}_N * x$, $\tilde{\varphi}_{N+1} * (x - x_N) = \tilde{\varphi}_{N+1} * x - \tilde{\varphi}_{N+1} * \tilde{\varphi}_N * x$. We then deduce

$$\|x - x_N\|_{B_{p,q}^\alpha}^q \leq 2 \sum_{k \geq N} 2^{qk\alpha} \|\tilde{\varphi}_k * x\|_p^q \rightarrow 0 \text{ as } N \rightarrow \infty.$$

(iv) In this part, we view $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ as $B_{p,c_0}^\alpha(\mathbb{T}_\theta^d)$ when $q = \infty$. Let $y \in B_{p',q'}^{-\alpha}(\mathbb{T}_\theta^d)$. Define $\ell_y(x) = \tau(xy)$ for $x \in \mathcal{P}_\theta$. Then

$$\begin{aligned} |\ell_y(x)| &= |\widehat{x}(0)\widehat{y}(0) + \sum_{k \geq 0} \tau(\tilde{\varphi}_k * x \sum_{j=k-1}^{k+1} \tilde{\varphi}_j * y)| \\ &\leq |\widehat{x}(0)\widehat{y}(0)| + \sum_{k \geq 0} \|\tilde{\varphi}_k * x\|_p \left\| \sum_{j=k-1}^{k+1} \tilde{\varphi}_j * y \right\|_{p'} \\ &\lesssim \|x\|_{B_{p,q}^\alpha} \|y\|_{B_{p',q'}^{-\alpha}}. \end{aligned}$$

Thus by the density of $\mathcal{P}_\theta$ in $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$, ℓ_y defines a continuous functional on $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$.

To prove the converse, we need a more notation. Given a Banach space X , let $\ell_q^\alpha(X)$ be the weighted direct sum of $(\mathbb{C}, X, X, \dots)$ in the ℓ_q -sense, that is, this is the space of all sequences $(a, x_0, x_1, \dots)$ with $a \in \mathbb{C}$ and $x_k \in X$, equipped with the norm

$$\left(|a|^q + \sum_{k \geq 0} 2^{qk\alpha} \|x_k\|^q \right)^{\frac{1}{q}}.$$

If $q = \infty$, we replace $\ell_q^\alpha(X)$ by its subspace $c_0^\alpha(X)$ consisting of sequences $(a, x_0, x_1, \dots)$ such that $2^{k\alpha}\|x_k\| \rightarrow 0$ as $k \rightarrow \infty$. Recall that the dual space of $\ell_q^\alpha(X)$ is $\ell_{q'}^{-\alpha}(X^*)$. By definition, $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ embeds into $\ell_q^\alpha(L_p(\mathbb{T}_\theta^d))$ via $x \mapsto (\widehat{x}(0), \tilde{\varphi}_0 * x, \tilde{\varphi}_1 * x, \dots)$. Now let ℓ be a continuous functional on $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ for $p < \infty$. Then by the Hahn-Banach theorem, ℓ extends to a continuous functional on $\ell_q^\alpha(L_p(\mathbb{T}_\theta^d))$ of unit norm, so there exists a unit element $(b, y_0, y_1, \dots)$ belonging to $\ell_{q'}^{-\alpha}(L_{p'}(\mathbb{T}_\theta^d))$ such that

$$\ell(x) = b\widehat{x}(0) + \sum_{k \geq 0} \tau(y_k \tilde{\varphi}_k * x), \quad x \in B_{p,q}^\alpha(\mathbb{T}_\theta^d).$$

Let

$$y = b + \sum_{k \geq 0} (\tilde{\varphi}_{k-1} * y_k + \tilde{\varphi}_k * y_k + \tilde{\varphi}_{k+1} * y_k).$$

Then clearly $y \in B_{p',q'}^{-\alpha}(\mathbb{T}_\theta^d)$ and $\ell = \ell_y$ when $p < \infty$. The same argument works for $p = \infty$ too. Indeed, for ℓ as above, there exists a unit element $(b, y_0, y_1, \dots)$ belonging to $\ell_{q'}^{-\alpha}(L_\infty(\mathbb{T}_\theta^d)^*)$ such that

$$\ell(x) = b\widehat{x}(0) + \sum_{k \geq 0} \langle y_k, \tilde{\varphi}_k * x \rangle, \quad x \in B_{p,q}^\alpha(\mathbb{T}_\theta^d).$$

Let y be defined as above. Then y is still a distribution and

$$\left(|\widehat{y}(0)|^{q'} + \sum_{k \geq 0} 2^{q'k\alpha} \|\tilde{\varphi}_k * y\|_{L_\infty(\mathbb{T}_\theta^d)^*}^{q'} \right)^{\frac{1}{q'}} < \infty.$$

Since it is a polynomial, $\tilde{\varphi}_k * y$ belongs to $L_1(\mathbb{T}_\theta^d)$; and we have

$$\|\tilde{\varphi}_k * y\|_{L_\infty(\mathbb{T}_\theta^d)^*} = \|\tilde{\varphi}_k * y\|_{L_1(\mathbb{T}_\theta^d)}.$$

Thus we are done for $p = \infty$ too. $\square$

To proceed further, we require some elementary lemmas. Recall that $J_\alpha(\xi) = (1+|\xi|^2)^{\frac{\alpha}{2}}$ and $I_\alpha(\xi) = |\xi|^\alpha$.

Lemma 3.4. *Let $\alpha \in \mathbb{R}$ and $k \in \mathbb{N}_0$. Then*

$$\|\mathcal{F}^{-1}(J_\alpha \varphi_k)\|_1 \lesssim 2^{\alpha k} \quad \text{and} \quad \|\mathcal{F}^{-1}(I_\alpha \varphi_k)\|_1 \lesssim 2^{\alpha k}.$$

where the constants depend only on φ , α and d . Consequently, for $x \in L_p(\mathbb{T}_\theta^d)$ with $1 \leq p \leq \infty$,

$$\|J^\alpha(\tilde{\varphi}_k * x)\|_p \lesssim 2^{\alpha k} \|\tilde{\varphi}_k * x\|_p \quad \text{and} \quad \|I^\alpha(\tilde{\varphi}_k * x)\|_p \lesssim 2^{\alpha k} \|\tilde{\varphi}_k * x\|_p.$$

Proof. The first part is well-known and easy to check. Indeed,

$$\|\mathcal{F}^{-1}(J_\alpha \varphi_k)\|_1 = 2^{\alpha k} \|\mathcal{F}^{-1}((4^{-k} + |\cdot|^2)^{\frac{\alpha}{2}} \varphi)\|_1;$$

the function $(4^{-k} + |\cdot|^2)^{\frac{\alpha}{2}} \varphi$ is a Schwartz function supported by $\{\xi : 2^{-1} \leq |\xi| \leq 2\}$, whose all partial derivatives, up to a fixed order, are bounded uniformly in k , so

$$\sup_{k \geq 0} \|\mathcal{F}^{-1}((4^{-k} + |\cdot|^2)^{\frac{\alpha}{2}} \varphi)\|_1 < \infty.$$

Similarly,

$$\|\mathcal{F}^{-1}(I_\alpha \varphi_k)\|_1 = 2^{\alpha k} \|\mathcal{F}^{-1}(I_\alpha \varphi)\|_1.$$

Since $\varphi_k = \varphi_k(\varphi_{k-1} + \varphi_k + \varphi_{k+1})$, by Lemma 1.7, we obtain the second part. $\square$

Given $a \in \mathbb{R}_+$, we define $D_{i,a}(\xi) = (2\pi i \xi_i)^a$ for $\xi \in \mathbb{R}^d$, and D_i^a to be the associated Fourier multiplier on $\mathbb{T}_\theta^d$. We set $D_a = D_{1,a_1} \cdots D_{d,a_d}$ and $D^a = D_1^{a_1} \cdots D_d^{a_d}$ for any $a = (a_1, \dots, a_d) \in \mathbb{R}_+^d$. Note that if a is a positive integer, $D_i^a = \partial_i^a$, so there does not exist any conflict of notation. The following lemma is well-known. We include a sketch of proof for the reader's convenience (see the proof of Remark 1 in Section 2.4.1 of [74]).

Lemma 3.5. *Let ρ be a compactly supported infinitely differentiable function on $\mathbb{R}^d$. Assume $\sigma, \beta \in \mathbb{R}_+$ and $a \in \mathbb{R}_+^d$ such that $\sigma > \frac{d}{2}$, $\beta > \sigma - \frac{d}{2}$ and $|a|_1 > \sigma - \frac{d}{2}$. Then the functions $I_\beta \rho$ and $D_a \rho$ belong to $H_2^\sigma(\mathbb{R}^d)$.*

Proof. If σ is a positive integer, the assertion clearly holds in view of $H_2^\sigma(\mathbb{R}^d) = W_2^\sigma(\mathbb{R}^d)$. On the other hand, $I_\beta \rho \in L_2(\mathbb{R}^d) = H_2^0(\mathbb{R}^d)$ for $\beta > -\frac{d}{2}$. The general case follows by complex interpolation. Indeed, under the assumption on σ and β , we can choose $\sigma_1 \in \mathbb{N}$, $\beta_1, \beta_0 \in \mathbb{R}$ and $\eta \in (0, 1)$ such that

$$\sigma_1 > \sigma, \quad \beta_1 > \sigma_1 - \frac{d}{2}, \quad \beta_0 > -\frac{d}{2}, \quad \sigma = \eta \sigma_1, \quad \beta = (1 - \eta) \beta_0 + \eta \beta_1.$$

For a complex number z in the strip $\{z \in \mathbb{C} : 0 \leq \operatorname{Re}(z) \leq 1\}$ define

$$F_z(\xi) = e^{(z-\eta)^2} |\xi|^{\beta_0(1-z) + \beta_1 z} \rho(\xi).$$

Then

$$\sup_{b \in \mathbb{R}} \|F_{ib}\|_{L_2} \lesssim \|I_{\beta_0} \rho\|_{L_2} \quad \text{and} \quad \sup_{b \in \mathbb{R}} \|F_{1+ib}\|_{H_2^{\sigma_1}} \lesssim \|I_{\beta_1} \rho\|_{H_2^{\sigma_1}}.$$

It thus follows that

$$I_\beta \rho = F_\eta \in (L_2(\mathbb{R}^d), H_2^{\sigma_1}(\mathbb{R}^d))_\eta.$$

The second assertion is proved in the same way. $\square$

The usefulness of the previous lemma relies upon the following well-known fact.

Remark 3.6. Let $\sigma > \frac{d}{2}$ and $f \in H_2^\sigma(\mathbb{R}^d)$. Then

$$\|\mathcal{F}^{-1}(f)\|_1 \lesssim \|f\|_{H_2^\sigma}.$$

The verification is extremely easy:

$$\begin{aligned} \|\mathcal{F}^{-1}(f)\|_1 &= \int_{|s| \leq 1} |\mathcal{F}^{-1}(f)(s)| ds + \sum_{k \geq 0} \int_{2^k < |s| \leq 2^{k+1}} |\mathcal{F}^{-1}(f)(s)| ds \\ &\lesssim \left(\int_{|s| \leq 1} |\mathcal{F}^{-1}(f)(s)|^2 ds + \sum_{k \geq 0} 2^{2k\sigma} \int_{2^k < |s| \leq 2^{k+1}} |\mathcal{F}^{-1}(f)(s)|^2 ds \right)^{\frac{1}{2}} \\ &\approx \|f\|_{H_2^\sigma}. \end{aligned}$$

The following is the so-called reduction (or lifting) theorem of Besov spaces.

Theorem 3.7. Let $1 \leq p, q \leq \infty$, $\alpha \in \mathbb{R}$.

- (i) For any $\beta \in \mathbb{R}$, both J^β and I^β are isomorphisms between $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ and $B_{p,q}^{\alpha-\beta}(\mathbb{T}_\theta^d)$.
- (ii) Let $a \in \mathbb{R}_+^d$. If $x \in B_{p,q}^\alpha(\mathbb{T}_\theta^d)$, then $D^a x \in B_{p,q}^{\alpha-|a|_1}(\mathbb{T}_\theta^d)$ and

$$\|D^a x\|_{B_{p,q}^{\alpha-|a|_1}} \lesssim \|x\|_{B_{p,q}^\alpha}.$$

- (iii) Let $\beta > 0$. Then $x \in B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ iff $D_i^\beta x \in B_{p,q}^{\alpha-\beta}(\mathbb{T}_\theta^d)$ for all $i = 1, \dots, d$. Moreover, in this case,

$$\|x\|_{B_{p,q}^\alpha} \approx |\hat{x}(0)| + \sum_{i=1}^d \|D_i^\beta x\|_{B_{p,q}^{\alpha-\beta}}.$$

Proof. (i) Let $x \in B_{p,q}^s(\mathbb{T}_\theta^d)$ with $\hat{x}(0) = 0$. Then by Lemma 3.4,

$$\begin{aligned} \|J^\beta x\|_{B_{p,q}^{\alpha-\beta}} &= \left(\sum_{k \geq 0} (2^{k(\alpha-\beta)} \|J^\beta(\tilde{\varphi}_k * x)\|_p)^q \right)^{\frac{1}{q}} \\ &\lesssim \left(\sum_{k \geq 0} (2^{k\alpha} \|\tilde{\varphi}_k * x\|_p)^q \right)^{\frac{1}{q}} = \|x\|_{B_{p,q}^\alpha}. \end{aligned}$$

Thus J^β is bounded from $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ to $B_{p,q}^{\alpha-\beta}(\mathbb{T}_\theta^d)$, its inverse, which is $J^{-\beta}$, is bounded too. The case of I^β is treated similarly.

(ii) By Lemma 3.5 and Remark 3.6, we have

$$\|\mathcal{F}^{-1}(D_a \varphi_k)\|_1 = 2^{k|a|_1} \|\mathcal{F}^{-1}(D_a \varphi)\|_1 \lesssim 2^{k|a|_1}.$$

Consequently, by Lemma 1.7,

$$\|\tilde{\varphi}_k * D^a x\|_p \lesssim 2^{k|a|_1} \|\tilde{\varphi}_k * x\|_p, \forall j \geq 0,$$

whence

$$\|D^a x\|_{B_{p,q}^{\alpha-|a|_1}} \lesssim \|x\|_{B_{p,q}^\alpha}.$$

(iii) One implication is contained in (ii). To show the other, choose an infinitely differentiable function $\chi : \mathbb{R} \rightarrow \mathbb{R}_+$ such that $\chi(s) = 0$ if $|s| < \frac{1}{4\sqrt{d}}$ and $\chi(s) = 1$ if $|s| \geq \frac{1}{2\sqrt{d}}$. For $i = 1, \dots, d$, let χ_i on $\mathbb{R}^d$ be defined by

$$\chi_i(\xi) = \frac{1}{\chi(\xi_1)|\xi_1|^\beta + \dots + \chi(\xi_d)|\xi_d|^\beta} \frac{\chi(\xi_i)|\xi_i|^\beta}{(2\pi i \xi_i)^\beta}$$

whenever the first denominator is positive, which is the case when $|\xi| \geq 2^{-1}$. Then for any $k \geq 0$, $\chi_i \varphi_k$ is a well-defined infinitely differentiable function on $\mathbb{R}^d \setminus \{\xi : \xi_i = 0\}$. We have

$$\|\mathcal{F}^{-1}(\chi_i \varphi_k)\|_1 = 2^{-k\beta} \|\mathcal{F}^{-1}(\psi \varphi)\|_1,$$

where

$$\psi(\xi) = \frac{1}{\chi(2^k \xi_1)|\xi_1|^\beta + \cdots + \chi(2^k \xi_d)|\xi_d|^\beta} \frac{\chi(2^k \xi_i)|\xi_i|^\beta}{(2\pi i \xi_i)^\beta}.$$

The function $\psi \varphi$ is supported in $\{\xi : 2^{-1} \leq |\xi| \leq 2\}$. An inspection reveals that all its partial derivatives of order less than a fixed integer are bounded uniformly in k . It then follows that the L_1 -norm of $\mathcal{F}^{-1}(\psi \varphi)$ is majorized by a constant independent of k , so

$$\|\mathcal{F}^{-1}(\chi_i \varphi_k)\|_1 \lesssim 2^{-k\beta},$$

and by Lemma 1.7,

$$\|\tilde{\chi}_i * \tilde{\varphi}_k * D_i^\beta x\|_p \lesssim 2^{-k\beta} \|\tilde{\varphi}_k * D_i^\beta x\|_p.$$

Since

$$\varphi_k = \sum_{i=1}^d \chi_i D_{i,\beta} \varphi_k,$$

we deduce

$$\|\tilde{\varphi}_k * x\| \lesssim 2^{-k\beta} \sum_{i=1}^d \|\tilde{\varphi}_k * D_i^\beta x\|_p,$$

which implies

$$\|x\|_{B_{p,q}^\alpha} \lesssim |\hat{x}(0)| + \sum_{i=1}^d \|D_i^\beta x\|_{B_{p,q}^{\alpha-\beta}}.$$

Thus (iii) is proved. $\square$

The following result relates the Besov and potential Sobolev spaces.

Theorem 3.8. *Let $1 \leq p \leq \infty$ and $\alpha \in \mathbb{R}^d$. Then we have the following continuous inclusions:*

$$B_{p,\min(p,2)}^\alpha(\mathbb{T}_\theta^d) \subset H_p^\alpha(\mathbb{T}_\theta^d) \subset B_{p,\max(p,2)}^\alpha(\mathbb{T}_\theta^d).$$

Proof. By Propositions 2.7 and 3.7, we can assume $\alpha = 0$. In this case, $H_p^a(\mathbb{T}_\theta^d) = L_p(\mathbb{T}_\theta^d)$. Let x be a distribution on $\mathbb{T}_\theta^d$ with $\hat{x}(0) = 0$. Since

$$x = \sum_{k \geq 0} \tilde{\varphi}_k * x,$$

we see that the inclusion $B_{p,1}^0(\mathbb{T}_\theta^d) \subset L_p(\mathbb{T}_\theta^d)$ follows immediately from triangular inequality. On the other hand, the inequality

$$\|\tilde{\varphi}_k * x\|_p \lesssim \|x\|_p, \quad k \geq 0$$

yields the inclusion $L_p(\mathbb{T}_\theta^d) \subset B_{p,\infty}^0(\mathbb{T}_\theta^d)$. Both inclusions can be improved in the range $1 < p < \infty$.

Let us consider only the case $2 \leq p < \infty$. Then the inclusion $L_p(\mathbb{T}_\theta^d) \subset B_{p,p}^0(\mathbb{T}_\theta^d)$ can be easily proved by interpolation. Indeed, the two spaces coincide isometrically when $p = 2$. The other extreme case $p = \infty$ has been already proved. We then deduce the case $2 < p < \infty$ by complex interpolation and by embedding $B_{p,\infty}^0(\mathbb{T}_\theta^d)$ into $\ell_\infty(L_p(\mathbb{T}_\theta^d))$.

The converse inclusion $B_{p,2}^0(\mathbb{T}_\theta^d) \subset L_p(\mathbb{T}_\theta^d)$ is subtler. To show it, we use Hardy spaces and the equality $L_p(\mathbb{T}_\theta^d) = \mathcal{H}_p(\mathbb{T}_\theta^d)$ (see Lemma 1.9). Then we must show

$$\max(\|x\|_{\mathcal{H}_p^c}, \|x\|_{\mathcal{H}_p^r}) \lesssim \|x\|_{B_{p,2}^0}.$$

To this end, we appeal to Lemma 1.10. The function ψ there is now equal to φ . The associated square function of x is thus given by

$$s_\varphi^c(x) = \left(\sum_{k \geq 0} |\tilde{\varphi}_k * x|^2 \right)^{\frac{1}{2}}.$$

Recall the following well-known inequality

$$\left\| \left(\sum_{k \geq 0} |x_k|^2 \right)^{\frac{1}{2}} \right\|_p \leq \left(\sum_{k \geq 0} \|x_k\|_p^2 \right)^{\frac{1}{2}}$$

for $x_k \in L_p(\mathbb{T}_\theta^d)$ and $2 \leq p \leq \infty$. Note that this inequality is proved simply by the triangular inequality in $L_{\frac{p}{2}}(\mathbb{T}_\theta^d)$. Thus

$$\|x\|_{\mathcal{H}_p^c} \approx \|s_\varphi^c(x)\|_p \leq \left(\sum_{k \geq 0} \|\tilde{\varphi}_k * x\|_p^2 \right)^{\frac{1}{2}} = \|x\|_{B_{p,2}^0}.$$

Passing to adjoints, we get $\|x\|_{\mathcal{H}_p^r} \lesssim \|x\|_{B_{p,2}^0}$. Therefore, the desired inequality follows. $\square$

3.2 A general characterization

In this and next sections we extend some characterizations of the classical Besov spaces to the quantum setting. Our treatment follows Triebel [74] rather closely.

We give a general characterization in this section. We have observed in the previous section that the definition of the Besov spaces is independent of the choice of φ satisfying (3.1). We now show that φ can be replaced by more general functions. To state the required conditions, we introduce an auxiliary Schwartz function h such that

$$\text{supp } h \subset \{\xi \in \mathbb{R}^d : |\xi| \leq 4\} \quad \text{and} \quad h = 1 \text{ on } \{\xi \in \mathbb{R}^d : |\xi| \leq 2\}. \quad (3.3)$$

Let $\alpha_0, \alpha_1 \in \mathbb{R}$. Let ψ be an infinitely differentiable function on $\mathbb{R}^d \setminus \{0\}$ satisfying the following conditions

$$\begin{cases} |\psi| > 0 \text{ on } \{\xi : 2^{-1} \leq |\xi| \leq 2\}, \\ \mathcal{F}^{-1}(\psi h I_{-\alpha_1}) \in L_1(\mathbb{R}^d), \\ \sup_{j \in \mathbb{N}_0} 2^{-\alpha_0 j} \|\mathcal{F}^{-1}(\psi(2^j \cdot) \varphi)\|_1 < \infty. \end{cases} \quad (3.4)$$

The first nonvanishing condition above on ψ is a Tauberian condition. The integrability of the inverse Fourier transforms can be reduced to a handier criterion by means of the potential Sobolev space $H_2^\sigma(\mathbb{R}^d)$ with $\sigma > \frac{d}{2}$; see Remark 3.6.

We will use the same notation for ψ as for φ . In particular, ψ_k is the inverse Fourier transform of $\psi(2^{-k} \cdot)$ and $\tilde{\psi}_k$ is the Fourier multiplier on $\mathbb{T}_\theta^d$ with symbol $\psi(2^{-k} \cdot)$. It is to note that compared with [74, Theorem 2.5.1], we need not require $\alpha_1 > 0$ in the following theorem. This will have interesting consequences in the next section.

Theorem 3.9. *Let $1 \leq p, q \leq \infty$ and $\alpha \in \mathbb{R}$. Assume $\alpha_0 < \alpha < \alpha_1$. Let ψ satisfy the above assumption. Then a distribution x on $\mathbb{T}_\theta^d$ belongs to $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ iff*

$$\left(\sum_{k \geq 0} (2^{k\alpha} \|\tilde{\psi}_k * x\|_p)^q \right)^{\frac{1}{q}} < \infty.$$

If this is the case, then

$$\|x\|_{B_{p,q}^\alpha} \approx \left(|\hat{x}(0)|^q + \sum_{k \geq 0} (2^{k\alpha} \|\tilde{\psi}_k * x\|_p)^q \right)^{\frac{1}{q}} \quad (3.5)$$

with relevant constants depending only on $\varphi, \psi, \alpha, \alpha_0, \alpha_1$ and d .

Proof. We will follow the pattern of the proof of [74, Theorem 2.4.1]. Given a function f on $\mathbb{R}^d$, we will use the notation that $f^{(k)} = f(2^{-k} \cdot)$ for $k \geq 0$ and $f^{(k)} = 0$ for $k \leq -1$. Recall that f_k is the inverse Fourier transform of $f^{(k)}$ and $\tilde{f}_k$ is the 1-periodization of f_k :

$$\tilde{f}_k(s) = \sum_{m \in \mathbb{Z}^d} f_k(s + m).$$

In the following, we will fix a distribution x on $\mathbb{T}_\theta^d$. Without loss of generality, we assume $\hat{x}(0) = 0$. We will denote the right-hand side of (3.5) by $\|x\|_{B_{p,q}^{\alpha,\psi}}$ when it is finite. For clarity, we divide the proof into several steps.

Step 1. In the first two steps, we assume $x \in B_{p,q}^\alpha(\mathbb{T}_\theta^d)$. Let K be a positive integer to be determined later in step 3. By (3.1), we have

$$\psi^{(j)} = \sum_{k=0}^{\infty} \psi^{(j)} \varphi^{(k)} = \sum_{k=-\infty}^K \psi^{(j)} \varphi^{(j+k)} + \sum_{k=K}^{\infty} \psi^{(j)} \varphi^{(j+k)} \quad \text{on } \{\xi : |\xi| \geq 1\}.$$

Then

$$\tilde{\psi}_j * x = \sum_{k \leq K} \tilde{\psi}_j * \tilde{\varphi}_{j+k} * x + \sum_{k > K} \tilde{\psi}_j * \tilde{\varphi}_{j+k} * x. \quad (3.6)$$

For the moment, we do not care about the convergence issue of the second series above, which is postponed to the last step. Let $a_{j,k} = 2^{j\alpha} \|\tilde{\psi}_j * \tilde{\varphi}_{j+k} * x\|_p$. Then

$$\|x\|_{B_{p,q}^{\alpha,\psi}} \leq \left(\sum_{j=0}^{\infty} \left[\sum_{k \leq K} a_{j,k} \right]^q \right)^{\frac{1}{q}} + \left(\sum_{j=0}^{\infty} \left[\sum_{k > K} a_{j,k} \right]^q \right)^{\frac{1}{q}}. \quad (3.7)$$

We will treat the two sums on the right-hand side separately. For the first one, by the support assumption on φ and h , for $k \leq K$, we can write

$$\begin{aligned} \psi^{(j)}(\xi) \varphi^{(j+k)}(\xi) &= 2^{k\alpha_1} \frac{\psi^{(j)}(\xi)}{|2^{-j}\xi|^{\alpha_1}} h^{(j+K)}(\xi) |2^{-j-k}\xi|^{\alpha_1} \varphi^{(j+k)}(\xi) \\ &= 2^{k\alpha_1} \eta^{(j)}(\xi) \rho^{(j+k)}(\xi), \end{aligned} \quad (3.8)$$

where η and ρ are defined by

$$\eta(\xi) = \frac{\psi(\xi)}{|\xi|^{\alpha_1}} h^{(K)}(\xi) \quad \text{and} \quad \rho(\xi) = |\xi|^{\alpha_1} \varphi(\xi).$$

Note that $\mathcal{F}^{-1}(\eta)$ is integrable on $\mathbb{R}^d$. Indeed, write

$$\eta(\xi) = \frac{\psi(\xi)}{|\xi|^{\alpha_1}} h(\xi) + \frac{\psi(\xi)}{|\xi|^{\alpha_1}} (h^{(K)}(\xi) - h(\xi)). \quad (3.9)$$

By (3.4), the inverse Fourier transform of the first function on the right-hand side is integrable. The second one is an infinitely differentiable function with compact support, so its inverse Fourier transform is also integrable with L_1 -norm controlled by a constant depending only on ψ , h , α_1 and K . Therefore, Lemma 1.7 implies that each $\eta^{(j)}$ is a Fourier multiplier on $L_p(\mathbb{T}_\theta^d)$ for all $1 \leq p \leq \infty$ with norm controlled by a constant c_1 , depending only ψ , h , α_1 and K . Therefore,

$$a_{j,k} \leq c_1 2^{j\alpha+k\alpha_1} \|\tilde{\rho}_{j+k} * x\|_p = c_1 2^{k(\alpha_1-\alpha)} (2^{(j+k)\alpha} \|\tilde{\rho}_{j+k} * x\|_p). \quad (3.10)$$

Thus by triangular inequality and Lemma 3.4, we deduce

$$\begin{aligned} \left(\sum_{j=0}^{\infty} \left[\sum_{k \leq K} a_{j,k} \right]^q \right)^{\frac{1}{q}} &\leq c_1 \sum_{k \leq K} 2^{k(\alpha_1-\alpha)} \left(\sum_{j=-\infty}^{\infty} (2^{(j+k)\alpha} \|\tilde{\rho}_{j+k} * x\|_p)^q \right)^{\frac{1}{q}} \\ &= c_1 \sum_{k \leq K} 2^{k(\alpha_1-\alpha)} \left(\sum_{j=0}^{\infty} (2^{j\alpha} 2^{-j\alpha_1} \|I^{\alpha_1} \tilde{\varphi}_j * x\|_p)^q \right)^{\frac{1}{q}} \\ &\leq c'_1 \|x\|_{B_{p,q}^{\alpha}}, \end{aligned}$$

where c'_1 depends only ψ , h , K , α and α_1 .

Step 2. The second sum on the right-hand side of (3.7) is treated similarly. Like in step 1 and by (3.2), we write

$$\begin{aligned} \psi^{(j)}(\xi) \varphi^{(j+k)}(\xi) &= \frac{\psi^{(j)}(\xi)}{|2^{-j-k}\xi|^{\alpha_0}} (\varphi^{(j+k-1)} + \varphi^{(j+k)} + \varphi^{(j+k+1)})(\xi) |2^{-j-k}\xi|^{\alpha_0} \varphi^{(j+k)}(\xi) \\ &= \left[\frac{\psi(2^{-j-k} \cdot 2^k \xi)}{|2^{-j-k}\xi|^{\alpha_0}} H(2^{-j-k}\xi) \right] \rho^{(j+k)}(\xi), \end{aligned} \quad (3.11)$$

where $H = \varphi^{(-1)} + \varphi + \varphi^{(1)}$, and where ρ is now defined by

$$\rho(\xi) = |\xi|^{\alpha_0} \varphi(\xi).$$

The L_1 -norm of the inverse Fourier transform of the function

$$\frac{\psi(2^{-j-k} \cdot 2^k \xi)}{|2^{-j-k}\xi|^{\alpha_0}} H(2^{-j-k}\xi)$$

is equal to $\|\mathcal{F}^{-1}(I_{-\alpha_0} H \psi(2^k \cdot))\|_1$. Using Lemma 3.4, we see that the last norm is majorized by

$$\|\mathcal{F}^{-1}(\psi(2^k \cdot) H)\|_1 \lesssim \|\mathcal{F}^{-1}(\psi(2^k \cdot) \varphi)\|_1.$$

Then, using (3.4), for $k > K$ we get

$$a_{j,k} \leq c_2 2^{k(\alpha_0-\alpha)} (2^{(j+k)\alpha} \|\tilde{\rho}_{j+k} * x\|_p), \quad (3.12)$$

where c_2 depends only on φ , α_0 and the supremum in (3.4). Thus as before, we get

$$\left(\sum_{j=0}^{\infty} \left[\sum_{k > K} a_{j,k} \right]^q \right)^{\frac{1}{q}} \leq c'_2 \frac{2^{K(\alpha_0-\alpha)}}{1 - 2^{\alpha_0-\alpha}} \|x\|_{B_{p,q}^{\alpha}},$$

which, together with the inequality obtained in step 1, yields

$$\|x\|_{B_{p,q}^{\alpha,\psi}} \lesssim \|x\|_{B_{p,q}^{\alpha}}.$$

Step 3. Now we prove the inequality reverse to the previous one. We first assume that x is a polynomial. We write

$$\varphi^{(j)} = \varphi^{(j)} h^{(j+K)} = \frac{\varphi^{(j)}}{\psi^{(j)}} h^{(j+K)} \psi^{(j)}. \quad (3.13)$$

The function $\varphi\psi^{-1}$ is an infinitely differentiable function with compact support, so its inverse Fourier transform belongs to $L_1(\mathbb{R}^d)$. Thus by Lemma 1.7,

$$\|\tilde{\varphi}_j * x\|_p \leq c_3 \|\tilde{h}_{j+K} * \tilde{\psi}_j * x\|_p,$$

where $c_3 = \|\mathcal{F}^{-1}(\varphi\psi^{-1})\|_1$. Hence,

$$\|x\|_{B_{p,q}^\alpha} \leq c_3 \left(\sum_{j=0}^{\infty} (2^{j\alpha} \|\tilde{h}_{j+K} * \tilde{\psi}_j * x\|_p)^q \right)^{\frac{1}{q}}.$$

To handle the right-hand side, we let $\lambda = 1 - h$ and write $h^{(j+K)}\psi^{(j)} = \psi^{(j)} - \lambda^{(j+K)}\psi^{(j)}$. Then

$$\left(\sum_{j=0}^{\infty} (2^{j\alpha} \|\tilde{h}_{j+K} * \tilde{\psi}_j * x\|_p)^q \right)^{\frac{1}{q}} \leq \|x\|_{B_{p,q}^{\alpha,\psi}} + \left(\sum_{j=0}^{\infty} (2^{j\alpha} \|\tilde{\lambda}_{j+K} * \tilde{\psi}_j * x\|_p)^q \right)^{\frac{1}{q}}.$$

Thus it remains to deal with the last sum. We do this as in the previous steps with ψ replaced by $\lambda\psi$, by writing

$$\lambda^{(j+K)}\psi^{(j)} = \sum_{k=-\infty}^{\infty} \lambda^{(j+K)}\psi^{(j)}\varphi^{(j+k)}.$$

The crucial point now is the fact that $\lambda^{(j+K)}\varphi^{(j+k)} = 0$ for all $k \leq K$ and all j . So

$$\lambda^{(j+K)}\psi^{(j)} = \sum_{k>K} \lambda^{(j+K)}\psi^{(j)}\varphi^{(j+k)},$$

that is, only the second sum on the right-hand side of (3.7) survives now:

$$\left(\sum_{j=0}^{\infty} (2^{j\alpha} \|\tilde{\lambda}_{j+K} * \tilde{\psi}_j * x\|_p)^q \right)^{\frac{1}{q}} \leq \left(\sum_{j=0}^{\infty} [2^{j\alpha} \sum_{k>K} \|\tilde{\lambda}_{j+K} * \tilde{\psi}_j * \tilde{\varphi}_{j+k} * x\|_p]^q \right)^{\frac{1}{q}}.$$

Let us reexamine the argument of step 2 and formulate (3.11) with $\lambda^{(K)}\psi$ instead of ψ . We then arrive at majorizing the norm $\|\mathcal{F}^{-1}(\lambda(2^{k-K}\cdot)\psi(2^k\cdot))\|_1$:

$$\|\mathcal{F}^{-1}(\lambda(2^{k-K}\cdot)\psi(2^k\cdot))\|_1 \leq \|\mathcal{F}^{-1}(\lambda)\|_1 \|\mathcal{F}^{-1}(\psi(2^k\cdot))\|_1.$$

Keeping the notation of step 2 and as for (3.12), we get

$$\|\tilde{\lambda}_{j+K} * \tilde{\psi}_j * \tilde{\varphi}_{j+k} * x\|_p \leq cc_2 2^{k(\alpha_0 - \alpha)} (2^{(j+k)\alpha} \|\tilde{\rho}_{j+k} * x\|_p),$$

where $c = \|\mathcal{F}^{-1}(\lambda)\|_1$. Thus

$$\left(\sum_{j=0}^{\infty} [2^{\alpha j} \sum_{k>K} \|\tilde{\lambda}_{j+K} * \tilde{\psi}_j * \tilde{\varphi}_{j+k} * x\|_p]^q \right)^{\frac{1}{q}} \leq cc'_2 \frac{2^{K(\alpha_0 - \alpha)}}{1 - 2^{\alpha_0 - \alpha}} \|x\|_{B_{p,q}^\alpha}.$$

Combining the preceding inequalities, we obtain

$$\|x\|_{B_{p,q}^\alpha} \leq c_3 \|x\|_{B_{p,q}^{\alpha,\psi}} + cc'_2 \frac{2^{K(\alpha_0-\alpha)}}{1-2^{\alpha_0-\alpha}} \|x\|_{B_{p,q}^\alpha}.$$

Choosing K so that

$$cc'_2 \frac{2^{K(\alpha_0-\alpha)}}{1-2^{\alpha_0-\alpha}} \leq \frac{1}{2},$$

we finally deduce

$$\|x\|_{B_{p,q}^\alpha} \leq 2c_3 \|x\|_{B_{p,q}^{\alpha,\psi}},$$

which shows (3.5) in case x is a polynomial.

The general case can be easily reduced to this special one. Indeed, assume $\|x\|_{B_{p,q}^{\alpha,\psi}} < \infty$. Then using the Fejér means F_N as in the proof of Proposition 2.7, we see that

$$\|F_N(x)\|_{B_{p,q}^{\alpha,\psi}} \leq \|x\|_{B_{p,q}^{\alpha,\psi}}.$$

Applying the above part already proved for polynomials yields

$$\|F_N(x)\|_{B_{p,q}^\alpha} \leq 2c_3 \|F_N(x)\|_{B_{p,q}^{\alpha,\psi}} \leq 2c_3 \|x\|_{B_{p,q}^{\alpha,\psi}},$$

However, it is easy to check that

$$\lim_{N \rightarrow \infty} \|F_N(x)\|_{B_{p,q}^\alpha} = \|x\|_{B_{p,q}^\alpha}.$$

We thus deduce (3.5) for general x , modulo the convergence problem on the second series of (3.6).

Step 4. We now settle up the convergence issue left previously. Each term $\tilde{\psi}_j * \tilde{\varphi}_{j+k} * x$ is a polynomial, so a distribution on $\mathbb{T}_\theta^d$. We must show that the series converges in $\mathcal{S}'(\mathbb{T}_\theta^d)$. By (3.12), for any $L > K$, by the Hölder inequality (with q' the conjugate index of q), we get

$$\begin{aligned} 2^{j\alpha} \sum_{k=K+1}^L \|\tilde{\psi}_j * \tilde{\varphi}_{j+k} * x\|_p &\leq c'_2 \sum_{k=K+1}^L 2^{k(\alpha_0-\alpha)} (2^{(j+k)\alpha} \|\tilde{\varphi}_{j+k} * x\|_p) \\ &\leq c'_2 R_{K,L} \left(\sum_{k=K+1}^L (2^{(j+k)\alpha} \|\tilde{\varphi}_{j+k} * x\|_p)^q \right)^{\frac{1}{q}} \\ &\lesssim c'_2 R_{K,L} \|x\|_{B_{p,q}^\alpha}, \end{aligned}$$

where

$$R_{K,L} = \left(\sum_{k=K+1}^L 2^{q'k(\alpha_0-\alpha)} \right)^{\frac{1}{q'}}.$$

Since $\alpha_0 < \alpha$, $R_{K,L} \rightarrow 0$ as K tends to ∞ . Thus the series $\sum_{k>K} \tilde{\psi}_j * \tilde{\varphi}_{j+k} * x$ converges in $L_p(\mathbb{T}_\theta^d)$, so in $\mathcal{S}'(\mathbb{T}_\theta^d)$ too. In the same way, we show that the series also converges in $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$. Hence, the proof of the theorem is complete. $\square$

Remark 3.10. The infinite differentiability of ψ can be substantially relaxed without changing the proof. Indeed, we have used this condition only once to insure that the inverse Fourier transform of the second term on the right-hand side of (3.9) is integrable. This integrability is guaranteed when ψ is continuously differentiable up to order $[\frac{d}{2}] + 1$. The latter condition can be replaced by the following slightly weaker one: there exists $\sigma > \frac{d}{2} + 1$ such that $\psi\eta \in H_2^\sigma(\mathbb{R}^d)$ for any compactly supported infinite differentiable function η which vanishes in a neighborhood of the origin.

The following is the continuous version of Theorem 3.9. We will use similar notation for continuous parameters: given $\varepsilon > 0$, ψ_ε denotes the function with Fourier transform $\psi^{(\varepsilon)} = \psi(\varepsilon \cdot)$, and $\tilde{\psi}_\varepsilon$ denotes the Fourier multiplier on $\mathbb{T}_\theta^d$ associated to $\psi^{(\varepsilon)}$.

Theorem 3.11. *Keep the assumption of the previous theorem. Then for any distribution x on $\mathbb{T}_\theta^d$,*

$$\|x\|_{B_{p,q}^\alpha} \approx \left(|\hat{x}(0)|^q + \int_0^1 \varepsilon^{-q\alpha} \|\tilde{\psi}_\varepsilon * x\|_p^q \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}}. \quad (3.14)$$

The above equivalence is understood in the sense that if one side is finite, so is the other, and the two are then equivalent with constants independent of x .

Proof. This proof is very similar to the previous one. Keeping the notation there, we will point out only the necessary changes. Let us first discretize the integral on the right-hand side of (3.14):

$$\int_0^1 (\varepsilon^{-\alpha} \|\tilde{\psi}_\varepsilon * x\|_p)^q \frac{d\varepsilon}{\varepsilon} \approx \sum_{j=0}^{\infty} 2^{jq\alpha} \int_{2^{-j-1}}^{2^{-j}} \|\tilde{\psi}_\varepsilon * x\|_p^q \frac{d\varepsilon}{\varepsilon}.$$

Now for $j \geq 0$ and $2^{-j-1} < \varepsilon \leq 2^{-j}$, we transfer (3.8) to the present setting:

$$\psi^{(\varepsilon)}(\xi) \varphi^{(j+k)}(\xi) = 2^{\alpha_1 k} \left[\frac{\psi(2^{-j} \cdot 2^j \varepsilon \xi)}{|2^{-j} \varepsilon \xi|^{\alpha_1}} h^{(j+K)}(\xi) \right] \rho^{(j+k)}(\xi).$$

We then must estimate the L_1 -norm of the inverse Fourier transform of the function in the brackets. It is equal to

$$\|\mathcal{F}^{-1}(I_{-\alpha_1} \psi(2^j \varepsilon \cdot) h^{(K)})\|_1 = \delta^{-\alpha_1} \|\mathcal{F}^{-1}(I_{-\alpha_1} \psi h(\delta 2^{-K} \cdot))\|_1,$$

where $\delta = 2^{-j} \varepsilon^{-1}$. The last norm is estimated as follows:

$$\begin{aligned} \|\mathcal{F}^{-1}(I_{-\alpha_1} \psi h(\delta 2^{-K} \cdot))\|_1 &\leq \|\mathcal{F}^{-1}(I_{-\alpha_1} \psi h)\|_1 + \|\mathcal{F}^{-1}(I_{-\alpha_1} \psi [h - h(\delta 2^{-K} \cdot)])\|_1 \\ &\leq \|\mathcal{F}^{-1}(I_{-\alpha_1} \psi h)\|_1 + \sup_{1 \leq \delta \leq 2} \|\mathcal{F}^{-1}(I_{-\alpha_1} \psi [h - h(\delta 2^{-K} \cdot)])\|_1. \end{aligned}$$

Note that the above supremum is finite since $I_{-\alpha_1} \psi [h - h(\delta 2^{-K} \cdot)]$ is a compactly supported infinite differentiable function and its inverse Fourier transform depends continuously on δ . It follows that for $k \leq K$ and $2^{-j-1} \leq \varepsilon \leq 2^{-j}$

$$2^{j\alpha} \|\tilde{\psi}_\varepsilon * \tilde{\varphi}_{j+k} * x\|_p \lesssim 2^{k(\alpha_1 - \alpha)} (2^{(j+k)\alpha} \|\tilde{\rho}_{j+k} * x\|_p), \quad (3.15)$$

which is the analogue of (3.10). Thus, we get the continuous analogue of the final inequality of step 1 in the preceding proof.

We can make similar modifications in step 2, and then show the second part. Hence, we have proved

$$\left(\int_0^1 (\varepsilon^{-\alpha} \|\tilde{\psi}_\varepsilon * x\|_p)^q \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}} \lesssim \|x\|_{B_{p,q}^\alpha}.$$

To show the converse inequality, we proceed as in step 3 above. By (3.4), there exists a constant $a > 2$ such that $\psi > 0$ on $\{\xi : a^{-1} \leq |\xi| \leq a\}$. Assume also $a \leq 2\sqrt{2}$. For $j \geq 0$ let $R_j = (a^{-1} 2^{-j-1}, a 2^{-j+1}]$. The R_j 's are disjoint subintervals of $(0, 1]$. Now we slightly modify (3.13) as follows:

$$\varphi^{(j)} = \varphi^{(j)} h^{(j+K)} = \frac{\varphi^{(j)}}{\psi^{(\varepsilon)}} h^{(j+K)} \psi^{(\varepsilon)}, \quad \varepsilon \in R_j. \quad (3.16)$$

Then

$$\|\mathcal{F}^{-1}\left(\frac{\varphi^{(j)}}{\psi^{(\varepsilon)}}\right)\|_1 = \|\mathcal{F}^{-1}\left(\frac{\varphi(2^{-j}\varepsilon^{-1}\cdot)}{\psi}\right)\|_1 \leq \sup_{2a^{-1} \leq \delta \leq a2^{-1}} \|\mathcal{F}^{-1}\left(\frac{\varphi^{(\delta)}}{\psi}\right)\|_1 < \infty.$$

Like in step 3, we deduce

$$\begin{aligned} \|x\|_{B_{p,q}^\alpha} &\lesssim \left(\sum_{j=0}^{\infty} (2^{j\alpha} \int_{R_j} \|\tilde{h}_{j+K} * \tilde{\psi}_\varepsilon * x\|_p)^q \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}} \\ &\lesssim \left(\int_0^1 (\varepsilon^{-\alpha} \|\tilde{\psi}_\varepsilon * x\|_p)^q \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}} + \left(\sum_{j=0}^{\infty} (2^{j\alpha} \int_{R_j} \|\tilde{h}_{j+K} * \tilde{\psi}_\varepsilon * x\|_p)^q \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}}. \end{aligned}$$

The remaining of the proof follows step 3 with necessary modifications as in the first part. $\square$

Remark 3.12. Theorems 3.9 and 3.11 admit analogous characterizations for $B_{p,c_0}^\alpha(\mathbb{T}_\theta^d)$ too. For example, a distribution x on $\mathbb{T}_\theta^d$ belongs to $B_{p,c_0}^\alpha(\mathbb{T}_\theta^d)$ iff

$$\lim_{\varepsilon \rightarrow 0} \frac{\|\tilde{\psi}_\varepsilon * x\|_p}{\varepsilon^\alpha} = 0.$$

This easily follows from Theorem 3.11 for $q = \infty$. The same remark applies to the characterizations by the Poisson, heat semigroups and differences in the next two sections.

3.3 The characterizations by Poisson and heat semigroups

We now concretize the general characterization in the previous section to the case of Poisson and heat kernels. We begin with the Poisson case. Recall that P denotes the Poisson kernel of $\mathbb{R}^d$ and

$$\tilde{P}_\varepsilon(x) = \tilde{P}_\varepsilon * x = \sum_{m \in \mathbb{Z}^d} e^{-2\pi\varepsilon|m|} \hat{x}(m) U^m.$$

So for any positive integer k , the k th derivation relative to ε is given by

$$\frac{\partial^k}{\partial \varepsilon^k} \tilde{P}_\varepsilon(x) = \sum_{m \in \mathbb{Z}^d} (-2\pi|m|)^k e^{-2\pi\varepsilon|m|} \hat{x}(m) U^m.$$

The inverse of the k th derivation is the k th integration $\mathcal{I}^k$ defined for x with $\hat{x}(0) = 0$ by

$$\begin{aligned} \mathcal{I}_\varepsilon^k \tilde{P}_\varepsilon(x) &= \int_\varepsilon^\infty \int_{\varepsilon_k}^\infty \cdots \int_{\varepsilon_2}^\infty \tilde{P}_{\varepsilon_1}(x) d\varepsilon_1 \cdots d\varepsilon_{k-1} d\varepsilon_k \\ &= \sum_{m \in \mathbb{Z}^d \setminus \{0\}} (2\pi|m|)^{-k} e^{-2\pi\varepsilon|m|} \hat{x}(m) U^m. \end{aligned}$$

In order to simplify the presentation, for any $k \in \mathbb{Z}$, we define

$$\mathcal{J}_\varepsilon^k = \frac{\partial^k}{\partial \varepsilon^k} \text{ for } k \geq 0 \quad \text{and} \quad \mathcal{J}_\varepsilon^k = \mathcal{I}_\varepsilon^{-k} \text{ for } k < 0.$$

It is worth to point out that all concrete characterizations in this section in terms of integration operators are new even in the classical case. Also, compare the following theorem with [74, Section 2.6.4], in which k is assumed to be a positive integer in the Poisson characterization, and a nonnegative integer in the heat characterization.

Theorem 3.13. *Let $1 \leq p, q \leq \infty$, $\alpha \in \mathbb{R}$ and $k \in \mathbb{Z}$ such that $k > \alpha$. Then for any distribution x on $\mathbb{T}_\theta^d$, we have*

$$\|x\|_{B_{p,q}^\alpha} \approx \left(|\widehat{x}(0)|^q + \int_0^1 \varepsilon^{q(k-\alpha)} \|\mathcal{J}_\varepsilon^k \widetilde{P}_\varepsilon(x)\|_p^q \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}}.$$

Proof. Recall that $P = P_1$. Thanks to $\widehat{P}(\xi) = e^{-2\pi|\xi|}$, we introduce the function $\psi(\xi) = (-\operatorname{sgn}(k)2\pi|\xi|)^k e^{-2\pi|\xi|}$. Then

$$\psi(\varepsilon\xi) = \varepsilon^k \mathcal{J}_\varepsilon^k e^{-2\pi\varepsilon|\xi|} = \varepsilon^k \mathcal{J}_\varepsilon^k \widehat{P}_\varepsilon(\xi).$$

It follows that for $x \in B_{p,q}^\alpha(\mathbb{T}_\theta^d)$,

$$\widetilde{\psi}_\varepsilon * x = \varepsilon^k \mathcal{J}_\varepsilon^k \widetilde{P}_\varepsilon * x = \varepsilon^k \mathcal{J}_\varepsilon^k \widetilde{P}_\varepsilon(x).$$

Thus by Theorem 3.11, it remains to check that ψ satisfies (3.4) for some $\alpha_0 < \alpha < \alpha_1$. It is clear that the last condition there is verified for any α_0 . For the second one, choosing $k = \alpha_1 > \alpha$, we have $I_{-\alpha_1} h \psi = (-\operatorname{sgn}(k)2\pi)^k h \widehat{P}$. So

$$\|\mathcal{F}^{-1}(I_{k-\alpha_1} h \psi)\|_1 \leq (2\pi)^k \|\mathcal{F}^{-1}(h)\|_1 \|P\|_1 < \infty.$$

The theorem is thus proved. $\square$

There exists an analogous characterization in terms of the heat kernel. Let W_ε be the heat semigroup of $\mathbb{R}^d$:

$$W_\varepsilon(s) = \frac{1}{(4\pi\varepsilon)^{\frac{d}{2}}} e^{-\frac{|s|^2}{4\varepsilon}}.$$

As usual, let $\widetilde{W}_\varepsilon$ be the periodization of W_ε . Given a distribution x on $\mathbb{T}_\theta^d$ let

$$\widetilde{W}_\varepsilon(x) = \widetilde{W}_\varepsilon * x = \sum_{m \in \mathbb{Z}^d} \widehat{W}(\sqrt{\varepsilon} m) \widehat{x}(m) U^m,$$

where $W = W_1$. Recall that

$$\widehat{W}(\xi) = e^{-4\pi^2|\xi|^2}.$$

Theorem 3.14. *Let $1 \leq p, q \leq \infty$, $\alpha \in \mathbb{R}$ and $k \in \mathbb{Z}$ such that $k > \frac{\alpha}{2}$. Then for any distribution x on $\mathbb{T}_\theta^d$,*

$$\|x\|_{B_{p,q}^\alpha} \approx \left(|\widehat{x}(0)|^q + \int_0^1 \varepsilon^{q(k-\frac{\alpha}{2})} \|\mathcal{J}_\varepsilon^k \widetilde{W}_\varepsilon(x)\|_p^q \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}}.$$

Proof. This proof is similar to and simpler than the previous one. This time, the function ψ is defined by $\psi(\xi) = (-\operatorname{sgn}(k)4\pi^2|\xi|^2)^k e^{-4\pi^2|\xi|^2}$. Clearly, it satisfies (3.4) for $2k = \alpha_1 > \alpha$ and any $\alpha_0 < \alpha$. Thus Theorem 3.11 holds for this choice of ψ . Note that a simple change of variables shows that the integral in (3.14) is equal to

$$2^{-\frac{1}{q}} \left(\int_0^1 \varepsilon^{-\frac{\alpha q}{2}} \|\widetilde{\psi}_{\sqrt{\varepsilon}} * x\|_p^q \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}}.$$

Then using the identity

$$\psi(\sqrt{\varepsilon}\xi) = \varepsilon^k \mathcal{J}_\varepsilon^k \widehat{W}_\varepsilon(\xi),$$

we obtain the desired assertion. $\square$

Now we wish to formulate Theorems 3.13 and 3.14 directly in terms of the circular Poisson and heat semigroups of $\mathbb{T}^d$. Recall that $\mathbb{P}_r$ denote the circular Poisson kernel of $\mathbb{T}^d$ introduced by (1.6) and the Poisson integral of a distribution x on $\mathbb{T}_\theta^d$ is defined by

$$\mathbb{P}_r(x) = \sum_{m \in \mathbb{Z}^d} \hat{x}(m) r^{|m|} U^m, \quad 0 \leq r < 1.$$

Accordingly, we introduce the circular heat kernel $\mathbb{W}$ of $\mathbb{T}^d$:

$$\mathbb{W}_r(z) = \sum_{m \in \mathbb{Z}^d} r^{|m|^2} z^m, \quad z \in \mathbb{T}^d, \quad 0 \leq r < 1. \quad (3.17)$$

Then for $x \in \mathcal{S}'(\mathbb{T}_\theta^d)$ we put

$$\mathbb{W}_r(x) = \sum_{m \in \mathbb{Z}^d} \hat{x}(m) r^{|m|^2} U^m, \quad 0 \leq r < 1.$$

As before, $\mathcal{J}_r^k$ denotes the k th derivation $\frac{\partial^k}{\partial r^k}$ if $k \geq 0$ and the $(-k)$ th integration $\mathcal{I}_r^{-k}$ if $k < 0$:

$$\mathcal{J}_r^k \mathbb{P}_r(x) = \sum_{m \in \mathbb{Z}^d} C_{m,k} \hat{x}(m) r^{|m|-k} U^m,$$

where

$$C_{m,k} = |m| \cdots (|m| - k + 1) \quad \text{if } k \geq 0 \quad \text{and} \quad C_{m,k} = \frac{1}{(|m| + 1) \cdots (|m| - k)} \quad \text{if } k < 0.$$

$\mathcal{J}_r^k \mathbb{W}_r(x)$ is defined similarly. Since $|m|$ is not necessarily an integer, the coefficient $C_{m,k}$ may not vanish for $|m| < k$ and $k \geq 2$. In this case, the corresponding term in $\mathcal{J}_r^k \mathbb{P}_r(x)$ above cause a certain problem of integrability since $r^{(|m|-k)q}$ is integrable on $(0, 1)$ only when $(|m| - k)q > -1$. This explains why we will remove all these terms from $\mathcal{J}_r^k \mathbb{P}_r(x)$ in the following theorem. However, this difficulty does not occur for the heat semigroup.

The following is new even in the classical case, that is, in the case of $\theta = 0$.

Theorem 3.15. *Let $1 \leq p, q \leq \infty$, $\alpha \in \mathbb{R}$ and $k \in \mathbb{Z}$. Let x be a distribution on $\mathbb{T}_\theta^d$.*

(i) *If $k > \alpha$, then*

$$\|x\|_{B_{p,q}^\alpha} \approx \left(\max_{|m| < k} |\hat{x}(m)|^q + \int_0^1 (1-r)^{q(k-\alpha)} \|\mathcal{J}_r^k \mathbb{P}_r(x_k)\|_p^q \frac{dr}{1-r} \right)^{\frac{1}{q}},$$

$$\text{where } x_k = x - \sum_{|m| < k} \hat{x}(m) U^m.$$

(ii) *If $k > \frac{\alpha}{2}$, then*

$$\|x\|_{B_{p,q}^\alpha} \approx \left(\max_{|m|^2 < k} |\hat{x}(m)|^q + \int_0^1 (1-r)^{q(k-\frac{\alpha}{2})} \|\mathcal{J}_r^k \mathbb{W}_r(x)\|_p^q \frac{dr}{1-r} \right)^{\frac{1}{q}}.$$

Proof. We consider only the case of the Poisson kernel. Fix $x \in B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ with $\hat{x}(0) = 0$. We first claim that for any $0 < \varepsilon_0 < 1$,

$$\left(\int_0^1 \varepsilon^{q(k-\alpha)} \|\mathcal{J}_\varepsilon^k \tilde{\mathbb{P}}_\varepsilon(x)\|_p^q \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}} \approx \left(\int_0^{\varepsilon_0} \varepsilon^{q(k-\alpha)} \|\mathcal{J}_\varepsilon^k \tilde{\mathbb{P}}_\varepsilon(x)\|_p^q \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}}.$$

Indeed, since

$$\mathcal{J}_\varepsilon^k \tilde{\mathbf{P}}_\varepsilon(x) = \sum_{m \in \mathbb{Z}^d \setminus \{0\}} (-\operatorname{sgn}(k) 2\pi|m|)^k e^{-2\pi\varepsilon|m|} \hat{x}(m) U^m,$$

we have

$$\begin{aligned} \left(\int_0^{\varepsilon_0} \varepsilon^{q(k-\alpha)} \|\mathcal{J}_\varepsilon^k \tilde{\mathbf{P}}_\varepsilon(x)\|_p^q \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}} &\geq \sup_{m \in \mathbb{Z}^d \setminus \{0\}} (2\pi|m|)^k |\hat{x}(m)| \left(\int_0^{\varepsilon_0} \varepsilon^{q(k-\alpha)} e^{-2\pi\varepsilon|m|q} \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}} \\ &\gtrsim \sup_{m \in \mathbb{Z}^d \setminus \{0\}} |m|^\alpha |\hat{x}(m)|. \end{aligned}$$

On the other hand, by triangular inequality,

$$\begin{aligned} \left(\int_{\varepsilon_0}^1 \varepsilon^{q(k-\alpha)} \|\mathcal{J}_\varepsilon^k \tilde{\mathbf{P}}_\varepsilon(x)\|_p^q \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}} &\leq \sum_{m \in \mathbb{Z}^d \setminus \{0\}} (2\pi|m|)^k |\hat{x}(m)| \left(\int_{\varepsilon_0}^1 \varepsilon^{q(k-\alpha)} e^{-2\pi\varepsilon|m|q} \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}} \\ &\lesssim \sup_{m \in \mathbb{Z}^d \setminus \{0\}} |m|^\alpha |\hat{x}(m)| \sum_{m \in \mathbb{Z}^d} |m|^{k-\alpha} e^{-2\pi\varepsilon_0|m|} \\ &\lesssim \sup_{m \in \mathbb{Z}^d \setminus \{0\}} |m|^\alpha |\hat{x}(m)|. \end{aligned}$$

We then deduce the claim.

Similarly, we can show that for any $0 < r_0 < 1$,

$$\left(\int_0^1 (1-r)^{q(k-\alpha)} \|\mathcal{J}_r^k \mathbb{P}_r(x_k)\|_p^q \frac{dr}{1-r} \right)^{\frac{1}{q}} \approx \left(\int_{r_0}^1 (1-r)^{q(k-\alpha)} \|\mathcal{J}_r^k \mathbb{P}_r(x_k)\|_p^q \frac{dr}{1-r} \right)^{\frac{1}{q}}.$$

Now we use the relation $r = e^{-2\pi\varepsilon}$. If $\varepsilon_0 > 0$ is sufficiently small, then

$$1 - r \approx \varepsilon \quad \text{for } \varepsilon \in (0, \varepsilon_0).$$

So

$$\left(\int_0^{\varepsilon_0} \varepsilon^{q(k-\alpha)} \|\mathcal{J}_\varepsilon^k \tilde{\mathbf{P}}_\varepsilon(x)\|_p^q \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}} \approx \left(\sup_{0 < |m| < k} |\hat{x}(m)|^q + \int_{r_0}^1 (1-r)^{q(k-\alpha)} \|\mathcal{J}_r^k \mathbb{P}_r(x_k)\|_p^q \frac{dr}{1-r} \right)^{\frac{1}{q}}.$$

Combining this with Theorem 3.13, we get the desired assertion. $\square$

3.4 The characterization by differences

In this section we show the quantum analogue of the classical characterization of Besov spaces by differences. Recall that $\omega_p^k(x, \varepsilon)$ is the L_p -modulus of smoothness of x introduced in Definition 2.18. The result of this section is the following:

Theorem 3.16. *Let $1 \leq p, q \leq \infty$ and $0 < \alpha < n$ with $n \in \mathbb{N}$. Then for any distribution x on $\mathbb{T}_\theta^d$,*

$$\|x\|_{B_{p,q}^\alpha} \approx \left(|\hat{x}(0)|^q + \int_0^1 \varepsilon^{-q\alpha} \omega_p^n(x, \varepsilon)^q \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}}. \quad (3.18)$$

Proof. We will derive the result from Theorem 3.11, or more precisely, from its proof. Since $\alpha > 0$, we take $\alpha_0 = 0$ and $\alpha_1 = n$ in that theorem. Recall that $d_u(\xi) = e^{2\pi i u \cdot \xi} - 1$. Then the last condition of (3.4) with $\psi = d_u^n$ is satisfied uniformly in u since

$$\|\mathcal{F}^{-1}(d_u(2^j \cdot)^n \varphi)\|_1 = \|\Delta_{2^j u}^n \mathcal{F}^{-1}(\varphi)\|_1 \leq 2^n \|\mathcal{F}^{-1}(\varphi)\|_1.$$

We will use a variant of the second one (which is not necessarily verified). To this end, let us come back to (3.8) and rewrite it as follows:

$$\begin{aligned}\psi^{(j)}(\xi)\varphi^{(j+k)}(\xi) &= 2^{nk} \frac{\psi^{(j)}(\xi)}{(2^{-j}u \cdot \xi)^n} h^{(j+k)}(\xi) (2^{-j-k}u \cdot \xi)^n \varphi^{(j+k)}(\xi) \\ &= 2^{nk} \eta^{(j)}(\xi) \rho^{(j+k)}(\xi),\end{aligned}$$

where η and ρ are now defined by

$$\eta(\xi) = \frac{\psi(\xi)}{(u \cdot \xi)^n} h^{(K)}(\xi) \quad \text{and} \quad \rho(\xi) = (u \cdot \xi)^n \varphi(\xi).$$

The second condition of (3.4) becomes the requirement that

$$\sup_{u \in \mathbb{R}^d, |u| \leq 1} \|\mathcal{F}^{-1}(\eta)\|_1 < \infty.$$

The crucial point here is that $\psi(\xi) = d_u^n(\xi) = (u \cdot \xi)^n \zeta(u \cdot \xi)$, where ζ is an analytic function on $\mathbb{R}$. This shows that the above supremum is finite.

However, the first condition of (3.4), the Tauberian condition is not verified for a single d_u^n . We will return back to this point later. For the moment, we just observe that the Tauberian condition has not been used in steps 1 and 2 of the proof of Theorem 3.9. Reexamining those two steps with $\psi = d_u^n$, we see that all estimates there can be made independent of u . For instance, (3.15) now becomes (with $\alpha_1 = n$)

$$2^{j\alpha} \|\Delta_{\varepsilon u}^n \tilde{\varphi}_{j+k} * x\|_p \lesssim 2^{(\alpha_1 - \alpha)k} (2^{(j+k)\alpha} \|\tilde{\rho}_{j+k} * x\|_p),$$

where the new function ρ is defined as above. Thus taking the supremum over all u with $|u| \leq 1$, we get

$$2^{j\alpha} \omega_p^n(x, \varepsilon) \lesssim 2^{k(\alpha_1 - \alpha)} (2^{(j+k)\alpha} \|\tilde{\rho}_{(j+k)} * x\|_p).$$

Therefore, by Lemma 3.4, we obtain

$$\left(\int_0^1 \varepsilon^{-q\alpha} \omega_p^n(x, \varepsilon)^q \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}} \lesssim \|x\|_{B_{p,q}^\alpha}.$$

The reverse inequality requires necessarily a Tauberian-type condition. Although a single d_u^n does not satisfy it, a finite family of d_u^n 's does satisfy this condition, which we precise below. Choose a $\frac{1}{2}$ -net $\{v_\ell\}_{1 \leq \ell \leq L}$ of the unit sphere of $\mathbb{R}^d$. Let $u_\ell = 4^{-1}v_\ell$ and

$$\Omega_\ell = \{\xi : 2^{-1} \leq |\xi| \leq 2, \left| \frac{\xi}{|\xi|} - v_\ell \right| \leq 2^{-1}\}.$$

Then the union of the Ω_ℓ 's is equal to $\{\xi : 2^{-1} \leq |\xi| \leq 2\}$ and $|d_{u_\ell}^n| > 0$ on Ω_ℓ . So the family $\{d_{u_\ell}^n\}_{1 \leq \ell \leq L}$ satisfies the following Tauberian-type condition:

$$\sum_{\ell=1}^L |d_{u_\ell}^n| > 0 \quad \text{on} \quad \{\xi : 2^{-1} \leq |\xi| \leq 2\}.$$

Now we reexamine step 3 of the proof of Theorem 3.9. To adapt it to the present setting, by an appropriate partition of unity, we decompose φ into a sum of infinitely differentiable functions, $\varphi = \varphi_1 + \dots + \varphi_L$ such that $\text{supp } \varphi_\ell \subset \Omega_\ell$. Accordingly, for every $j \geq 0$, let

$$\varphi^{(j)} = \sum_{\ell=1}^L \varphi_\ell^{(j)}.$$

Then we write the corresponding (3.16) with $(\varphi_\ell, d_{u_\ell}^n)$ in place of (φ, ψ) for every $\ell \in \{1, \dots, L\}$. Arguing as in step 3 of the proof of Theorem 3.9, we get

$$\|x\|_{B_{p,q}^\alpha} \lesssim \left(|\widehat{x}(0)|^q + \sum_{\ell=1}^L \int_0^1 \varepsilon^{-q\alpha} \|(\widetilde{d_{u_\ell}^n})_\varepsilon * x\|_p^q \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}}.$$

Since $(\widetilde{d_{u_\ell}^n})_\varepsilon * x = \Delta_{\varepsilon u_\ell}^n x$, we deduce

$$\begin{aligned} \|x\|_{B_{p,q}^\alpha} &\lesssim \left(|\widehat{x}(0)|^q + \int_0^1 \varepsilon^{-q\alpha} \sup_{1 \leq \ell \leq L} \|\Delta_{\varepsilon u_\ell}^n x\|_p^q \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}} \\ &\lesssim \left(|\widehat{x}(0)|^q + \int_0^1 \varepsilon^{-q\alpha} \omega_p^n(x, \varepsilon)^q \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}}. \end{aligned}$$

Thus the theorem is proved. $\square$

As a byproduct, the preceding theorem implies that the right-hand side of (3.18) does not depend on n with $n > \alpha$, up to equivalence. This fact admits a direct simple proof and is an immediate consequence of the following analogue of Marchaud's classical inequality which is of interest in its own right.

Proposition 3.17. *For any positive integers n and N with $n < N$ and for any $\varepsilon > 0$, we have*

$$2^{n-N} \omega_p^N(x, \varepsilon) \leq \omega_p^n(x, \varepsilon) \lesssim \varepsilon^n \int_\varepsilon^\infty \frac{\omega_p^N(x, \delta)}{\delta^n} \frac{d\delta}{\delta}.$$

Proof. The argument below is standard. Using the identity $\Delta_u^N = \Delta_u^{N-n} \Delta_u^n$, we get

$$\|\Delta_u^N(x)\|_p \leq 2^{N-n} \|\Delta_u^n(x)\|_p,$$

whence the lower estimate. The upper one is less obvious. By elementary calculations, for any $u \in \mathbb{R}^d$, we have

$$d_{2u}^n = 2^n d_u^n + \left[\sum_{j=0}^n \binom{n}{j} e_{ju} - 2^n \right] d_u^n = 2^n d_u^n + \sum_{j=0}^n \binom{n}{j} \sum_{i=0}^{j-1} e_{iu} d_u^{n+1}.$$

In terms of Fourier multipliers, this means

$$\Delta_{2u}^n = 2^n \Delta_u^n + \sum_{j=0}^n \binom{n}{j} \sum_{i=0}^{j-1} T_{iu} \Delta_u^{n+1}.$$

It then follows that

$$\omega_p^n(x, \varepsilon) \leq \frac{n}{2} \omega_p^{n+1}(x, \varepsilon) + 2^{-n} \omega_p^n(x, 2\varepsilon).$$

Iterating this inequality yields

$$\omega_p^n(x, \varepsilon) \leq \frac{n}{2} \sum_{j=1}^{J-1} \omega_p^{n+1}(x, 2^j \varepsilon) + 2^{-Jn} \omega_p^n(x, 2^J \varepsilon),$$

from which we deduce the desired inequality for $N = n + 1$ as $J \rightarrow \infty$. Another iteration argument then yields the general case. $\square$

Remark 3.18. Theorem 3.16 shows that if $\alpha > 0$, $B_{\infty,\infty}^\alpha(\mathbb{T}_\theta^d)$ coincides with the quantum analogue of the classical Zygmund class of order α . In particular, for $0 < \alpha < 1$, $B_{\infty,\infty}^\alpha(\mathbb{T}_\theta^d)$ and $B_{\infty,c_0}^\alpha(\mathbb{T}_\theta^d)$ are the Lipschitz and little Lipschitz classes of order α , already studied by Weaver [79]. Note that like in the classical setting, if k is a positive integer, $B_{\infty,\infty}^k(\mathbb{T}_\theta^d)$ is closely related to but different from the Lipschitz class $W_\infty^k(\mathbb{T}_\theta^d)$ discussed in section 2.4.

3.5 Limits of Besov norms

In this section we consider the behavior of the right-hand side of (3.18) as $\alpha \rightarrow n$. The study of this behavior is the subject of several recent publications in the classical setting; see, for instance, [4, 5, 38, 41, 75]. It originated from [14] in which Bourgain, Brézis and Mironescu proved that for any $1 \leq p < \infty$ and any $f \in C_0^\infty(\mathbb{R}^d)$

$$\lim_{\alpha \rightarrow 1} \left((1 - \alpha) \int_{\mathbb{R}^d \times \mathbb{R}^d} \frac{|f(s) - f(t)|^p}{|s - t|^{\alpha p + d}} ds dt \right)^{\frac{1}{p}} = C_{p,d} \|\nabla f(t)\|_p.$$

It is well known that

$$\left(\int_{\mathbb{R}^d \times \mathbb{R}^d} \frac{|f(s) - f(t)|^p}{|s - t|^{\alpha p + d}} ds dt \right)^{\frac{1}{p}} \approx \left(\int_0^\infty \sup_{u \in \mathbb{R}^d, |u| \leq \varepsilon} \|\Delta_u f\|_p^p \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{p}}.$$

The right-hand side is the norm of f in the Besov space $B_{p,p}^1(\mathbb{R}^d)$. Higher order extensions have been obtained in [38, 75].

The main result of the present section is the following quantum extension of these results. Let

$$\|x\|_{B_{p,q}^{\alpha,\omega}} = \left(\int_0^1 \varepsilon^{-\alpha q} \omega_p^k(x, \varepsilon)^q \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{q}}. \quad (3.19)$$

Theorem 3.19. *Let $1 \leq p \leq \infty$, $1 \leq q < \infty$ and $0 < \alpha < k$ with k an integer. Then for $x \in W_p^k(\mathbb{T}_\theta^d)$,*

$$\lim_{\alpha \rightarrow k} (k - \alpha)^{\frac{1}{q}} \|x\|_{B_{p,q}^{\alpha,\omega}} \approx q^{-\frac{1}{q}} |x|_{W_p^k}$$

with relevant constants depending only on d and k .

Proof. The proof is easy by using the results of section 2.4. Let $x \in W_p^k(\mathbb{T}_\theta^d)$ with $\hat{x}(0) = 0$. Let A denote the limit in Lemma 2.21. Then

$$\int_0^1 \varepsilon^{-\alpha q} \omega_p^k(x, \varepsilon)^q \frac{d\varepsilon}{\varepsilon} \leq A^q \int_0^1 \varepsilon^{(k-\alpha)q} \frac{d\varepsilon}{\varepsilon},$$

whence

$$\limsup_{\alpha \rightarrow k} (k - \alpha) \int_0^1 \varepsilon^{-\alpha q} \omega_p^k(x, \varepsilon)^q \frac{d\varepsilon}{\varepsilon} \leq \frac{A^q}{q}.$$

Conversely, for any $\eta > 0$, choose $\delta \in (0, 1)$ such that

$$\frac{\omega_p^k(x, \varepsilon)}{\varepsilon^k} \geq A - \eta, \quad \forall \varepsilon \leq \delta.$$

Then

$$(k - \alpha) \int_0^1 \varepsilon^{-\alpha q} \omega_p^k(x, \varepsilon)^q \frac{d\varepsilon}{\varepsilon} \geq \frac{(A - \eta)^q}{q} \delta^{(k-\alpha)q},$$

which implies

$$\liminf_{\alpha \rightarrow k} (k - \alpha) \int_0^1 \varepsilon^{-\alpha q} \omega_p^k(x, \varepsilon)^q \frac{d\varepsilon}{\varepsilon} \geq \frac{(A - \eta)^q}{q}.$$

Therefore,

$$\lim_{\alpha \rightarrow k} (k - \alpha) \int_0^1 \varepsilon^{-\alpha q} \omega_p^k(x, \varepsilon)^q \frac{d\varepsilon}{\varepsilon} = \frac{1}{q} \lim_{\varepsilon \rightarrow 0} \frac{\omega_p^k(x, \varepsilon)}{\varepsilon^k}.$$

So Theorem 2.20 implies the desired assertion. $\square$

Remark 3.20. We will determine later the behavior of $\|x\|_{B_{p,q}^{\alpha,\omega}}$ when $\alpha \rightarrow 0$; see Corollary 5.20 below.

3.6 The link with the classical Besov spaces

Like for the Sobolev spaces on $\mathbb{T}_\theta^d$, there exists a strong link between $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ and the classical vector-valued Besov spaces on $\mathbb{T}^d$. Let us give a precise definition of the latter spaces. We maintain the assumption and notation on φ in section 3.1. In particular, $f \mapsto \tilde{\varphi}_k * f$ is the Fourier multiplier on $\mathbb{T}^d$ associated to $\varphi(2^{-k}\cdot)$:

$$\tilde{\varphi}_k * f = \sum_{m \in \mathbb{Z}^d} \varphi(2^{-k}m) \hat{f}(m) z^m$$

for any $f \in \mathcal{S}'(\mathbb{T}^d; X)$. Here X is a Banach space.

Definition 3.21. Let $1 \leq p, q \leq \infty$ and $\alpha \in \mathbb{R}$. Define

$$B_{p,q}^\alpha(\mathbb{T}^d; X) = \{f \in \mathcal{S}'(\mathbb{T}^d; X) : \|f\|_{B_{p,q}^\alpha} < \infty\},$$

where

$$\|f\|_{B_{p,q}^\alpha} = \left(\|\hat{f}(0)\|_X^q + \sum_{k \geq 0} 2^{\alpha k q} \|\tilde{\varphi}_k * f\|_{L_p(\mathbb{T}^d; X)}^q \right)^{\frac{1}{q}}.$$

These vector-valued Besov spaces have been largely studied in literature. Note that almost all publications concern only the case of $\mathbb{R}^d$, but the periodic theory is parallel (see, for instance, [24, 71]; see also [3] for the vector-valued case). $B_{p,q}^\alpha(\mathbb{R}^d; X)$ is defined in the same way with the necessary modifications among them the main difference concerns the term $\|\hat{f}(0)\|_X$ above which is replaced by $\|\phi * f\|_{L_p(\mathbb{R}^d; X)}$, where ϕ is the function whose Fourier transform is equal to $1 - \sum_{k \geq 0} \varphi(2^{-k}\cdot)$.

All results proved in the previous sections remain valid in the present vector-valued setting with essentially the same proofs for any Banach space X , except Theorem 3.8 whose vector-valued version holds only if X is isomorphic to a Hilbert space. On the other hand, the duality assertion in Proposition 3.3 should be slightly modified by requiring that X^* have the Radon-Nikodym property.

Let us state the vector-valued analogue of Theorem 3.15. As said before, it is new even in the scalar case. The circular Poisson and heat semigroups are extended to the present case too. For any $f \in \mathcal{S}'(\mathbb{T}^d; X)$,

$$\mathbb{P}_r(f)(z) = \sum_{m \in \mathbb{Z}^d} r^{|m|} \hat{f}(m) z^m \quad \text{and} \quad \mathbb{W}_r(f)(z) = \sum_{m \in \mathbb{Z}^d} r^{|m|^2} \hat{f}(m) z^m, \quad z \in \mathbb{T}^d, \quad 0 \leq r < 1.$$

The operator $\mathcal{J}_r^k$ has the same meaning as before, for instance, in the Poisson case, we have

$$\mathcal{J}_r^k \mathbb{P}_r(f) = \sum_{m \in \mathbb{Z}^d} C_{m,k} \hat{f}(m) r^{|m|-k} z^m,$$

where

$$C_{m,k} = |m| \cdots (|m| - k + 1) \quad \text{if } k \geq 0 \quad \text{and} \quad C_{m,k} = \frac{1}{(|m| + 1) \cdots (|m| - k)} \quad \text{if } k < 0.$$

Theorem 3.22. Let $1 \leq p, q \leq \infty$, $\alpha \in \mathbb{R}$ and $k \in \mathbb{Z}$. Let X be a Banach space.

(i) If $k > \alpha$, then for any $f \in B_{p,q}^\alpha(\mathbb{T}^d; X)$,

$$\|f\|_{B_{p,q}^\alpha} \approx \left(\sup_{|m| < k} \|\hat{f}(m)\|_X^q + \int_0^1 (1-r)^{(k-\alpha)q} \|\mathcal{J}_r^k \mathbb{P}_r(f_k)\|_{L_p(\mathbb{T}^d; X)}^q \frac{dr}{1-r} \right)^{\frac{1}{q}},$$

$$\text{where } f_k = f - \sum_{|m| < k} \hat{f}(m) z^m.$$

(ii) If $k > \frac{\alpha}{2}$, then for any $f \in B_{p,q}^\alpha(\mathbb{T}^d; X)$,

$$\|f\|_{B_{p,q}^\alpha} \approx \left(\sup_{|m|^2 < k} \|\widehat{f}(m)\|_X^q + \int_0^1 (1-r)^{(k-\frac{\alpha}{2})q} \|\mathcal{J}_r^k \mathbb{W}_r(f)\|_{L_p(\mathbb{T}^d; X)}^q \frac{dr}{1-r} \right)^{\frac{1}{q}}.$$

The following transference result from $\mathbb{T}_\theta^d$ to $\mathbb{T}^d$ is clear. It can be used to prove a majority of the preceding results on $\mathbb{T}_\theta^d$, under the hypothesis that the corresponding results in the vector-valued case on $\mathbb{T}^d$ are known.

Proposition 3.23. *Let $1 \leq p, q \leq \infty$ and $\alpha \in \mathbb{R}$. The map $x \mapsto \tilde{x}$ in Corollary 1.2 is an isometric embedding from $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ into $B_{p,q}^\alpha(\mathbb{T}^d; L_p(\mathbb{T}_\theta^d))$ with 1-complemented range.*

Remark 3.24. As a subspace of $\ell_q^\alpha(L_p(\mathbb{T}_\theta^d))$ (see the proof of Proposition 3.3 for the definition of this space), $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ can be equipped with a natural operator space structure in Pisier's sense [54]. Moreover, in the spirit of the preceding vector-valued case, we can also introduce the vector-valued quantum Besov spaces. Given an operator space E , $B_{p,q}^\alpha(\mathbb{T}_\theta^d; E)$ is defined exactly as in the scalar case; it is a subspace of $\ell_q^\alpha(L_p(\mathbb{T}_\theta^d; E))$. Then all results of this chapter are extended to this vector-valued case, except the duality in Proposition 3.3 and Theorem 3.15.

Chapter 4

Triebel-Lizorkin spaces

This chapter is devoted to the study of Triebel-Lizorkin spaces. These spaces are much subtler than Besov spaces even in the classical setting. Like Besov spaces, the classical Triebel-Lizorkin spaces $F_{p,q}^\alpha(\mathbb{R}^d)$ have three parameters, p, q and α . The difference is that the ℓ_q -norm is now taken before the L_p -norm. Namely, $f \in F_{p,q}^\alpha(\mathbb{R}^d)$ iff $\|(\sum_{k \geq 0} 2^{k\alpha q} |\varphi_k * f|^q)^{\frac{1}{q}}\|_p$ is finite. Besides the usual subtlety just mentioned, more difficulties appear in the noncommutative setting. For instance, a first elementary one concerns the choice of the internal ℓ_q -norm. It is a well-known fact in the noncommutative integration that the simple replacement of the usual absolute value by the modulus of operators does not give a norm except for $q = 2$. Alternately, one could use Pisier's definition of ℓ_q -valued noncommutative L_p -spaces in the category of operator spaces. However, we will not study the latter choice and will confine ourselves only to the case $q = 2$, by considering the column and row norms (and their mixture) for the internal ℓ_2 -norms. This choice is dictated by the theory of noncommutative Hardy spaces. In fact, we will show that the so-defined Triebel-Lizorkin spaces on $\mathbb{T}_\theta^d$ are isomorphic to the Hardy spaces developed in [17].

Another difficulty is linked with the frequent use of maximal functions in the commutative case. These functions play a crucial role in the classical theory. However, the pointwise analogue of maximal functions is no longer available in the present setting, which makes our study harder than the classical case. We have already encountered this difficulty in the study of Besov spaces. It is much more substantial now. Instead, our development will rely heavily on the theory of Hardy spaces developed in [81] through a Fourier multiplier theorem that is proved in the first section. It is this multiplier theorem which clears the obstacles on our route. After the definitions and basic properties, we prove some characterizations of the quantum Triebel-Lizorkin spaces. Like in the Besov case, they are better than the classical ones even in the commutative case. We conclude the chapter with a short section on the operator-valued Triebel-Lizorkin spaces on $\mathbb{T}^d$ (or $\mathbb{R}^d$). These spaces are interesting in view of the theory of operator-valued Hardy spaces.

Throughout this chapter, we will use the notation introduced in the previous one. In particular, φ is a function satisfying (3.1), $\varphi^{(k)} = \varphi(2^{-k}\cdot)$ and $\widehat{\varphi}_k = \varphi^{(k)}$.

4.1 A multiplier theorem

The following multiplier result will play a crucial role in this chapter.

Theorem 4.1. *Let $\sigma \in \mathbb{R}$ with $\sigma > \frac{d}{2}$. Assume that $(\phi_j)_{j \geq 0}$ and $(\rho_j)_{j \geq 0}$ are two sequences of continuous functions on $\mathbb{R}^d \setminus \{0\}$ such that*

$$\begin{cases} \text{supp}(\phi_j \rho_j) \subset \{\xi : 2^{j-1} \leq |\xi| \leq 2^{j+1}\}, \forall j \geq 0, \\ \sup_{j \geq 0} \|\phi_j(2^j \cdot) \varphi\|_{H_2^\sigma(\mathbb{R}^d)} < \infty. \end{cases} \quad (4.1)$$

(i) *Let $1 < p < \infty$. Then for any distribution x on $\mathbb{T}_\theta^d$,*

$$\left\| \left(\sum_{j \geq 0} 2^{2j\alpha} |\tilde{\phi}_j * \tilde{\rho}_j * x|^2 \right)^{\frac{1}{2}} \right\|_{L_p(\mathbb{T}_\theta^d)} \lesssim \sup_{j \geq 0} \|\phi_j(2^j \cdot) \varphi\|_{H_2^\sigma} \left\| \left(\sum_{j \geq 0} 2^{2j\alpha} |\tilde{\rho}_j * x|^2 \right)^{\frac{1}{2}} \right\|_{L_p(\mathbb{T}_\theta^d)},$$

where the constant depends only on p, σ, d and φ .

(ii) *Assume, in addition, that $\rho_j = \widehat{\rho(2^{-j} \cdot)}$ for some Schwartz function with $\text{supp}(\rho) = \{\xi : 2^{-1} \leq |\xi| \leq 2\}$. Then the above inequality also holds for $p = 1$ with relevant constant depending additionally on ρ .*

The remainder of this section is devoted to the proof of the above theorem. As one can guess, the proof is based on the Calderón-Zygmund theory. We require several lemmas. The first one is an elementary inequality.

Lemma 4.2. *Assume that $f : \mathbb{R}^d \rightarrow \ell_2$ and $g : \mathbb{R}^d \rightarrow \mathbb{C}$ satisfy*

$$f \in H_2^\sigma(\mathbb{R}^d; \ell_2) \quad \text{and} \quad \int_{\mathbb{R}^d} (1 + |s|^2)^\sigma |\mathcal{F}^{-1}(g)(s)| ds < \infty.$$

Then $fg \in H_2^\sigma(\mathbb{R}^d; \ell_2)$ and

$$\|fg\|_{H_2^\sigma(\mathbb{R}^d; \ell_2)} \leq \|f\|_{H_2^\sigma(\mathbb{R}^d; \ell_2)} \int_{\mathbb{R}^d} (1 + |s|^2)^\sigma |\mathcal{F}^{-1}(g)(s)| ds.$$

Proof. The norm $\|\cdot\|$ below is that of ℓ_2 . By the Cauchy-Schwarz inequality, for $s \in \mathbb{R}^d$, we have

$$\|\mathcal{F}^{-1}(fg)(s)\|^2 = \|\mathcal{F}^{-1}(f) * \mathcal{F}^{-1}(g)(s)\|^2 \leq \|\mathcal{F}^{-1}(g)\|_1 \int_{\mathbb{R}^d} \|\mathcal{F}^{-1}(f)(s-t)\|^2 |\mathcal{F}^{-1}(g)(t)| dt.$$

It then follows that

$$\begin{aligned} \|fg\|_{H_2^\sigma(\mathbb{R}^d; \ell_2)}^2 &= \int_{\mathbb{R}^d} (1 + |s|^2)^\sigma \|\mathcal{F}^{-1}(fg)(s)\|^2 ds \\ &\leq \|\mathcal{F}^{-1}(g)\|_1 \int_{\mathbb{R}^d} (1 + |s|^2)^\sigma \int_{\mathbb{R}^d} \|\mathcal{F}^{-1}(f)(s-t)\|^2 |\mathcal{F}^{-1}(g)(t)| dt ds \\ &\leq \|\mathcal{F}^{-1}(g)\|_1 \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} (1 + |s-t|^2)^\sigma \|\mathcal{F}^{-1}(f)(s-t)\|^2 ds (1 + |t|^2)^\sigma |\mathcal{F}^{-1}(g)(t)| dt \\ &\leq \|f\|_{H_2^\sigma(\mathbb{R}^d; \ell_2)}^2 \left(\int_{\mathbb{R}^d} (1 + |t|^2)^\sigma |\mathcal{F}^{-1}(g)(t)| dt \right)^2. \end{aligned}$$

Thus the assertion is proved. $\square$

The following lemma is a well-known result in harmonic analysis, which asserts that every Hörmander multiplier is a Calderón-Zygmund operator. Note that the usual Hörmander condition is expressed in terms of partial derivatives up to order $[\frac{d}{2}] + 1$, while the condition below, in terms of the potential Sobolev space $H_2^\sigma(\mathbb{R}^d)$, is not commonly used (it is explicitly stated on page 263 of [68]). Combined with the previous lemma, the standard argument as described in [25, p. 211-214], [68, p. 245-247] or [72, p. 161-165] can be easily adapted to the present setting. We include a proof for the convenience of the reader.

Lemma 4.3. *Let $\phi = (\phi_j)_{j \geq 0}$ be a sequence of continuous functions on $\mathbb{R}^d \setminus \{0\}$, viewed as a function from $\mathbb{R}^d$ to ℓ_2 . Assume that*

$$\sup_{k \in \mathbb{Z}} \|\phi(2^k \cdot) \varphi\|_{H_2^\sigma(\mathbb{R}^d; \ell_2)} < \infty. \quad (4.2)$$

Let $\mathbf{k} = (\mathbf{k}_j)_{j \geq 0}$ with $\mathbf{k}_j = \mathcal{F}^{-1}(\phi_j)$. Then $\mathbf{k}$ is a Calderón-Zygmund kernel with values in ℓ_2 , more precisely,

- $\|\widehat{\mathbf{k}}\|_{L_\infty(\mathbb{R}^d; \ell_2)} \lesssim \sup_{k \in \mathbb{Z}} \|\phi(2^k \cdot) \varphi\|_{H_2^\sigma(\mathbb{R}^d; \ell_2)};$
- $\sup_{t \in \mathbb{R}^d} \int_{|s| > 2|t|} \|\mathbf{k}(s-t) - \mathbf{k}(s)\|_{\ell_2} ds \lesssim \sup_{k \in \mathbb{Z}} \|\phi(2^k \cdot) \varphi\|_{H_2^\sigma(\mathbb{R}^d; \ell_2)}.$

The relevant constants depend only on φ , σ and d .

Proof. We will follow standard arguments; see, for instance, [25, p. 211-214], [68, p. 245-247] or [72, p. 161-165] where, instead of (4.2), the usual Hörmander condition in terms of partial derivatives up to order $[\frac{d}{2}] + 1$ is assumed. Let $\|\phi\|_{2,\sigma}$ denote the supremum in (4.2), which can be reformulated as

$$\|\phi\|_{2,\sigma}^2 = \sup_{k \in \mathbb{Z}} \int_{\mathbb{R}^d} (1 + |s|^2)^\sigma \|\mathcal{F}^{-1}(\phi(2^k \cdot) \varphi)(s)\|^2 ds < \infty.$$

Then for $\xi \in \mathbb{R}^d$ and $k \in \mathbb{Z}$, by the Cauchy-Schwarz inequality and the condition $\sigma > \frac{d}{2}$, we have

$$\begin{aligned} \|\phi(2^k \xi) \varphi(\xi)\| &= \left\| \int_{\mathbb{R}^d} \mathcal{F}^{-1}(\phi(2^k \cdot) \varphi)(s) e^{-2\pi i s \cdot \xi} ds \right\| \\ &\leq \|\phi(2^k \cdot) \varphi\|_{H_2^\sigma(\mathbb{R}^d; \ell_2)} \left(\int_{\mathbb{R}^d} (1 + |s|^2)^{-\sigma} ds \right)^{\frac{1}{2}} \lesssim \|\phi\|_{2,\sigma}. \end{aligned}$$

This clearly implies

$$\|\widehat{\mathbf{k}}\|_{L_\infty(\mathbb{R}^d; \ell_2)} \lesssim \|\phi\|_{2,\sigma}.$$

To show the second property of the kernel $\mathbf{k}$, we decompose ϕ as

$$\phi = \sum_{k \in \mathbb{Z}} \phi \varphi^{(k)}.$$

Here and in the rest of the proof, we do not worry about the convergence issue of the above series; more rigorously, one should first consider its partial sums, then use a limit procedure, which is quite formal. For each k , by (3.2), we write

$$\phi \varphi^{(k)} = [\phi(\varphi^{(k-1)} + \varphi^{(k)} + \varphi^{(k+1)})] \varphi^{(k)} \stackrel{\text{def}}{=} \phi^k \varphi^{(k)}.$$

Then for $s \in \mathbb{R}^d$,

$$\mathcal{F}^{-1}(\phi \varphi^{(k)})(s) = \mathcal{F}^{-1}(\phi^k) * \mathcal{F}^{-1}(\varphi^{(k)})(s) = 2^{kd} \mathcal{F}^{-1}(\phi^k(2^k \cdot)) * \mathcal{F}^{-1}(\varphi)(2^k s).$$

So by Lemma 4.2

$$\left(\int_{\mathbb{R}^d} (1 + |2^k s|^2)^\sigma \|\mathcal{F}^{-1}(\phi \varphi^{(k)})(s)\|^2 ds \right)^{\frac{1}{2}} \lesssim 2^{\frac{kd}{2}} \|\phi\|_{2,\sigma},$$

where the constant depends only on φ, d and σ . Thus for any $t \in \mathbb{R}^d \setminus \{0\}$,

$$\int_{|s|>|t|} \|\mathcal{F}^{-1}(\phi\varphi^{(k)})(s)\| ds \lesssim 2^{\frac{kd}{2}} \|\phi\|_{2,\sigma} \left(\int_{|s|>|t|} (1 + |2^k s|^2)^{-\sigma} ds \right)^{\frac{1}{2}} \lesssim (2^k |t|)^{\frac{d}{2}-\sigma} \|\phi\|_{2,\sigma}.$$

Consequently,

$$\int_{|s|>2|t|} \|\mathcal{F}^{-1}(\phi\varphi^{(k)})(s) - \mathcal{F}^{-1}(\phi\varphi^{(k)})(s-t)\| ds \lesssim (2^k |t|)^{\frac{d}{2}-\sigma} \|\phi\|_{2,\sigma}.$$

This estimate is good only when $2^k |t| \geq 1$. Otherwise, we need another one by using the cancellation condition. We have (recalling that $e_t(\xi) = e^{2\pi i t \cdot \xi}$)

$$\begin{aligned} \mathcal{F}^{-1}(\phi\varphi^{(k)})(s) - \mathcal{F}^{-1}(\phi\varphi^{(k)})(s-t) &= \mathcal{F}^{-1}(\phi^k \varphi^{(k)}(1 - e_t)) \\ &= 2^{kd} \mathcal{F}^{-1}(\phi^k(2^k \cdot)) * \mathcal{F}^{-1}(\varphi(1 - e_{2^k t}))(2^k s) \\ &= 2^{kd} \mathcal{F}^{-1}(\phi^k(2^k \cdot)) * [\mathcal{F}^{-1}(\varphi) - \mathcal{F}^{-1}(\varphi)(\cdot - 2^k t)](2^k s). \end{aligned}$$

Thus

$$\left(\int_{\mathbb{R}^d} (1 + |2^k s|^2)^\sigma \|\mathcal{F}^{-1}(\phi\varphi^{(k)})(s) - \mathcal{F}^{-1}(\phi\varphi^{(k)})(s-t)\|^2 ds \right)^{\frac{1}{2}} \lesssim 2^{\frac{kd}{2}} 2^k |t| \|\phi\|_{2,\sigma}.$$

Then as before, for $2^k |t| < 1$, we get

$$\int_{|s|>2|t|} \|\mathcal{F}^{-1}(\phi\varphi^{(k)})(s) - \mathcal{F}^{-1}(\phi\varphi^{(k)})(s-t)\| ds \lesssim 2^k |t| \|\phi\|_{2,\sigma}.$$

Combining the previous estimates, we deduce

$$\begin{aligned} \sup_{t \in \mathbb{R}^d} \int_{|s|>2|t|} \|k(s-t) - k(s)\| ds &\leq \sum_{k \in \mathbb{Z}} \int_{|s|>2|t|} \|\mathcal{F}^{-1}(\phi\varphi^{(k)})(s) - \mathcal{F}^{-1}(\phi\varphi^{(k)})(s-t)\| ds \\ &\lesssim \|\phi\|_{2,\sigma} \sum_{k \in \mathbb{Z}} \min((2^k |t|)^{\frac{d}{2}-\sigma}, 2^k |t|) \lesssim \|\phi\|_{2,\sigma}. \end{aligned}$$

Hence, the lemma is proved. $\square$

The above kernel k defines a Calderón-Zygmund operator on $\mathbb{R}^d$. But we will consider only the periodic case, so we need to periodize k :

$$\tilde{k}(s) = \sum_{m \in \mathbb{Z}^d} k(s+m).$$

By a slight abuse of notation, we use $\tilde{k}_j$ to denote the Calderón-Zygmund operator on $\mathbb{T}^d$ associated to $\tilde{k}_j$ too:

$$\tilde{k}_j(f)(s) = \int_{\mathbb{T}^d} \tilde{k}_j(s-t) f(t) dt,$$

where we have identified $\mathbb{T}$ with $\mathbb{I} = [0, 1)$. $\tilde{k}_j$ is the Fourier multiplier on $\mathbb{T}^d$ with symbol ϕ_j : $f \mapsto \tilde{\phi}_j * f$.

We have $\widehat{\tilde{k}} = \widehat{k}|_{\mathbb{Z}^d}$. If ϕ satisfies (4.2), then Lemma 4.3 implies

$$\begin{cases} \|\widehat{\tilde{k}}\|_{\ell_\infty(\mathbb{Z}^d; \ell_2)} < \infty, \\ \sup_{t \in \mathbb{I}^d} \int_{\{s \in \mathbb{I}^d: |s|>2|t|\}} \|\tilde{k}(s-t) - \tilde{k}(s)\|_{\ell_2} ds < \infty. \end{cases} \quad (4.3)$$

Now let $\mathcal{M}$ be a von Neumann algebra equipped with a normal faithful tracial state τ , and let $\mathcal{N} = L_\infty(\mathbb{T}^d) \bar{\otimes} \mathcal{M}$, equipped with the tensor trace. The following lemma should be known to experts; it is closely related to similar results of [29, 46, 51], notably to [35, Lemma 2.3]. Note that the sole difference between the following condition (4.4) and (4.2) is that the supremum in (4.2) runs over all integers while the one below is restricted to nonnegative integers.

Lemma 4.4. *Let $\phi = (\phi_j)_{j \geq 0}$ be a sequence of continuous functions on $\mathbb{R}^d \setminus \{0\}$ such that*

$$\|\phi\|_{h_2^\sigma} = \sup_{k \geq 0} \|\phi(2^k \cdot) \varphi\|_{H_2^\sigma(\mathbb{R}^d; \ell_2)} < \infty. \quad (4.4)$$

Then for $1 < p < \infty$ and any finite sequence $(f_j) \subset L_p(\mathcal{N})$,

$$\left\| \left(\sum_{j \geq 0} |\tilde{\phi}_j * f_j|^2 \right)^{\frac{1}{2}} \right\|_p \lesssim \|\phi\|_{h_2^\sigma} \left\| \left(\sum_{j \geq 0} |f_j|^2 \right)^{\frac{1}{2}} \right\|_p.$$

The relevant constant depends only on p , φ , σ and d .

Proof. The argument below is standard. First, note that the Fourier multiplier on $\mathbb{T}^d$ with symbol ϕ_j does not depend on the values of ϕ_j in the open unit ball of $\mathbb{R}^d$. So letting η be an infinite differentiable function on $\mathbb{R}^d$ such that $\eta(\xi) = 0$ for $|\xi| \leq \frac{1}{2}$ and $\eta(\xi) = 1$ for $|\xi| \geq 1$, we see that ϕ_j and $\eta\phi_j$ induce the same Fourier multiplier on $\mathbb{T}^d$ (restricted to distributions with vanishing Fourier coefficients at the origin). On the other hand, it is easy to see that (4.4) implies that the sequence $(\eta\phi_j)_{j \geq 0}$ satisfies (4.2) with $(\eta\phi_j)_{j \geq 0}$ in place of ϕ . Thus replacing ϕ_j by $\eta\phi_j$ if necessary, we will assume that ϕ satisfies the stronger condition (4.2).

We will use the Calderón-Zygmund theory and consider $\tilde{\mathbf{k}}$ as a diagonal matrix with diagonal entries $(\mathbf{k}_j)_{j \geq 1}$. The Calderón-Zygmund operator associated to $\tilde{\mathbf{k}}$ is thus the convolution operator:

$$\tilde{\mathbf{k}}(f)(s) = \int_{\mathbb{I}^d} \tilde{\mathbf{k}}(s-t) f(t) dt$$

for any finite sequence $f = (f_j)$ (viewed as a column matrix). Then the assertion to prove amounts to the boundedness of $\tilde{\mathbf{k}}$ on $L_p(\mathcal{N}; \ell_2^c)$.

We first show that $\tilde{\mathbf{k}}$ is bounded from $L_\infty(\mathcal{N}; \ell_2^c)$ into $\text{BMO}^c(\mathbb{T}^d, B(\ell_2) \bar{\otimes} \mathcal{M})$. Let f be a finite sequence in $L_\infty(\mathcal{N}; \ell_2^c)$, and let Q be a cube of $\mathbb{I}^d$ whose center is denoted by c . We decompose f as $f = g + h$ with $g = f \mathbf{1}_{\tilde{Q}}$, where $\tilde{Q} = 2Q$, the cube with center c and twice the side length of Q . Setting

$$a = \int_{\mathbb{I}^d \setminus \tilde{Q}} \tilde{\mathbf{k}}(c-t) f(t) dt,$$

we have

$$\tilde{\mathbf{k}}(f)(s) - a = \tilde{\mathbf{k}}(g)(s) + \int_{\mathbb{I}^d} (\tilde{\mathbf{k}}(s-t) - \tilde{\mathbf{k}}(c-t)) h(t) dt.$$

Thus

$$\frac{1}{|Q|} \int_Q |\tilde{\mathbf{k}}(f)(s) - a|^2 ds \leq 2(A + B),$$

where

$$\begin{aligned} A &= \frac{1}{|Q|} \int_Q |\tilde{\mathbf{k}}(g)(s)|^2 ds, \\ B &= \frac{1}{|Q|} \int_Q \left| \int_{\mathbb{I}^d} (\tilde{\mathbf{k}}(s-t) - \tilde{\mathbf{k}}(c-t)) h(t) dt \right|^2 ds. \end{aligned}$$

The first term A is easy to estimate. Indeed, by (4.3) and the Plancherel formula,

$$\begin{aligned}
|Q|A &\leq \int_{\mathbb{I}^d} |\tilde{\mathbf{k}}(g)(s)|^2 ds = \sum_{m \in \mathbb{Z}^d} |\hat{\mathbf{k}}(m) \hat{g}(m)|^2 \\
&= \sum_{m \in \mathbb{Z}^d} \hat{g}(m)^* [\hat{\mathbf{k}}(m)^* \hat{\mathbf{k}}(m)] \hat{g}(m) \leq \sum_{m \in \mathbb{Z}^d} \|\hat{\mathbf{k}}(m)\|_{B(\ell_2)}^2 |\hat{g}(m)|^2 \\
&\leq \|\hat{\mathbf{k}}\|_{\ell_\infty(\mathbb{Z}^d; \ell_\infty)} \int_{\mathbb{I}^d} |g(s)|^2 ds \\
&\leq \|\hat{\mathbf{k}}\|_{\ell_\infty(\mathbb{Z}^d; \ell_2)} \int_{\mathbb{I}^d} |g(s)|^2 ds \lesssim |\tilde{Q}| \|f\|_{L_\infty(\mathcal{N}; \ell_2^c)}^2,
\end{aligned}$$

whence

$$\|A\|_{B(\ell_2) \otimes \mathcal{M}} \lesssim \|f\|_{L_\infty(\mathcal{N}; \ell_2^c)}^2.$$

To estimate B , let $h = (h_j)$. Then by (4.3), for any $s \in Q$ we have

$$\begin{aligned}
\left| \int_{\mathbb{I}^d} (\tilde{\mathbf{k}}(s-t) - \tilde{\mathbf{k}}(c-t)) h(t) dt \right|^2 &= \sum_j \left| \int_{\mathbb{I}^d} (\tilde{\mathbf{k}}_j(s-t) - \tilde{\mathbf{k}}_j(c-t)) h_j(t) dt \right|^2 \\
&\lesssim \sum_j \int_{\mathbb{I}^d \setminus \tilde{Q}} |\tilde{\mathbf{k}}_j(s-t) - \tilde{\mathbf{k}}_j(c-t)| |h_j(t)|^2 dt \\
&\lesssim \int_{\mathbb{I}^d \setminus \tilde{Q}} \|\tilde{\mathbf{k}}(s-t) - \tilde{\mathbf{k}}(c-t)\|_{\ell_\infty} \sum_j |h_j(t)|^2 dt \\
&\lesssim \|f\|_{L_\infty(\mathcal{N}; \ell_2^c)}^2 \int_{\mathbb{I}^d \setminus \tilde{Q}} \|\tilde{\mathbf{k}}(s-t) - \tilde{\mathbf{k}}(c-t)\|_{\ell_2} dt \\
&\lesssim \|f\|_{L_\infty(\mathcal{N}; \ell_2^c)}^2.
\end{aligned}$$

Thus

$$\begin{aligned}
\|B\|_{B(\ell_2) \otimes \mathcal{M}} &\leq \frac{1}{|Q|} \int_Q \left\| \int_{\mathbb{I}^d} (\tilde{\mathbf{k}}(s-t) - \tilde{\mathbf{k}}(c-t)) h(t) dt \right\|_{B(\ell_2) \otimes \mathcal{M}}^2 ds \\
&= \frac{1}{|Q|} \int_Q \left\| \left| \int_{\mathbb{I}^d} (\tilde{\mathbf{k}}(s-t) - \tilde{\mathbf{k}}(c-t)) h(t) dt \right|^2 \right\|_{B(\ell_2) \otimes \mathcal{M}} ds \\
&\lesssim \|f\|_{L_\infty(\mathcal{N}; \ell_2^c)}^2.
\end{aligned}$$

Therefore, $\tilde{\mathbf{k}}$ is bounded from $L_\infty(\mathcal{N}; \ell_2^c)$ into $\text{BMO}^c(\mathbb{T}^d, B(\ell_2) \otimes \mathcal{M})$.

We next show that $\tilde{\mathbf{k}}$ is bounded from $L_\infty(\mathcal{N}; \ell_2^c)$ into $\text{BMO}^r(\mathbb{T}^d, B(\ell_2) \otimes \mathcal{M})$. Let f, Q and a be as above. Now we have to estimate

$$\left\| \frac{1}{|Q|} \int_Q |[\tilde{\mathbf{k}}(f)(s) - a]^*|^2 ds \right\|_{B(\ell_2) \otimes \mathcal{M}}.$$

We will use the same decomposition $f = g + h$. Then

$$\frac{1}{|Q|} \int_Q |[\tilde{\mathbf{k}}(f)(s) - a]^*|^2 ds \leq 2(A' + B'),$$

where

$$\begin{aligned}
A' &= \frac{1}{|Q|} \int_Q |[\tilde{\mathbf{k}}(g)(s)]^*|^2 ds, \\
B' &= \frac{1}{|Q|} \int_Q \left| \int_{\mathbb{I}^d} [(\tilde{\mathbf{k}}(s-t) - \tilde{\mathbf{k}}(c-t)) h(t)]^* dt \right|^2 ds.
\end{aligned}$$

The estimate of B' can be reduced to that of B before. Indeed,

$$\begin{aligned} \|B'\|_{B(\ell_2)\overline{\otimes}\mathcal{M}} &\leq \frac{1}{|Q|} \int_Q \left\| \int_{\mathbb{I}^d} [\tilde{\mathbf{k}}(s-t) - \tilde{\mathbf{k}}(c-t))h(t)]^* dt \right\|_{B(\ell_2)\overline{\otimes}\mathcal{M}}^2 ds \\ &= \frac{1}{|Q|} \int_Q \left\| \left[\int_{\mathbb{I}^d} (\tilde{\mathbf{k}}(s-t) - \tilde{\mathbf{k}}(c-t))h(t) dt \right]^* \right\|_{B(\ell_2)\overline{\otimes}\mathcal{M}}^2 ds \\ &= \frac{1}{|Q|} \int_Q \left\| \int_{\mathbb{I}^d} (\tilde{\mathbf{k}}(s-t) - \tilde{\mathbf{k}}(c-t))h(t) dt \right\|_{B(\ell_2)\overline{\otimes}\mathcal{M}}^2 ds \\ &\lesssim \|f\|_{L_\infty(\mathcal{N};\ell_2^c)}^2. \end{aligned}$$

However, A' needs a different argument. Setting $g = (g_j)$, we have

$$\|A'\|_{B(\ell_2)\overline{\otimes}\mathcal{M}} = \sup \left\{ \frac{1}{|Q|} \int_Q \tau \left[\sum_{i,j} \tilde{\mathbf{k}}_i(g_i)(s) \tilde{\mathbf{k}}_j(g_j)(s)^* a_j^* a_i \right] ds \right\},$$

where the supremum runs over all $a = (a_i)$ in the unit ball of $\ell_2(L_2(\mathcal{M}))$. Considering a_i as a constant function on $\mathbb{I}^d$, we can write

$$a_i \tilde{\mathbf{k}}_i(g_i) = \tilde{\mathbf{k}}_i(a_i g_i).$$

Thus

$$\int_Q \tau \left[\sum_{i,j} \tilde{\mathbf{k}}_i(g_i)(s) \tilde{\mathbf{k}}_j(g_j)(s)^* a_j^* a_i \right] ds = \int_Q \left\| \sum_i \tilde{\mathbf{k}}_i(a_i g_i)(s) \right\|_{L_2(\mathcal{M})}^2 ds.$$

So by the Plancherel formula,

$$\begin{aligned} \int_Q \left\| \sum_i \tilde{\mathbf{k}}_i(a_i g_i)(s) \right\|_{L_2(\mathcal{M})}^2 ds &\leq \int_{\mathbb{I}^d} \left\| \sum_i \tilde{\mathbf{k}}_i(a_i g_i)(s) \right\|_{L_2(\mathcal{M})}^2 ds \\ &= \sum_{m \in \mathbb{Z}^d} \left\| \sum_i \widehat{\tilde{\mathbf{k}}_i}(m) a_i \widehat{g_i}(m) \right\|_{L_2(\mathcal{M})}^2. \end{aligned}$$

On the other hand, by the Cauchy-Schwarz inequality, (4.3) and the Plancherel formula once more, we have

$$\begin{aligned} \sum_{m \in \mathbb{Z}^d} \left\| \sum_i \widehat{\tilde{\mathbf{k}}_i}(m) a_i \widehat{g_i}(m) \right\|_{L_2(\mathcal{M})}^2 &\leq \sum_{m \in \mathbb{Z}^d} \|\widehat{\tilde{\mathbf{k}}}(m)\|_{\ell_2}^2 \sum_i \tau(|a_i \widehat{g_i}(m)|^2) \\ &\lesssim \sum_i \tau \left[a_i \sum_{m \in \mathbb{Z}^d} \widehat{g_i}(m) \widehat{g_i}(m)^* a_i^* \right] \\ &= \sum_i \tau \left[a_i \int_{\mathbb{I}^d} g_i(s) g_i(s)^* ds a_i^* \right] \\ &= \sum_i \tau \left[a_i \int_{\tilde{Q}} f_i(s) f_i(s)^* ds a_i^* \right] \\ &\leq |\tilde{Q}| \sum_i \tau[a_i \|f_i\|_{L_\infty(\mathcal{N})}^2 a_i^*] \\ &\lesssim |Q| \|f\|_{L_\infty(\mathcal{N};\ell_2^c)}^2 \sum_i \tau(|a_i|^2) \\ &\leq |Q| \|f\|_{L_\infty(\mathcal{N};\ell_2^c)}^2. \end{aligned}$$

Combining the above estimates, we get the desired estimate of A' :

$$\|A'\|_{B(\ell_2)\overline{\otimes}\mathcal{M}} \lesssim \|f\|_{L_\infty(\mathcal{N};\ell_2^c)}^2.$$

Thus, $\tilde{k}$ is bounded from $L_\infty(\mathcal{N}; \ell_2^c)$ into $\text{BMO}^r(\mathbb{T}^d, B(\ell_2) \otimes \mathcal{M})$, so is it from $L_\infty(\mathcal{N}; \ell_2^c)$ into $\text{BMO}(\mathbb{T}^d, B(\ell_2) \otimes \mathcal{M})$.

It is clear that k is bounded from $L_2(\mathcal{N}; \ell_2^c)$ into $L_2(B(\ell_2) \otimes \mathcal{N})$. Hence, by interpolation via (1.2) and Lemma 1.8, $\tilde{k}$ is bounded from $L_p(\mathcal{N}; \ell_2^c)$ into $L_p(B(\ell_2) \otimes \mathcal{N})$ for any $2 < p < \infty$. This is the announced assertion for $2 \leq p < \infty$. The case $1 < p < 2$ is obtained by duality. $\square$

Remark 4.5. In the commutative case, i.e., $\mathcal{M} = \mathbb{C}$, it is well known that the conclusion of the preceding lemma holds under the following weaker assumption on ϕ :

$$\sup_{k \geq 0} \left(\int_{\mathbb{R}^d} (1 + |s|^2)^\sigma \|\mathcal{F}^{-1}(\phi(2^k \cdot) \varphi)(s)\|_{\ell_\infty}^2 ds \right)^{\frac{1}{2}} < \infty. \quad (4.5)$$

Like at the beginning of the preceding proof, this assumption can be strengthened to

$$\sup_{k \in \mathbb{Z}} \left(\int_{\mathbb{R}^d} (1 + |s|^2)^\sigma \|\mathcal{F}^{-1}(\phi(2^k \cdot) \varphi)(s)\|_{\ell_\infty}^2 ds \right)^{\frac{1}{2}} < \infty.$$

Then if we consider $k = (k_j)_{j \geq 0}$ as a kernel with values in ℓ_∞ , Lemma 4.3 admits the following ℓ_∞ -analogue:

- $\|\widehat{k}\|_{L_\infty(\mathbb{R}^d; \ell_\infty)} < \infty$;
- $\sup_{t \in \mathbb{R}^d} \int_{|s| > 2|t|} \|k(s-t) - k(s)\|_{\ell_\infty} ds < \infty$.

Transferring this to the periodic case, we have

- $\|\widehat{k}\|_{\ell_\infty(\mathbb{Z}^d; \ell_\infty)} < \infty$;
- $\sup_{t \in \mathbb{I}^d} \int_{\{s \in \mathbb{I}^d : |s| > 2|t|\}} \|\tilde{k}(s-t) - \tilde{k}(s)\|_{\ell_\infty} ds < \infty$.

The last two properties of the kernel $\tilde{k}$ are exactly what is needed for the estimates of A and B in the proof of Lemma 4.4, so the conclusion holds when $\mathcal{M} = \mathbb{C}$. However, we do not know whether Lemma 4.4 remains true when (4.4) is weakened to (4.5).

Lemma 4.6. *Let $\phi = (\phi_j)_{j \geq 0}$ be a sequence of continuous functions on $\mathbb{R}^d \setminus \{0\}$ satisfying (4.4). Then for $1 \leq p \leq 2$ and any $f \in \mathcal{H}_p^c(\mathbb{T}^d, \mathcal{M})$,*

$$\left\| \left(\sum_{j \geq 0} |\tilde{\phi}_j * f|^2 \right)^{\frac{1}{2}} \right\|_{L_p(\mathcal{N})} \lesssim \|\phi\|_{h_2^\sigma} \|f\|_{\mathcal{H}_p^c}.$$

The relevant constant depends only on φ , σ and d .

Proof. Like in the proof of Lemma 4.4, we can assume, without loss of generality, that ϕ satisfies (4.2). We use again the Calderón-Zygmund theory. Now we view $\tilde{k} = (\tilde{k}_j)_{j \geq 0}$ as a column matrix and the associated Calderón-Zygmund operator $\tilde{k}$ as defined on $L_p(\mathcal{N})$:

$$\tilde{k}(f)(s) = \int_{\mathbb{I}^d} \tilde{k}(s-t) f(t) dt.$$

Thus $\tilde{k}$ maps functions to sequences of functions. We have to show that $\tilde{k}$ is bounded from $\mathcal{H}_p^c(\mathbb{T}^d, \mathcal{M})$ to $L_p(\mathcal{N}; \ell_2^c)$ for $1 \leq p \leq 2$. This is trivial for $p = 2$. So by Lemma 1.9 via interpolation, it suffices to consider the case $p = 1$. The argument below is based on the atomic decomposition of $\mathcal{H}_1^c(\mathbb{T}^d, \mathcal{M})$ obtained in [17] (see also [44]). Recall that an $\mathcal{M}^c$ -atom is a function $a \in L_1(\mathcal{M}; L_2^c(\mathbb{T}^d))$ such that

- a is supported by a cube $Q \subset \mathbb{T}^d \approx \mathbb{I}^d$;
- $\int_Q a(s)ds = 0$;
- $\tau[(\int_Q |a(s)|^2 ds)^{\frac{1}{2}}] \leq |Q|^{-\frac{1}{2}}$.

Thus we need only to show that for any atom a

$$\|\tilde{\mathbf{k}}(a)\|_{L_1(\mathcal{N}; \ell_2^c)} \lesssim 1.$$

Let Q be the supporting cube of a . By translation invariance of the operator $\tilde{\mathbf{k}}$, we can assume that Q is centered at the origin. Set $\tilde{Q} = 2Q$ as before. Then

$$\|\tilde{\mathbf{k}}(a)\|_{L_1(\mathcal{N}; \ell_2^c)} \leq \|\tilde{\mathbf{k}}(a)\mathbf{1}_{\tilde{Q}}\|_{L_1(\mathcal{N}; \ell_2^c)} + \|\tilde{\mathbf{k}}(a)\mathbf{1}_{\mathbb{I}^d \setminus \tilde{Q}}\|_{L_1(\mathcal{N}; \ell_2^c)}. \quad (4.6)$$

The operator convexity of the square function $x \mapsto |x|^2$ implies

$$\int_{\tilde{Q}} |\mathbf{k}(a)(s)| ds \leq |\tilde{Q}|^{\frac{1}{2}} \left(\int_{\tilde{Q}} |\tilde{\mathbf{k}}(a)(s)|^2 ds \right)^{\frac{1}{2}}.$$

However, by the Plancherel formula,

$$\begin{aligned} \int_{\tilde{Q}} |\tilde{\mathbf{k}}(a)(s)|^2 ds &\leq \int_{\mathbb{I}^d} |\tilde{\mathbf{k}}(a)(s)|^2 ds = \sum_{m \in \mathbb{Z}^d} |\widehat{\tilde{\mathbf{k}}(a)}(m)|^2 = \sum_{m \in \mathbb{Z}^d} |\widehat{\mathbf{k}}(m)\widehat{a}(m)|^2 \\ &\leq \|\widehat{\mathbf{k}}\|_{\ell_\infty(\mathbb{Z}^d; \ell_2)} \sum_{m \in \mathbb{Z}^d} |\widehat{a}(m)|^2 = \|\widehat{\mathbf{k}}\|_{\ell_\infty(\mathbb{Z}^d; \ell_2)} \int_Q |a(s)|^2 ds. \end{aligned}$$

Therefore, by (4.3)

$$\|\tilde{\mathbf{k}}(a)\mathbf{1}_{\tilde{Q}}\|_{L_1(\mathcal{N}; \ell_2^c)} = \tau \int_{\tilde{Q}} |\tilde{\mathbf{k}}(a)(s)| ds \lesssim |\tilde{Q}|^{\frac{1}{2}} \tau[(\int_Q |a(s)|^2 ds)^{\frac{1}{2}}] \lesssim 1.$$

This is the desired estimate of the first term of the right-hand side of (4.6). For the second, since a is of vanishing mean, for every $s \notin \tilde{Q}$ we can write

$$\tilde{\mathbf{k}}(a)(s) = \int_Q [\tilde{\mathbf{k}}(s-t) - \tilde{\mathbf{k}}(s)]a(t)dt.$$

Then by the Cauchy-Schwarz inequality via the operator convexity of the square function $x \mapsto |x|^2$, we have

$$|\tilde{\mathbf{k}}(a)(s)|^2 \leq \int_Q \|\tilde{\mathbf{k}}(s-t) - \tilde{\mathbf{k}}(s)\|_{\ell_2}^2 dt \cdot \int_Q \|\tilde{\mathbf{k}}(s-t) - \tilde{\mathbf{k}}(s)\|_{\ell_2}^2 |a(t)|^2 dt.$$

Thus by (4.3),

$$\begin{aligned} \|\tilde{\mathbf{k}}(a)\mathbf{1}_{\mathbb{I}^d \setminus \tilde{Q}}\|_{L_1(\mathcal{N}; \ell_2^c)} &= \tau \int_{\mathbb{I}^d \setminus \tilde{Q}} |\tilde{\mathbf{k}}(a)(s)| ds \\ &\lesssim \tau \left[\left(\int_Q \int_{\mathbb{I}^d \setminus \tilde{Q}} \|\tilde{\mathbf{k}}(s-t) - \tilde{\mathbf{k}}(s)\|_{\ell_2}^2 ds dt \right)^{\frac{1}{2}} \cdot \left(\int_Q \int_{\mathbb{I}^d \setminus \tilde{Q}} \|\tilde{\mathbf{k}}(s-t) - \tilde{\mathbf{k}}(s)\|_{\ell_2}^2 |a(t)|^2 ds dt \right)^{\frac{1}{2}} \right] \\ &\lesssim |Q|^{1/2} \tau[(\int_Q |a(s)|^2 ds)^{\frac{1}{2}}] \lesssim 1. \end{aligned}$$

Hence the desired assertion is proved. $\square$

By transference, the previous lemmas imply the following. According to our convention used in the previous chapters, the map $x \mapsto \tilde{\phi} * x$ denotes the Fourier multiplier associated to ϕ on $\mathbb{T}_\theta^d$.

Lemma 4.7. *Let $\phi = (\phi_j)_j$ satisfy (4.4).*

(i) *For $1 < p < \infty$ we have*

$$\left\| \left(\sum_{j \geq 0} |\tilde{\phi}_j * x_j|^2 \right)^{\frac{1}{2}} \right\|_p \lesssim \|\phi\|_{h_2^\sigma} \left\| \left(\sum_{j \geq 0} |x_j|^2 \right)^{\frac{1}{2}} \right\|_p, \quad x_j \in L_p(\mathbb{T}_\theta^d)$$

with relevant constant depending only on p, φ, σ and d .

(ii) *For $1 \leq p \leq 2$ we have*

$$\left\| \left(\sum_{j \geq 0} |\tilde{\phi}_j * x|^2 \right)^{\frac{1}{2}} \right\|_p \lesssim \|\phi\|_{h_2^\sigma} \|x\|_{\mathcal{H}_p^c}, \quad x \in \mathcal{H}_p^c(\mathbb{T}_\theta^d)$$

with relevant constant depending only on φ, σ and d .

We are now ready to prove Theorem 4.1.

Proof of Theorem 4.1. Let $\zeta_j = \phi_j(\varphi^{(j-1)} + \varphi^{(j)} + \varphi^{(j+1)})$. By (3.2) and the support assumption on $\phi_j \rho_j$, we have

$$\phi_j \rho_j = \zeta_j \bar{\rho}_j, \quad \text{so} \quad \tilde{\phi}_j * \tilde{\rho}_j * x = \tilde{\zeta}_j * \tilde{\rho}_j * x$$

for any distribution x on $\mathbb{T}_\theta^d$. We claim that $\zeta = (\zeta_j)_{j \geq 0}$ satisfies (4.4) in place of ϕ . Indeed, given $k \in \mathbb{N}_0$, by the support assumption on φ in (3.1), the sequence $\zeta(2^k \cdot) \varphi = (\zeta_j(2^k \cdot) \varphi)_{j \geq 0}$ has at most five nonzero terms of indices j such that $k-2 \leq j \leq k+2$. Thus

$$\|\zeta(2^k \cdot) \varphi\|_{H_2^\sigma(\mathbb{R}^d; \ell_2)} \leq \sum_{j=k-2}^{k+2} \|\zeta_j(2^k \cdot) \varphi\|_{H_2^\sigma(\mathbb{R}^d)}.$$

However, by Lemma 4.2,

$$\|\zeta_j(2^k \cdot) \varphi\|_{H_2^\sigma(\mathbb{R}^d)} \lesssim \|\phi_j(2^j \cdot) \varphi\|_{H_2^\sigma(\mathbb{R}^d)}, \quad k-2 \leq j \leq k+2,$$

where the relevant constant depends only on d, σ and φ . Therefore, the second condition of (4.1) yields the claim.

Now applying Lemma 4.7 (i) with ζ_j instead of ϕ_j and $x_j = 2^{j\alpha} \tilde{\varphi}_j * x$, we prove part (i) of the theorem.

To show part (ii), we need the characterization of $\mathcal{H}_1^c(\mathbb{T}_\theta^d)$ by discrete square functions stated in Lemma 1.10 with $\psi = I_{-\alpha} \rho$. Let x be a distribution on $\mathbb{T}_\theta^d$ with $\hat{x}(0) = 0$ such that

$$\left\| \left(\sum_{j \geq 0} 2^{j\alpha} |\tilde{\rho}_j * x|^2 \right)^{\frac{1}{2}} \right\|_1 < \infty.$$

Let $y = I^\alpha(x)$. Then the discrete square function of y associated to ψ is given by

$$s_\psi^c(y)^2 = \sum_{j \geq 0} |\tilde{\psi}_j * y|^2 = \sum_{j \geq 0} 2^{j\alpha} |\tilde{\rho}_j * x|^2.$$

So $y \in \mathcal{H}_1^c(\mathbb{T}_\theta^d)$ and

$$\|y\|_{\mathcal{H}_1^c} \approx \left\| \left(\sum_{j \geq 0} 2^{j\alpha} |\tilde{\rho}_j * x|^2 \right)^{\frac{1}{2}} \right\|_1.$$

We want to apply Lemma 4.7 (ii) to y but with a different multiplier in place of ϕ . To that end, let $\eta_j = 2^{j\alpha} I_{-\alpha} \phi_j$ and $\eta = (\eta_j)_{j \geq 0}$. We claim that η satisfies (4.1) too. The support condition of (4.1) is obvious for η . To prove the second one, by (3.2), we write

$$\eta_j(2^j \xi) \varphi(\xi) = |\xi|^{-\alpha} \varphi(\xi) \phi_j(2^j \xi) = |\xi|^{-\alpha} [\varphi(2^{-1} \xi) + \varphi(\xi) + \varphi(2\xi)] \varphi(\xi) \phi_j(2^j \xi).$$

Since $I_{-\alpha}(\varphi^{(-1)} + \varphi + \varphi^{(1)})$ is an infinitely differentiable function with compact support,

$$\int_{\mathbb{R}^d} (1 + |s|^2)^\sigma |\mathcal{F}^{-1}(I_{-\alpha}(\varphi^{(-1)} + \varphi + \varphi^{(1)}))(s)| ds < \infty.$$

Thus by Lemma 4.2,

$$\|\eta_j(2^j \cdot) \varphi\|_{H_2^s(\mathbb{R}^d)} \lesssim \|\phi_j(2^j \cdot) \varphi\|_{H_2^s(\mathbb{R}^d)},$$

whence the claim.

As in the first part of the proof, we define a new sequence ζ by setting $\zeta_j = \eta_j \rho_j$. Then the new sequence ζ satisfies (4.4) too and

$$\sup_{k \geq 0} \|\zeta_j(2^k \cdot) \varphi\|_{H_2^s(\mathbb{R}^d)} \lesssim \sup_{j \geq 0} \|\eta_j(2^j \cdot) \varphi\|_{H_2^s(\mathbb{R}^d)} \lesssim \sup_{j \geq 0} \|\phi_j(2^j \cdot) \varphi\|_{H_2^s(\mathbb{R}^d)}.$$

On the other hand, we have

$$2^{j\alpha} \tilde{\phi}_j * \tilde{\rho}_j * x = \tilde{\zeta}_j * y.$$

Thus we can apply Lemma 4.7 (ii) to y with this new ζ instead of ϕ , and as before, we get

$$\begin{aligned} \left\| \left(\sum_{j \geq 0} 2^{2j\alpha} |\tilde{\phi}_j * \tilde{\rho}_j * x|^2 \right)^{\frac{1}{2}} \right\|_1 &= \left\| \left(\sum_{j \geq 0} |\tilde{\zeta}_j * y|^2 \right)^{\frac{1}{2}} \right\|_1 \\ &\lesssim \sup_{k \geq 0} \|\zeta(2^k \cdot) \varphi\|_{H_2^s(\mathbb{R}^d; \ell_2)} \|y\|_{\mathcal{H}_p^c} \\ &\lesssim \sup_{j \geq 0} \|\phi_j(2^j \cdot) \varphi\|_{H_2^s(\mathbb{R}^d)} \left\| \left(\sum_{j \geq 0} 2^{j\alpha} |\tilde{\rho}_j * x|^2 \right)^{\frac{1}{2}} \right\|_1. \end{aligned}$$

Hence the proof of the theorem is complete. $\square$

4.2 Definitions and basic properties

As said at the beginning of this chapter, we consider the Triebel-Lizorkin spaces on $\mathbb{T}_\theta^d$ only for $q = 2$. In this case, there exist three different families of spaces according to the three choices of the internal ℓ_2 -norms.

Definition 4.8. Let $1 \leq p < \infty$ and $\alpha \in \mathbb{R}$.

- (i) The column Triebel-Lizorkin space $F_p^{\alpha, c}(\mathbb{T}_\theta^d)$ is defined by

$$F_p^{\alpha, c}(\mathbb{T}_\theta^d) = \{x \in \mathcal{S}'(\mathbb{T}_\theta^d) : \|x\|_{F_p^{\alpha, c}} < \infty\},$$

where

$$\|x\|_{F_p^{\alpha, c}} = |\hat{x}(0)| + \left\| \left(\sum_{k \geq 0} 2^{2k\alpha} |\tilde{\varphi}_k * x|^2 \right)^{\frac{1}{2}} \right\|_p.$$

- (ii) The row space $F_p^{\alpha, r}(\mathbb{T}_\theta^d)$ consists of all x such that $x^* \in F_p^{\alpha, c}(\mathbb{T}_\theta^d)$, equipped with the norm $\|x\|_{F_p^{\alpha, r}} = \|x^*\|_{F_p^{\alpha, c}}$.

(iii) The mixture space $F_p^\alpha(\mathbb{T}_\theta^d)$ is defined to be

$$F_p^\alpha(\mathbb{T}_\theta^d) = \begin{cases} F_p^{\alpha,c}(\mathbb{T}_\theta^d) + F_p^{\alpha,r}(\mathbb{T}_\theta^d) & \text{if } 1 \leq p < 2, \\ F_p^{\alpha,c}(\mathbb{T}_\theta^d) \cap F_p^{\alpha,r}(\mathbb{T}_\theta^d) & \text{if } 2 \leq p < \infty, \end{cases}$$

equipped with

$$\|x\|_{F_p^\alpha} = \begin{cases} \inf \{ \|y\|_{F_p^{\alpha,c}} + \|z\|_{F_p^{\alpha,r}} : x = y + z \} & \text{if } 1 \leq p < 2, \\ \max(\|x\|_{F_p^{\alpha,c}}, \|x\|_{F_p^{\alpha,r}}) & \text{if } 2 \leq p < \infty. \end{cases}$$

In the sequel, we will concentrate our study only on the column Triebel-Lizorkin spaces. All results will admit the row and mixture analogues. The following shows that $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ is independent of the choice of the function φ .

Proposition 4.9. *Let ψ be a Schwartz function satisfying the same condition (3.1) as φ . Let $\hat{\psi}_k = \psi^{(k)} = \psi(2^{-k}\cdot)$. Then*

$$\|x\|_{F_p^{\alpha,c}} \approx |\hat{x}(0)| + \left\| \left(\sum_{k \geq 0} 2^{2k\alpha} |\tilde{\psi}_k * x|^2 \right)^{\frac{1}{2}} \right\|_p.$$

Proof. Fix a distribution x on $\mathbb{T}_\theta^d$ with $\hat{x}(0) = 0$. By the support assumption on $\psi^{(k)}$ and (3.2), we have (with $\tilde{\varphi}_{-1} = 0$)

$$\tilde{\psi}_k * x = \sum_{j=-1}^1 \tilde{\psi}_k * \tilde{\varphi}_{k+j} * x.$$

Thus by Theorem 4.1,

$$\begin{aligned} \left\| \left(\sum_{k \geq 0} 2^{2k\alpha} |\tilde{\psi}_k * x|^2 \right)^{\frac{1}{2}} \right\|_p &\leq \sum_{j=-1}^1 \left\| \left(\sum_{k \geq 0} 2^{2k\alpha} |\tilde{\psi}_k * \tilde{\varphi}_{k+j} * x|^2 \right)^{\frac{1}{2}} \right\|_p \\ &\lesssim \left\| \left(\sum_{k \geq 0} 2^{2k\alpha} |\tilde{\varphi}_k * x|^2 \right)^{\frac{1}{2}} \right\|_p. \end{aligned}$$

Changing the role of φ and ψ , we get the reverse inequality. $\square$

Proposition 4.10. *Let $1 \leq p < \infty$ and $\alpha \in \mathbb{R}$.*

- (i) $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ is a Banach space.
- (ii) $F_p^{\alpha,c}(\mathbb{T}_\theta^d) \subset F_p^{\beta,c}(\mathbb{T}_\theta^d)$ for $\beta < \alpha$.
- (iii) $\mathcal{P}_\theta$ is dense in $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$.
- (iv) $F_p^{0,c}(\mathbb{T}_\theta^d) = \mathcal{H}_p^c(\mathbb{T}_\theta^d)$.
- (v) $B_{p, \min(p,2)}^\alpha(\mathbb{T}_\theta^d) \subset F_p^{\alpha,c}(\mathbb{T}_\theta^d) \subset B_{p, \max(p,2)}^\alpha(\mathbb{T}_\theta^d)$.

Proof. (i) is proved as in the case of Besov spaces; see the corresponding proof of Proposition 3.3. (ii) is obvious. To show (iii), we use the Fejér means as in the proof of Proposition 2.7. We need one more property of those means, that is, they are completely contractive. So they are also contractive on $L_p(B(\ell_2) \otimes \mathbb{T}_\theta^d)$, in particular, on the column subspace $L_p(\mathbb{T}_\theta^d; \ell_2^c)$ too. We then deduce that F_N is contractive on $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ and $\lim_{N \rightarrow \infty} F_N(x) = x$ for every $x \in F_p^{\alpha,c}(\mathbb{T}_\theta^d)$.

(iv) has been already observed during the proof of Theorem 4.1. Indeed, for any distribution x on $\mathbb{T}_\theta^d$, the square function associated to φ defined in Lemma 1.10 is given by

$$s_\varphi^c(x) = \left(\sum_{k \geq 0} |\tilde{\varphi}_k * x|^2 \right)^{\frac{1}{2}}.$$

Thus $\|x\|_{\mathcal{H}_p^c} \approx \|x\|_{F_p^{0,c}}$.

(v) follows from the following well-known property:

$$\ell_2(L_p(\mathbb{T}_\theta^d)) \subset L_p(\mathbb{T}_\theta^d; \ell_2^c) \subset \ell_p(L_p(\mathbb{T}_\theta^d))$$

are contractive inclusions for $2 \leq p \leq \infty$; both inclusions are reversed for $1 \leq p \leq 2$. Note that the first inclusion is an immediate consequence of the triangular inequality of $L_{\frac{p}{2}}(\mathbb{T}_\theta^d)$, the second is proved by complex interpolation. $\square$

The following is the Triebel-Lizorkin analogue of Theorem 3.7. We keep the notation introduced before that theorem.

Theorem 4.11. *Let $1 \leq p < \infty$ and $\alpha \in \mathbb{R}$.*

(i) *For any $\beta \in \mathbb{R}$, both J^β and I^β are isomorphisms between $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ and $F_p^{\alpha-\beta,c}(\mathbb{T}_\theta^d)$. In particular, J^α and I^α are isomorphisms between $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ and $\mathcal{H}_p^c(\mathbb{T}_\theta^d)$.*

(ii) *Let $a \in \mathbb{R}_+^d$. If $x \in F_p^{\alpha,c}(\mathbb{T}_\theta^d)$, then $D^a x \in F_p^{\alpha-|a|_1,c}(\mathbb{T}_\theta^d)$ and*

$$\|D^a x\|_{F_p^{\alpha-|a|_1,c}} \lesssim \|x\|_{F_p^{\alpha,c}}.$$

(iii) *Let $\beta > 0$. Then $x \in F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ iff $D_i^\beta x \in F_p^{\alpha-\beta,c}(\mathbb{T}_\theta^d)$ for all $i = 1, \dots, d$. Moreover, in this case,*

$$\|x\|_{F_p^{\alpha,c}} \approx |\hat{x}(0)| + \sum_{i=1}^d \|D_i^\beta x\|_{F_p^{\alpha-\beta,c}}.$$

Proof. (i) Let $x \in F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ with $\hat{x}(0) = 0$. By Theorem 4.1,

$$\begin{aligned} \|J^\beta x\|_{F_p^{\alpha-\beta,c}} &= \left\| \left(\sum_{k \geq 0} 2^{2k(\alpha-\beta)} |J^\beta * \tilde{\varphi}_k * x|^2 \right)^{\frac{1}{2}} \right\|_p \\ &\lesssim \sup_{k \geq 0} 2^{-k\beta} \|J_\beta(2^k \cdot) \varphi\|_{H_2^c(\mathbb{R}^d)} \left\| \left(\sum_{k \geq 0} 2^{2k\alpha} |\tilde{\varphi}_k * x|^2 \right)^{\frac{1}{2}} \right\|_p. \end{aligned}$$

However, it is easy to see that all partial derivatives of the function $2^{-k\beta} J_\beta(2^k \cdot) \varphi$, of order less than a fixed integer, are bounded uniformly in k . It then follows that

$$\sup_{k \geq 0} 2^{-k\beta} \|J_\beta(2^k \cdot) \varphi\|_{H_2^c(\mathbb{R}^d)} < \infty.$$

Thus $\|J^\beta x\|_{F_p^{\alpha-\beta,c}} \lesssim \|x\|_{F_p^{\alpha,c}}$. So J^β is bounded from $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ to $F_p^{\alpha-\beta,c}(\mathbb{T}_\theta^d)$, its inverse, which is $J^{-\beta}$, is bounded too. I^β is handled similarly.

If $\beta = \alpha$, then $F_p^{\alpha-\beta,c}(\mathbb{T}_\theta^d) = F_p^{0,c}(\mathbb{T}_\theta^d) = \mathcal{H}_p^c(\mathbb{T}_\theta^d)$ by Proposition 4.10 (iv).

(ii) This proof is similar to the previous one by replacing J^β by D^a and using Lemma 3.5.

(iii) One implication is contained in (ii). To show the other, we follow the proof of Theorem 3.7 (iii) and keep the notation there. Since

$$\varphi_k = \sum_{i=1}^d \chi_i D_{i,\beta} \varphi_k,$$

by Theorem 4.1,

$$\begin{aligned} \|x\|_{F_p^{\alpha,c}} &\leq \sum_{i=1}^d \left\| \left(\sum_{k \geq 0} 2^{2k\alpha} |\tilde{\chi}_i * \tilde{\varphi}_k * D_i^\beta x|^2 \right)^{\frac{1}{2}} \right\|_p \\ &\lesssim \sum_{i=1}^d \sup_{k \geq 0} 2^{k\beta} \|\chi_i(2^k \cdot) \varphi\|_{H_2^\sigma(\mathbb{R}^d)} \left\| \left(\sum_{k \geq 0} 2^{2k(\alpha-\beta)} |\tilde{\varphi}_k * D_i^\beta x|^2 \right)^{\frac{1}{2}} \right\|_p. \end{aligned}$$

However,

$$2^{k\beta} \|\chi_i(2^k \cdot) \varphi\|_{H_2^\sigma(\mathbb{R}^d)} = \|\psi \varphi\|_{H_2^\sigma(\mathbb{R}^d)},$$

where

$$\psi(\xi) = \frac{1}{\chi(2^k \xi_1) |\xi_1|^\beta + \cdots + \chi(2^k \xi_d) |\xi_d|^\beta} \frac{\chi(2^k \xi_i) |\xi_i|^\beta}{(2\pi i \xi_i)^\beta}.$$

As all partial derivatives of $\psi \varphi$, of order less than a fixed integer, are bounded uniformly in k , the norm of $\psi \varphi$ in $H_2^\sigma(\mathbb{R}^d)$ are controlled by a constant independent of k . We then deduce

$$\|x\|_{F_p^{\alpha,c}} \lesssim \sum_{i=1}^d \left\| \left(\sum_{k \geq 0} 2^{2k(\alpha-\beta)} |\tilde{\varphi}_k * D_i^\beta x|^2 \right)^{\frac{1}{2}} \right\|_p = \sum_{i=1}^d \|D_i^\beta x\|_{F_p^{\alpha-\beta,c}}.$$

The theorem is thus completely proved. $\square$

Corollary 4.12. *Let $1 < p < \infty$ and $\alpha \in \mathbb{R}$. Then $F_p^\alpha(\mathbb{T}_\theta^d) = H_p^\alpha(\mathbb{T}_\theta^d)$ with equivalent norms.*

Proof. Since J^α is an isomorphism from $F_p^\alpha(\mathbb{T}_\theta^d)$ onto $F_p^0(\mathbb{T}_\theta^d)$, and from $H_p^\alpha(\mathbb{T}_\theta^d)$ onto $H_p^0(\mathbb{T}_\theta^d)$, it suffices to consider the case $\alpha = 0$. But then $H_p^0(\mathbb{T}_\theta^d) = L_p(\mathbb{T}_\theta^d)$ by definition, and $F_p^0(\mathbb{T}_\theta^d) = \mathcal{H}_p(\mathbb{T}_\theta^d)$ by Proposition 4.10. It remains to apply Lemma 1.9 to conclude $F_p^0(\mathbb{T}_\theta^d) = H_p^0(\mathbb{T}_\theta^d)$. $\square$

We now discuss the duality of $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$. For this we need to define $F_\infty^{\alpha,c}(\mathbb{T}_\theta^d)$ that is excluded from the definition at the beginning of the present section. Let ℓ_2^α denote the Hilbert space of all complex sequences $a = (a_k)_{k \geq 0}$ such that

$$\|a\| = \left(\sum_{k \geq 0} 2^{2k\alpha} |a_k|^2 \right)^{\frac{1}{2}} < \infty.$$

Thus $L_p(\mathbb{T}_\theta^d; \ell_2^{\alpha,c})$ is the column subspace of $L_p(B(\ell_2^\alpha) \otimes \mathbb{T}_\theta^d)$.

Definition 4.13. For $\alpha \in \mathbb{R}$ we define $F_\infty^{\alpha,c}(\mathbb{T}_\theta^d)$ as the space of all distributions x on $\mathbb{T}_\theta^d$ that admit a representation of the form

$$x = \sum_{k \geq 0} \tilde{\varphi}_k * x_k \quad \text{with} \quad (x_k)_{k \geq 0} \in L_\infty(\mathbb{T}_\theta^d; \ell_2^{\alpha,c}),$$

and endow it with the norm

$$\|x\|_{F_\infty^{\alpha,c}} = |\hat{x}(0)| + \inf \left\{ \left\| \left(\sum_{k \geq 0} 2^{2k\alpha} |\tilde{\varphi}_k * x_k|^2 \right)^{\frac{1}{2}} \right\|_\infty \right\},$$

where the infimum runs over all representations of x as above.

Proposition 4.14. *Let $1 \leq p < \infty$ and $\alpha \in \mathbb{R}$. Then the dual space of $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ coincides isomorphically with $F_{p'}^{-\alpha,c}(\mathbb{T}_\theta^d)$.*

Proof. For simplicity, we will consider only distributions with vanishing Fourier coefficients at $m = 0$. We view $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ as an isometric subspace of $L_p(\mathbb{T}_\theta^d; \ell_2^{\alpha,c})$ via $x \mapsto (\tilde{\varphi}_k * x)_{k \geq 0}$. Then the dual space of $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ is identified with the following quotient of the latter:

$$G_{p'} = \{y = \sum_{k \geq 0} \tilde{\varphi}_k * y_k : (y_k)_{k \geq 0} \in L_{p'}(\mathbb{T}_\theta^d; \ell_2^{-\alpha,c})\},$$

equipped with the quotient norm

$$\|y\| = \inf \{ \|(y_k)_{k \geq 0}\|_{L_{p'}(\mathbb{T}_\theta^d; \ell_2^{-\alpha,c})} : y = \sum_{k \geq 0} \tilde{\varphi}_k * y_k \}.$$

The duality bracket is given by $\langle x, y \rangle = \tau(xy^*)$. If $p = 1$, then $G_{p'} = F_\infty^{-\alpha,c}(\mathbb{T}_\theta^d)$ by definition. It remains to show that $G_{p'} = F_{p'}^{-\alpha,c}(\mathbb{T}_\theta^d)$ for $1 < p < \infty$. It is clear that $F_{p'}^{-\alpha,c}(\mathbb{T}_\theta^d) \subset G_{p'}$, a contractive inclusion. Conversely, let $y \in G_{p'}$ and $y = \sum \tilde{\varphi}_k * y_k$ for some $(y_k)_{k \geq 0} \in L_{p'}(\mathbb{T}_\theta^d; \ell_2^{-\alpha,c})$. Then

$$\tilde{\varphi}_k * y = \tilde{\varphi}_k * \tilde{\varphi}_{k-1} * y_{k-1} + \tilde{\varphi}_k * \tilde{\varphi}_k * y_k + \tilde{\varphi}_k * \tilde{\varphi}_{k+1} * y_{k+1}.$$

Therefore, by Lemma 4.7,

$$\begin{aligned} \left\| \left(\sum_{k \geq 0} 2^{2k\alpha} |\tilde{\varphi}_k * y|^2 \right)^{\frac{1}{2}} \right\|_{p'} &\leq \sum_{j=-1}^1 \left\| \left(\sum_{k \geq 0} 2^{-2k\alpha} |\tilde{\varphi}_k * \tilde{\varphi}_{k+j} * y_{k+j}|^2 \right)^{\frac{1}{2}} \right\|_{p'} \\ &\lesssim \left\| \left(\sum_{k \geq 0} 2^{-2k\alpha} |y_k|^2 \right)^{\frac{1}{2}} \right\|_{p'}. \end{aligned}$$

Thus $y \in F_{p'}^{-\alpha,c}(\mathbb{T}_\theta^d)$ and $\|y\|_{F_{p'}^{-\alpha,c}} \lesssim \|y\|_{G_{p'}}$. $\square$

Remark 4.15. (i) The above proof shows that $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ is a complemented subspace of $L_p(\mathbb{T}_\theta^d; \ell_2^{\alpha,c})$ for $1 < p < \infty$.

(ii) By duality, Propositions 4.9, 4.10 and Theorem 4.11 remain valid for $p = \infty$, except the density of $\mathcal{P}_\theta$. In particular, $F_\infty^{0,c}(\mathbb{T}_\theta^d) = \text{BMO}^c(\mathbb{T}_\theta^d)$.

We conclude this section with the following Fourier multiplier theorem, which is an immediate consequence of Theorem 4.1 for $p < \infty$. The case $p = \infty$ is obtained by duality. In the case of $\alpha = 0$, this result is to be compared with Lemma 1.7 where more smoothness of ϕ is assumed.

Theorem 4.16. *Let ϕ be a continuous function on $\mathbb{R}^d \setminus \{0\}$ such that*

$$\sup_{k \geq 0} \|\phi(2^k \cdot) \varphi\|_{H_2^\sigma(\mathbb{R}^d)} < \infty$$

for some $\sigma > \frac{d}{2}$. Then ϕ is a bounded Fourier multiplier on $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ for all $1 \leq p \leq \infty$ and $\alpha \in \mathbb{R}$.

In particular, ϕ is a bounded Fourier multiplier on $\mathcal{H}_p^c(\mathbb{T}_\theta^d)$ for $1 \leq p < \infty$ and on $\text{BMO}^c(\mathbb{T}_\theta^d)$.

4.3 A general characterization

In this section we give a general characterization of Triebel-Lizorkin spaces on $\mathbb{T}_\theta^d$ in the same spirit as that given in section 3.2 for Besov spaces.

Let $\alpha_0, \alpha_1, \sigma \in \mathbb{R}$ with $\sigma > \frac{d}{2}$. Let h be a Schwartz function satisfying (3.3). Assume that ψ is an infinitely differentiable function on $\mathbb{R}^d \setminus \{0\}$ such that

$$\begin{cases} |\psi| > 0 \text{ on } \{\xi : 2^{-1} \leq |\xi| \leq 2\}, \\ \int_{\mathbb{R}^d} (1 + |s|^2)^\sigma |\mathcal{F}^{-1}(\psi h I_{-\alpha_1})(s)| ds < \infty, \\ \sup_{k \in \mathbb{N}_0} 2^{-k\alpha_0} \|\mathcal{F}^{-1}(\psi(2^k \cdot) \varphi)\|_{H_2^\sigma(\mathbb{R}^d)} < \infty. \end{cases} \quad (4.7)$$

Writing $\varphi = \varphi^{(-1)} + \varphi + \varphi^{(1)}$ and using Lemma 4.2, we have

$$\|\mathcal{F}^{-1}(\psi(2^k \cdot) \varphi)\|_{H_2^\sigma(\mathbb{R}^d)} \lesssim \int_{\mathbb{R}^d} (1 + |s|^2)^\sigma |\mathcal{F}^{-1}(\psi(2^k \cdot) \varphi)(s)| ds.$$

So the third condition of (4.7) is weaker than the corresponding one assumed in [74, Theorem 2.4.1]. On the other hand, consistent with Theorem 3.9 but contrary to [74, Theorem 2.4.1], our following theorem does not require that $\alpha_1 > 0$.

Theorem 4.17. *Let $1 \leq p < \infty$ and $\alpha \in \mathbb{R}$. Assume that $\alpha_0 < \alpha < \alpha_1$ and ψ satisfies (4.7). Then for any distribution x on $\mathbb{T}_\theta^d$, we have*

$$\|x\|_{F_p^{\alpha,c}} \approx |\widehat{x}(0)| + \left\| \left(\sum_{k \geq 0} 2^{2k\alpha} |\widetilde{\psi}_k * x|^2 \right)^{\frac{1}{2}} \right\|_p. \quad (4.8)$$

The equivalence is understood in the sense that whenever one side is finite, so is the other, and the two are then equivalent with constants independent of x .

Proof. Although it resembles, in form, the proof of Theorem 3.9, the proof given below is harder and subtler than the Besov space case. The key new ingredient is Theorem 4.1. The main differences will already appear in the first part of the proof, which is an adaptation of step 1 of the proof of Theorem 3.9. In the following, we will fix x with $\widehat{x}(0) = 0$. By approximation, we can assume that x is a polynomial. We will denote the right-hand side of (4.8) by $\|x\|_{F_{p,\psi}^{\alpha,c}}$.

Given a positive integer K , we write, as before

$$\psi^{(j)} = \sum_{k=0}^{\infty} \psi^{(j)} \varphi^{(k)} = \sum_{k=-\infty}^K \psi^{(j)} \varphi^{(j+k)} + \sum_{k=K}^{\infty} \psi^{(j)} \varphi^{(j+k)}.$$

Then

$$\|x\|_{F_{p,\psi}^{\alpha,c}} \leq \text{I} + \text{II}, \quad (4.9)$$

where

$$\begin{aligned} \text{I} &= \sum_{k \leq K} \left\| \left(\sum_j 2^{2j\alpha} |\widetilde{\psi}_j * \widetilde{\varphi}_{j+k} * x|^2 \right)^{\frac{1}{2}} \right\|_p, \\ \text{II} &= \sum_{k > K} \left\| \left(\sum_j 2^{2j\alpha} |\widetilde{\psi}_j * \widetilde{\varphi}_{j+k} * x|^2 \right)^{\frac{1}{2}} \right\|_p. \end{aligned}$$

The estimate of the term I corresponds to step 1 of the proof of Theorem 3.9. We use again (3.8) with η and ρ defined there. Then applying Theorem 4.1 twice, we have

$$\begin{aligned}
\text{I} &= \sum_{k \leq K} 2^{k(\alpha_1 - \alpha)} \left\| \left(\sum_j 2^{2(j+k)\alpha} |\tilde{\eta}_j * \tilde{\rho}_{j+k} * x|^2 \right)^{\frac{1}{2}} \right\|_p \\
&= \sum_{k \leq K} 2^{k(\alpha_1 - \alpha)} \left\| \left(\sum_j 2^{2j\alpha} |\tilde{\eta}_{j-k} * \tilde{\rho}_j * x|^2 \right)^{\frac{1}{2}} \right\|_p \\
&\lesssim \sum_{k \leq K} 2^{k(\alpha_1 - \alpha)} \|\eta^{(-k)}\varphi\|_{H_2^\sigma} \left\| \left(\sum_j 2^{2j\alpha} |\tilde{\rho}_j * x|^2 \right)^{\frac{1}{2}} \right\|_p \\
&\lesssim \|I_{\alpha_1}\varphi\|_{H_2^\sigma} \sum_{k \leq K} 2^{k(\alpha_1 - \alpha)} \|\eta^{(-k)}\varphi\|_{H_2^\sigma} \left\| \left(\sum_j 2^{2j\alpha} |\tilde{\varphi}_j * x|^2 \right)^{\frac{1}{2}} \right\|_p \\
&= \|I_{\alpha_1}\varphi\|_{H_2^\sigma} \sum_{k \leq K} 2^{k(\alpha_1 - \alpha)} \|\eta^{(-k)}\varphi\|_{H_2^\sigma} \|x\|_{F_p^{\alpha, c}}.
\end{aligned}$$

Being an infinitely differentiable function with compact support, $I_{\alpha_1}\varphi$ belongs to $H_2^\sigma(\mathbb{R}^d)$, that is, $\|I_{\alpha_1}\varphi\|_{H_2^\sigma} < \infty$. Next, we must estimate $\|\eta^{(-k)}\varphi\|_{H_2^\sigma}$ uniformly in k . To that end, for $s \in \mathbb{R}^d$, using

$$\begin{aligned}
|\mathcal{F}^{-1}(\eta^{(-k)}\varphi)(s)|^2 &= \left| \int_{\mathbb{R}^d} \mathcal{F}^{-1}(\eta)(t) * \mathcal{F}^{-1}(\varphi)(s - 2^k t) dt \right|^2 \\
&\leq \|\mathcal{F}^{-1}(\eta)\|_1 \int_{\mathbb{R}^d} |\mathcal{F}^{-1}(\eta)(t)| |\mathcal{F}^{-1}(\varphi)(s - 2^k t)|^2 dt,
\end{aligned}$$

for $k \leq K$, we have

$$\begin{aligned}
\|\eta^{(-k)}\varphi\|_{H_2^\sigma}^2 &= \int_{\mathbb{R}^d} (1 + |s|^2)^\sigma |\mathcal{F}^{-1}(\eta^{(-k)}\varphi)(s)|^2 ds \\
&\leq \|\mathcal{F}^{-1}(\eta)\|_1 \int_{\mathbb{R}^d} (1 + |s|^2)^\sigma \int_{\mathbb{R}^d} |\mathcal{F}^{-1}(\eta)(t)| |\mathcal{F}^{-1}(\varphi)(s - 2^k t)|^2 dt ds \\
&\lesssim \|\mathcal{F}^{-1}(\eta)\|_1 \int_{\mathbb{R}^d} (1 + |2^k t|^2)^\sigma |\mathcal{F}^{-1}(\eta)(t)| \int_{\mathbb{R}^d} (1 + |s - 2^k t|^2)^\sigma |\mathcal{F}^{-1}(\varphi)(s - 2^k t)|^2 ds dt \\
&\leq 2^{K\sigma} \|\mathcal{F}^{-1}(\eta)\|_1 \int_{\mathbb{R}^d} (1 + |t|^2)^\sigma |\mathcal{F}^{-1}(\eta)(t)| dt \int_{\mathbb{R}^d} (1 + |s|^2)^\sigma |\mathcal{F}^{-1}(\varphi)(s)|^2 ds \\
&\leq c_{\varphi, \sigma, K} \left(\int_{\mathbb{R}^d} (1 + |t|^2)^\sigma |\mathcal{F}^{-1}(\eta)(t)| dt \right)^2.
\end{aligned}$$

In order to return back from η to ψ , write

$$\eta = I_{-\alpha_1}\psi h + I_{-\alpha_1}\psi(h^{(K)} - h).$$

Note that

$$\int_{\mathbb{R}^d} (1 + |t|^2)^\sigma |\mathcal{F}^{-1}(I_{-\alpha_1}\psi(h^{(K)} - h))(t)| dt = c_{\psi, h, \alpha_1, \sigma, K} < \infty \quad (4.10)$$

since $I_{-\alpha_1}\psi(h^{(K)} - h)$ is an infinitely differentiable function with compact support. We then deduce

$$\int_{\mathbb{R}^d} (1 + |t|^2)^\sigma |\mathcal{F}^{-1}(\eta)(t)| dt \lesssim \int_{\mathbb{R}^d} (1 + |t|^2)^\sigma |\mathcal{F}^{-1}(I_{-\alpha_1}\psi h)(t)| dt.$$

The term on the right-hand side is the second condition of (4.7). Combining the preceding inequalities, we obtain

$$\text{I} \lesssim \int_{\mathbb{R}^d} (1 + |t|^2)^\sigma |\mathcal{F}^{-1}(I_{-\alpha_1}\psi h)(t)| dt \|x\|_{F_p^{\alpha, c}}.$$

The second term II on the right-hand side of (4.9) is easier to estimate. Using (3.11), Theorem 4.1 and arguing as in the preceding part for the term I, we obtain

$$\begin{aligned} \text{II} &\lesssim \|I_{\alpha_0}\varphi\|_{H_2^\sigma} \sum_{k>K} 2^{-2k\alpha} \|I_{-\alpha_0}\psi(2^k\cdot)H\varphi\|_{H_2^\sigma} \|x\|_{F_p^{\alpha,c}} \\ &\lesssim \sum_{k>K} 2^{-2k\alpha} \|I_{-\alpha_0}\psi(2^k\cdot)H\varphi\|_{H_2^\sigma} \|x\|_{F_p^{\alpha,c}}, \end{aligned}$$

where $H = \varphi(2^{-1}\cdot) + \varphi + \varphi(2\cdot)$. To treat the last Sobolev norm, noting that $I_{-\alpha_0}H$ is an infinitely differentiable function with compact support, by Lemma 4.2, we have

$$\|I_{-\alpha_0}\psi(2^k\cdot)H\varphi\|_{H_2^\sigma} \leq \|\psi(2^k\cdot)\varphi\|_{H_2^\sigma} \int_{\mathbb{R}^d} (1+|t|^2)^\sigma |\mathcal{F}^{-1}(I_{-\alpha_0}H)(t)| dt \lesssim \|\psi(2^k\cdot)\varphi\|_{H_2^\sigma}.$$

Therefore,

$$\begin{aligned} \text{II} &\lesssim \sup_{k>K} 2^{-k\alpha_0} \|\psi(2^k\cdot)\varphi\|_{H_2^\sigma} \sum_{k>K} 2^{2k(\alpha_0-\alpha)} \|x\|_{F_p^{\alpha,c}} \\ &\leq c \sup_{k>K} 2^{-k\alpha_0} \|\psi(2^k\cdot)\varphi\|_{H_2^\sigma} \frac{2^{(\alpha_0-\alpha)K}}{1-2^{\alpha_0-\alpha}} \|x\|_{F_p^{\alpha,c}} \end{aligned} \quad (4.11)$$

with some constant c independent of K . Putting this estimate together with that of I, we finally get

$$\|x\|_{F_{p,\psi}^{\alpha,c}} \lesssim \|x\|_{F_p^{\alpha,c}}.$$

Now we show the reverse inequality by following step 3 of the proof of Theorem 3.9 (recalling that $\lambda = 1 - h$). By (3.13) and Theorem 4.1,

$$\begin{aligned} \|x\|_{F_p^{\alpha,c}} &\lesssim \|\psi^{-1}\varphi^2\|_{H_2^\sigma} \left\| \left(\sum_{j=0}^{\infty} 2^{2j\alpha} |\tilde{h}_{j+K} * \tilde{\psi}_j * x|^2 \right)^{\frac{1}{2}} \right\|_p \\ &\lesssim \left\| \left(\sum_{j=0}^{\infty} 2^{2j\alpha} |\tilde{h}_{j+K} * \tilde{\psi}_j * x|^2 \right)^{\frac{1}{2}} \right\|_p \\ &\leq \|x\|_{F_{p,\psi}^{\alpha,c}} + \left\| \left(\sum_{j=0}^{\infty} 2^{2j\alpha} |\tilde{\lambda}_{j+K} * \tilde{\psi}_j * x|^2 \right)^{\frac{1}{2}} \right\|_p. \end{aligned}$$

Then combining the arguments in step 3 of the proof of Theorem 3.9 and (4.11) with $\lambda^{(K)}\psi$ in place of ψ , we deduce

$$\left\| \left(\sum_{j=0}^{\infty} 2^{2j\alpha} |\tilde{\lambda}_{j+K} * \tilde{\psi}_j * x|^2 \right)^{\frac{1}{2}} \right\|_p \leq c \sup_{k>K} 2^{-k\alpha_0} \|\lambda(2^{k-K}\cdot)\psi(2^k\cdot)\varphi\|_{H_2^\sigma} \frac{2^{(\alpha_0-\alpha)K}}{1-2^{\alpha_0-\alpha}} \|x\|_{F_p^{\alpha,c}}.$$

To remove $\lambda(2^{k-K}\cdot)$ from the above Sobolev norm, by triangular inequality, we have

$$\|\lambda(2^{k-K}\cdot)\psi(2^k\cdot)\varphi\|_{H_2^\sigma} \leq \|\psi(2^k\cdot)\varphi\|_{H_2^\sigma} + \|h(2^{k-K}\cdot)\psi(2^k\cdot)\varphi\|_{H_2^\sigma}.$$

By the support assumption on h and φ , $h(2^{k-K}\cdot)\varphi \neq 0$ only for $k \leq K+2$, so the second term on the right hand side above matters only for $k = K+1$ and $k = K+2$. But for these two values of k , by Lemma 4.2, we have

$$\|h(2^{k-K}\cdot)\psi(2^k\cdot)\varphi\|_{H_2^\sigma} \leq c' \|\psi(2^k\cdot)\varphi\|_{H_2^\sigma},$$

where c' depends only on h . Thus

$$\|\lambda(2^{k-K}\cdot)\psi(2^k\cdot)\varphi\|_{H_2^\sigma} \leq (1+c')\|\psi(2^k\cdot)\varphi\|_{H_2^\sigma}.$$

Putting together all estimates so far obtained, we deduce

$$\|x\|_{F_p^{\alpha,c}} \leq \|x\|_{F_{p,\psi}^{\alpha,c}} + c(1+c') \sup_{k \geq K} 2^{-k\alpha_0} \|\psi(2^k\cdot)\varphi\|_{H_2^\sigma} \frac{2^{(\alpha_0-\alpha)K}}{1-2^{\alpha_0-\alpha}} \|x\|_{F_p^{\alpha,c}}.$$

So if K is chosen sufficiently large, we finally obtain

$$\|x\|_{F_p^{\alpha,c}} \lesssim \|x\|_{F_{p,\psi}^{\alpha,c}},$$

which finishes the proof of the theorem. $\square$

Remark 4.18. Note that we have used the infinite differentiability of ψ only to insure (4.10), which holds whenever ψ is continuously differentiable up to order $[\frac{3d}{2}] + 1$. More generally, we need only to assume that there exists $\sigma > \frac{3d}{2} + 1$ such that $\psi\eta \in H_2^\sigma(\mathbb{R}^d)$ for any compactly supported infinite differentiable function η which vanishes in a neighborhood of the origin.

Like in the case of Besov spaces, Theorem 4.17 admits the following continuous version.

Theorem 4.19. *Under the assumption of the previous theorem, for any distribution x on $\mathbb{T}_\theta^d$,*

$$\|x\|_{F_p^{\alpha,c}} \approx |\widehat{x}(0)| + \left\| \left(\int_0^1 \varepsilon^{-2\alpha} |\widetilde{\psi}_\varepsilon * x|^2 \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{2}} \right\|_p.$$

Proof. This proof is very similar to that of Theorem 4.17. The main idea is, of course, to discretize the continuous square function:

$$\int_0^1 \varepsilon^{-2\alpha} |\widetilde{\psi}_\varepsilon * x|^2 \frac{d\varepsilon}{\varepsilon} \approx \sum_{j=0}^{\infty} 2^{2j\alpha} \int_{2^{-j-1}}^{2^{-j}} |\widetilde{\psi}_\varepsilon * x|^2 \frac{d\varepsilon}{\varepsilon}.$$

We can further discretize the internal integrals on the right-hand side. Indeed, by approximation and assuming that x is a polynomial, each internal integral can be approximated uniformly by discrete sums. Then we follow the proof of Theorem 3.11 with necessary modifications as in the preceding proof. The only difference is that when Theorem 4.1 is applied, the L_1 -norm of the inverse Fourier transforms of the various functions in consideration there must be replaced by the two norms of these functions appearing in (4.7). We omit the details. $\square$

4.4 Concrete characterizations

This section concretizes the general characterization in the previous one in terms of the Poisson and heat kernels. We keep the notation introduced in section 3.3.

The following result improves [74, Section 2.6.4] at two aspects even in the classical case: Firstly, in addition to derivation operators, it can also use integration operators (corresponding to negative k); secondly, [74, Section 2.6.4] requires $k > d + \max(\alpha, 0)$ for the Poisson characterization while we only need $k > \alpha$.

Theorem 4.20. *Let $1 \leq p < \infty$ and $\alpha \in \mathbb{R}$.*

(i) Let $k \in \mathbb{Z}$ such that $k > \alpha$. Then for any distribution x on $\mathbb{T}_\theta^d$,

$$\|x\|_{F_p^{\alpha,c}} \approx |\widehat{x}(0)| + \left\| \left(\int_0^1 \varepsilon^{2(k-\alpha)} |\mathcal{J}_\varepsilon^k \widetilde{P}_\varepsilon(x)|^2 \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{2}} \right\|_p.$$

(ii) Let $k \in \mathbb{Z}$ such that $k > \frac{\alpha}{2}$. Then for any distribution x on $\mathbb{T}_\theta^d$,

$$\|x\|_{F_p^{\alpha,c}} \approx |\widehat{x}(0)| + \left\| \left(\int_0^1 \varepsilon^{2(k-\frac{\alpha}{2})} |\mathcal{J}_\varepsilon^k \widetilde{W}_\varepsilon(x)|^2 \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{2}} \right\|_p.$$

The preceding theorem can be formulated directly in terms of the circular Poisson and heat semigroups of $\mathbb{T}_\theta^d$. The proof of the following result is similar to that of Theorem 3.15, and is left to the reader.

Theorem 4.21. Let $1 \leq p < \infty$, $\alpha \in \mathbb{R}$ and $k \in \mathbb{Z}$.

(i) If $k > \alpha$, then for any distribution x on $\mathbb{T}_\theta^d$,

$$\|x\|_{F_p^{\alpha,c}} \approx \max_{|m| < k} |\widehat{x}(m)| + \left\| \left(\int_0^1 (1-r)^{2(k-\alpha)} |\mathcal{J}_r^k \mathbb{P}_r(x_k)|^2 \frac{dr}{1-r} \right)^{\frac{1}{2}} \right\|_p,$$

$$\text{where } x_k = x - \sum_{|m| < k} \widehat{x}(m) U^m.$$

(ii) If $k > \frac{\alpha}{2}$, then for any any distribution x on $\mathbb{T}_\theta^d$,

$$\|x\|_{F_p^{\alpha,c}} \approx \max_{|m|^2 < k} |\widehat{x}(m)| + \left\| \left(\int_0^1 (1-r)^{2(k-\frac{\alpha}{2})} |\mathcal{J}_r^k \mathbb{W}_r(x)|^2 \frac{dr}{1-r} \right)^{\frac{1}{2}} \right\|_p.$$

The proof of Theorem 4.20. Similar to the Besov case, the proof of (ii) is done by choosing $\alpha_1 = 2k > \alpha$. But (i) is much subtler. We will first prove (i) under the stronger assumption that $k > d + \alpha$, the remaining case being postponed. The proof in this case is similar to and a little bit harder than the proof of Theorem 3.13. Let again $\psi(\xi) = (-\operatorname{sgn}(k)2\pi|\xi|)^k e^{-2\pi|\xi|}$. As in that proof, it remains to show that ψ satisfies the second condition of (4.7) for some $\alpha_1 > \alpha$ and $\sigma > \frac{d}{2}$. Since $k > d + \alpha$, we can choose α_1 such that $\alpha < \alpha_1 < k - d$. We claim that $I_{k-\alpha_1} h \widehat{P} \in H_2^{\sigma_1}(\mathbb{R}^d)$ for every $\sigma_1 \in (\frac{d}{2}, k - \alpha_1 + \frac{d}{2})$. Indeed, this is a variant of Lemma 3.5 with $a = k - \alpha_1$ and $\rho = h \widehat{P}$. The difference is that this function ρ is not infinitely differentiable at the origin. However, the claim is true if σ_1 is an integer. Then by complex interpolation as in the proof of that lemma, we deduce the claim in the general case. Now choose σ such that $\frac{d}{2} < \sigma < \frac{1}{2}(\sigma_1 - \frac{d}{2})$ and set $\eta = \sigma_1 - 2\sigma$. Then $\eta > \frac{d}{2}$, and by the Cauchy-Schwarz inequality, we have

$$\begin{aligned} \int_{\mathbb{R}^d} (1+|s|^2)^\sigma |\mathcal{F}^{-1}(I_{k-\alpha_1} h \widehat{P})(s)| ds &\leq \left(\int_{\mathbb{R}^d} (1+|s|^2)^{2\sigma+\eta} |\mathcal{F}^{-1}(I_{k-\alpha_1} h \widehat{P})(s)|^2 ds \right)^{\frac{1}{2}} \\ &\lesssim \|\mathcal{F}^{-1}(I_{k-\alpha_1} h \widehat{P})\|_{H_2^{\sigma_1}(\mathbb{R}^d)}. \end{aligned}$$

Therefore, the second condition of (4.7) is verified. This shows part (i) in the case $k > d + \alpha$. $\square$

To deal with the remaining case $k > \alpha$, we need the following:

Lemma 4.22. *Let $1 \leq p < \infty$ and $k, \ell \in \mathbb{Z}$ such that $\ell > k > \alpha$. Then for any distribution x on $\mathbb{T}_\theta^d$ with $\hat{x}(0) = 0$,*

$$\left\| \left(\int_0^1 \varepsilon^{2(k-\alpha)} |\mathcal{J}_\varepsilon^k \tilde{P}_\varepsilon(x)|^2 \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{2}} \right\|_p \approx \left\| \left(\int_0^1 \varepsilon^{2(\ell-\alpha)} |\mathcal{J}_\varepsilon^\ell \tilde{P}_\varepsilon(x)|^2 \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{2}} \right\|_p.$$

Proof. By induction, it suffices to consider the case $\ell = k + 1$. We first show the lower estimate:

$$\left\| \left(\int_0^1 \varepsilon^{2(k-\alpha)} |\mathcal{J}_\varepsilon^k \tilde{P}_\varepsilon(x)|^2 \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{2}} \right\|_p \lesssim \left\| \left(\int_0^1 \varepsilon^{2(k+1-\alpha)} |\mathcal{J}_\varepsilon^{k+1} \tilde{P}_\varepsilon(x)|^2 \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{2}} \right\|_p.$$

To that end, we use

$$\mathcal{J}_\varepsilon^k \tilde{P}_\varepsilon(x) = -\operatorname{sgn}(k) \int_\varepsilon^\infty \mathcal{J}_\delta^{k+1} \tilde{P}_\delta(x) d\delta.$$

Choose $\beta \in (0, k - \alpha)$. By the Cauchy-Schwarz inequality via the operator convexity of the function $t \mapsto t^2$, we obtain

$$|\mathcal{J}_\varepsilon^k \tilde{P}_\varepsilon(x)|^2 \leq \frac{\varepsilon^{-2\beta}}{2\beta} \int_\varepsilon^\infty \delta^{2(1+\beta)} |\mathcal{J}_\delta^{k+1} \tilde{P}_\delta(x)|^2 \frac{d\delta}{\delta}.$$

It then follows that

$$\begin{aligned} \int_0^1 \varepsilon^{2(k-\alpha)} |\mathcal{J}_\varepsilon^k \tilde{P}_\varepsilon(x)|^2 \frac{d\varepsilon}{\varepsilon} &\leq \frac{1}{2\beta} \int_0^\infty \delta^{2(1+\beta)} |\mathcal{J}_\delta^{k+1} \tilde{P}_\delta(x)|^2 \frac{d\delta}{\delta} \int_0^\delta \varepsilon^{2(k-\alpha-\beta)} \frac{d\varepsilon}{\varepsilon} \\ &= \frac{1}{4\beta(k-\alpha-\beta)} \int_0^\infty \delta^{2(k+1-\alpha)} |\mathcal{J}_\delta^{k+1} \tilde{P}_\delta(x)|^2 \frac{d\delta}{\delta}. \end{aligned}$$

Therefore,

$$\begin{aligned} \left\| \left(\int_0^1 \varepsilon^{2(k-\alpha)} |\mathcal{J}_\varepsilon^k \tilde{P}_\varepsilon(x)|^2 \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{2}} \right\|_p &\lesssim \left\| \left(\int_0^\infty \delta^{2(k+1-\alpha)} |\mathcal{J}_\delta^{k+1} \tilde{P}_\delta(x)|^2 \frac{d\delta}{\delta} \right)^{\frac{1}{2}} \right\|_p \\ &\lesssim \left\| \left(\int_0^1 \delta^{2(k+1-\alpha)} |\mathcal{J}_\delta^{k+1} \tilde{P}_\delta(x)|^2 \frac{d\delta}{\delta} \right)^{\frac{1}{2}} \right\|_p, \end{aligned}$$

as desired.

The upper estimate is harder. This time, writing $\tilde{P}_{\varepsilon_1+\varepsilon_2} = \tilde{P}_{\varepsilon_1} * \tilde{P}_{\varepsilon_2}$, we have

$$\begin{aligned} \left(\delta^{k+1} \mathcal{J}_\delta^{k+1} \tilde{P}_\delta \right) \Big|_{\delta=2\varepsilon} &= \operatorname{sgn}(k) 2^{k+1} \varepsilon^{k+1} \frac{\partial}{\partial \varepsilon} \tilde{P}_\varepsilon * \mathcal{J}_\varepsilon^k \tilde{P}_\varepsilon \\ &= \operatorname{sgn}(k) 2^{k+1} \varepsilon^k \tilde{\phi}_\varepsilon * \mathcal{J}_\varepsilon^k \tilde{P}_\varepsilon, \end{aligned}$$

where $\phi(\xi) = -2\pi|\xi| e^{-2\pi|\xi|}$. Thus

$$\begin{aligned} \left\| \left(\int_0^1 \varepsilon^{2(k+1-\alpha)} |\mathcal{J}_\varepsilon^{k+1} \tilde{P}_\varepsilon(x)|^2 \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{2}} \right\|_p &= \left\| \left(\int_0^{\frac{1}{2}} \left| \left(\delta^{k+1-\alpha} \mathcal{J}_\delta^{k+1} \tilde{P}_\delta \right) \Big|_{\delta=2\varepsilon} \right|^2 \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{2}} \right\|_p \\ &= 2^{k+1-\alpha} \left\| \left(\int_0^{\frac{1}{2}} \varepsilon^{2(k-\alpha)} |\tilde{\phi}_\varepsilon * \mathcal{J}_\varepsilon^k \tilde{P}_\varepsilon(x)|^2 \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{2}} \right\|_p \\ &\leq 2^{k+1-\alpha} \left\| \left(\int_0^1 \varepsilon^{2(k-\alpha)} |\tilde{\phi}_\varepsilon * \mathcal{J}_\varepsilon^k \tilde{P}_\varepsilon(x)|^2 \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{2}} \right\|_p. \end{aligned}$$

Now our task is to remove ϕ_ε from the integrand on the right-hand side in the spirit of Theorem 4.1. To that end, we will use a multiplier theorem analogous to Lemma 4.7. Let $H = L_2((0, 1), \frac{d\varepsilon}{\varepsilon})$ and define the H -valued kernel k on $\mathbb{R}^d$ by $k(s) = (\phi_\varepsilon(s))_{0 < \varepsilon < 1}$. It is a well-known elementary fact that this is a Calderón-Zygmund kernel, namely,

- $\|\widehat{k}\|_{L_\infty(\mathbb{R}^d; H)} < \infty$;
- $\sup_{t \in \mathbb{R}^d} \int_{|s| > 2|t|} \|k(s-t) - k(s)\|_H ds < \infty$.

Thus by Lemma 4.7 (i) (more exactly, following its proof), we obtain that the singular integral operator associated to k is bounded on $L_p(\mathbb{T}_\theta^d; H^c)$ for any $1 < p < \infty$; consequently,

$$\left\| \left(\int_0^1 \varepsilon^{2(k-\alpha)} |\widetilde{\phi}_\varepsilon * \mathcal{J}_\varepsilon^k \widetilde{P}_\varepsilon(x)|^2 \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{2}} \right\|_p \lesssim \left\| \left(\int_0^1 \varepsilon^{2(k-\alpha)} |\mathcal{J}_\varepsilon^k \widetilde{P}_\varepsilon(x)|^2 \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{2}} \right\|_p, \quad (4.12)$$

whence

$$\left\| \left(\int_0^1 \varepsilon^{2(k+1-\alpha)} |\mathcal{J}_\varepsilon^{k+1} \widetilde{P}_\varepsilon(x)|^2 \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{2}} \right\|_p \lesssim \left\| \left(\int_0^1 \varepsilon^{2(k-\alpha)} |\mathcal{J}_\varepsilon^k \widetilde{P}_\varepsilon(x)|^2 \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{2}} \right\|_p.$$

Thus the lemma is proved for $1 < p < \infty$.

The case $p = 1$ necessitates a separate argument like Lemma 4.7. We will require a more characterization of $\mathcal{H}_1^c(\mathbb{T}_\theta^d)$ which is a complement to Lemma 1.10. It is the following equivalence proved in [81]:

Let $\beta > 0$. Then for a distribution x on $\mathbb{T}_\theta^d$ with $\widehat{x}(0) = 0$, we have

$$\|x\|_{\mathcal{H}_1^c} \approx \left\| \left(\int_0^1 |(I^\beta P)_\varepsilon * x|^2 \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{2}} \right\|_1. \quad (4.13)$$

Armed with this characterization, we can easily complete the proof of the lemma. Indeed,

$$(\widetilde{I^{k-\alpha}P})_\varepsilon * (I^\alpha x) = (-\operatorname{sgn}(k)2\pi)^{-k} \varepsilon^{k-\alpha} \mathcal{J}_\varepsilon^k \widetilde{P}_\varepsilon(x).$$

Thus by (4.13) with $\beta = k - \alpha$,

$$\left\| \left(\int_0^1 \varepsilon^{2(k-\alpha)} |\mathcal{J}_\varepsilon^k \widetilde{P}_\varepsilon(x)|^2 \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{2}} \right\|_1 \approx \|I^\alpha x\|_{\mathcal{H}_1^c}.$$

It then remains to apply Lemma 4.7 (ii) to $I^\alpha x$ to conclude that (4.12) holds for $p = 1$ too, so the proof of the lemma is complete. $\square$

End of the proof of Theorem 4.20. The preceding lemma shows that the norm in the right-hand side of the equivalence in part (i) is independent of k with $k > \alpha$. As (i) has been already proved to be true for $k > d + \alpha$, we deduce the assertion in full generality. $\square$

We end this section with a Littlewood-Paley type characterizations of Sobolev spaces. The following is an immediate consequence of Corollary 4.12 and the characterizations proved previously in this chapter.

Proposition 4.23. *Let ψ satisfy (4.7), $k > \alpha$ and $1 < p < \infty$. Then for any distribution on $\mathbb{T}_\theta^d$,*

$$\|x\|_{H_p^\alpha} \approx |\widehat{x}(0)| + \begin{cases} \inf \left\{ \left\| \left(\sum_{k \geq 0} (2^{2k\alpha} |\widetilde{\psi}_k * y|^2)^{\frac{1}{2}} \right) \right\|_p + \left\| \left(\sum_{k \geq 0} (2^{2k\alpha} |(\widetilde{\psi}_k * z)^*|^2)^{\frac{1}{2}} \right) \right\|_p \right\} & \text{if } 1 < p < 2, \\ \max \left\{ \left\| \left(\sum_{k \geq 0} (2^{2k\alpha} |\widetilde{\psi}_k * x|^2)^{\frac{1}{2}} \right) \right\|_p, \left\| \left(\sum_{k \geq 0} (2^{2k\alpha} |(\widetilde{\psi}_k * x)^*|^2)^{\frac{1}{2}} \right) \right\|_p \right\} & \text{if } 2 \leq p < \infty; \end{cases}$$

and

$$\begin{aligned} \|x\|_{H_p^\alpha} &\approx |\widehat{x}(0)| + \\ &\begin{cases} \inf \left\{ \left\| \left(\int_0^1 \varepsilon^{2(k-\alpha)} |\mathcal{J}_\varepsilon^k \tilde{P}_\varepsilon(y)|^2 \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{2}} \right\|_p + \left\| \left(\int_0^1 \varepsilon^{2(k-\alpha)} |(\mathcal{J}_\varepsilon^k \tilde{P}_\varepsilon(z))^*|^2 \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{2}} \right\|_p \right\} & \text{if } 1 < p < 2, \\ \max \left\{ \left\| \left(\int_0^1 \varepsilon^{2(k-\alpha)} |\mathcal{J}_\varepsilon^k \tilde{P}_\varepsilon(x)|^2 \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{2}} \right\|_p, \left\| \left(\int_0^1 \varepsilon^{2(k-\alpha)} |(\mathcal{J}_\varepsilon^k \tilde{P}_\varepsilon(x))^*|^2 \frac{d\varepsilon}{\varepsilon} \right)^{\frac{1}{2}} \right\|_p \right\} & \text{if } 2 \leq p < \infty. \end{cases} \end{aligned}$$

The above infima are taken above all decompositions $x = y + z$.

4.5 Operator-valued Triebel-Lizorkin spaces

Unlike Sobolev and Besov spaces, the study of vector-valued Triebel-Lizorkin spaces in the classical setting does not allow one to handle their counterparts in quantum tori by means of transference. Given a Banach space X , a straightforward way of defining the X -valued Triebel-Lizorkin spaces on $\mathbb{T}^d$ is as follows: for $1 \leq p < \infty$, $1 \leq q \leq \infty$ and $\alpha \in \mathbb{R}$, an X -valued distribution f on $\mathbb{T}^d$ belongs to $F_{p,q}^\alpha(\mathbb{T}^d; X)$ if

$$\|f\|_{F_{p,q}^\alpha} = \|\widehat{f}(0)\|_X + \left\| \left(\sum_{k \geq 0} 2^{qk\alpha} \|\tilde{\varphi}_k * f\|_X^q \right)^{\frac{1}{q}} \right\|_{L_p(\mathbb{T}^d)} < \infty.$$

A majority of the classical results on Triebel-Lizorkin spaces can be proved to be true in this vector-valued setting with essentially the same methods. Contrary to the Sobolev or Besov case, the space $F_{p,2}^\alpha(\mathbb{T}^d; L_p(\mathbb{T}_\theta^d))$ is very different from the previously studied space $F_{p,2}^{\alpha,c}(\mathbb{T}_\theta^d)$. This explains why the transference method is not efficient here.

However, there exists another way of defining $F_{p,2}^\alpha(\mathbb{T}^d; X)$. Let (r_k) be a Rademacher sequence, that is, an independent sequence of random variables on a probability space (Ω, P) , taking only two values ± 1 with equal probability. We define $F_{p,\text{rad}}^\alpha(\mathbb{T}^d; X)$ to be the space of all X -valued distributions f on $\mathbb{T}^d$ such that

$$\|f\|_{F_{p,\text{rad}}^\alpha} = \|\widehat{f}(0)\|_X + \left\| \sum_{k \geq 0} r_k 2^{k\alpha} \tilde{\varphi}_k * f \right\|_{L_p(\Omega \times \mathbb{T}^d; X)} < \infty.$$

It seems that these spaces $F_{p,\text{rad}}^\alpha(\mathbb{T}^d; X)$ have never been studied so far in literature. They might be worth to be investigated. If X is a Banach lattice of finite concavity, then by the Khintchine inequality,

$$\|f\|_{F_{p,\text{rad}}^\alpha} \approx \|\widehat{f}(0)\|_X + \left\| \left(\sum_{k \geq 0} 2^{2k\alpha} |\tilde{\varphi}_k * f|^2 \right)^{\frac{1}{2}} \right\|_{L_p(\mathbb{T}^d; X)}.$$

This norm resembles, in form, more the previous one $\|f\|_{F_{p,2}^\alpha}$. Moreover, in this case, one can also define a similar space by replacing the internal ℓ_2 -norm by any ℓ_q -norm.

But what we are interested in here is the noncommutative case, where X is a noncommutative L_p -space, say, $X = L_p(\mathbb{T}_\theta^d)$. Then by the noncommutative Khintchine inequality [40], we can show that for $2 \leq p < \infty$ (assuming $\widehat{f}(0) = 0$),

$$\|f\|_{F_{p,\text{rad}}^\alpha} \approx \max \left\{ \left\| \left(\sum_{k \geq 0} 2^{2k\alpha} |\tilde{\varphi}_k * f|^2 \right)^{\frac{1}{2}} \right\|_p, \left\| \left(\sum_{k \geq 0} 2^{2k\alpha} |(\tilde{\varphi}_k * f)^*|^2 \right)^{\frac{1}{2}} \right\|_p \right\}.$$

Here $\|\cdot\|_p$ is the norm of $L_p(\mathbb{T}^d; L_p(\mathbb{T}_\theta^d))$. Thus the right hand-side is closely related to the norm of $F_p^\alpha(\mathbb{T}_\theta^d)$ defined in section 4.2. In fact, if $x \mapsto \tilde{x}$ denotes the transference map introduced in Corollary 1.2, then for $1 < p < \infty$, we have

$$\|x\|_{F_p^\alpha(\mathbb{T}_\theta^d)} \approx \|\tilde{x}\|_{F_{p,\text{rad}}^\alpha(\mathbb{T}^d; L_p(\mathbb{T}_\theta^d))}.$$

This shows that if one wishes to treat Triebel-Lizorkin spaces on $\mathbb{T}_\theta^d$ via transference, one should first investigate the spaces $F_{p,\text{rad}}^\alpha(\mathbb{T}^d; L_p(\mathbb{T}_\theta^d))$. The latter ones are as hard to deal with as $F_p^\alpha(\mathbb{T}_\theta^d)$.

We would like to point out, at this stage, that the method we have developed in this chapter applies as well to $F_{p,\text{rad}}^\alpha(\mathbb{T}^d; L_p(\mathbb{T}_\theta^d))$. In view of operator-valued Hardy spaces, we will call $F_{p,\text{rad}}^\alpha(\mathbb{T}^d; L_p(\mathbb{T}_\theta^d))$ an operator-valued Triebel-Lizorkin space on $\mathbb{T}^d$. We can define similarly its column and row counterparts. We will give below an outline of these operator-valued Triebel-Lizorkin spaces in the light of the development made in the previous sections. A systematic study will be given elsewhere. In the remainder of this section, $\mathcal{M}$ will denote a finite von Neumann algebra $\mathcal{M}$ with a faithful normal radial state τ and $\mathcal{N} = L_\infty(\mathbb{T}^d) \overline{\otimes} \mathcal{M}$.

Definition 4.24. Let $1 \leq p < \infty$ and $\alpha \in \mathbb{R}$. The column operator-valued Triebel-Lizorkin space $F_p^{\alpha,c}(\mathbb{T}^d, \mathcal{M})$ is defined to be

$$F_p^{\alpha,c}(\mathbb{T}^d, \mathcal{M}) = \{f \in \mathcal{S}'(\mathbb{T}^d; L_1(\mathcal{M})) : \|f\|_{F_p^{\alpha,c}} < \infty\},$$

where

$$\|f\|_{F_p^{\alpha,c}} = \|\widehat{f}(0)\|_{L_p(\mathcal{M})} + \left\| \left(\sum_{k \geq 0} 2^{2k\alpha} |\widetilde{\varphi}_k * f|^2 \right)^{\frac{1}{2}} \right\|_{L_p(\mathcal{N})}.$$

The main ingredient for the study of these spaces is still a multiplier result like Theorem 4.1 that is restated as follows:

Theorem 4.25. Assume that $(\phi_j)_{j \geq 0}$ and $(\rho_j)_{j \geq 0}$ satisfy (4.1) with some $\sigma > \frac{d}{2}$.

(i) Let $1 < p < \infty$. Then for any $f \in \mathcal{S}'(\mathbb{T}^d; L_1(\mathcal{M}))$,

$$\left\| \left(\sum_{j \geq 0} 2^{2j\alpha} |\widetilde{\phi}_j * \widetilde{\rho}_j * f|^2 \right)^{\frac{1}{2}} \right\|_{L_p(\mathcal{N})} \lesssim \sup_{j \geq 0} \|\phi_j(2^j \cdot) \varphi\|_{H_2^\sigma} \left\| \left(\sum_{j \geq 0} 2^{2j\alpha} |\widetilde{\rho}_j * f|^2 \right)^{\frac{1}{2}} \right\|_{L_p(\mathcal{N})}.$$

(ii) If $\rho_j = \widehat{\rho(2^{-j} \cdot)}$ for some Schwartz function ρ with $\text{supp}(\rho) = \{\xi : 2^{-1} \leq |\xi| \leq 2\}$. Then the above inequality holds for $p = 1$ too.

The proof of Theorem 4.1 already gives the above result. Armed with this multiplier theorem, we can check that all results proved in the previous sections admit operator-valued analogues with the same proofs. For instance, the dual space of $F_1^{\alpha,c}(\mathbb{T}^d, \mathcal{M})$ can be described as a space $F_\infty^{-\alpha,c}(\mathbb{T}^d, \mathcal{M})$ analogous to the one defined in Definition 4.13. However, following the H_1 -BMO duality developed in the theory of operator-valued Hardy spaces in [81], we can show the following nicer characterization of the latter space in the style of Carleson measures:

Theorem 4.26. A distribution $f \in \mathcal{S}'(\mathbb{T}^d; L_1(\mathcal{M}))$ with $\widehat{f}(0) = 0$ belongs to $F_\infty^{\alpha,c}(\mathbb{T}^d, \mathcal{M})$ iff

$$\sup_Q \left\| \frac{1}{|Q|} \int_Q \sum_{k \geq \log_2(l(Q))} 2^{2k\alpha} |\widetilde{\varphi}_k * f(s)|^2 ds \right\|_{\mathcal{M}} < \infty,$$

where the supremum runs over all cubes of $\mathbb{T}^d$, and where $l(Q)$ denotes the side length of Q .

The characterizations of Triebel-Lizorkin spaces given in the previous two sections can be transferred to the present setting too. Let us formulate only the analogue of Theorem 4.21.

Theorem 4.27. *Let $1 \leq p < \infty$, $\alpha \in \mathbb{R}$ and $k \in \mathbb{Z}$.*

(i) *If $k > \alpha$, then for any $f \in \mathcal{S}'(\mathbb{T}^d; L_1(\mathcal{M}))$,*

$$\|f\|_{F_p^{\alpha,c}} \approx \max_{|m| < k} \|\widehat{f}(m)\|_{L_p(\mathcal{M})} + \left\| \left(\int_0^1 (1-r)^{2(k-\alpha)} |\mathcal{J}_r^k \mathbb{P}_r(f_k)|^2 \frac{dr}{1-r} \right)^{\frac{1}{2}} \right\|_{L_p(\mathcal{N})},$$

$$\text{where } f_k = f - \sum_{|m| < k} \widehat{f}(m) U^m.$$

(ii) *If $k > \frac{\alpha}{2}$, then for any $f \in \mathcal{S}'(\mathbb{T}^d; L_1(\mathcal{M}))$,*

$$\|f\|_{F_p^{\alpha,c}} \approx \max_{|m|^2 < k} \|\widehat{f}(m)\|_{L_p(\mathcal{M})} + \left\| \left(\int_0^1 (1-r)^{2(k-\frac{\alpha}{2})} |\mathcal{J}_r^k \mathbb{W}_r(f)|^2 \frac{dr}{1-r} \right)^{\frac{1}{2}} \right\|_{L_p(\mathcal{N})}.$$

Chapter 5

Interpolation

Now we study the interpolation of the various spaces introduced in the preceding three chapters. We start with the interpolation of Besov and Sobolev spaces. Like in the classical case, the interpolation of Besov spaces on $\mathbb{T}_\theta^d$ is very simple. However, the situation of (fractional) Sobolev spaces is much more delicate. Recall that the complex interpolation problem of the classical couple $(W_1^k(\mathbb{R}^d), W_\infty^k(\mathbb{R}^d))$ remains always open (see [27, p. 173]). We show in the first section some partial results on the interpolation of $W_p^k(\mathbb{T}_\theta^d)$ and $H_p^\alpha(\mathbb{T}_\theta^d)$. The main result there concerns the Hardy-Sobolev spaces $W_{\mathcal{H}_1}^k(\mathbb{T}_\theta^d)$ and $H_{\mathcal{H}_1}^\alpha(\mathbb{T}_\theta^d)$, that is, when the L_1 -norm is replaced by the nicer $\mathcal{H}_1$ -norm on $\mathbb{T}_\theta^d$. The spaces $W_{\text{BMO}}^k(\mathbb{T}_\theta^d)$ and $H_{\text{BMO}}^\alpha(\mathbb{T}_\theta^d)$ are also considered. The most important problem left unsolved in the first section is to transfer DeVore and Scherer's theorem on the real interpolation of $(W_1^k(\mathbb{R}^d), W_\infty^k(\mathbb{R}^d))$ to the quantum setting. The main result of the second section characterizes the K-functional of the couple $(L_p(\mathbb{T}_\theta^d), W_p^k(\mathbb{T}_\theta^d))$ by the L_p -modulus of smoothness, thereby extending a theorem of Johnen and Scherer to the quantum tori. This result is closely related to the limit theorem of Besov spaces proved in section 3.5. The last short section contains some simple results on the interpolation of Triebel-Lizorkin spaces.

5.1 Interpolation of Besov and Sobolev spaces

This section collects some results on the interpolation of Besov and Sobolev spaces. We start with the Besov spaces.

Proposition 5.1. *Let $0 < \eta < 1$. Assume that $\alpha, \alpha_0, \alpha_1 \in \mathbb{R}$ and $p, p_0, p_1, q, q_0, q_1 \in [1, \infty]$ satisfy the constraints given in the formulas below. We have*

- (i) $(B_{p_0, q_0}^{\alpha_0}(\mathbb{T}_\theta^d), B_{p_1, q_1}^{\alpha_1}(\mathbb{T}_\theta^d))_{\eta, q} = B_{p, q}^\alpha(\mathbb{T}_\theta^d), \quad \alpha_0 \neq \alpha_1, \alpha = (1 - \eta)\alpha_0 + \eta\alpha_1;$
- (ii) $(B_{p, q_0}^\alpha(\mathbb{T}_\theta^d), B_{p, q_1}^\alpha(\mathbb{T}_\theta^d))_{\eta, q} = B_{p, q}^\alpha(\mathbb{T}_\theta^d), \quad \frac{1}{q} = \frac{1 - \eta}{q_0} + \frac{\eta}{q_1};$
- (iii) $(B_{p_0, q_0}^{\alpha_0}(\mathbb{T}_\theta^d), B_{p_1, q_1}^{\alpha_1}(\mathbb{T}_\theta^d))_{\eta, q} = B_{p, q}^\alpha(\mathbb{T}_\theta^d), \quad \alpha = (1 - \eta)\alpha_0 + \eta\alpha_1,$
 $\frac{1}{p} = \frac{1 - \eta}{p_0} + \frac{\eta}{p_1}, \frac{1}{q} = \frac{1 - \eta}{q_0} + \frac{\eta}{q_1}, p = q;$
- (iv) $(B_{p_0, q_0}^{\alpha_0}(\mathbb{T}_\theta^d), B_{p_1, q_1}^{\alpha_1}(\mathbb{T}_\theta^d))_\eta = B_{p, q}^\alpha(\mathbb{T}_\theta^d), \quad \alpha = (1 - \eta)\alpha_0 + \eta\alpha_1, \frac{1}{p} = \frac{1 - \eta}{p_0} + \frac{\eta}{p_1},$
 $\frac{1}{q} = \frac{1 - \eta}{q_0} + \frac{\eta}{q_1}, q < \infty.$

Proof. We will use the embedding of $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ into $\ell_q^\alpha(L_p(\mathbb{T}_\theta^d))$. Recall that given a Banach space X , $\ell_q^\alpha(X)$ denotes the weighted ℓ_q -direct sum of $(\mathbb{C}, X, X, \dots)$, equipped with the norm

$$\|(a, x_0, x_1, \dots)\| = \left(|a|^q + \sum_{k \geq 0} 2^{kq\alpha} \|x_k\|^q\right)^{\frac{1}{q}}.$$

Then $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ isometrically embeds into $\ell_q^\alpha(L_p(\mathbb{T}_\theta^d))$ via the map $\mathcal{I}$ defined by $\mathcal{I}x = (\widehat{x}(0), \widetilde{\varphi}_0 * x, \widetilde{\varphi}_1 * x, \dots)$. On the other hand, it is easy to check that the range of $\mathcal{I}$ is 1-complemented. Indeed, let $\mathcal{P} : \ell_q^\alpha(L_p(\mathbb{T}_\theta^d)) \rightarrow B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ be defined by (with $\widetilde{\varphi}_k = 0$ for $k \leq -1$)

$$\mathcal{P}(a, x_0, x_1, \dots) = a + \sum_{k \geq 0} (\widetilde{\varphi}_{k-1} + \widetilde{\varphi}_k + \widetilde{\varphi}_{k+1}) * x_k.$$

Then by (3.2), $\mathcal{P}\mathcal{I}x = x$ for all $x \in B_{p,q}^\alpha(\mathbb{T}_\theta^d)$. On the other hand, letting $y = \mathcal{P}(a, x_0, x_1, \dots)$, we have

$$\widetilde{\varphi}_j * y = \sum_{k=j-2}^{j+2} \varphi_j * (\widetilde{\varphi}_{k-1} + \widetilde{\varphi}_k + \widetilde{\varphi}_{k+1}) * x_k, \quad j \geq 0.$$

Thus we deduce that $\mathcal{P}$ is bounded with norm at most 15.

Therefore, the interpolation of the Besov spaces is reduced to that of the spaces $\ell_q^\alpha(L_p(\mathbb{T}_\theta^d))$, which is well-known and is treated in [8, Section 5.6]. Let us recall the results needed here. For a Banach space X and an interpolation couple (X_0, X_1) of Banach spaces, we have

- $(\ell_{q_0}^{\alpha_0}(X), \ell_{q_1}^{\alpha_1}(X))_{\eta,q} = \ell_q^\alpha(X), \quad \alpha_0 \neq \alpha_1, \alpha = (1-\eta)\alpha_0 + \eta\alpha_1;$
- $(\ell_{q_0}^{\alpha_0}(X_0), \ell_{q_1}^{\alpha_1}(X_1))_{\eta,q} = \ell_q^\alpha((X_0, X_1)_{\eta,q}), \quad \alpha = (1-\eta)\alpha_0 + \eta\alpha_1, \frac{1}{q} = \frac{1-\eta}{q_0} + \frac{\eta}{q_1};$
- $(\ell_{q_0}^{\alpha_0}(X_0), \ell_{q_1}^{\alpha_1}(X_1))_\eta = \ell_q^\alpha((X_0, X_1)_\eta), \quad \alpha = (1-\eta)\alpha_0 + \eta\alpha_1, \frac{1}{q} = \frac{1-\eta}{q_0} + \frac{\eta}{q_1}, \quad q < \infty.$

It is then clear that the interpolation formulas of the theorem follow from the above ones thanks to the complementation result proved previously. $\square$

Remark 5.2. If $q = \infty$, part (iv) holds for Calderón's second interpolation method, namely,

$$(B_{p_0,\infty}^{\alpha_0}(\mathbb{T}_\theta^d), B_{p_1,\infty}^{\alpha_1}(\mathbb{T}_\theta^d))^\eta = B_{p,\infty}^\alpha(\mathbb{T}_\theta^d), \quad \alpha = (1-\eta)\alpha_0 + \eta\alpha_1, \quad \frac{1}{p} = \frac{1-\eta}{p_0} + \frac{\eta}{p_1}.$$

On the other hand, if one wishes to stay with the first complex interpolation method in the case $q = \infty$, one should replace $B_{p,\infty}^\alpha(\mathbb{T}_\theta^d)$ by $B_{p,c_0}^\alpha(\mathbb{T}_\theta^d)$:

$$(B_{p_0,c_0}^{\alpha_0}(\mathbb{T}_\theta^d), B_{p_1,c_0}^{\alpha_1}(\mathbb{T}_\theta^d))_\eta = B_{p,c_0}^\alpha(\mathbb{T}_\theta^d).$$

Now we consider the potential Sobolev spaces. Since J^α is an isometry between $H_p^\alpha(\mathbb{T}_\theta^d)$ and $L_p(\mathbb{T}_\theta^d)$ for all $1 \leq p \leq \infty$, we get immediately the following

Remark 5.3. Let $0 < \eta < 1$, $\alpha \in \mathbb{R}$, $1 \leq p_0, p_1 \leq \infty$ and $\frac{1}{p} = \frac{1-\eta}{p_0} + \frac{\eta}{p_1}$. Then

$$(H_{p_0}^\alpha(\mathbb{T}_\theta^d), H_{p_1}^\alpha(\mathbb{T}_\theta^d))_\eta = H_p^\alpha(\mathbb{T}_\theta^d) \quad \text{and} \quad (H_{p_0}^\alpha(\mathbb{T}_\theta^d), H_{p_1}^\alpha(\mathbb{T}_\theta^d))_{\eta,p} = H_p^\alpha(\mathbb{T}_\theta^d).$$

The interpolation problem of the couple $(H_{p_0}^{\alpha_0}(\mathbb{T}_\theta^d), H_{p_1}^{\alpha_1}(\mathbb{T}_\theta^d))$ for $\alpha_0 \neq \alpha_1$ is delicate. At the time of this writing, we cannot, unfortunately, solve it completely. To our knowledge, it seems that even in the commutative case, its interpolation spaces by real or complex interpolation method have not been determined in full generality. We will prove some partial results.

Proposition 5.4. *Let $0 < \eta < 1$, $\alpha_0 \neq \alpha_1 \in \mathbb{R}$ and $1 \leq p, q \leq \infty$. Then*

$$(H_p^{\alpha_0}(\mathbb{T}_\theta^d), H_p^{\alpha_1}(\mathbb{T}_\theta^d))_{\eta, q} = B_{p, q}^\alpha(\mathbb{T}_\theta^d), \quad \alpha = (1 - \eta)\alpha_0 + \eta\alpha_1.$$

Proof. The assertion follows from Theorem 3.8, the reiteration theorem and Proposition 5.1 (i). $\square$

To treat the complex interpolation, we introduce the potential Hardy-Sobolev spaces.

Definition 5.5. For $\alpha \in \mathbb{R}$, define

$$H_{\mathcal{H}_1}^\alpha(\mathbb{T}_\theta^d) = \{x \in \mathcal{S}'(\mathbb{T}_\theta^d) : J^\alpha x \in \mathcal{H}_1(\mathbb{T}_\theta^d)\} \quad \text{with} \quad \|x\|_{H_{\mathcal{H}_1}^\alpha} = \|J^\alpha x\|_{\mathcal{H}_1}.$$

We define $H_{\text{BMO}}^\alpha(\mathbb{T}_\theta^d)$ similarly.

Theorem 5.6. *Let $\alpha_0, \alpha_1 \in \mathbb{R}$ and $1 < p < \infty$. Then*

$$(H_{\text{BMO}}^{\alpha_0}(\mathbb{T}_\theta^d), H_{\mathcal{H}_1}^{\alpha_1}(\mathbb{T}_\theta^d))_{\frac{1}{p}} = H_p^\alpha(\mathbb{T}_\theta^d), \quad \alpha = (1 - \frac{1}{p})\alpha_0 + \frac{\alpha_1}{p}.$$

We require the following result which extends Lemma 1.7(ii):

Lemma 5.7. *Let ϕ be a Mihlin multiplier in the sense of Definition 1.5. Then ϕ is a Fourier multiplier on both $\mathcal{H}_1(\mathbb{T}_\theta^d)$ and $\text{BMO}(\mathbb{T}_\theta^d)$ with norms majorized by $c_d \|\phi\|_M$.*

Proof. This is an immediate consequence of Lemma 4.7 (the sequence (ϕ_j) there becomes now the single function ϕ). Indeed, by that Lemma, ϕ is a bounded Fourier multiplier on $\mathcal{H}_1(\mathbb{T}_\theta^d)$, so by duality, it is bounded on $\text{BMO}(\mathbb{T}_\theta^d)$ too. $\square$

We will use Bessel potentials of complex order. For $z \in \mathbb{C}$, define $J_z(\xi) = (1 + |\xi|^2)^{\frac{z}{2}}$ and J^z to be the associated Fourier multiplier.

Lemma 5.8. *Let $t \in \mathbb{R}$. Then J^{it} is bounded on both $\mathcal{H}_1(\mathbb{T}_\theta^d)$ and $\text{BMO}(\mathbb{T}_\theta^d)$ with norms majorized by $c_d(1 + |t|)^d$.*

Proof. One easily checks that J_{it} is a Mihlin multiplier and $\|J_{it}\|_M \leq c_d(1 + |t|)^d$. Thus, the assertion follows from the previous lemma. $\square$

Proof of Theorem 5.6. Let $x \in H_p^\alpha(\mathbb{T}_\theta^d)$ with norm less than 1, that is, $J^\alpha x \in L_p(\mathbb{T}_\theta^d)$ and $\|J^\alpha x\|_p < 1$. By Lemma 1.9, and the definition of complex interpolation, there exists a continuous function f from the strip $S = \{z \in \mathbb{C} : 0 \leq \text{Re}(z) \leq 1\}$ to $\mathcal{H}_1(\mathbb{T}_\theta^d)$, analytic in the interior, such that $f(\frac{1}{p}) = J^\alpha x$,

$$\sup_{t \in \mathbb{R}} \|f(it)\|_{\text{BMO}} \leq c \quad \text{and} \quad \sup_{t \in \mathbb{R}} \|f(1 + it)\|_{\mathcal{H}_1} \leq c.$$

Define (with $\eta = \frac{1}{p}$)

$$F(z) = e^{(z-\eta)^2} J^{-(1-z)\alpha_0 - z\alpha_1} f(z), \quad z \in S.$$

Then for any $t \in \mathbb{R}$, by the preceding lemma,

$$\|F(it)\|_{H_{\text{BMO}}^{\alpha_0}} = e^{-t^2 + \eta^2} \|J^{\text{it}(\alpha_0 - \alpha_1)} f(it)\|_{\text{BMO}} \leq c'.$$

A similar estimate holds for the other extreme point $H_{\mathcal{H}_1}^{\alpha_1}(\mathbb{T}_\theta^d)$. Therefore,

$$x = F(\eta) \in (H_{\text{BMO}}^{\alpha_0}(\mathbb{T}_\theta^d), H_{\mathcal{H}_1}^{\alpha_1}(\mathbb{T}_\theta^d))_\eta \text{ with norm } \leq c'.$$

We have thus proved

$$H_p^\alpha(\mathbb{T}_\theta^d) \subset (H_{\text{BMO}}^{\alpha_0}(\mathbb{T}_\theta^d), H_{\mathcal{H}_1}^{\alpha_1}(\mathbb{T}_\theta^d))_\eta.$$

Since the dual space of $\mathcal{H}_1(\mathbb{T}_\theta^d)$ is $\text{BMO}(\mathbb{T}_\theta^d)$, we have

$$H_{\mathcal{H}_1}^{\alpha_1}(\mathbb{T}_\theta^d)^* = H_{\text{BMO}}^{-\alpha_1}(\mathbb{T}_\theta^d).$$

Thus dualizing the above inclusion (for appropriate α_i and p), we get

$$(H_{\text{BMO}}^{\alpha_0}(\mathbb{T}_\theta^d), H_{\text{BMO}}^{-\alpha_1}(\mathbb{T}_\theta^d)^*)^\eta \subset H_p^\alpha(\mathbb{T}_\theta^d),$$

where $(\cdot \cdot)^\eta$ denotes Calderón's second complex interpolation method. However, by [7]

$$(H_{\text{BMO}}^{\alpha_0}(\mathbb{T}_\theta^d), H_{\text{BMO}}^{-\alpha_1}(\mathbb{T}_\theta^d)^*)_\eta \subset (H_{\text{BMO}}^{\alpha_0}(\mathbb{T}_\theta^d), H_{\text{BMO}}^{-\alpha_1}(\mathbb{T}_\theta^d)^*)^\eta \text{ isometrically.}$$

Since

$$H_{\mathcal{H}_1}^{\alpha_1}(\mathbb{T}_\theta^d) \subset H_{\text{BMO}}^{-\alpha_1}(\mathbb{T}_\theta^d)^* \text{ isometrically,}$$

we finally deduce

$$(H_{\text{BMO}}^{\alpha_0}(\mathbb{T}_\theta^d), H_{\mathcal{H}_1}^{\alpha_1}(\mathbb{T}_\theta^d))_\eta \subset H_p^\alpha(\mathbb{T}_\theta^d),$$

which concludes the proof of the theorem. $\square$

Corollary 5.9. *Let $0 < \eta < 1$, $\alpha_0, \alpha_1 \in \mathbb{R}$ and $1 < p_0, p_1 < \infty$. Then*

$$(H_{p_0}^{\alpha_0}(\mathbb{T}_\theta^d), H_{p_1}^{\alpha_1}(\mathbb{T}_\theta^d))_\eta = H_p^\alpha(\mathbb{T}_\theta^d), \quad \alpha = (1 - \eta)\alpha_0 + \eta\alpha_1, \quad \frac{1}{p} = \frac{1 - \eta}{p_0} + \frac{\eta}{p_1}.$$

Proof. The preceding proof works equally for this corollary. Alternately, in the case $p_0 \neq p_1$, the corollary immediately follows from the previous theorem by reiteration. Indeed, if $p_0 \neq p_1$, then for any $\alpha_0, \alpha_1 \in \mathbb{R}$ there exist $\beta_0, \beta_1 \in \mathbb{R}$ such that

$$(1 - \frac{1}{p_0})\beta_0 + \frac{1}{p_0}\beta_1 = \alpha_0 \text{ and } (1 - \frac{1}{p_1})\beta_0 + \frac{1}{p_1}\beta_1 = \alpha_1.$$

Thus the previous theorem implies

$$(H_{\text{BMO}}^{\beta_0}(\mathbb{T}_\theta^d), H_{\mathcal{H}_1}^{\beta_1}(\mathbb{T}_\theta^d))_{\frac{1}{p_j}} = H_{p_j}^{\alpha_j}(\mathbb{T}_\theta^d), \quad j = 0, 1.$$

The corollary then follows by the reiteration theorem. $\square$

It is likely that the above corollary still holds for all $1 \leq p_0, p_1 \leq \infty$:

Conjecture 5.10. Let $\alpha_0, \alpha_1 \in \mathbb{R}$ and $1 < p < \infty$. Then

$$(H_\infty^{\alpha_0}(\mathbb{T}_\theta^d), H_1^{\alpha_1}(\mathbb{T}_\theta^d))_{\frac{1}{p}} = H_p^\alpha(\mathbb{T}_\theta^d), \quad \alpha = (1 - \frac{1}{p})\alpha_0 + \frac{\alpha_1}{p}.$$

By duality and Wolff's reiteration theorem [80], the conjecture is reduced to showing that for any $0 < \eta < 1$ and $1 < p_0 < \infty$,

$$(H_{p_0}^{\alpha_0}(\mathbb{T}_\theta^d), H_1^{\alpha_1}(\mathbb{T}_\theta^d))_\eta = H_p^\alpha(\mathbb{T}_\theta^d), \quad \alpha = (1 - \eta)\alpha_0 + \eta\alpha_1, \quad \frac{1}{p} = \frac{1 - \eta}{p_0} + \frac{\eta}{1}.$$

Since $H_{\mathcal{H}_1}^{\alpha_1}(\mathbb{T}_\theta^d) \subset H_1^{\alpha_1}(\mathbb{T}_\theta^d)$, Theorem 5.6 implies

$$H_p^\alpha(\mathbb{T}_\theta^d) \subset (H_{p_0}^{\alpha_0}(\mathbb{T}_\theta^d), H_1^{\alpha_1}(\mathbb{T}_\theta^d))_\eta.$$

So the conjecture is equivalent to the validity of the converse inclusion.

Remark 5.11. The proof of Theorem 5.6 shows that for $\alpha_0, \alpha_1 \in \mathbb{R}$ and $0 < \eta < 1$,

$$(H_{\mathcal{H}_1}^{\alpha_0}(\mathbb{T}_\theta^d), H_{\mathcal{H}_1}^{\alpha_1}(\mathbb{T}_\theta^d))_\eta = H_{\mathcal{H}_1}^\alpha(\mathbb{T}_\theta^d), \quad \alpha = (1 - \eta)\alpha_0 + \eta\alpha_1.$$

We do not know if this equality remains true for the couple $(H_1^{\alpha_0}(\mathbb{T}_\theta^d), H_1^{\alpha_1}(\mathbb{T}_\theta^d))$.

We conclude this section with a discussion on the interpolation of $(W_{p_0}^k(\mathbb{T}_\theta^d), W_{p_1}^k(\mathbb{T}_\theta^d))$. Here, the most interesting case is, of course, that where $p_0 = \infty$ and $p_1 = 1$. Recall that in the commutative case, the K-functional of $(W_\infty^k(\mathbb{T}^d), W_1^k(\mathbb{T}^d))$ is determined by DeVore and Scherer [21]; however, determining the complex interpolation spaces of this couple is a longstanding open problem.

Note that if $1 < p_0, p_1 < \infty$, $(W_{p_0}^k(\mathbb{T}_\theta^d), W_{p_1}^k(\mathbb{T}_\theta^d))$ reduces to $(H_{p_0}^k(\mathbb{T}_\theta^d), H_{p_1}^k(\mathbb{T}_\theta^d))$ by virtue of Theorem 2.9. So in this case, the interpolation problem is solved by the preceding results on potential Sobolev spaces. This reduction is, unfortunately, impossible when one of p_0 and p_1 is equal to 1 or ∞ . However, in the spirit of potential Hardy Sobolev spaces, it remains valid if we work with the Hardy Sobolev spaces $W_{\text{BMO}}^k(\mathbb{T}_\theta^d)$ and $W_{\mathcal{H}_1}^k(\mathbb{T}_\theta^d)$ instead of $W_\infty^k(\mathbb{T}_\theta^d)$ and $W_1^k(\mathbb{T}_\theta^d)$, respectively. Here, the Hardy Sobolev spaces are defined as they should be.

Using Lemma 5.7, we see that the proof of Theorem 2.9 remains valid for the Hardy Sobolev spaces too. Thus we have the following:

Lemma 5.12. For any $k \in \mathbb{N}$, $W_{\text{BMO}}^k(\mathbb{T}_\theta^d) = H_{\text{BMO}}^k(\mathbb{T}_\theta^d)$ and $W_{\mathcal{H}_1}^k(\mathbb{T}_\theta^d) = H_{\mathcal{H}_1}^k(\mathbb{T}_\theta^d)$.

Theorem 5.13. Let $k \in \mathbb{N}$ and $1 < p < \infty$. Then for $X = W_{\mathcal{H}_1}^k(\mathbb{T}_\theta^d)$ or $X = W_1^k(\mathbb{T}_\theta^d)$,

$$(W_{\text{BMO}}^k(\mathbb{T}_\theta^d), X)_{\frac{1}{p}} = W_p^k(\mathbb{T}_\theta^d) = (W_{\text{BMO}}^k(\mathbb{T}_\theta^d), X)_{\frac{1}{p}, p}.$$

Consequently, for any $0 < \eta < 1$ and $1 < p_0 < \infty$,

$$(W_{p_0}^k(\mathbb{T}_\theta^d), W_1^k(\mathbb{T}_\theta^d))_\eta = W_p^k(\mathbb{T}_\theta^d) = (W_{p_0}^k(\mathbb{T}_\theta^d), W_1^k(\mathbb{T}_\theta^d))_{\eta, p}, \quad \frac{1}{p} = \frac{1 - \eta}{p_0} + \frac{\eta}{1}.$$

Proof. The first part for $X = W_{\mathcal{H}_1}^k(\mathbb{T}_\theta^d)$ follows immediately from Remark 5.3, Theorem 5.6 and Lemma 5.12. Then by the reiteration theorem, for any $1 < p < \infty$ and $0 < \eta < 1$, we get

$$(W_{\text{BMO}}^k(\mathbb{T}_\theta^d), W_p^k(\mathbb{T}_\theta^d))_\eta = W_q^k(\mathbb{T}_\theta^d) \quad \text{and} \quad (W_p^k(\mathbb{T}_\theta^d), W_{\mathcal{H}_1}^k(\mathbb{T}_\theta^d))_\eta = W_r^k(\mathbb{T}_\theta^d),$$

where $\frac{1}{q} = \frac{1 - \eta}{\infty} + \frac{\eta}{p}$ and $\frac{1}{r} = \frac{1 - \eta}{p} + \frac{\eta}{1}$. On the other hand, by the continuous inclusion $\mathcal{H}_1(\mathbb{T}_\theta^d) \subset L_1(\mathbb{T}_\theta^d)$, we have

$$W_r^k(\mathbb{T}_\theta^d) = (W_p^k(\mathbb{T}_\theta^d), W_{\mathcal{H}_1}^k(\mathbb{T}_\theta^d))_\eta \subset (W_p^k(\mathbb{T}_\theta^d), W_1^k(\mathbb{T}_\theta^d))_\eta \subset W_r^k(\mathbb{T}_\theta^d),$$

the last inclusion above being trivial. Thus

$$(W_p^k(\mathbb{T}_\theta^d), W_1^k(\mathbb{T}_\theta^d))_\eta = W_r^k(\mathbb{T}_\theta^d).$$

Therefore, by Wolff's reiteration theorem [80], we deduce the first part for $X = W_1^k(\mathbb{T}_\theta^d)$. The second part follows from the first by the reiteration theorem. $\square$

Remark 5.14. The second part of the previous theorem had been proved by Marius Junge by a different method; he reduced it to the corresponding problem on $\mathcal{H}_1$ too.

The main problem left open at this stage is the following:

Problem 5.15. Does the second part of the previous theorem hold for $p_0 = \infty$?

5.2 The K-functional of (L_p, W_p^k)

In this section we characterize the K-functional of the couple $(L_p(\mathbb{T}_\theta^d), W_p^k(\mathbb{T}_\theta^d))$ for any $1 \leq p \leq \infty$ and $k \in \mathbb{N}$. First, recall the definition of the K-functional. For an interpolation couple (X_0, X_1) of Banach spaces, we define

$$K(x, \varepsilon; X_0, X_1) = \inf \{ \|x_0\|_{X_0} + \varepsilon \|x_1\|_{X_1} : x = x_0 + x_1, x_0 \in X_0, x_1 \in X_1 \}$$

for $\varepsilon > 0$ and $x \in X_0 + X_1$. Since $W_p^k(\mathbb{T}_\theta^d) \subset L_p(\mathbb{T}_\theta^d)$ contractively, $K(x, \varepsilon; L_p(\mathbb{T}_\theta^d), W_p^k(\mathbb{T}_\theta^d)) = \|x\|_p$ for $\varepsilon \geq 1$; so only the case $\varepsilon < 1$ is nontrivial. The following result is the quantum analogue of Johnen-Scherer's theorem for Sobolev spaces on $\mathbb{R}^d$ (see [30]; see also [6, Theorem 5.4.12]). Recall that $\omega_p^k(x, \varepsilon)$ denotes the k th order modulus of L_p -smoothness of x introduced in section 3.4.

Theorem 5.16. *Let $1 \leq p \leq \infty$ and $k \in \mathbb{N}$. Then*

$$K(x, \varepsilon^k; L_p(\mathbb{T}_\theta^d), W_p^k(\mathbb{T}_\theta^d)) \approx \varepsilon^k |\widehat{x}(0)| + \omega_p^k(x, \varepsilon), \quad 0 < \varepsilon \leq 1$$

with relevant constants depending only on d and k .

Proof. We will adapt the proof of [6, Theorem 5.4.12]. Denote $K(x, \varepsilon; L_p(\mathbb{T}_\theta^d), W_p^k(\mathbb{T}_\theta^d))$ simply by $K(x, \varepsilon)$. It suffices to consider the elements of $L_p(\mathbb{T}_\theta^d)$ whose Fourier coefficients vanish at $m = 0$. Fix such an element x . Let $x = y + z$ with $y \in L_p(\mathbb{T}_\theta^d)$ and $z \in W_p^k(\mathbb{T}_\theta^d)$ (with vanishing Fourier coefficients at 0). Then by Theorem 2.20,

$$\omega_p^k(x, \varepsilon) \leq \omega_p^k(y, \varepsilon) + \omega_p^k(z, \varepsilon) \lesssim \|y\|_p + \varepsilon^k |z|_{W_p^k},$$

which implies

$$\omega_p^k(x, \varepsilon) \lesssim K(x, \varepsilon^k).$$

The converse inequality is harder. We have to produce an appropriate decomposition of x . To this end, let $\mathbb{I} = [0, 1)^d$ and define the required decomposition by

$$y = (-1)^k \int_{\mathbb{I}} \cdots \int_{\mathbb{I}} \Delta_{\varepsilon u}^k(x) du_1 \cdots du_k \quad \text{and} \quad z = x - y,$$

where $u = u_1 + \cdots + u_k$. Then

$$\|y\|_p \leq \int_{\mathbb{I}} \cdots \int_{\mathbb{I}} \|\Delta_{\varepsilon u}^k(x)\|_p du_1 \cdots du_k \leq \omega_p^k(x, k\sqrt{d}\varepsilon) \lesssim \omega_p^k(x, \varepsilon).$$

To handle z , using the formula

$$\Delta_{\varepsilon u}^k = \sum_{j=0}^k (-1)^{k-j} \binom{k}{j} T_{j\varepsilon u},$$

we rewrite z as

$$z = (-1)^{k+1} \sum_{j=1}^k (-1)^{k-j} \binom{k}{j} \int_{\mathbb{I}} \cdots \int_{\mathbb{I}} T_{j\varepsilon u}(x) du_1 \cdots du_k.$$

All terms on the right-hand side are treated in the same way. Let us consider only the first one by setting

$$z_1 = \int_{\mathbb{I}} \cdots \int_{\mathbb{I}} T_{\varepsilon u}(x) du_1 \cdots du_k.$$

Write each u_i in the canonical basis of $\mathbb{R}^d$:

$$u_i = \sum_{j=1}^d u_{i,j} \mathbf{e}_j.$$

We compute $\partial_1 z_1$ explicitly, as example, in the spirit of (2.1):

$$\partial_1 z_1 = \frac{1}{\varepsilon} \sum_{i=1}^k \int_{\mathbb{I}} \cdots \int_{\mathbb{I}} \frac{\partial}{\partial u_{i,1}} T_{\varepsilon u}(x) du_1 \cdots du_k.$$

Integrating the partial derivative on the right-hand side with respect to $u_{i,1}$ yields:

$$\int_0^1 \frac{\partial}{\partial u_{i,1}} T_{\varepsilon u}(x) du_{i,1} = \Delta_{\varepsilon(\mathbf{e}_1 + u - u_{i,1}\mathbf{e}_1)}(x) = \int_0^1 \Delta_{\varepsilon(\mathbf{e}_1 + u - u_{i,1}\mathbf{e}_1)}(x) du_{i,1},$$

where for the second equality, we have used the fact that $\Delta_{\varepsilon(\mathbf{e}_1 + u - u_{i,1}\mathbf{e}_1)}(x)$ is constant in $u_{i,1}$. Thus

$$\partial_1 z_1 = \frac{1}{\varepsilon} \sum_{i=1}^k \int_{\mathbb{I}} \cdots \int_{\mathbb{I}} \Delta_{\varepsilon(\mathbf{e}_1 + u - u_{i,1}\mathbf{e}_1)}(x) du_1 \cdots du_k.$$

To iterate this formula, we use multi-index notation. For $n \in \mathbb{N}$ let

$$[[k]]^n = \{\underline{i} = (i_1, \dots, i_n) : 1 \leq i_\ell \leq k, \text{ all } i_\ell \text{'s are distinct}\}.$$

Then for any $m_1 \in \mathbb{N}$ with $m_1 \leq k$, we have

$$\partial_1^{m_1} z_1 = \varepsilon^{-m_1} \sum_{\underline{i}^1 \in [[k]]^{m_1}} \int_{\mathbb{I}} \cdots \int_{\mathbb{I}} \Delta_{\varepsilon u_{\underline{i}^1}}^{m_1}(x) du_1 \cdots du_k,$$

where

$$u_{\underline{i}^1} = \mathbf{e}_1 + u - (u_{i_1,1} + \cdots + u_{i_{m_1},1})\mathbf{e}_1.$$

Iterating this procedure, for any $m \in \mathbb{N}_0^d$ with $|m|_1 = k$, we get

$$D^m z_1 = \varepsilon^{-k} \sum_{\underline{i}^d \in [[k]]^{m_d}} \cdots \sum_{\underline{i}^1 \in [[k]]^{m_1}} \int_{\mathbb{I}^k} \Delta_{\varepsilon u_{\underline{i}^d}}^{m_d} \cdots \Delta_{\varepsilon u_{\underline{i}^1}}^{m_1}(x) du_1 \cdots du_k,$$

where the $u_{\underline{i}^j}$'s are defined by induction

$$u_{\underline{i}^j} = \mathbf{e}_j + u_{\underline{i}^{j-1}} - (u_{i_1^j,j} + \cdots + u_{i_{m_j}^j,j})\mathbf{e}_j, \quad j = 2, \dots, d.$$

Thus we are in a position of appealing Lemma 2.22 to conclude that

$$\|D^m z_1\|_p \lesssim \varepsilon^{-k} \omega_p^k(x, \varepsilon),$$

whence

$$|z|_{W_p^k} \lesssim \varepsilon^{-k} \omega_p^k(x, \varepsilon).$$

Therefore, $K(x, \varepsilon^k) \lesssim \omega_p^k(x, \varepsilon)$. $\square$

Remark 5.17. The preceding proof shows a little bit more: for any $x \in W_p^k(\mathbb{T}_\theta^d)$ with $\hat{x}(0) = 0$,

$$\omega_p^k(x, \varepsilon) \approx \{\|y\|_p + \varepsilon^k |z|_{W_p^k} : x = y + z, \hat{y}(0) = \hat{z}(0) = 0\}, \quad 0 < \varepsilon \leq 1.$$

In particular, this implies

$$\|x\|_p \lesssim \omega_p^k(x, \varepsilon),$$

which is the analogue for moduli of L_p -continuity of the inequality in Theorem 2.12 (the Poincaré inequality). On the other hand, together with Lemma 2.22, the above inequality provides an alternate proof of Theorem 2.12.

The preceding theorem, together with Theorem 3.16 and the reiteration theorem, implies the following

Corollary 5.18. *Let $0 < \eta < 1$, $\alpha > 0$, $k, k_0, k_1 \in \mathbb{N}$ and $1 \leq p, q, q_1 \leq \infty$. Then*

- (i) $(L_p(\mathbb{T}_\theta^d), W_p^k(\mathbb{T}_\theta^d))_{\eta, q} = B_{p, q}^{\eta k}(\mathbb{T}_\theta^d)$;
- (ii) $(W_p^k(\mathbb{T}_\theta^d), B_{p, q_1}^\alpha(\mathbb{T}_\theta^d))_{\eta, q} = B_{p, q}^\beta(\mathbb{T}_\theta^d)$, $k \neq \alpha$, $\beta = (1 - \eta)k + \eta\alpha$;
- (iii) $(W_p^{k_0}(\mathbb{T}_\theta^d), W_p^{k_1}(\mathbb{T}_\theta^d))_{\eta, q} = B_{p, q}^\alpha(\mathbb{T}_\theta^d)$, $k_0 \neq k_1$, $\alpha = (1 - \eta)k_0 + \eta k_1$.

We can also consider the complex interpolation of $(L_p(\mathbb{T}_\theta^d), W_p^k(\mathbb{T}_\theta^d))$. If $1 < p < \infty$, this is reduced to that of $(L_p(\mathbb{T}_\theta^d), H_p^k(\mathbb{T}_\theta^d))$; so by the result of the previous section, for any $0 < \eta < 1$,

$$(L_p(\mathbb{T}_\theta^d), W_p^k(\mathbb{T}_\theta^d))_\eta = H_p^{\eta k}(\mathbb{T}_\theta^d).$$

Problem 5.19. Does the above equality hold for $p = 1$? The problem is closely related to that in Remark 5.11.

We conclude this section with a remark on the link between Theorem 3.19 and Theorem 5.16. The former can be easily deduced from the latter, by using the following elementary fact (see [8] p. 40): for any couple (X_0, X_1) of Banach spaces and $x \in X_0 \cap X_1$

$$\begin{aligned} \lim_{\eta \rightarrow 1} (\eta(1 - \eta))^{\frac{1}{q}} \|x\|_{(X_0, X_1)_{\eta, q}} &= q^{-\frac{1}{q}} \|x\|_{X_1}, \\ \lim_{\eta \rightarrow 0} (\eta(1 - \eta))^{\frac{1}{q}} \|x\|_{(X_0, X_1)_{\eta, q}} &= q^{-\frac{1}{q}} \|x\|_{X_0}. \end{aligned}$$

Here the norm of $(X_0, X_1)_{\eta, q}$ is that defined by the K-functional. Then Theorem 3.19 follows from Theorem 5.16. and the first limit above. This is the approach adopted in [38, 47]. It also allows us to determine the other extreme case $\alpha = 0$ in Theorem 3.19, which was done by Maz'ya and Shaposhnikova [41] in the commutative case. Let us record this result here.

Corollary 5.20. *Let $1 \leq p \leq \infty$ and $1 \leq q < \infty$. Then for $x \in B_{p, q}^{\alpha_0}(\mathbb{T}_\theta^d)$ with $\hat{x}(0) = 0$ for some $\alpha_0 > 0$,*

$$\lim_{\alpha \rightarrow 0} \alpha^{\frac{1}{q}} \|x\|_{B_{p, q}^{\alpha, \omega}} \approx q^{-\frac{1}{q}} \|x\|_p.$$

5.3 Interpolation of Triebel-Lizorkin spaces

This short section contains some simple results on the interpolation of Triebel-Lizorkin spaces. They are similar to those for potential Sobolev spaces presented in section 5.1. It is surprising, however, that the real interpolation spaces of $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ for a fixed p do not depend on the column structure.

Proposition 5.21. *Let $1 \leq p, q \leq \infty$ and $\alpha_0, \alpha_1 \in \mathbb{R}$ with $\alpha_0 \neq \alpha_1$. Then*

$$(F_p^{\alpha_0,c}(\mathbb{T}_\theta^d), F_p^{\alpha_1,c}(\mathbb{T}_\theta^d))_{\eta,q} = B_{p,q}^\alpha(\mathbb{T}_\theta^d), \quad \alpha = (1-\eta)\alpha_0 + \eta\alpha_1.$$

Similar statements hold for the row and mixture Triebel-Lizorkin spaces.

Proof. The assertion is an immediate consequence of Proposition 4.10 (v) and Proposition 5.1 (i). Note, however, that Proposition 4.10 (v) is stated for $p < \infty$; but by duality via Proposition 4.14, it continues to hold for $p = \infty$. $\square$

On the other hand, the interpolation of $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ for a fixed α is reduced to that of Hardy spaces by virtue of Proposition 4.10 (iv) and Lemma 1.9.

Remark 5.22. Let $\alpha \in \mathbb{R}$ and $1 < p < \infty$. Then

$$(F_\infty^{\alpha,c}(\mathbb{T}_\theta^d), F_1^{\alpha,c}(\mathbb{T}_\theta^d))_{\frac{1}{p}} = F_p^{\alpha,c}(\mathbb{T}_\theta^d) = (F_\infty^{\alpha,c}(\mathbb{T}_\theta^d), F_1^{\alpha,c}(\mathbb{T}_\theta^d))_{\frac{1}{p},p}.$$

Proposition 5.23. *Let $\alpha_0, \alpha_1 \in \mathbb{R}$ and $1 < p < \infty$. Then*

$$(F_\infty^{\alpha_0,c}(\mathbb{T}_\theta^d), F_1^{\alpha_1,c}(\mathbb{T}_\theta^d))_{\frac{1}{p}} = F_p^{\alpha,c}(\mathbb{T}_\theta^d), \quad \alpha = (1 - \frac{1}{p})\alpha_0 + \frac{\alpha_1}{p}.$$

Proof. This proof is similar to that of Theorem 5.6. Let x be in the unit ball of $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$. Then by Proposition 4.10, $J^\alpha(x) \in \mathcal{H}_p^c(\mathbb{T}_\theta^d)$. Thus by Lemma 1.9, there exists a continuous function f from the strip $S = \{z \in \mathbb{C} : 0 \leq \operatorname{Re}(z) \leq 1\}$ to $\mathcal{H}_1^c(\mathbb{T}_\theta^d)$, analytic in the interior, such that $f(\frac{1}{p}) = J^\alpha(x)$ and such that

$$\sup_{t \in \mathbb{R}} \|f(it)\|_{\operatorname{BMO}^c} \leq c, \quad \sup_{t \in \mathbb{R}} \|f(1+it)\|_{\mathcal{H}_1^c} \leq c.$$

Define

$$F(z) = e^{(z-\frac{1}{p})^2} J^{-(1-z)\alpha_0 - z\alpha_1} f(z), \quad z \in S.$$

By Remark 4.15 and Lemma 5.8, for any $t \in \mathbb{R}$,

$$\|F(it)\|_{F_\infty^{\alpha,c}} \approx e^{-t^2 + \frac{1}{p^2}} \|J^{it(\alpha_0 - \alpha_1)} f(it)\|_{\operatorname{BMO}^c} \leq c'.$$

Similarly,

$$\|F(1+it)\|_{F_1^{\alpha,c}} \approx e^{-t^2 + (1-\frac{1}{p})^2} \|J^{it(\alpha_0 - \alpha_1)} f(1+it)\|_{\mathcal{H}_1^c} \leq c'.$$

Therefore,

$$x = F\left(\frac{1}{p}\right) \in (F_\infty^{\alpha_0,c}(\mathbb{T}_\theta^d), F_1^{\alpha_1,c}(\mathbb{T}_\theta^d))_{\frac{1}{p}},$$

whence

$$F_p^{\alpha,c}(\mathbb{T}_\theta^d) \subset (F_\infty^{\alpha_0,c}(\mathbb{T}_\theta^d), F_1^{\alpha_1,c}(\mathbb{T}_\theta^d))_{\frac{1}{p}}.$$

The converse inclusion is obtained by duality. $\square$

Chapter 6

Embedding

We consider the embedding problem in this chapter. We begin with Besov spaces, then pass to Sobolev spaces. Our embedding theorem for Besov spaces is complete; however, the embedding problem of $W_1^1(\mathbb{T}_\theta^d)$ is, unfortunately, left unsolved at the time of this writing. The last section deals with the compact embedding.

6.1 Embedding of Besov spaces

This section deals with the embedding of Besov spaces. We will follow the semigroup approach developed by Varopolous [76] (see also [20, 77]). This approach can be adapted to the noncommutative setting, which has been done by Junge and Mei [33]. Here we can use either the circular Poisson or heat semigroup of $\mathbb{T}_\theta^d$, already considered in section 3.3. We choose to work with the latter. Recall that for $x \in \mathcal{S}'(\mathbb{T}_\theta^d)$,

$$\mathbb{W}_r(x) = \sum_{m \in \mathbb{Z}^d} \hat{x}(m) r^{|m|^2} U^m, \quad 0 \leq r < 1.$$

The following elementary lemma will be crucial.

Lemma 6.1. *Let $1 \leq p \leq p_1 \leq \infty$. Then*

$$\|\mathbb{W}_r(x)\|_{p_1} \lesssim (1-r)^{\frac{d}{2}(\frac{1}{p_1} - \frac{1}{p})} \|x\|_p, \quad x \in L_p(\mathbb{T}_\theta^d), \quad 0 \leq r < 1. \quad (6.1)$$

Proof. Consider first the case $p = 1$ and $p_1 = \infty$. Then

$$\begin{aligned} \|\mathbb{W}_r(x)\|_\infty &\leq \sum_{m \in \mathbb{Z}^d} r^{|m|^2} |\hat{x}(m)| \leq \|x\|_1 \sum_{m \in \mathbb{Z}^d} r^{|m|^2} \\ &= \|x\|_1 \sum_{k \geq 0} r^k \sum_{|m|^2=k} 1 \lesssim \|x\|_1 \sum_{k \geq 0} (1+k)^{\frac{d}{2}} r^k \\ &\approx (1-r)^{-\frac{d}{2}} \|x\|_1. \end{aligned}$$

The general case easily follows from this special one by interpolation. Indeed, the inequality just proved means that $\mathbb{W}_r$ is bounded from $L_1(\mathbb{T}_\theta^d)$ to $L_\infty(\mathbb{T}_\theta^d)$ with norm controlled by $(1-r)^{-\frac{d}{2}}$. On the other hand, $\mathbb{W}_r$ is a contraction on $L_p(\mathbb{T}_\theta^d)$ for $1 \leq p \leq \infty$. Interpolating these two cases, we get (6.1) for $1 < p < p_1 = \infty$. The remaining case $p_1 < \infty$ is treated similarly. $\square$

The following is the main theorem of this section.

Theorem 6.2. *Assume that $1 \leq p < p_1 \leq \infty, 1 \leq q \leq q_1 \leq \infty$ and $\alpha, \alpha_1 \in \mathbb{R}$ such that $\alpha - \frac{d}{p} = \alpha_1 - \frac{d}{p_1}$. Then we have the following continuous inclusion:*

$$B_{p,q}^\alpha(\mathbb{T}_\theta^d) \subset B_{p_1,q_1}^{\alpha_1}(\mathbb{T}_\theta^d).$$

Proof. Since $B_{p_1,q}^{\alpha_1}(\mathbb{T}_\theta^d) \subset B_{p_1,q_1}^{\alpha_1}(\mathbb{T}_\theta^d)$, it suffices to consider the case $q = q_1$. On the other hand, by the lifting Theorem 3.7, we can assume $\max\{\alpha, \alpha_1\} < 0$, so that we can take $k = 0$ in Theorem 3.15. Thus, we are reduced to showing

$$\left(\int_0^1 (1-r)^{-\frac{q\alpha_1}{2}} \|\mathbb{W}_r(x)\|_{p_1}^q \frac{dr}{1-r} \right)^{\frac{1}{q}} \lesssim \left(\int_0^1 (1-r)^{-\frac{q\alpha}{2}} \|\mathbb{W}_r(x)\|_p^q \frac{dr}{1-r} \right)^{\frac{1}{q}}.$$

To this end, we write $\mathbb{W}_r(x) = \mathbb{W}_{\sqrt{r}}(\mathbb{W}_{\sqrt{r}}(x))$ and apply (6.1) to get

$$\|\mathbb{W}_r(x)\|_{p_1} \lesssim (1-\sqrt{r})^{\frac{d}{2}(\frac{1}{p_1}-\frac{1}{p})} \|\mathbb{W}_{\sqrt{r}}(x)\|_p.$$

Thus

$$\begin{aligned} \left(\int_0^1 (1-r)^{-\frac{q\alpha_1}{2}} \|\mathbb{W}_r(x)\|_{p_1}^q \frac{dr}{1-r} \right)^{\frac{1}{q}} &\lesssim \left(\int_0^1 (1-r)^{-\frac{q\alpha_1}{2}} (1-\sqrt{r})^{\frac{qd}{2}(\frac{1}{p_1}-\frac{1}{p})} \|\mathbb{W}_{\sqrt{r}}(x)\|_p^q \frac{dr}{1-r} \right)^{\frac{1}{q}} \\ &= \left(\int_0^1 (1-r^2)^{-\frac{q\alpha_1}{2}} (1-r)^{\frac{qd}{2}(\frac{1}{p_1}-\frac{1}{p})} \|\mathbb{W}_r(x)\|_p^q \frac{2rdr}{1-r^2} \right)^{\frac{1}{q}} \\ &\lesssim \left(\int_0^1 (1-r)^{-\frac{q\alpha}{2}} \|\mathbb{W}_r(x)\|_p^q \frac{dr}{1-r} \right)^{\frac{1}{q}}, \end{aligned}$$

as desired. $\square$

Corollary 6.3. *Assume that $1 \leq p < p_1 \leq \infty, 1 \leq q \leq \infty$ and $\alpha = d(\frac{1}{p} - \frac{1}{p_1})$. Then*

$$B_{p,q}^\alpha(\mathbb{T}_\theta^d) \subset L_{p_1,q}(\mathbb{T}_\theta^d) \text{ if } p_1 < \infty \text{ and } B_{p,1}^\alpha(\mathbb{T}_\theta^d) \subset L_\infty(\mathbb{T}_\theta^d) \text{ if } p_1 = \infty.$$

Proof. Applying the previous theorem to $\alpha_1 = 0$ and $q = q_1 = 1$, and by Theorem 3.8, we get

$$B_{p,1}^\alpha(\mathbb{T}_\theta^d) \subset B_{p_1,1}^0(\mathbb{T}_\theta^d) \subset L_{p_1}(\mathbb{T}_\theta^d).$$

This gives the assertion in the case $p_1 = \infty$. For $p_1 < \infty$, we fix p and choose two appropriate values of α (which give the two corresponding values of p_1); then we interpolate the resulting embeddings as above by real interpolation; finally, using (1.1) and Proposition 5.1, we obtain the announced embedding for $p_1 < \infty$. $\square$

The preceding corollary admits a self-improvement in terms of modulus of smoothness.

Corollary 6.4. *Assume that $1 \leq p < p_1 \leq \infty, \alpha = d(\frac{1}{p} - \frac{1}{p_1})$ and $k \in \mathbb{N}$ such that $k > \alpha$. Then*

$$\omega_{p_1}^k(x, \varepsilon) \lesssim \int_0^\varepsilon \delta^{-\alpha} \omega_p^k(x, \delta) \frac{d\delta}{\delta}, \quad 0 < \varepsilon \leq 1.$$

Proof. Without loss of generality, assume $\hat{x}(0) = 0$. Then by the preceding corollary and Theorem 3.16, we have

$$\|x\|_{p_1} \lesssim \int_0^1 \delta^{-\alpha} \omega_p^k(x, \delta) \frac{d\delta}{\delta}.$$

Now let $u \in \mathbb{R}^d$ with $|u| \leq \varepsilon$. Noting that

$$\omega_p^k(\Delta_u(x), \delta) \leq 2^k \min(\omega_p^k(x, \varepsilon), \omega_p^k(x, \delta)) \leq 2^k \omega_p^k(x, \min(\varepsilon, \delta)),$$

we obtain

$$\begin{aligned} \|\Delta_u(x)\|_{p_1} &\lesssim \int_0^\varepsilon \delta^{-\alpha} \omega_p^k(x, \delta) \frac{d\delta}{\delta} + \varepsilon^{-\alpha} \omega_p^k(x, \varepsilon) \\ &\lesssim \int_0^\varepsilon \delta^{-\alpha} \omega_p^k(x, \delta) \frac{d\delta}{\delta} + \int_{\frac{\varepsilon}{2}}^\varepsilon \delta^{-\alpha} \omega_p^k(x, \delta) \frac{d\delta}{\delta} \\ &\lesssim \int_0^\varepsilon \delta^{-\alpha} \omega_p^k(x, \delta) \frac{d\delta}{\delta}. \end{aligned}$$

Taking the supremum over all u with $|u| \leq \varepsilon$ yields the desired inequality. $\square$

Remark 6.5. We will discuss the optimal order of the best constant of the embedding in Corollary 6.3 at the end of the next section.

6.2 Embedding of Sobolev spaces

This section is devoted to the embedding of Sobolev spaces. The following is our main theorem. Recall that $B_{\infty, \infty}^{\alpha_1}(\mathbb{T}_\theta^d)$ in the second part below is the quantum analogue of the classical Zygmund class of order α_1 (see Remark 3.18).

Theorem 6.6. *Let $\alpha, \alpha_1 \in \mathbb{R}$ with $\alpha > \alpha_1$.*

- (i) *If $1 < p < p_1 < \infty$ are such that $\alpha - \frac{d}{p} = \alpha_1 - \frac{d}{p_1}$, then*

$$H_p^\alpha(\mathbb{T}_\theta^d) \subset H_{p_1}^{\alpha_1}(\mathbb{T}_\theta^d) \text{ continuously.}$$

In particular, if additionally $\alpha = k$ and $\alpha_1 = k_1$ are nonnegative integers, then

$$W_p^k(\mathbb{T}_\theta^d) \subset W_{p_1}^{k_1}(\mathbb{T}_\theta^d) \text{ continuously.}$$

- (ii) *If $1 \leq p < \infty$ is such that $p(\alpha - \alpha_1) > d$ and $\alpha_1 = \alpha - \frac{d}{p}$, then*

$$H_p^\alpha(\mathbb{T}_\theta^d) \subset B_{\infty, \infty}^{\alpha_1}(\mathbb{T}_\theta^d) \text{ continuously.}$$

In particular, if additionally $\alpha = k \in \mathbb{N}$, and if either $p > 1$ or $p = 1$ and k is even, then

$$W_p^k(\mathbb{T}_\theta^d) \subset B_{\infty, \infty}^{\alpha_1}(\mathbb{T}_\theta^d) \text{ continuously.}$$

Proof. (i) By Theorem 2.9, the embedding of $W_p^k(\mathbb{T}_\theta^d)$ is a special case of that of $H_p^\alpha(\mathbb{T}_\theta^d)$. Thus we just deal with the potential spaces $H_p^\alpha(\mathbb{T}_\theta^d)$. On the other hand, by the lifting property of potential Sobolev spaces, we can assume $\alpha_1 = 0$. By Theorem 3.8 and Corollary 6.3, we have

$$H_p^\alpha(\mathbb{T}_\theta^d) \subset L_{q, \infty}(\mathbb{T}_\theta^d).$$

Now choose $0 < \eta < 1$ and two indices s_0, s_1 with $1 < s_0, s_1 < \frac{d}{\alpha}$ such that

$$\frac{1}{p} = \frac{1-\eta}{s_0} + \frac{\eta}{s_1}.$$

Let

$$\frac{1}{t_j} = \frac{1}{s_j} - \frac{\alpha}{d}, \quad j = 0, 1.$$

Then interpolating the above inclusions with s_j in place of p for $j = 0, 1$, using Remark 5.3 and (1.1), we get

$$H_p^\alpha(\mathbb{T}_\theta^d) = (H_{s_0}^\alpha(\mathbb{T}_\theta^d), H_{s_1}^\alpha(\mathbb{T}_\theta^d))_{\eta, p} \subset (L_{t_0, \infty}(\mathbb{T}_\theta^d), L_{t_1, \infty}(\mathbb{T}_\theta^d))_{\eta, p} = L_{p_1, p}(\mathbb{T}_\theta^d) \subset L_{p_1}(\mathbb{T}_\theta^d).$$

(ii) By Theorems 3.8 and 6.2, we obtain

$$H_p^\alpha(\mathbb{T}_\theta^d) \subset B_{p, \infty}^\alpha(\mathbb{T}_\theta^d) \subset B_{\infty, \infty}^{\alpha_1}(\mathbb{T}_\theta^d).$$

If k is even, $W_1^k(\mathbb{T}_\theta^d) \subset H_1^k(\mathbb{T}_\theta^d)$. Thus the theorem is proved. $\square$

Remark 6.7. The case $p\alpha = d$ with $\alpha_1 = 0$ is excluded from the preceding theorem. In this case, it is easy to see that $H_p^\alpha(\mathbb{T}_\theta^d) \subset L_q(\mathbb{T}_\theta^d)$ for any $q < \infty$. It is well known in the classical case that this embedding is false for $q = \infty$. Consider, for instance, the ball $B = \{s \in \mathbb{R}^d : |s| \leq \frac{1}{4}\}$ and the function f defined by $f(s) = \log \log(1 + \frac{1}{|s|})$. Then f belongs to $W_1^d(B)$ but is unbounded on B . Now extending f to a 1-periodic function on $\mathbb{R}^d$ which is infinitely differentiable in $[-\frac{1}{2}, \frac{1}{2}]^d \setminus B$, we obtain a function in $W_1^d(\mathbb{T}^d)$ but unbounded on $\mathbb{T}^d$.

Remark 6.8. Part (ii) of the preceding theorem implies $W_p^d(\mathbb{T}_\theta^d) \subset L_\infty(\mathbb{T}_\theta^d)$ for all $p > 1$. In the commutative case, representing a function as an indefinite integral of its derivatives, one easily checks that this embedding remains true for $p = 1$. However, we do not know how to prove it in the noncommutative case. A related question concerns the embedding $W_p^k(\mathbb{T}_\theta^d) \subset B_{\infty, \infty}^{\alpha_1}(\mathbb{T}_\theta^d)$ in the case of odd k which is not covered by the same part (ii).

The quantum analogue of the Gagliardo-Nirenberg inequality can be also proved easily by interpolation.

Proposition 6.9. *Let $k \in \mathbb{N}$, $1 < r, p < \infty$, $1 \leq q < \infty$ and $\beta \in \mathbb{N}_0^d$ with $0 < |\beta|_1 < k$. If*

$$\eta = \frac{|\beta|_1}{k} \quad \text{and} \quad \frac{1}{r} = \frac{1-\eta}{q} + \frac{\eta}{p},$$

then for every $x \in W_p^k(\mathbb{T}_\theta^d) \cap L_q(\mathbb{T}_\theta^d)$,

$$\|D^\beta x\|_r \lesssim \|x\|_q^{1-\eta} \left(\sum_{|m|=k} \|D^m x\|_p \right)^\eta.$$

Proof. This inequality immediately follows from Theorem 5.13 and the well-known relation between real and complex interpolations:

$$(L_q(\mathbb{T}_\theta^d), W_p^k(\mathbb{T}_\theta^d))_{\eta, 1} \subset (L_q(\mathbb{T}_\theta^d), W_p^k(\mathbb{T}_\theta^d))_\eta = W_r^{|\beta|_1}(\mathbb{T}_\theta^d).$$

It then follows that

$$\|x\|_{W_r^{|\beta|_1}} \lesssim \|x\|_q^{1-\eta} \|x\|_{W_p^k}^\eta.$$

Applying this inequality to $x - \widehat{x}(0)$ instead of x and using Theorem 2.12, we get the desired Gagliardo-Nirenberg inequality. $\square$

An alternate approach to Sobolev embedding. Note that the preceding proof of Theorem 6.6 is based on Theorem 6.2, which is, in its turn, proved by Varopolous' semigroup approach. Varopolous initially developed his method for the Sobolev embedding, which was transferred to the noncommutative setting by Junge and Mei [33]. Our argument for the embedding of Besov spaces has followed this route. Let us now give an alternate proof of Theorem 6.6 (i) by the same way. We state its main part as the following lemma that is of interest in its own right.

Lemma 6.10. *Let $1 \leq p < q < \infty$ such that $\frac{1}{q} = \frac{1}{p} - \frac{1}{d}$. Then*

$$W_p^1(\mathbb{T}_\theta^d) \subset L_{q,\infty}(\mathbb{T}_\theta^d).$$

Proof. We will use again the heat semigroup $\mathbb{W}_r$ of $\mathbb{T}_\theta^d$. Recall that $\mathbb{W}_r = \widetilde{\mathbb{W}}_\varepsilon$ with $r = e^{-4\pi^2\varepsilon}$, where $\widetilde{\mathbb{W}}_\varepsilon$ is the periodization of the usual heat kernel $\mathbb{W}_\varepsilon$ of $\mathbb{R}^d$ (see section 3.3). It is more convenient to work with $\widetilde{\mathbb{W}}_\varepsilon$. In the following, we assume $x \in \mathcal{S}(\mathbb{T}_\theta^d)$ and $\widehat{x}(0) = 0$. Let $\Delta_j = \Delta^{-1}\partial_j$, $1 \leq j \leq d$. Then

$$\Delta^{-1}x = 4\pi^2 \int_0^\infty \widetilde{\mathbb{W}}_\varepsilon(x) d\varepsilon \quad \text{and} \quad \Delta_j x = 4\pi^2 \int_0^\infty \widetilde{\mathbb{W}}_\varepsilon(\partial_j x) d\varepsilon.$$

We claim that for any $1 \leq p \leq \infty$

$$\|\widetilde{\mathbb{W}}_\varepsilon(\partial_j x)\|_p \lesssim \varepsilon^{-\frac{1}{2}} \|x\|_p \quad \text{and} \quad \|\widetilde{\mathbb{W}}_\varepsilon(\partial_j x)\|_\infty \lesssim \varepsilon^{-\frac{1}{2}(\frac{d}{p}+1)} \|x\|_p, \quad \varepsilon > 0. \quad (6.2)$$

Indeed, in order to prove the first inequality, by the transference method, it suffices to show a similar one for the Banach space valued heat semigroup of the usual d -torus. The latter immediately follows from the following standard estimate on the heat kernel $\mathbb{W}_\varepsilon$ of $\mathbb{R}^d$:

$$\sup_{\varepsilon>0} \varepsilon^{\frac{1}{2}} \int_{\mathbb{R}^d} |\nabla \mathbb{W}_\varepsilon(s)| ds < \infty.$$

The second inequality of (6.2) is proved in the same way as (6.1). First, for the case $p = 1$, we have (recalling that $\widehat{x}(0) = 0$)

$$\begin{aligned} \|\widetilde{\mathbb{W}}_\varepsilon(\partial_j x)\|_\infty &\leq 2\pi \sum_{m \in \mathbb{Z}^d \setminus \{0\}} |m_j| e^{-\varepsilon|m|^2} |\widehat{x}(m)| \\ &\leq 2\pi \|x\|_1 \sum_{m \in \mathbb{Z}^d \setminus \{0\}} |m_j| e^{-\varepsilon|m|^2} \\ &\lesssim e^{-\varepsilon} (1 - e^{-\varepsilon})^{-\frac{d+1}{2}} \|x\|_1 \lesssim \varepsilon^{-\frac{d+1}{2}} \|x\|_1. \end{aligned}$$

Interpolating this with the first inequality for $p = \infty$, we get the second one in the general case.

Now let $\varepsilon > 0$ and decompose $\Delta_j x$ into the following two parts:

$$y = 4\pi^2 \int_\varepsilon^\infty \widetilde{\mathbb{W}}_\delta(\partial_j x) d\delta \quad \text{and} \quad z = 4\pi^2 \int_0^\varepsilon \widetilde{\mathbb{W}}_\delta(\partial_j x) d\delta.$$

Then by (6.2),

$$\|y\|_\infty \lesssim \|x\|_p \int_\varepsilon^\infty \delta^{-\frac{1}{2}(\frac{d}{p}+1)} d\delta \approx \varepsilon^{-\frac{1}{2}(\frac{d}{p}-1)} \|x\|_p$$

and

$$\|z\|_p \lesssim \|x\|_p \int_0^\varepsilon \delta^{-\frac{1}{2}} d\delta \approx \varepsilon^{\frac{1}{2}} \|x\|_p.$$

Thus for any $t > 0$, choosing ε such that $\varepsilon^{-\frac{d}{2p}} = t$, we deduce

$$\|y\|_\infty + t\|z\|_p \lesssim t^{1-\frac{p}{d}}\|x\|_p = t^\eta\|x\|_p,$$

where $\eta = 1 - \frac{p}{d}$. It then follows that

$$\|\Delta_j x\|_{q,\infty} \approx \|\Delta_j x\|_{(L^\infty(\mathbb{T}_\theta^d), L_p(\mathbb{T}_\theta^d))_{\eta,\infty}} \lesssim \|x\|_p.$$

Since

$$x = -\sum_{j=1}^d \Delta_j \partial_j x,$$

we finally get

$$\|x\|_{q,\infty} \lesssim \sum_{j=1}^d \|\Delta_j \partial_j x\|_{q,\infty} \lesssim \sum_{j=1}^d \|\partial_j x\|_p = \|\nabla x\|_p.$$

Thus the lemma is proved. $\square$

Alternate proof of Theorem 6.6 (i). For $1 < p < d$, choose p_0, p_1 such that $1 < p_0 < p < p_1 < d$. Let $\frac{1}{q_i} = \frac{1}{p_i} - \frac{1}{d}$ for $i = 0, 1$. Then by the previous lemma,

$$W_{p_i}^1(\mathbb{T}_\theta^d) \subset L_{q_i,\infty}(\mathbb{T}_\theta^d), \quad i = 0, 1.$$

Interpolating these two inclusions by real method, we obtain

$$W_p^1(\mathbb{T}_\theta^d) \subset L_{q,p}(\mathbb{T}_\theta^d).$$

This is the embedding of Sobolev spaces in Theorem 6.6 (i) for $k = 1$. The case $k > 1$ immediately follows by iteration. Then using real interpolation, we deduce the embedding of potential Sobolev spaces. $\square$

Sobolev embedding for $p = 1$. Now we discuss the case $p = 1$ which is not covered by Theorem 6.6. The main problem concerns the following:

$$W_1^1(\mathbb{T}_\theta^d) \subset L_{\frac{d}{d-1}}(\mathbb{T}_\theta^d). \quad (6.3)$$

At the time of this writing, we are unable, unfortunately, to prove it. However, Lemma 6.10 provides a weak substitute, namely,

$$W_1^1(\mathbb{T}_\theta^d) \subset L_{\frac{d}{d-1},\infty}(\mathbb{T}_\theta^d). \quad (6.4)$$

In the classical case, one can rather easily deduce (6.3) from (6.4). Let us explain the idea coming from [77, page 58]. It was kindly pointed out to us by Marius Junge. Let f be a nice real function on $\mathbb{T}^d$ with $\hat{f}(0) = 0$. For any $t \in \mathbb{R}$ let f_t be the indicator function of the subset $\{f > t\}$. Then f can be decomposed as an integral of the f_t 's:

$$f = \int_{-\infty}^{+\infty} f_t dt. \quad (6.5)$$

By triangular inequality (with $q = \frac{d}{d-1}$),

$$\|f\|_q \leq \int_{-\infty}^{+\infty} \|f_t\|_q dt.$$

However,

$$\|f_t\|_q = \|f_t\|_{q,\infty} \quad \forall t \in \mathbb{R}.$$

Thus by (6.4) for $\theta = 0$,

$$\|f_t\|_q \lesssim \|f_t\|_1 + \|\nabla f_t\|_1.$$

It comes now the crucial point which is the following

$$\int_{-\infty}^{+\infty} \|\nabla f_t\|_1 dt \lesssim \|\nabla f\|_1. \quad (6.6)$$

In fact, the two sides are equal in view of Sard's theorem. We then get the strong embedding (6.3) in the case $\theta = 0$. Note that this proof yields a stronger embedding:

$$W_1^1(\mathbb{T}^d) \subset L_{\frac{d}{d-1},1}(\mathbb{T}^d). \quad (6.7)$$

The above decomposition of f is not smooth in the sense that f_t is not derivable even though f is nice. In his proof of Hardy's inequality in Sobolev spaces, Bourgain [11] discovered independently the same decomposition but using nicer functions f_t (see also [60]). Using (6.7) and the Hausdorff-Young inequality, Bourgain derived the following Hardy type inequality (assuming $d \geq 3$):

$$\sum_{m \in \mathbb{Z}^d} \frac{|\widehat{f}(m)|}{(1 + |m|)^{d-1}} \lesssim \|f\|_{W_1^1(\mathbb{T}^d)}.$$

We have encountered difficulties in the attempt of extending this approach to the noncommutative case. Let us formulate the corresponding open problems explicitly as follows:

Problem 6.11. Let $d \geq 2$.

(i) Does one have the following embedding

$$W_1^1(\mathbb{T}_\theta^d) \subset L_{\frac{d}{d-1}}(\mathbb{T}_\theta^d) \quad \text{or} \quad W_1^1(\mathbb{T}_\theta^d) \subset L_{\frac{d}{d-1},1}(\mathbb{T}_\theta^d)?$$

(ii) Does one have the following inequality

$$\sum_{m \in \mathbb{Z}^d} \frac{|\widehat{x}(m)|}{(1 + |m|)^{d-1}} \lesssim \|x\|_{W_1^1(\mathbb{T}_\theta^d)}?$$

By the previous discussion, part (i) is reduced to a decomposition for operators in $W_1^1(\mathbb{T}_\theta^d)$ of the form (6.5) and satisfying (6.6). One could be attempted to do this by transference by first considering operator-valued functions on $\mathbb{R}^d$. With this in mind, the following observation, due to Marius Junge, might be helpful.

Given an interval $I = [s, t] \subset \mathbb{R}$ and an element $a \in L_1(\mathbb{T}_\theta^d)$, we have

$$\partial(\mathbf{1}_I \otimes a) = \delta_s \otimes a - \delta_t \otimes a,$$

where ∂ denotes the distribution derivative relative to $\mathbb{R}$. Let $\|\cdot\|_L$ denote the norm of the dual space $C_0(\mathbb{R}; \mathcal{A}_\theta)^*$, which contains $L_1(\mathbb{R}; L_1(\mathbb{T}_\theta^d))$ isometrically. If f is a (nice) linear combination of $\mathbf{1}_I \otimes a$'s, then we have the desired decomposition of f . Indeed, assume

$f = \sum_{i=1}^n \alpha_i \mathbb{1}_{I_i} \otimes e_i$, where $\alpha_i \in \mathbb{R}_+$ and the $\mathbb{1}_{I_i} \otimes e_i$'s are pairwise disjoint projections of $L_\infty(\mathbb{R}) \overline{\otimes} \mathbb{T}_\theta^d$. Let $f_t = \mathbb{1}_{(t,\infty)}(f)$. Then

$$f = \int_0^\infty f_t dt.$$

So for any $q \geq 1$,

$$\|f\|_q \leq \int_0^\infty \|f_t\|_q dt.$$

On the other hand, by writing explicitly f_t for every t , one easily checks

$$\|\partial f\|_L = \int_0^\infty \|\partial f_t\|_L dt.$$

By iteration, the above decomposition can be extended to higher dimensional case for all functions f of the form $\sum_{i=1}^n \alpha_i \mathbb{1}_{R_i} \otimes e_i$, where $\alpha_i \in \mathbb{R}_+$, R_i 's are rectangles (with sides parallel to the axes) and $\mathbb{1}_{R_i} \otimes e_i$'s are pairwise disjoint projections of $L_\infty(\mathbb{R}^d) \overline{\otimes} \mathbb{T}_\theta^d$.

The next idea would be to apply Lemma 6.10 to these special functions. Then two difficulties come up to us, even in the commutative case. The first is that these functions do not belong to W_1^1 ; this difficulty can be resolved quite easily by regularization. The second one, substantial, is the density of these functions, more precisely, of suitable regularizations of them, in W_1^1 .

Uniform Besov embedding. We end this section with a discussion on the link between a certain uniform embedding of Besov spaces and the embedding of Sobolev spaces. Let $0 < \alpha < 1$, $1 \leq p < \infty$ with $\alpha p < d$ and $\frac{1}{r} = \frac{1}{p} - \frac{\alpha}{d}$. Then

$$\|x\|_r^p \leq c_{d,p} \frac{\alpha(1-\alpha)}{(d-\alpha p)} \|x\|_{B_{p,p}^{\alpha,\omega}}^p, \quad x \in B_{p,p}^\alpha(\mathbb{T}_\theta^d), \quad (6.8)$$

where $\|x\|_{B_{p,p}^{\alpha,\omega}}$ is the Besov norm defined by (3.19). In the commutative case, this inequality is proved in [13] for α close to 1 and in [41] for general α . One can show that (6.8) is essentially equivalent to the embedding of $W_p^1(\mathbb{T}_\theta^d)$ into $L_q(\mathbb{T}_\theta^d)$ (or $L_{q,p}(\mathbb{T}_\theta^d)$) for $d > p$ and $\frac{1}{q} = \frac{1}{p} - \frac{1}{d}$. Indeed, assume (6.8). Then taking limit in both sides of (6.8) as $\alpha \rightarrow 1$, by Theorem 3.19, we get

$$\|x\|_q \lesssim |x|_{W_p^1}$$

for all $x \in W_p^1(\mathbb{T}_\theta^d)$ with $\widehat{x}(0) = 0$. Conversely, if $W_p^1(\mathbb{T}_\theta^d) \subset L_q(\mathbb{T}_\theta^d)$, then

$$(L_p(\mathbb{T}_\theta^d), W_p^1(\mathbb{T}_\theta^d))_{\alpha,p} \subset (L_p(\mathbb{T}_\theta^d), L_q(\mathbb{T}_\theta^d))_{\alpha,p}.$$

Theorem 5.16 implies that

$$(L_p(\mathbb{T}_\theta^d), W_p^1(\mathbb{T}_\theta^d))_{\alpha,p} \subset B_{p,p}^\alpha(\mathbb{T}_\theta^d)$$

with relevant constant depending only on d , here $B_{p,p}^\alpha(\mathbb{T}_\theta^d)$ being equipped with the norm $\|\cdot\|_{B_{p,p}^{\alpha,\omega}}$. On the other hand, By a classical result of Holmstedt [28] on real interpolation of L_p -spaces (which readily extends to the noncommutative case, as observed in [37, Lemma 3.7]),

$$(L_p(\mathbb{T}_\theta^d), L_q(\mathbb{T}_\theta^d))_{\alpha,p} \subset L_{r,p}(\mathbb{T}_\theta^d)$$

with the inclusion constant uniformly controlled by $\alpha^{\frac{1}{p}}(1-\alpha)^{\frac{1}{q}}$. We then deduce

$$\|x\|_{r,p} \lesssim \alpha^{\frac{1}{p}}(1-\alpha)^{\frac{1}{q}} \|x\|_{B_{p,p}^{\alpha,\omega}}.$$

This implies a variant of (6.8) since $L_{r,p}(\mathbb{T}_\theta^d) \subset L_r(\mathbb{T}_\theta^d)$.

Since we have proved the embedding $W_p^1(\mathbb{T}_\theta^d) \subset L_q(\mathbb{T}_\theta^d)$ for $p > 1$, (6.8) holds for $p > 1$. Let us record this explicitly as follows:

Proposition 6.12. *Let $0 < \alpha < 1$, $1 < p < \infty$ with $\alpha p < d$ and $\frac{1}{r} = \frac{1}{p} - \frac{\alpha}{d}$. Then*

$$\|x\|_r \lesssim (\alpha(1-\alpha))^{\frac{1}{p}} \|x\|_{B_{p,p}^{\alpha,\omega}}^p, \quad x \in B_{p,p}^\alpha(\mathbb{T}_\theta^d)$$

with relevant constant independent of α .

In the case $p = 1$, Problem 6.11 (i) is equivalent to (6.8) for $p = 1$ and α close to 1.

6.3 Compact embedding

This section deals with the compact embedding. The case $p = 2$ for potential Sobolev spaces was solved by Spera [65]:

Lemma 6.13. *The embedding $H_2^{\alpha_1}(\mathbb{T}_\theta^d) \hookrightarrow H_2^{\alpha_2}(\mathbb{T}_\theta^d)$ is compact for $\alpha_1 > \alpha_2 \geq 0$.*

We will require the following real interpolation result on compact operators, due to Cwikel [19].

Lemma 6.14. *Let (X_0, X_1) and (Y_0, Y_1) be two interpolation couples of Banach spaces, and let $T : X_j \rightarrow Y_j$ be a bounded linear operator, $j = 0, 1$. If $T : X_0 \rightarrow Y_0$ is compact, then $T : (X_0, X_1)_{\eta,p} \rightarrow (Y_0, Y_1)_{\eta,p}$ is compact too for any $0 < \eta < 1$ and $1 \leq p \leq \infty$.*

Theorem 6.15. *Assume that $1 \leq p < p_1 \leq \infty$, $1 \leq p^* < p_1$, $1 \leq q \leq q_1 \leq \infty$ and $\alpha - \frac{d}{p} = \alpha_1 - \frac{d}{p_1}$. Then the embedding $B_{p,q}^\alpha(\mathbb{T}_\theta^d) \hookrightarrow B_{p^*,q_1}^{\alpha_1}(\mathbb{T}_\theta^d)$ is compact.*

Proof. Without loss of generality, we can assume $q = q_1$. First consider the case $p = 2$. Choose t sufficiently close to q and $0 < \eta < 1$ such that

$$\frac{1}{q} = \frac{1-\eta}{2} + \frac{\eta}{t}.$$

Then by Proposition 5.1,

$$B_{2,q}^\alpha(\mathbb{T}_\theta^d) = (B_{2,2}^\alpha(\mathbb{T}_\theta^d), B_{2,t}^\alpha(\mathbb{T}_\theta^d))_{\eta,q}.$$

By Lemma 6.13, $B_{2,2}^\alpha(\mathbb{T}_\theta^d) \hookrightarrow B_{2,2}^{\alpha_1}(\mathbb{T}_\theta^d)$ is compact. On the other hand, by Theorem 6.2, $B_{2,t}^\alpha(\mathbb{T}_\theta^d) \hookrightarrow B_{p_1,t}^{\alpha_1}(\mathbb{T}_\theta^d)$ is continuous. So by Lemma 6.14,

$$B_{2,q}^\alpha(\mathbb{T}_\theta^d) \hookrightarrow (B_{2,2}^{\alpha_1}(\mathbb{T}_\theta^d), B_{p_1,t}^{\alpha_1}(\mathbb{T}_\theta^d))_{\eta,q} \text{ is compact.}$$

However, by the proof of Proposition 5.1 and (1.1), we have

$$(B_{2,2}^{\alpha_1}(\mathbb{T}_\theta^d), B_{p_1,t}^{\alpha_1}(\mathbb{T}_\theta^d))_{\eta,q} \subset \ell_q^{\alpha_1}((L_2(\mathbb{T}_\theta^d), L_{p_1}(\mathbb{T}_\theta^d))_{\eta,q}) = \ell_q^{\alpha_1}(L_{s,q}(\mathbb{T}_\theta^d)),$$

where s is determined by

$$\frac{1}{s} = \frac{1-\eta}{2} + \frac{\eta}{p_1} = \frac{1}{p_1} + \frac{(1-\eta)(\alpha - \alpha_1)}{d}.$$

Note that η tends to 1 as t tends to q , so we can choose η such that $s > p^*$. Then $L_{s,q}(\mathbb{T}_\theta^d) \subset L_{p^*}(\mathbb{T}_\theta^d)$. Thus the desired assertion for $p = 2$ follows.

The case $p \neq 2$ but $p > 1$ is dealt with similarly. Let t and η be as above. Choose $r < p$ (r close to p). Then

$$(B_{2,2}^\alpha(\mathbb{T}_\theta^d), B_{r,t}^\alpha(\mathbb{T}_\theta^d))_{\eta,q} \subset \ell_q^\alpha((L_2(\mathbb{T}_\theta^d), L_r(\mathbb{T}_\theta^d))_{\eta,q}) = \ell_q^\alpha(L_{p_0,q}(\mathbb{T}_\theta^d)),$$

where p_0 is determined by

$$\frac{1}{p_0} = \frac{1-\eta}{2} + \frac{\eta}{r}.$$

If η is sufficiently close to 1, then $p_0 < p$ that we will assume. Thus $L_p(\mathbb{T}_\theta^d) \subset L_{p_0,q}(\mathbb{T}_\theta^d)$. It then follows that

$$B_{p,q}^\alpha(\mathbb{T}_\theta^d) \subset (B_{2,2}^\alpha(\mathbb{T}_\theta^d), B_{r,t}^\alpha(\mathbb{T}_\theta^d))_{\eta,q}.$$

The rest of the proof is almost the same as the case $p = 2$, so is omitted.

The remaining case $p = 1$ can be easily reduced to the previous one. Indeed, first embed $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ into $B_{p_2,q}^{\alpha_2}(\mathbb{T}_\theta^d)$ for some $\alpha_2 \in (\alpha, \alpha_1)$ (α_2 close to α) and p_2 determined by $\alpha - \frac{d}{p} = \alpha_2 - \frac{d}{p_2}$. Then by the previous case, the embedding $B_{p_2,q}^{\alpha_2}(\mathbb{T}_\theta^d) \hookrightarrow B_{p^*,q_1}^{\alpha_1}(\mathbb{T}_\theta^d)$ is compact, so we are done. $\square$

Theorem 6.16. *Let $1 < p < p_1 < \infty$ and $\alpha, \alpha_1 \in \mathbb{R}$.*

- (i) *If $\alpha - \frac{d}{p} = \alpha_1 - \frac{d}{p_1}$, then $H_p^\alpha(\mathbb{T}_\theta^d) \hookrightarrow H_{p^*}^{\alpha_1}(\mathbb{T}_\theta^d)$ is compact for $p^* < p_1$. In particular, if additionally $\alpha = k$ and $\alpha_1 = k_1$ are nonnegative integers, then $W_p^k(\mathbb{T}_\theta^d) \hookrightarrow W_{p^*}^{k_1}(\mathbb{T}_\theta^d)$ is compact.*
- (ii) *If $p(\alpha - \alpha_1) > d$ and $\alpha^* < \alpha_1 = \alpha - \frac{d}{p}$, then $H_p^\alpha(\mathbb{T}_\theta^d) \hookrightarrow B_{\infty,\infty}^{\alpha^*}(\mathbb{T}_\theta^d)$ is compact. In particular, if additionally $\alpha = k \in \mathbb{N}$, then $W_p^k(\mathbb{T}_\theta^d) \hookrightarrow B_{\infty,\infty}^{\alpha^*}(\mathbb{T}_\theta^d)$ is compact.*

Proof. Based on the preceding theorem, this proof is similar to that of Theorem 6.6 and left to the reader. $\square$

Chapter 7

Fourier multiplier

This chapter deals with Fourier multipliers on Sobolev, Besov and Triebel-Lizorkin spaces on $\mathbb{T}_\theta^d$. The first section concerns the Sobolev spaces. Its main result is the analogue for $W_p^k(\mathbb{T}_\theta^d)$ of [17, Theorem 7.3] (see also Lemma 1.3) on c.b. Fourier multipliers on $L_p(\mathbb{T}_\theta^d)$; so the space of c.b. Fourier multipliers on $W_p^k(\mathbb{T}_\theta^d)$ is independent of θ . The second section turns to Besov spaces on which Fourier multipliers behave better. We extend some classical results to the present setting. We show that the space of c.b. Fourier multipliers on $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ does not depend θ (nor on q or α). We also prove that a function on $\mathbb{Z}^d$ is a Fourier multiplier on $B_{1,q}^\alpha(\mathbb{T}_\theta^d)$ iff it is the Fourier transform of an element of $B_{1,\infty}^0(\mathbb{T}^d)$. The last section deals with Fourier multipliers on Triebel-Lizorkin spaces.

7.1 Fourier multipliers on Sobolev spaces

We now investigate Fourier multipliers on Sobolev spaces. We refer to [59, 9] for the study of Fourier multipliers on the classical Sobolev spaces. If X is a Banach space of distributions on $\mathbb{T}_\theta^d$, we denote by $\mathbf{M}(X)$ the space of bounded Fourier multipliers on X ; if X is further equipped with an operator space structure, $\mathbf{M}_{\text{cb}}(X)$ is the space of c.b. Fourier multipliers on X . These spaces are endowed with their natural norms. Recall that the Sobolev spaces $W_p^k(\mathbb{T}_\theta^d)$, $H_p^\alpha(\mathbb{T}_\theta^d)$ and the Besov $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ are equipped with their natural operator space structures as defined in Remarks 2.29 and 3.24.

The aim of this section is to extend [17, Theorem 7.3] (see also Lemma 1.3) on c.b. Fourier multipliers on $L_p(\mathbb{T}_\theta^d)$ to Sobolev spaces. Inspired by Neuwirth and Ricard's transference theorem [48], we will relate Fourier multipliers with Schur multipliers. Given a distribution x on $\mathbb{T}_\theta^d$, we write its matrix in the basis $(U^m)_{m \in \mathbb{Z}^d}$:

$$[x] = \left(\langle x U^n, U^m \rangle \right)_{m,n \in \mathbb{Z}^d} = \left(\widehat{x}(m-n) e^{in\tilde{\theta}(m-n)^t} \right)_{m,n \in \mathbb{Z}^d}.$$

Here k^t denotes the transpose of $k = (k_1, \dots, k_d)$ and $\tilde{\theta}$ is the following $d \times d$ -matrix deduced from the skew symmetric matrix θ :

$$\tilde{\theta} = -2\pi \begin{pmatrix} 0 & \theta_{12} & \theta_{13} & \dots & \theta_{1d} \\ 0 & 0 & \theta_{23} & \dots & \theta_{2d} \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & \dots & \theta_{d-1,d} \\ 0 & 0 & 0 & \dots & 0 \end{pmatrix}.$$

Now let $\phi : \mathbb{Z}^d \rightarrow \mathbb{C}$ and M_ϕ be the associated Fourier multiplier on $\mathbb{T}_\theta^d$. Set $\mathring{\phi} = (\phi_{m-n})_{m,n \in \mathbb{Z}^d}$. Then

$$[M_\phi x] = (\phi_{m-n} \hat{x}(m-n) e^{in\tilde{\theta}(m-n)^t})_{m,n \in \mathbb{Z}^d} = S_{\mathring{\phi}}([x]), \quad (7.1)$$

where $S_{\mathring{\phi}}$ is the Schur multiplier with symbol $\mathring{\phi}$.

According to the definition of $W_p^k(\mathbb{T}_\theta^d)$, for any matrix $a = (a_{m,n})_{m,n \in \mathbb{Z}^d}$ and $\ell \in \mathbb{N}_0^d$ define

$$D^\ell a = ((2\pi i(m-n))^\ell a_{m,n})_{m,n \in \mathbb{Z}^d}.$$

If x is a distribution on $\mathbb{T}_\theta^d$, then clearly

$$[M_\phi D^\ell x] = S_{\mathring{\phi}}(D^\ell [x]).$$

We introduce the space

$$S_p^k = \left\{ a = (a_{m,n})_{m,n \in \mathbb{Z}^d} : D^\ell a \in S_p(\ell_2(\mathbb{Z}^d)), \forall \ell \in \mathbb{N}_0^d, 0 \leq |\ell|_1 \leq k \right\}$$

and endow it with the norm

$$\|a\|_{S_p^k} = \left(\sum_{0 \leq |\ell|_1 \leq k} \|D^\ell a\|_{S_p}^p \right)^{\frac{1}{p}}.$$

Then S_p^k is a closed subspace of the ℓ_p -direct sum of L copies of $S_p(\ell_2(\mathbb{Z}^d))$ with $L = \sum_{0 \leq |\ell|_1 \leq k} 1$. The latter direct sum is equipped with its natural operator space structure, which induces an operator space structure on S_p^k too.

If $\psi = (\psi_{m,n})_{m,n \in \mathbb{Z}^d}$ is a complex matrix, its associated Schur multiplier S_ψ on S_p^k is defined by $S_\psi a = (\psi_{m,n} a_{m,n})_{m,n \in \mathbb{Z}^d}$. Let $\mathbf{M}_{\text{cb}}(S_p^k)$ denote the space of all c.b. Schur multipliers on S_p^k , equipped with the natural norm.

Theorem 7.1. *Let $1 \leq p \leq \infty$ and $k \in \mathbb{N}$. Then*

$$\mathbf{M}_{\text{cb}}(W_p^k(\mathbb{T}_\theta^d)) = \mathbf{M}_{\text{cb}}(S_p^k) \text{ with equal norms.}$$

Consequently,

$$\mathbf{M}_{\text{cb}}(W_p^k(\mathbb{T}_\theta^d)) = \mathbf{M}_{\text{cb}}(W_p^k(\mathbb{T}^d)) \text{ with equal norms.}$$

Proof. This proof is an adaptation of that of [17, Theorem 7.3]. We start with an elementary observation. Let $V = \text{diag}(\dots, U^n, \dots)_{n \in \mathbb{Z}^d}$. For any $a = (a_{m,n})_{m,n \in \mathbb{Z}^d} \in B(\ell_2(\mathbb{Z}^d))$, let $x = V(a \otimes 1_{\mathbb{T}_\theta^d})V^* \in B(\ell_2(\mathbb{Z}^d)) \bar{\otimes} \mathbb{T}_\theta^d$, where $1_{\mathbb{T}_\theta^d}$ denotes the unit of $\mathbb{T}_\theta^d$. Then

$$x = (U^m a_{mn} U^{-n})_{m,n \in \mathbb{Z}^d} = \sum_{m,n} a_{m,n} e_{m,n} \otimes U^m U^{-n} = \sum_{m,n} a_{m,n} e_{m,n} \otimes e^{-in\tilde{\theta}m^t} U^{m-n},$$

where $(e_{m,n})$ are the canonical matrix units of $B(\ell_2(\mathbb{Z}^d))$. So,

$$[x] = (a_{m,n} e_{m,n})_{m,n \in \mathbb{Z}^d},$$

a matrix with entries in $B(\ell_2(\mathbb{Z}^d))$. Since V is unitary, we have

$$\|x\|_{L_p(B(\ell_2(\mathbb{Z}^d)) \bar{\otimes} \mathbb{T}_\theta^d)} = \|a \otimes 1_{\mathbb{T}_\theta^d}\|_{L_p(B(\ell_2(\mathbb{Z}^d)) \bar{\otimes} \mathbb{T}_\theta^d)} = \|a\|_{S_p(\ell_2(\mathbb{Z}^d))}.$$

Similarly, for $\ell \in \mathbb{N}_0^d$,

$$\|D^\ell x\|_{L_p(B(\ell_2(\mathbb{Z}^d)) \otimes \mathbb{T}_\theta^d)} = \|D^\ell a\|_{S_p(\ell_2(\mathbb{Z}^d))}. \quad (7.2)$$

Now suppose that $\phi \in \mathbf{M}_{\text{cb}}(W_p^k(\mathbb{T}_\theta^d))$. For $a = (a_{m,n})_{m,n \in \mathbb{Z}^d} \in B(\ell_2(\mathbb{Z}^d))$, define $x = V(a \otimes 1_{\mathbb{T}_\theta^d})V^*$ as above. Then by (7.1), for $\ell \in \mathbb{N}_0^d$,

$$(\text{Id}_{B(\ell_2(\mathbb{Z}^d))} \otimes M_\phi)(D^\ell x) = V(S_\phi^\circ(D^\ell a) \otimes 1_{\mathbb{T}_\theta^d})V^*.$$

It then follows from (7.2) that

$$\begin{aligned} \|S_\phi^\circ(a)\|_{S_p^k} &= \left[\sum_{|\ell|_1 \leq k} \|(\text{Id}_{B(\ell_2(\mathbb{Z}^d))} \otimes M_\phi)(D^\ell x)\|_{L_p(B(\ell_2(\mathbb{Z}^d)) \otimes \mathbb{T}_\theta^d)}^p \right]^{\frac{1}{p}} \\ &\leq \|\phi\|_{\mathbf{M}_{\text{cb}}(W_p^k(\mathbb{T}_\theta^d))} \left[\sum_{|\ell|_1 \leq k} \|D^\ell x\|_{L_p(B(\ell_2(\mathbb{Z}^d)) \otimes \mathbb{T}_\theta^d)}^p \right]^{\frac{1}{p}} \\ &= \|\phi\|_{\mathbf{M}_{\text{cb}}(W_p^k(\mathbb{T}_\theta^d))} \|a\|_{S_p^k}. \end{aligned}$$

Therefore, ϕ° is a bounded Schur multiplier on S_p^k . Considering matrices $a = (a_{m,n})_{m,n \in \mathbb{Z}^d}$ with entries in S_p , we show in the same way that M_ϕ° is c.b. on S_p^k , so ϕ° is a c.b. Schur multiplier on S_p^k and

$$\|\phi^\circ\|_{\mathbf{M}_{\text{cb}}(S_p^k)} \leq \|\phi\|_{\mathbf{M}_{\text{cb}}(W_p^k(\mathbb{T}_\theta^d))}.$$

To show the converse direction, introducing the following Følner sequence of $\mathbb{Z}^d$:

$$Z_N = \{-N, \dots, -1, 0, 1, \dots, N\}^d \subset \mathbb{Z}^d,$$

we define two maps A_N and B_N as follows:

$$A_N : \mathbb{T}_\theta^d \rightarrow B(\ell_2^{|Z_N|}) \quad \text{with} \quad x \mapsto P_N([x]),$$

where $P_N : B(\ell_2(\mathbb{Z}^d)) \rightarrow B(\ell_2^{|Z_N|})$ with $(a_{m,n}) \mapsto (a_{m,n})_{m,n \in Z_N}$; and

$$B_N : B(\ell_2^{|Z_N|}) \rightarrow \mathbb{T}_\theta^d \quad \text{with} \quad e_{m,n} \mapsto \frac{1}{|Z_N|} e^{-in\tilde{\theta}(m-n)t} U^{m-n}.$$

Here $B(\ell_2^{|Z_N|})$ is endowed with the normalized trace. Both A_N, B_N are unital, completely positive and trace preserving, so extend to complete contractions between the corresponding L_p -spaces. Moreover,

$$\lim_{N \rightarrow \infty} B_N \circ A_N(x) = x \quad \text{in} \quad L_p(\mathbb{T}_\theta^d), \quad \forall x \in L_p(\mathbb{T}_\theta^d).$$

If we define $S_p^k(\ell_2^{|Z_N|})$ as before for S_p^k just replacing $S_p(\ell_2(\mathbb{Z}^d))$ by $S_p(\ell_2^{|Z_N|})$, we see that A_N extends to a complete contraction from $W_p^k(\mathbb{T}_\theta^d)$ into $S_p^k(\ell_2^{|Z_N|})$, while B_N a complete contraction from $S_p^k(\ell_2^{|Z_N|})$ into $W_p^k(\mathbb{T}_\theta^d)$.

Now assume that ϕ° is a c.b. Schur multiplier on S_p^k , then it is also a c.b. Schur multiplier on $S_p^k(\ell_2^{|Z_N|})$. We want to prove that M_ϕ is c.b. on $W_p^k(\mathbb{T}_\theta^d)$. For any $x \in$

$$L_p(B(\ell_2(\mathbb{Z}^d)) \overline{\otimes} \mathbb{T}_\theta^d),$$

$$\begin{aligned} \|\text{Id} \otimes M_\phi(x)\|_{L_p(B(\ell_2(\mathbb{Z}^d)) \overline{\otimes} \mathbb{T}_\theta^d)} &= \lim_N \|(\text{Id} \otimes B_N) \circ (\text{Id} \otimes A_N)(\text{Id} \otimes M_\phi(x))\|_{L_p(B(\ell_2(\mathbb{Z}^d)) \overline{\otimes} \mathbb{T}_\theta^d)} \\ &= \lim_N \|(\text{Id} \otimes B_N) \circ (\text{Id} \otimes S_\phi^\circ)(\text{Id} \otimes A_N(x))\|_{L_p(B(\ell_2(\mathbb{Z}^d)) \overline{\otimes} \mathbb{T}_\theta^d)} \\ &\leq \limsup_N \|\text{Id} \otimes S_\phi^\circ(\text{Id} \otimes A_N(x))\|_{S_p(\ell_2(\mathbb{Z}^d); S_p^k(\ell_2^{|Z_N|}))} \\ &\leq \limsup_N \|S_\phi^\circ\|_{\text{cb}} \|\text{Id} \otimes A_N(x)\|_{S_p(\ell_2(\mathbb{Z}^d); S_{p,k}(\ell_2^{|Z_N|}))} \\ &\leq \|S_\phi^\circ\|_{\text{cb}} \|x\|_{L_p(B(\ell_2(\mathbb{Z}^d)) \overline{\otimes} \mathbb{T}_\theta^d)}, \end{aligned}$$

where in the second equality we have used the fact that

$$\text{Id} \otimes A_N(\text{Id} \otimes M_\phi(x)) = \text{Id} \otimes S_\phi^\circ(\text{Id} \otimes A_N(x)),$$

which follows from (7.1). Therefor, M_ϕ is c.b. on $W_p^k(\mathbb{T}_\theta^d)$ and

$$\|\phi\|_{\text{M}_{\text{cb}}(W_p^k(\mathbb{T}_\theta^d))} \leq \|\phi^\circ\|_{\text{M}_{\text{cb}}(S_p^k)}.$$

The theorem is thus proved. $\square$

Remark 7.2. Let $1 \leq p \leq \infty$ and $\alpha \in \mathbb{R}$. Since J^α is a complete isometry from $H_p^\alpha(\mathbb{T}_\theta^d)$ onto $L_p(\mathbb{T}_\theta^d)$, we have

$$\text{M}_{\text{cb}}(H_p^\alpha(\mathbb{T}_\theta^d)) = \text{M}_{\text{cb}}(L_p(\mathbb{T}_\theta^d)) \text{ with equal norms.}$$

Thus, by Lemma 1.3

$$\text{M}_{\text{cb}}(H_p^\alpha(\mathbb{T}_\theta^d)) = \text{M}_{\text{cb}}(H_p^\alpha(\mathbb{T}^d)) \text{ with equal norms.}$$

Note that the proof of Theorem 2.9 shows that $W_p^k(\mathbb{T}_\theta^d) = H_p^k(\mathbb{T}_\theta^d)$ holds completely isomorphically for $1 < p < \infty$. Thus the above remark implies

Corollary 7.3. *Let $1 < p < \infty$ and $k \in \mathbb{N}$.*

$$\text{M}_{\text{cb}}(W_p^k(\mathbb{T}_\theta^d)) = \text{M}_{\text{cb}}(L_p(\mathbb{T}^d)) \text{ with equivalent norms.}$$

Clearly, the above equality still holds for $p = 1$ or $p = \infty$ if $d = 1$ (the commutative case) since then $W_p^k(\mathbb{T}) = L_p(\mathbb{T})$ for all $1 \leq p \leq \infty$ by the (complete) isomorphism

$$L_p(\mathbb{T}) \ni x \mapsto \hat{x}(0) + \sum_{m \in \mathbb{Z} \setminus \{0\}} \frac{1}{(2\pi i m)^k} \hat{x}(m) z^m \in W_p^k(\mathbb{T}).$$

However, this is no longer the case as soon as $d \geq 2$, as proved by Poornima [59] in the commutative case for $\mathbb{R}^d$. Poornima's example comes from Ornstein [50] which is still valid for our setting. Indeed, by [50], there exists a distribution T on $\mathbb{T}^2$ which is not a measure and such that $T = \partial_1 \mu_0$, $\partial_1 T = \partial_2 \mu_1$ and $\partial_2 T = \partial_1 \mu_2$ for three measures μ_i on $\mathbb{T}^2$. T induces a Fourier multiplier on $\mathbb{T}_\theta^2$, which is defined by the Fourier transform of T and is denoted by $x \mapsto T * x$. Then for any $x \in W_1^1(\mathbb{T}_\theta^2)$,

$$\begin{aligned} T * x &= \partial_1 \mu_0 * x = \mu_0 * \partial_1 x \in L_1(\mathbb{T}_\theta^2), \\ \partial_1 T * x &= \partial_1 \mu_1 * x = \mu_1 * \partial_2 x \in L_1(\mathbb{T}_\theta^2), \\ \partial_2 T * x &= \partial_2 \mu_2 * x = \mu_2 * \partial_1 x \in L_1(\mathbb{T}_\theta^2). \end{aligned}$$

Thus $T * x \in W_1^1(\mathbb{T}_\theta^2)$, so the Fourier multiplier induced by T is bounded on $W_1^1(\mathbb{T}_\theta^2)$. We show in the same way that it is c.b. too. Since T is not a measure, it does not belong to $\text{M}(L_1(\mathbb{T}^2))$.

7.2 Fourier multipliers on Besov spaces

It is well known that in the classical setting, Fourier multipliers behave better on Besov spaces than on L_p -spaces. We will see that this fact remains true in the quantum case. We maintain the notation introduced in section 3.1. In particular, φ is a function satisfying (3.1) and $\varphi^{(k)}(\xi) = \varphi(2^{-k}\xi)$ for $k \in \mathbb{N}_0$. As usual, $\varphi^{(k)}$ is viewed as a function on $\mathbb{Z}^d$ too.

The following is the main result of this section. Compared with the corresponding result in the classical case (see, for instance, Section 2.6 of [73]), our result is more precise since it gives a characterization of Fourier multipliers on $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ in terms of those on $L_p(\mathbb{T}_\theta^d)$.

Theorem 7.4. *Let $\alpha \in \mathbb{R}$ and $1 \leq p, q \leq \infty$. Let $\phi : \mathbb{Z}^d \rightarrow \mathbb{C}$. Then ϕ is a Fourier multiplier on $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ iff the $\phi\varphi^{(k)}$'s are Fourier multipliers on $L_p(\mathbb{T}_\theta^d)$ uniformly in k . In this case, we have*

$$\|\phi\|_{\mathbf{M}(B_{p,q}^\alpha(\mathbb{T}_\theta^d))} \approx |\phi(0)| + \sup_{k \geq 0} \|\phi\varphi^{(k)}\|_{\mathbf{M}(L_p(\mathbb{T}_\theta^d))}$$

with relevant constants depending only on α . A similar c.b. version holds too.

Proof. Without loss of generality, we assume that $\phi(0) = 0$ and all elements x considered below have vanishing Fourier coefficients at the origin. Let $\phi \in \mathbf{M}(B_{p,q}^\alpha(\mathbb{T}_\theta^d))$ and $x \in L_p(\mathbb{T}_\theta^d)$. Then $y = (\tilde{\varphi}_{k-1} + \tilde{\varphi}_k + \tilde{\varphi}_{k+1}) * x \in B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ and

$$\|y\|_{B_{p,q}^\alpha} \leq c_\alpha 2^{k\alpha} \|x\|_p \quad \text{with } c_\alpha = 9 \cdot 4^{|\alpha|}.$$

So

$$\|M_\phi(y)\|_{B_{p,q}^\alpha} \leq \|\phi\|_{\mathbf{M}(B_{p,q}^\alpha(\mathbb{T}_\theta^d))} \|y\|_{B_{p,q}^\alpha} \leq c_\alpha 2^{k\alpha} \|\phi\|_{\mathbf{M}(B_{p,q}^\alpha(\mathbb{T}_\theta^d))} \|x\|_p.$$

On the other hand, by (3.2), $\tilde{\varphi}_k * M_\phi(y) = M_{\phi\varphi^{(k)}}(x)$ and

$$\|M_\phi(y)\|_{B_{p,q}^\alpha} \geq 2^{k\alpha} \|\tilde{\varphi}_k * M_\phi(y)\|_p = 2^{k\alpha} \|M_{\phi\varphi^{(k)}}(x)\|_p.$$

It then follows that

$$\|M_{\phi\varphi^{(k)}}(x)\|_p \leq c_\alpha \|\phi\|_{\mathbf{M}(B_{p,q}^\alpha(\mathbb{T}_\theta^d))} \|x\|_p,$$

whence

$$\sup_{k \geq 0} \|\phi\varphi^{(k)}\|_{\mathbf{M}(L_p(\mathbb{T}_\theta^d))} \leq c_\alpha \|\phi\|_{\mathbf{M}(B_{p,q}^\alpha(\mathbb{T}_\theta^d))}.$$

Conversely, for $x \in B_{p,q}^\alpha(\mathbb{T}_\theta^d)$,

$$\begin{aligned} \|\tilde{\varphi}_k * M_\phi(x)\|_p &= \|M_{\phi\varphi^{(k)}}((\tilde{\varphi}_{k-1} + \tilde{\varphi}_k + \tilde{\varphi}_{k+1}) * x)\|_p \\ &\leq \|\phi\varphi^{(k)}\|_{\mathbf{M}(L_p(\mathbb{T}_\theta^d))} \|(\tilde{\varphi}_{k-1} + \tilde{\varphi}_k + \tilde{\varphi}_{k+1}) * x\|_p. \end{aligned}$$

We then deduce

$$\|M_\phi(x)\|_{B_{p,q}^\alpha} \leq 3 \cdot 2^{|\alpha|} \sup_{k \geq 0} \|\phi\varphi^{(k)}\|_{\mathbf{M}(L_p(\mathbb{T}_\theta^d))} \|x\|_{B_{p,q}^\alpha},$$

which implies

$$\|\phi\|_{\mathbf{M}(B_{p,q}^\alpha(\mathbb{T}_\theta^d))} \leq 3 \cdot 2^{|\alpha|} \sup_{k \geq 0} \|\phi\varphi^{(k)}\|_{\mathbf{M}(L_p(\mathbb{T}_\theta^d))}.$$

Thus the assertion concerning bounded multipliers is proved.

The preceding argument can be modified to work in the c.b. case too. First note that for $k \geq 0$, $\varphi^{(k)}$ is a c.b. Fourier multiplier on $L_p(\mathbb{T}_\theta^d)$ for all $1 \leq p \leq \infty$ with c.b. norm 1, that is, the map $x \mapsto \tilde{\varphi}_k * x$ is c.b. on $L_p(\mathbb{T}_\theta^d)$. So for any $x \in S_q[L_p(\mathbb{T}_\theta^d)]$ (the $L_p(\mathbb{T}_\theta^d)$ -valued Schatten q -class),

$$\|(\text{Id}_{S_q} \otimes M_{\varphi^{(k)}})(x)\|_{S_q[L_p(\mathbb{T}_\theta^d)]} \leq \|x\|_{S_q[L_p(\mathbb{T}_\theta^d)]}.$$

Now let $\phi \in \mathbf{M}_{\text{cb}}(B_{p,q}^\alpha(\mathbb{T}_\theta^d))$ and $x \in S_q[L_p(\mathbb{T}_\theta^d)]$. Define y as above: $y = (\tilde{\varphi}_{k-1} + \tilde{\varphi}_k + \tilde{\varphi}_{k+1}) * x$. Then for $k-2 \leq j \leq k+2$,

$$\|\tilde{\varphi}_j * y\|_{S_q[L_p(\mathbb{T}_\theta^d)]} \leq 3 \|x\|_{S_q[L_p(\mathbb{T}_\theta^d)]}.$$

It thus follows that

$$\begin{aligned} \|(\text{Id}_{S_q} \otimes M_\phi)(y)\|_{S_q[B_{p,q}^\alpha(\mathbb{T}_\theta^d)]} &\leq \|\phi\|_{\mathbf{M}_{\text{cb}}(B_{p,q}^\alpha(\mathbb{T}_\theta^d))} \|y\|_{S_q[B_{p,q}^\alpha(\mathbb{T}_\theta^d)]} \\ &\leq \|\phi\|_{\mathbf{M}_{\text{cb}}(B_{p,q}^\alpha(\mathbb{T}_\theta^d))} \sum_{j=k-2}^{k+2} 2^{j\alpha} \|\tilde{\varphi}_j * y\|_{S_q[L_p(\mathbb{T}_\theta^d)]} \\ &\leq c_\alpha 2^{k\alpha} \|\phi\|_{\mathbf{M}_{\text{cb}}(B_{p,q}^\alpha(\mathbb{T}_\theta^d))} \|x\|_{S_q[L_p(\mathbb{T}_\theta^d)]}. \end{aligned}$$

Then as before, we deduce

$$\sup_{k \geq 0} \|\phi \varphi^{(k)}\|_{\mathbf{M}_{\text{cb}}(L_p(\mathbb{T}_\theta^d))} \leq c_\alpha \|\phi\|_{\mathbf{M}_{\text{cb}}(B_{p,q}^\alpha(\mathbb{T}_\theta^d))}.$$

To show the converse inequality, assume

$$\sup_{k \geq 0} \|\phi \varphi^{(k)}\|_{\mathbf{M}_{\text{cb}}(L_p(\mathbb{T}_\theta^d))} \leq 1.$$

Then for $x \in S_q[B_{p,q}^\alpha(\mathbb{T}_\theta^d)]$,

$$\begin{aligned} \|\tilde{\varphi}_k * M_\phi(x)\|_{S_q[L_p(\mathbb{T}_\theta^d)]} &\leq \|\phi \varphi^{(k)}\|_{\mathbf{M}_{\text{cb}}(L_p(\mathbb{T}_\theta^d))} \|(\tilde{\varphi}_{k-1} + \tilde{\varphi}_k + \tilde{\varphi}_{k+1}) * x\|_{S_q[L_p(\mathbb{T}_\theta^d)]} \\ &\leq \|(\tilde{\varphi}_{k-1} + \tilde{\varphi}_k + \tilde{\varphi}_{k+1}) * x\|_{S_q[L_p(\mathbb{T}_\theta^d)]}. \end{aligned}$$

Therefore,

$$\begin{aligned} \|M_\phi(x)\|_{S_q[B_{p,q}^\alpha(\mathbb{T}_\theta^d)]} &= \left(\sum_{k \geq 0} (2^{k\alpha} \|\tilde{\varphi}_k * M_\phi(x)\|_{S_q[L_p(\mathbb{T}_\theta^d)]})^q \right)^{\frac{1}{q}} \\ &\leq \left(\sum_{k \geq 0} (2^{k\alpha} \|(\tilde{\varphi}_{k-1} + \tilde{\varphi}_k + \tilde{\varphi}_{k+1}) * x\|_{S_q[L_p(\mathbb{T}_\theta^d)]})^q \right)^{\frac{1}{q}} \\ &\leq 3 \cdot 2^{|\alpha|} \|x\|_{S_q[B_{p,q}^\alpha(\mathbb{T}_\theta^d)]}. \end{aligned}$$

We thus get the missing converse inequality, so the theorem is proved. $\square$

The following is an immediate consequence of the preceding theorem.

Corollary 7.5. (i) $\mathbf{M}(B_{p,q}^\alpha(\mathbb{T}_\theta^d))$ is independent of α and q , up to equivalent norms.

(ii) $\mathbf{M}(B_{p,\infty}^0(\mathbb{T}_\theta^d)) = \mathbf{M}(B_{p',\infty}^0(\mathbb{T}_\theta^d))$, where p' is the conjugate index of p .

(iii) $\mathbf{M}(B_{p_0,\infty}^0(\mathbb{T}_\theta^d)) \subset \mathbf{M}(B_{p_1,\infty}^0(\mathbb{T}_\theta^d))$ for $1 \leq p_0 < p_1 \leq 2$.

(iv) $M(L_p(\mathbb{T}_\theta^d)) \subset M(B_{p,q}^\alpha(\mathbb{T}_\theta^d))$.

Similar statements hold for the spaces $M_{cb}(B_{p,q}^\alpha(\mathbb{T}_\theta^d))$.

Theorem 7.4 and Lemma 1.3 imply the following:

Corollary 7.6. $M_{cb}(B_{p,q}^\alpha(\mathbb{T}_\theta^d)) = M_{cb}(B_{p,q}^\alpha(\mathbb{T}^d))$ with equivalent norms.

Let $\mathcal{F}(B_{1,\infty}^0(\mathbb{T}^d))$ be the space of all Fourier transforms of functions in $B_{1,\infty}^0(\mathbb{T}^d)$ (a commutative Besov space), equipped with the norm $\|\hat{f}\| = \|f\|_{B_{1,\infty}^0}$.

Corollary 7.7. $M_{cb}(B_{1,q}^\alpha(\mathbb{T}_\theta^d)) = \mathcal{F}(B_{1,\infty}^0(\mathbb{T}^d))$ with equivalent norms.

Proof. Let $\phi \in M_{cb}(B_{1,\infty}^0(\mathbb{T}_\theta^d))$ and f be the distribution on $\mathbb{T}^d$ such that $\hat{f} = \phi$. By Theorem 7.4 and Lemma 1.3, we have

$$\sup_{k \geq 0} \|\phi \varphi^{(k)}\|_{M(L_1(\mathbb{T}^d))} < \infty.$$

Recall that the Fourier transform of φ_k is $\varphi^{(k)}$ and $\tilde{\varphi}_k$ is the periodization of φ_k . So

$$\|\tilde{\varphi}_k\|_{L_1(\mathbb{T}^d)} = \|\varphi_k\|_{L_1(\mathbb{R}^d)} = \|\varphi\|_{L_1(\mathbb{R}^d)}.$$

Noting that by (3.2), $\tilde{\varphi}_k * f = M_{\phi \varphi^{(k)}}(\tilde{\varphi}_{k-1} + \tilde{\varphi}_k + \tilde{\varphi}_{k+1})$, we get

$$\|\tilde{\varphi}_k * f\|_1 \leq \|\phi \varphi^{(k)}\|_{M(L_1(\mathbb{T}^d))} \|\tilde{\varphi}_{k-1} + \tilde{\varphi}_k + \tilde{\varphi}_{k+1}\|_1 \leq 3 \|\varphi\|_{L_1(\mathbb{R}^d)} \|\phi \varphi^{(k)}\|_{M(L_1(\mathbb{T}^d))},$$

whence

$$\|f\|_{B_{1,\infty}^0} \leq 3 \|\varphi\|_{L_1(\mathbb{R}^d)} \sup_{k \geq 0} \|\phi \varphi^{(k)}\|_{M(L_1(\mathbb{T}^d))}.$$

Conversely, assume $\phi = \hat{f}$ with $f \in B_{1,\infty}^0(\mathbb{T}^d)$. Let $g \in B_{1,\infty}^0(\mathbb{T}^d)$. Then

$$\begin{aligned} \|\tilde{\varphi}_k * M_\phi(g)\|_1 &= \|\tilde{\varphi}_k * f * g\|_1 = \|\tilde{\varphi}_k * f * (\tilde{\varphi}_{k-1} + \tilde{\varphi}_k + \tilde{\varphi}_{k+1}) * g\|_1 \\ &\leq \|\tilde{\varphi}_k * f\|_1 \|(\tilde{\varphi}_{k-1} + \tilde{\varphi}_k + \tilde{\varphi}_{k+1}) * g\|_1 \\ &\leq 3 \|f\|_{B_{1,\infty}^0} \|g\|_{B_{1,\infty}^0}. \end{aligned}$$

Thus $M_\phi(g) \in B_{1,\infty}^0(\mathbb{T}^d)$ and

$$\|M_\phi(g)\|_{B_{1,\infty}^0} \leq 3 \|f\|_{B_{1,\infty}^0} \|g\|_{B_{1,\infty}^0},$$

which implies that ϕ is a Fourier multiplier on $B_{1,\infty}^0(\mathbb{T}_\theta^d)$ and

$$\|\phi\|_{M(B_{1,\infty}^0(\mathbb{T}_\theta^d))} \leq 3 \|f\|_{B_{1,\infty}^0}.$$

Considering g with values in S_∞ , we show that ϕ is c.b. too. Alternately, since $M(L_1(\mathbb{T}^d)) = M_{cb}(L_1(\mathbb{T}^d))$, Theorem 7.4 yields $M(B_{1,\infty}^0(\mathbb{T}^d)) = M_{cb}(B_{1,\infty}^0(\mathbb{T}^d))$, which allows us to conclude the proof too. $\square$

We have seen previously that every bounded (c.b.) Fourier multiplier on $L_p(\mathbb{T}_\theta^d)$ is a bounded (c.b.) Fourier multiplier on $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$. Corollary 7.7 shows that the converse is false for $p = 1$. We now show that it also is false for any $p \neq 2$.

Proposition 7.8. *There exists a Fourier multiplier ϕ which is c.b. on $B_{p,q}^\alpha(\mathbb{T}_\theta^d)$ for any p, q and α but never belongs to $M(L_p(\mathbb{T}_\theta^d))$ for any $p \neq 2$ and any θ .*

Proof. The example constructed by Stein and Zygmund [70] for a similar circumstance can be shown to work for our setting too. Their example is a distribution on $\mathbb{T}$ defined as follows:

$$\mu(z) = \sum_{n=2}^{\infty} \frac{1}{\log n} (w_n z)^{2^n} D_n(w_n z)$$

for some appropriate $w_n \in \mathbb{T}$, where

$$D_n(z) = \sum_{j=0}^n z^j, \quad z \in \mathbb{T}.$$

Since $\|D_n\|_{L_1(\mathbb{T})} \approx \log n$, we see that $\mu \in B_{1,\infty}^0(\mathbb{T})$. Considered as a distribution on $\mathbb{T}^d$, $\mu \in B_{1,\infty}^0(\mathbb{T}^d)$ too. Thus by Corollaries 7.5 and 7.7, $\phi = \hat{\mu}$ belongs to $M_{cb}(B_{p,p}^\alpha(\mathbb{T}_\theta^d))$ for any p, q and α . However, Stein and Zygmund proved that ϕ is not a Fourier multiplier on $L_p(\mathbb{T})$ for any $p \neq 2$ if the w_n 's are appropriately chosen. Consequently, ϕ cannot be a Fourier multiplier on $L_p(\mathbb{T}_\theta^d)$ for any $p \neq 2$ and any θ since $L_p(\mathbb{T})$ isometrically embeds into $L_p(\mathbb{T}_\theta^d)$ by an embedding that is also a c.b. Fourier multiplier. $\square$

We conclude this section with some comments on the vector-valued case. The proof of Theorem 7.4 works equally for vector-valued Besov spaces. Recall that for an operator space E , $B_{p,q}^\alpha(\mathbb{T}_\theta^d; E)$ denotes the E -valued Besov space on $\mathbb{T}_\theta^d$ (see Remark 3.24).

Proposition 7.9. *For any operator space E ,*

$$\|\phi\|_{M(B_{p,q}^\alpha(\mathbb{T}_\theta^d; E))} \approx |\phi(0)| + \sup_{k \geq 0} \|\phi\varphi^{(k)}\|_{M(L_p(\mathbb{T}_\theta^d; E))}$$

with equivalence constants depending only on α .

If $\theta = 0$, we go back to the classical vector-valued case. The above proposition explains the well-known fact mentioned at the beginning of this section that Fourier multipliers behave better on Besov spaces than on L_p -spaces. To see this, it is more convenient to write the above right-hand side in a different form:

$$\|\phi\varphi^{(k)}\|_{M(L_p(\mathbb{T}_\theta^d; E))} = \|\phi(2^k \cdot)\varphi\|_{M(L_p(\mathbb{T}_\theta^d; E))}.$$

Thus if ϕ is homogeneous, the above multiplier norm is independent of k , so ϕ is a Fourier multiplier on $B_{p,q}^\alpha(\mathbb{T}^d; E)$ for any p, q, α and any Banach space E . In particular, the Riesz transform is bounded on $B_{p,q}^\alpha(\mathbb{T}^d; E)$.

The preceding characterization of Fourier multipliers is, of course, valid for $\mathbb{R}^d$ in place of $\mathbb{T}^d$. Let us record this here:

Proposition 7.10. *Let E be a Banach space. Then for any $\phi : \mathbb{R}^d \rightarrow \mathbb{C}$,*

$$\|\phi\|_{M(B_{p,q}^\alpha(\mathbb{R}^d; E))} \approx \|\phi\psi\|_{M(L_p(\mathbb{R}^d; E))} + \sup_{k \geq 0} \|\phi(2^k \cdot)\varphi\|_{M(L_p(\mathbb{R}^d; E))},$$

where ψ is defined by

$$\psi(\xi) = \sum_{k \geq 1} \varphi(2^k \xi), \quad \xi \in \mathbb{R}^d.$$

7.3 Fourier multipliers on Triebel-Lizorkin spaces

As we have seen in the chapter on Triebel-Lizorkin spaces, Fourier multipliers on such spaces are subtler than those on Sobolev and Besov spaces. Similarly to the previous two sections, our target here is to show that the c.b. Fourier multipliers on $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ are independent of θ . By definition, $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ can be viewed as a subspace of the column space $L_p(\mathbb{T}_\theta^d; \ell_2^{\alpha,c})$, the latter is the column subspace of $L_p(B(\ell_2^\alpha) \bar{\otimes} \mathbb{T}_\theta^d)$. Thus $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ inherits the natural operator space structure of $L_p(B(\ell_2^\alpha) \bar{\otimes} \mathbb{T}_\theta^d)$. Similarly, the row Triebel-Lizorkin space $F_p^{\alpha,r}(\mathbb{T}_\theta^d)$ carries a natural operator space structure too. Finally, the mixture space $F_p^\alpha(\mathbb{T}_\theta^d)$ is equipped with the sum or intersection operator space structure according to $p < 2$ or $p \geq 2$. Note that according to this definition, even though it is a commutative function space, the space $F_p^\alpha(\mathbb{T}^d)$ (corresponding to $\theta = 0$) is endowed with three different operator space structures, the first two being defined by its embedding into $L_p(\mathbb{T}^d; \ell_2^{\alpha,c})$ and $L_p(\mathbb{T}^d; \ell_2^{\alpha,r})$, the third one being the mixture of these two. The resulting operator spaces are denoted by $F_p^{\alpha,c}(\mathbb{T}^d)$, $F_p^{\alpha,r}(\mathbb{T}^d)$ and $F_p^\alpha(\mathbb{T}^d)$, respectively. Similarly, we introduce operator space structures on the Hardy space $\mathcal{H}_p^c(\mathbb{T}_\theta^d)$, its row and mixture versions too.

The main result of this section is the following:

Theorem 7.11. *Let $1 \leq p \leq \infty$ and $\alpha \in \mathbb{R}$. Then*

$$\begin{aligned} \mathbf{M}_{\text{cb}}(F_p^{\alpha,c}(\mathbb{T}_\theta^d)) &= \mathbf{M}_{\text{cb}}(F_p^{\alpha,c}(\mathbb{T}^d)) \quad \text{with equal norms,} \\ \mathbf{M}_{\text{cb}}(F_p^{\alpha,c}(\mathbb{T}_\theta^d)) &= \mathbf{M}_{\text{cb}}(F_p^{0,c}(\mathbb{T}^d)) \quad \text{with equivalent norms.} \end{aligned}$$

Similar statements hold for the row and mixture spaces.

We will show the theorem only in the case $p < \infty$. The proof presented below can be easily modified to work for $p = \infty$ too. Alternately, the case $p = \infty$ can be also obtained by duality from the case $p = 1$. Note, however, that this duality argument yields only the first equality of the theorem with equivalent norms for $p = \infty$.

We adapt the proof of Theorem 7.1 to the present situation, by introducing the space

$$S_p^{\alpha,c} = \left\{ a = (a_{m,n})_{m,n \in \mathbb{Z}^d} : \left(\sum_{k \geq 0} 2^{2k\alpha} |\tilde{\varphi}_k * a|^2 \right)^{\frac{1}{2}} \in S_p(\ell_2(\mathbb{Z}^d)) \right\}$$

and endow it with the norm

$$\|a\|_{S_p^{\alpha,c}} = \left\| \left(\sum_{k \geq 0} 2^{2k\alpha} |\tilde{\varphi}_k * a|^2 \right)^{\frac{1}{2}} \right\|_{S_p},$$

where

$$\tilde{\varphi}_k * a = (\varphi(2^{-k}(m-n)) a_{m,n})_{m,n \in \mathbb{Z}^d}.$$

Then $S_p^{\alpha,c}$ is a closed subspace of the column subspace of $S_p(\ell_2^\alpha \otimes \ell_2(\mathbb{Z}^d))$, which introduces a natural operator space structure on $S_p^{\alpha,c}$. Let $\mathbf{M}_{\text{cb}}(S_p^{\alpha,c})$ denote the space of all c.b. Schur multipliers on $S_p^{\alpha,c}$, equipped with the natural norm.

Lemma 7.12. *Let $1 \leq p < \infty$ and $\alpha \in \mathbb{R}$. Then*

$$\mathbf{M}_{\text{cb}}(F_p^{\alpha,c}(\mathbb{T}_\theta^d)) = \mathbf{M}_{\text{cb}}(S_p^{\alpha,c}) \quad \text{with equal norms.}$$

Proof. This proof is similar to the one of Theorem 7.1; we point out the necessary changes. Keeping the notation there, we have for $a = (a_{m,n})_{m,n \in \mathbb{Z}^d} \in S_p^{\alpha,c}$ and $x = V(a \otimes 1_{\mathbb{T}_\theta^d})V^* \in B(\ell_2(\mathbb{Z}^d)) \overline{\otimes} \mathbb{T}_\theta^d$

$$\left\| \left(\sum_k 2^{2k\alpha} |\tilde{\varphi}_k * x|^2 \right)^{\frac{1}{2}} \right\|_{L_p(B(\ell_2(\mathbb{Z}^d)) \overline{\otimes} \mathbb{T}_\theta^d)} = \|a\|_{S_p^{\alpha,c}}.$$

Suppose that $\phi \in \mathbf{M}_{\text{cb}}(F_p^{\alpha,c}(\mathbb{T}_\theta^d))$. It then follows that

$$\begin{aligned} \|S_{\dot{\phi}}(a)\|_{S_p^{\alpha,c}} &= \left\| \left(\sum_k 2^{2k\alpha} |\tilde{\varphi}_k * (V(S_{\dot{\phi}}(a) \otimes 1_{\mathbb{T}_\theta^d})V^*)|^2 \right)^{\frac{1}{2}} \right\|_{L_p(B(\ell_2(\mathbb{Z}^d)) \overline{\otimes} \mathbb{T}_\theta^d)} \\ &= \left\| \left(\sum_k 2^{2k\alpha} |M_\phi(\tilde{\varphi}_k * (V(a \otimes 1_{\mathbb{T}_\theta^d})V^*))|^2 \right)^{\frac{1}{2}} \right\|_{L_p(B(\ell_2(\mathbb{Z}^d)) \overline{\otimes} \mathbb{T}_\theta^d)} \\ &\leq \|\phi\|_{\mathbf{M}_{\text{cb}}(F_p^{\alpha,c}(\mathbb{T}_\theta^d))} \|x\|_{S_p[F_p^{\alpha,c}(\mathbb{T}_\theta^d)]} \\ &= \|\phi\|_{\mathbf{M}_{\text{cb}}(F_p^{\alpha,c}(\mathbb{T}_\theta^d))} \|a\|_{S_p^{\alpha,c}}. \end{aligned}$$

Therefore, $\dot{\phi}$ is a bounded Schur multiplier on $S_p^{\alpha,c}$. Considering matrices $a = (a_{m,n})_{m,n \in \mathbb{Z}^d}$ with entries in $B(\ell_2)$, we show in the same way that $S_{\dot{\phi}}$ is c.b. on $S_p^{\alpha,c}$, so $\dot{\phi}$ is a c.b. Schur multiplier on $S_p^{\alpha,c}$ and

$$\|\dot{\phi}\|_{\mathbf{M}_{\text{cb}}(S_p^{\alpha,c})} \leq \|\phi\|_{\mathbf{M}_{\text{cb}}(F_p^{\alpha,c}(\mathbb{T}_\theta^d))}.$$

To show the opposite inequality, we just note that the contractive and convergence properties of the maps A_N and B_N introduced in the proof of Theorem 7.1 also hold on the corresponding $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ or $S_p^{\alpha,c}$ spaces. To see this, we take A_N for example. Since it is c.b. between the corresponding L_p -spaces, it is also c.b. from $L_p(B(\ell_2) \overline{\otimes} \mathbb{T}_\theta^d)$ to $L_p(B(\ell_2) \overline{\otimes} B(\ell_2^{\mathbb{Z}^N}))$. Applying this to the elements of the form

$$\begin{pmatrix} \tilde{\varphi}_0 * x & 0 & 0 & \dots \\ 2^\alpha \tilde{\varphi}_1 * x & 0 & 0 & \dots \\ 2^{2\alpha} \tilde{\varphi}_2 * x & 0 & 0 & \dots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix}$$

we see that A_N is completely contractive from $F_p^{\alpha,c}(\mathbb{T}_\theta^d)$ to $S_p^{\alpha,c}(B(\ell_2^{\mathbb{Z}^N}))$, the latter space being the finite dimensional analogue of $S_p^{\alpha,c}$. We then argue as in the proof of Theorem 7.1 to deduce the desired opposite inequality. $\square$

Proof of Theorem 7.11. The first part is an immediate consequence of the previous lemma. For the second, we need the c.b. version of Theorem 4.11 (i), whose proof is already contained in section 4.1. To see this, we just note that, letting $\mathcal{M} = B(\ell_2(\mathbb{Z}^d)) \overline{\otimes} \mathbb{T}_\theta^d$ and $\mathcal{N} = B(\ell_2(\mathbb{Z}^d)) \overline{\otimes} L_\infty(\mathbb{T}^d) \overline{\otimes} \mathbb{T}_\theta^d$ in Lemma 4.6 we obtain the c.b. version of Lemma 4.7, and that, in the same way, the c.b. version of Lemma 1.10 holds, i.e., for $x \in S_p[\mathcal{H}_p^c(\mathbb{T}_\theta^d)]$,

$$\|x\|_{S_p[\mathcal{H}_p^c(\mathbb{T}_\theta^d)]} \approx \|\hat{x}(0)\|_{S_p} + \|s_\psi^c(x)\|_{S_p[L_p(\mathbb{T}_\theta^d)]}.$$

Finally, the previous lemma and the c.b. version of Theorem 4.11 (i) yield the desired conclusion. $\square$

Remark 7.13. The preceding theorem and the c.b. version of Theorem 4.11 (i) show that

$$\mathbf{M}_{\text{cb}}(\mathcal{H}_p^c(\mathbb{T}_\theta^d)) = \mathbf{M}_{\text{cb}}(\mathcal{H}_p^c(\mathbb{T}^d)) \text{ with equivalent norms.}$$

In fact, using arguments similar to the proof of the preceding theorem, we can show that the above equality holds with equal norms.

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